

REGULARITY FOR EIGENVALUE-TYPE EQUATIONS IN SETS WITH THICK COMPLEMENT

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ABSTRACT. We prove global higher integrability of the gradient for solutions to nonlinear PDEs with homogeneous Dirichlet condition. We work with general open sets (i.e. not necessarily bounded), under a minimal thickness assumption of the complement. The main focus is on improving the global summability of the gradient, by preserving the “zero trace” Sobolev class. We pay due attention to providing precise a priori estimates. As an application, we get universal global Hölder estimates for solutions of the super-homogeneous Lane-Emden equation for the N -Laplacian, in contractible open sets with finite inradius.

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1. INTRODUCTION

1.1. Initial motivation: an old-fashioned result. Let us consider an open set $\Omega \subsetneq \mathbb{R}^N$, for $0 < \alpha \leq 1$ we recall the definition of the following space

$$C_0^{0,\alpha}(\overline{\Omega}) = \left\{ u \in C_{\text{bound}}(\overline{\Omega}) : [u]_{C^{0,\alpha}(\overline{\Omega})} < +\infty, u = 0 \text{ on } \partial\Omega \right\},$$

of Hölder continuous functions, vanishing at the boundary. Here $C_{\text{bound}}(\overline{\Omega})$ denotes the Banach space of continuous and bounded functions on $\overline{\Omega}$, while as usual

$$[u]_{C^{0,\alpha}(\overline{\Omega})} = \sup_{x \neq y \in \overline{\Omega}} \frac{|u(x) - u(y)|}{|x - y|^\alpha}.$$

We also recall the definition of *inradius* of Ω , given by

$$r_\Omega = \sup \left\{ r > 0 : \exists B_r(x_0) \subseteq \Omega \right\}.$$

The original intent of the present paper was that of proving the following *global* regularity result for eigenfunctions of the Dirichlet-Laplacian in the two-dimensional case. We highlight the two main points of interest: the minimality of the assumptions and the universality of the estimate.

Theorem. *There exists a universal exponent $\delta > 0$ such that for every simply connected open set $\Omega \subseteq \mathbb{R}^2$ with $r_\Omega < +\infty$ and every eigenfunction $u \in W_0^{1,2}(\Omega) \setminus \{0\}$ of the Dirichlet-Laplacian with eigenvalue λ , we have*

$$u \in W_0^{1,2(1+\delta)}(\Omega) \quad \text{and} \quad u \in C_0^{0, \frac{\delta}{1+\delta}}(\overline{\Omega}).$$

Moreover, it holds

$$\|\nabla u\|_{L^{2(1+\delta)}(\Omega)} \leq C \left(\sqrt{\lambda} \right)^{\frac{1+2\delta}{1+\delta}} \|u\|_{L^2(\Omega)},$$

and

$$|u|_{C_0^{0, \frac{\delta}{1+\delta}}(\overline{\Omega})} \leq C \left(\sqrt{\lambda} \right)^{\frac{1+2\delta}{1+\delta}} \|u\|_{L^2(\Omega)},$$

for a universal constant $C > 0$.

We make a couple of comments about the optimality of the previous result.

Remark 1.1 (A two-dimensional example). Under the above assumptions, eigenfunctions can not be more than globally Hölder continuous. Indeed, take for example the *slit disk*

$$\Omega_0 = \left\{ (x_1, x_2) \in \mathbb{R}^2 : x_1^2 + x_2^2 < 1 \right\} \setminus \left([0, 1] \times \{0\} \right).$$

It is not difficult to see that the function given by (in polar coordinates)

$$v(\varrho, \vartheta) = \frac{\sin(\pi \varrho)}{\sqrt{\pi \varrho}} \sin\left(\frac{\vartheta}{2}\right), \quad \text{for } 0 < \varrho < 1, 0 < \vartheta < 2\pi,$$

is an eigenfunction (actually, the first one) of the Dirichlet-Laplacian on Ω_0 , see for example [38, Exercise 6.8.9]. Observe that for this function we have

$$v \in C_0^{0, \frac{1}{2}}(\overline{\Omega_0}) \quad \text{and} \quad \nabla v \in L^q(\Omega_0), \text{ for every } 2 \leq q < 4.$$

This example also shows that, without additional assumptions, we must have $\delta \leq 1$ in the Theorem above.

Remark 1.2 (Higher dimensions). No direct analogue of the above theorem holds in higher dimensions, for the Laplacian. Indeed, for $N \geq 3$ no topological assumptions on Ω can guarantee a global improvement of regularity. For example, consider the following open set (which is contractible and even starshaped)

$$\Omega_1 = B_1(0) \setminus \Sigma \quad \text{where } \Sigma = \{(t, 0, \dots, 0) : t \in [0, 1]\},$$

and use that a segment is *removable* for the Laplacian in dimensions $N \geq 3$. Thus, every eigenfunction $u \in W_0^{1,2}(\Omega_1)$ of the Dirichlet-Laplacian on Ω_1 “does not see” the boundary part Σ . In particular, it is not true that

$$u \in C_0^{0,\alpha}(\overline{\Omega_1}),$$

i.e. u does not vanish on Σ .

As we will see in Corollary 6.4, the Theorem above will be a consequence of a much more general statement, obtained by suitably adapting the proof of [26, Theorem 1.1], by Kilpelainen and Koskela. This general statement will be valid for solutions of the quasilinear elliptic equations

$$(1.1) \quad -\operatorname{div} \nabla G(\nabla u) = f, \quad \text{in } \Omega,$$

under suitable assumption on f , G and the open set Ω . We insist from the very beginning that our open sets are not necessarily bounded. Moreover, we will try to be as much *quantitative* as possible, by focusing on providing precise a priori estimates and discussing explicit classes of open sets to which our result could be applied.

1.2. Description of the main result. We start by introducing the class of open sets we wish to consider. They satisfy a sort of “uniform density condition of the complement”, whose relevance in boundary regularity has been first singled out by [26]. Nowadays, this is quite common, see for example [2, 3, 17, 27] and [34], among others. Here “density” has to be intended in terms of *capacity* (we refer to Section 2 for the notation and definition).

Definition 1.3. Let $1 \leq p < \infty$ and let $\Omega \subsetneq \mathbb{R}^N$ be an open set. If there exists $c_0 > 0$ and $r_0 > 0$ such that

$$\operatorname{cap}_p(\overline{Q_r(x_0)} \setminus \Omega; Q_{2r}(x_0)) \geq c_0 \operatorname{cap}_p(\overline{Q_r(x_0)}; Q_{2r}(x_0)), \quad \text{for every } x_0 \in \mathbb{R}^N \setminus \Omega, \ 0 < r \leq r_0,$$

then we say that $\mathbb{R}^N \setminus \Omega$ is *uniformly p -thick with constants (c_0, r_0)* . We denote the collection of all open sets with this property by

$$\operatorname{UT}_p(c_0, r_0).$$

If the previous estimate holds for every $r > 0$, we will simply say that $\mathbb{R}^N \setminus \Omega$ is *uniformly p -thick with constant c_0* . Accordingly, we will use the symbol

$$\operatorname{UT}_p(c_0),$$

for the collection of all open sets having this property.

In what follows, for $1 < p < \infty$ we will use the following notation

$$(1.2) \quad \lambda_p(\Omega) = \inf_{\varphi \in C_0^\infty(\Omega)} \left\{ \int_\Omega |\nabla \varphi|^p dx : \|\varphi\|_{L^p(\Omega)} = 1 \right\}.$$

Observe that we have

$$\lambda_p(\Omega) > 0 \quad \Longleftrightarrow \quad \mathcal{D}_0^{1,p}(\Omega) \hookrightarrow L^p(\Omega),$$

and it holds by definition

$$\lambda_p(\Omega) \int_{\Omega} |\varphi|^p dx \leq \int_{\Omega} |\nabla \varphi|^p dx, \quad \text{for every } \varphi \in C_0^\infty(\Omega).$$

The symbol $\mathcal{D}_0^{1,p}(\Omega)$ stands for the homogeneous Sobolev space, given by the completion of $C_0^\infty(\Omega)$ with respect to the norm

$$\varphi \mapsto \|\nabla \varphi\|_{L^p(\Omega)}.$$

As already anticipated, our main result is concerned with weak solutions to the equation (1.1). We will work under the following list of assumptions on the open set $\Omega \subseteq \mathbb{R}^N$, the right-hand side f and the function G .

Assumptions 1. We will assume that:

(E_f) $f \in L^{\gamma_0}(\Omega)$, with $\gamma_0 \leq p'$ such that

$$\begin{cases} (p^*)' \leq \gamma_0 & \text{if } 1 < p < N, \\ 1 < \gamma_0, & \text{if } p = N, \\ 1 \leq \gamma_0, & \text{if } p > N; \end{cases} \quad \text{where } p^* = \frac{Np}{N-p};$$

(E_Ω) $\Omega \subseteq \mathbb{R}^N$ is an open set such that $\lambda_p(\Omega) > 0$;

(E_G) $G : \mathbb{R}^N \rightarrow \mathbb{R}$ is a C^1 strictly convex function, such that

$$g_1 |z|^p \leq G(z) \leq g_2 |z|^p,$$

for two constants $g_2 \geq g_1 > 0$.

Assumptions 2. In order to improve the regularity of solutions, we will take the following further set of assumptions:

(R_f) $f \in L^\gamma(\Omega)$ for

$$\gamma_0 < \gamma \leq +\infty,$$

and γ_0 as in Assumptions 1;

(R_Ω) $\Omega \in \text{UT}_{p_0}(c_0)$, for some $1 \leq p_0 < p$ and $c_0 > 0$.

We are now ready to state the main result of this paper. We will denote by $W_0^{1,p}(\Omega)$ the closure of $C_0^\infty(\Omega)$ in the usual Sobolev space $W^{1,p}(\Omega)$.

Main Theorem. *Under the Assumptions 1 and 2, let $u \in W_0^{1,p}(\Omega)$ be the weak solution to the equation (1.1). Then, there exists $\delta = \delta(N, p, p_0, \gamma_0, \gamma, c_0, g_1, g_2) > 0$ with $\delta \leq (\gamma - \gamma_0)/\gamma_0$, such that*

$$u \in W_0^{1,p(1+\delta)}(\Omega).$$

Moreover, the following scale invariant estimate holds, for $C_1 = C_1(N, p, p_0, \gamma_0, \gamma, c_0, g_1, g_2) > 0$

$$\|\nabla u\|_{L^{p(1+\delta)}(\Omega)}^p \leq C_1 \left[\int_{\Omega} |\nabla u|^p dx \right]^{1-\Theta} \left[(\|f\|_{L^{\gamma_0}(\Omega)})^{p' - \frac{\delta}{1+\delta} \frac{\gamma \gamma_0}{\gamma - \gamma_0}} (\|f\|_{L^\gamma(\Omega)})^{\frac{\delta}{1+\delta} \frac{\gamma \gamma_0}{\gamma - \gamma_0}} \right]^\Theta,$$

where $\Theta = \Theta(N, p, \gamma_0, \delta) \in (0, 1)$ is given by

$$\frac{1}{\Theta} = 1 + \frac{1+\delta}{\delta} p' \left(\frac{1}{\gamma_0'} - \frac{N-p}{Np} \right).$$

1.3. Comments on the assumptions. We can not expect the reader to immediately understand both the necessity and the subtleties hidden in the (long) list of assumptions we made. Of course, some of them are needed in order to consider unbounded open sets. In this subsection, we want to thoroughly comment these assumptions.

- *The operator:* differently from [26], our equation (1.1) has a variational nature. In other words, it is the Euler-Lagrange equation of the (strictly) convex functional

$$(1.3) \quad \mathfrak{F}(\varphi) = \int_{\Omega} G(\nabla \varphi) dx - \int_{\Omega} f \varphi dx, \quad \text{for every } \varphi \in W_0^{1,p}(\Omega).$$

Accordingly, the solution u is the minimizer of this functional. This fact is exploited in order to run a certain approximation argument (see Section 4), that we will explain in a moment. In doing this, the minimality of u will be useful.

- *The function G :* if G is not differentiable, the equation (1.1) would not make sense. However, in this case it would be possible to get the Main Theorem for the minimizer u of (1.3). Indeed, the desired Caccioppoli inequality, which is the cornerstone of the result, could be obtained in this case by using the classical trick of Giaquinta and Giusti (see [19]), circumventing the lack of an Euler-Lagrange equation. This is the case treated by Granlund in [22, Theorem 1.4], for example, under stronger assumptions on the open set. Here, we preferred to keep the C^1 assumption, in order to simplify a bit some of the estimates.
- *The assumptions (E_f) and (R_f) on the right-hand side f :* since we are not assuming that Ω has finite volume, the two integrability requirements $f \in L^{\gamma_0}(\Omega)$ and $f \in L^{\gamma}(\Omega)$ are not redundant. Both of them play a role: the former, in conjunction with (E_{Ω}) , is needed in order to assure *existence* of a solution; the latter is responsible for the *gain of integrability*. We also point out the requirement $\gamma_0 \leq p'$: this is crucially needed to get a local improvement of integrability, which can be extended to the whole Ω by using a covering with cubes having *fixed* radius. We refer to *Step 5* and *Step 6* of the proof of the Main Theorem.

If Ω has finite volume, this restriction is immaterial, since for $\gamma_0 > p'$ we would have $f \in L^{\gamma_0}(\Omega) \subseteq L^{p'}(\Omega)$. However, when Ω has infinite volume and $\gamma_0 > p'$, it seems that a finer knowledge on the “geometry at infinity” of the open set Ω would be needed. We plan to come back on this issue in the future.

- *The assumption (E_{Ω}) on the open set Ω :* we notice that the assumption $\lambda_p(\Omega) > 0$ guarantees the continuity of the *super-homogeneous* embeddings

$$\mathcal{D}_0^{1,p}(\Omega) \hookrightarrow L^{\gamma'_0}(\Omega), \quad \text{for every } p \leq \gamma'_0 \begin{cases} \leq p^* & \text{if } 1 < p < N, \\ < \infty, & \text{if } p = N, \\ \leq \infty, & \text{if } p > N, \end{cases}$$

and the identification $\mathcal{D}_0^{1,p}(\Omega) = W_0^{1,p}(\Omega)$ (see Proposition 2.1). These embeddings are perfectly consistent with the assumption (E_f) on f and they assure that we have

$$f \varphi \in L^1(\Omega), \quad \text{for every } \varphi \in W_0^{1,p}(\Omega).$$

We recall that the assumption $\lambda_p(\Omega) > 0$ comes for free if Ω has finite volume, as it is well-known. In general, necessary and sufficient conditions for this condition to hold are contained in [32, Chapter 15].

We point out that the case $1 < p < N$ and $\gamma_0 = (p^*)'$ is somehow exceptional: we could get a result similar to that of the Main Theorem by deleting (E_Ω) and only requiring the thickness condition (R_Ω) . In this case, one should work with the homogeneous Sobolev space $\mathcal{D}_0^{1,p}(\Omega)$, in place of $W_0^{1,p}(\Omega)$ (we briefly report this variant in Theorem B.1 in Appendix B, for completeness).

However, for $p \geq N$ this would cause some serious issues: the lack of a suitable surrogate for the Sobolev inequality makes the space $\mathcal{D}_0^{1,p}(\Omega)$ quite nasty to work with. In general, this is not even a functional space (see for example [10, 15] or [32, Chapter 15]). This explains our choice of taking the assumption (E_Ω) ;

- *The thickness condition (R_Ω)* : the attentive reader could notice that our assumption on the thickness of $\mathbb{R}^N \setminus \Omega$ is apparently stronger than that taken in [26]. Actually, this is not exactly true: in [26] it is assumed that $\mathbb{R}^N \setminus \Omega$ is uniformly p -thick, but then this condition is used to obtain that the same property holds for *some* $p_0 < p$, in light of a sophisticated result by Lewis (see [30, Theorem 1]). In the proof of [26], it is really the p_0 -thickness condition which is used. Accordingly, all the estimates in [26] depends on this exponent p_0 , which is known to exist but is quite implicit. For this reason, we decided to directly assume the condition on some exponent $p_0 < p$, in order to have a better control on the a priori estimate;
- *The thickness condition (R_Ω) ...reprise*: the same attentive reader could also notice that in [26] it is assumed that $\Omega \in \text{UT}_p(c_0, r_0)$, for some *finite* $r_0 > 0$. This requirement is apparently less restrictive than ours. Again, we show in Appendix A that this is not the case: the assumption (E_Ω) automatically entails that if $\mathbb{R}^N \setminus \Omega$ is uniformly p -thick for a range of radius $0 < r \leq r_0$, then it actually has the same property *for every* $r > 0$ (possibly by spoiling the constant c_0 , in a controlled way). This is a straightforward consequence of our Theorem A.2, which we believe is interesting in itself. Observe that the sets considered in [26] were *bounded*, thus the assumption (E_Ω) was automatically satisfied.

1.4. Comments on the proof. Our strategy of proof is exactly the same as in the paper by Kilpelainen and Koskela [26], in turn inspired by Granlund's paper [22]. The idea is to join the Caccioppoli and Sobolev-Poincaré inequalities, in order to obtain an “unnatural” reverse Hölder inequality on cubes, for the gradient of the solution. By iterating the estimate at every scale, a suitable variant of Gehring's Lemma implies the desired integrability gain on the gradient. This is super classical nowadays, as far as *local interior* estimates are concerned. But the main point here and in [22, 26] is to get the estimate up to the boundary. Indeed, the crucial point is the following one: for cubes intersecting $\mathbb{R}^N \setminus \Omega$, it is possible to use a capacity variant of the Sobolev-Poincaré inequality due to Maz'ya, exactly because of the thickness condition.

However, some care is needed. This procedure would only lead to the global integrability improvement

$$\nabla u \in L^{p(1+\delta)}(\Omega),$$

for some $\delta > 0$ depending on the data, as in [26]. In other words, this integrability improvement would not keep the¹ “zero trace” Sobolev class, i.e. one could hope for the (slightly) stronger result

$$u \in W_0^{1,p(1+\delta)}(\Omega).$$

¹The reader should not be fooled by our informal use of the word “trace”. Actually, in this paper we are not using any trace theory and our functions “vanish at the boundary” in the weak sense prescribed by the membership to the closure space $W_0^{1,p}$.

This slight improvement would be more natural, since it would be a true result “up to the boundary”, in a sense, as it really keeps track of the boundary datum. But more important, such an improvement would be particularly interesting in the borderline case $p = N$, since the Sobolev–Morrey–type embedding into globally Hölder continuous functions

$$W_0^{1,N(1+\delta)}(\Omega) \hookrightarrow C_0^{0, \frac{\delta}{1+\delta}}(\overline{\Omega}),$$

holds *for every open set*, on the scale of $W_0^{1,p}$ spaces, differently from the case of $W^{1,p}$.

We thus devote the rest of this subsection to explain how we get the higher integrability with “zero trace”:

- at first, one can observe that it is sufficient to find a sequence $\{u_n\}_{n \in \mathbb{N}} \subseteq W_0^{1,p(1+\delta)}(\Omega)$, for some δ independent on n , such that

$$(1.4) \quad u_n \rightharpoonup u, \quad \text{in } W^{1,p(1+\delta)}(\Omega), \quad \text{as } n \rightarrow \infty,$$

thanks to the fact that $W_0^{1,p(1+\delta)}(\Omega)$ is closed both in the strong and weak topologies (see for example [7, Theorem 3.7.4]);

- we construct the sequence $\{u_n\}_{n \in \mathbb{N}}$ by variational methods: for every $n \in \mathbb{N}$, we take the minimizer of the following “fake” double-phase functional with $1 < p < q$

$$\varphi \mapsto \int_{\Omega} G(\nabla \varphi) dx + \frac{\alpha_n}{q} \int_{\Omega} |\nabla \varphi|^q dx - \int_{\Omega} f \varphi dx,$$

where $q > p$ is fixed and $\{\alpha_n\}_{n \in \mathbb{N}}$ is a sequence of *positive* real numbers, converging to 0 as n goes to ∞ . This functional is well-defined on $W_0^{1,p}(\Omega) \cap W_0^{1,q}(\Omega)$. Thus, by a simple interpolation argument each u_n lies in particular in $W_0^{1,t}(\Omega)$, for every $p \leq t \leq q$. Moreover, it is not difficult to show that

$$u_n \rightharpoonup u, \quad \text{in } W^{1,p}(\Omega), \quad \text{as } n \rightarrow \infty;$$

- thanks to our hypothesis on Ω , we have that for every $p \leq t \leq q$

$$\varphi \mapsto \|\nabla \varphi\|_{L^t(\Omega)},$$

is an equivalent norm on $W_0^{1,t}(\Omega)$. Therefore, in order to improve the previous weak convergence and get (1.4), it is sufficient to get

$$\nabla u_n \in L^{p(1+\delta)}(\Omega), \quad \text{for every } n \in \mathbb{N},$$

with a uniform δ and a uniform control on the relevant norms. Indeed, recall that $u_n \in W_0^{1,p(1+\delta)}(\Omega)$ just by construction (at least for δ small enough);

- the previous global higher integrability result can be obtained by repeating the above scheme of proof: joining Caccioppoli and Maz’ya–Sobolev–Poincaré inequalities to get a reverse Hölder inequality. This goes exactly as in [26], up to some technical modifications needed to cope with the presence of the right-hand side f and the possible unboundedness of Ω .

At a technical level, we proceed as follows: we get uniform reverse Hölder inequalities not for $|\nabla u_n|^p$, but for the following non-homogeneous quantity

$$H_n(|\nabla u_n|) := |\nabla u_n|^p + \alpha_n |\nabla u_n|^q,$$

which is naturally attached to the perturbed problem. The idea of considering this quantity is taken from [14, Theorem 1.1], which deals with regularity issues for double-phase problems

(see also [1, 13, 16]). The above described scheme will then produce the following estimate

$$\int_{\Omega} H_n(|\nabla u_n|)^{1+\delta} \leq C,$$

with C and δ uniform in n , implying the same uniform bound for the $L^{p(1+\delta)}(\Omega)$ norm of ∇u_n , as claimed.

Remark 1.4. After the completion of the paper, we became aware of the reference [25] by Hedberg and Kilpelainen. In particular, [25, Corollary 3.5] shows that

$$W^{1,p(1+\delta)}(\Omega) \cap W_0^{1,p}(\Omega) = W_0^{1,p(1+\delta)}(\Omega),$$

whenever $\Omega \subseteq \mathbb{R}^N$ is an open bounded set, whose complement $\mathbb{R}^N \setminus \Omega$ is uniformly p -thick. Apart from being stated for *bounded* sets, the proof of their result is based on the fine properties of Sobolev functions, i.e. it uses p -quasi continuous representatives of Sobolev functions and the sophisticated *Havin-Bagby Theorem* for characterizing $W_0^{1,p}$ spaces on open sets.

We believe that our proof, based on constructing in an elementary way an approximating sequence and relying only on very basic properties of capacities and Sobolev spaces, is interesting in itself. The argument is simple and robust enough to be quite easily extended to more general situations (i.e. double-phase operators, $p(x)$ -Laplacian, orthotropic equations with nonstandard growth, just to list some examples), without the necessity to exploit the fine properties of the underlying Sobolev spaces.

1.5. Plan of the paper. We start with Section 2, containing all the notation and definitions needed throughout the paper. We also collect some properties of open sets with thick complement, which will be needed for the proof of the Main Theorem. In Section 3 we construct the approximating sequence $\{u_n\}_{n \in \mathbb{N}}$, through the “fake” double-phase functional previously discussed. We also obtain the two main Caccioppoli inequalities for these minimizers, needed to get the desired reverse Hölder inequalities. The proof of the Main Theorem is contained in Section 4: it will be divided in various steps, in order to assist the reader. We will get the result at first under a normalization condition; then we will use a scaling argument to reach the desired scale invariant estimate. The main purpose of Section 5 is that of offering a simple and sufficiently rich class of open sets to which our main result applies: we will show that *contractible* open sets satisfy the thickness conditions, under suitable restrictions on p and the dimension. We will also comment on the geometric aspects of the assumption (E_{Ω}) , for this class of sets (see Proposition 5.5). Finally, in Section 6 we will come to our original intent and show how our Main Theorem can be applied to the case of Lane-Emden equations (including the eigenvalue equation).

The paper is complemented by two appendices: while Appendix B is very brief and simply shows how the Main Theorem could be further extended in the case $1 < p < N$, in Appendix A we prove a property of open sets in the class $UT_p(c_0, r_0)$, that we believe to be of independent interest (see Theorem A.2 below).

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2. PRELIMINARIES

2.1. Notation. Throughout the whole paper, we will always assume that the dimension $N \geq 2$. We denote by $Q_r(x_0)$ the open cube centered at $x_0 \in \mathbb{R}^N$ with side length $2r > 0$, that is

$$Q_r(x_0) = \{x \in \mathbb{R}^N : \|x - x_0\|_{\ell^\infty} < r\},$$

and by $B_r(x_0)$ the open ball centered at $x_0 \in \mathbb{R}^N$ with radius $r > 0$. For a measurable set $E \subseteq \mathbb{R}^N$ such that $0 < |E| < +\infty$ we denote by

$$\bar{u}_E = \oint_E u \, dx = \frac{1}{|E|} \int u \, dx,$$

the averaged integral of the summable function u over E . The symbol $|E|$ stands for the N -dimensional Lebesgue measure of E .

2.2. Basics. The next result collects some properties of the open sets satisfying Assumptions 1. Recall that the quantity $\lambda_p(\Omega)$ has been defined in (1.2).

Proposition 2.1. *Let $1 < p < \infty$ and let $\gamma_0 \leq p'$ be such that*

$$\begin{cases} (p^*)' \leq \gamma_0 & \text{if } 1 < p < N, \\ 1 < \gamma_0, & \text{if } p = N, \\ 1 \leq \gamma_0, & \text{if } p > N. \end{cases}$$

Let $\Omega \subseteq \mathbb{R}^N$ be an open set such that $\lambda_p(\Omega) > 0$, then:

(1) *Ω has finite inradius, i.e. $r_\Omega < +\infty$;*

(2) *$\mathcal{D}_0^{1,p}(\Omega) = W_0^{1,p}(\Omega)$ and*

$$\varphi \mapsto \|\nabla \varphi\|_{L^p(\Omega)},$$

is an equivalent norm on $W_0^{1,p}(\Omega)$;

(3) *we also have*

$$(2.1) \quad \lambda_{p,\gamma'_0}(\Omega) := \inf_{\varphi \in C_0^\infty(\Omega)} \left\{ \int_\Omega |\nabla \varphi|^p \, dx : \|\varphi\|_{L^{\gamma'_0}(\Omega)} = 1 \right\} > 0.$$

Proof. The first result is a well-known necessary condition for the validity of Poincaré inequalities, see for example [35, Proposition 2.1].

The point (2) easily follows from the fact that $\lambda_p(\Omega) > 0$, since this is equivalent to the validity of the Poincaré inequality

$$\lambda_p(\Omega) \int_\Omega |\varphi|^p \, dx \leq \int_\Omega |\nabla \varphi|^p \, dx, \quad \text{for every } \varphi \in C_0^\infty(\Omega),$$

as already said.

As for the last point, we can suppose $\gamma_0 < p'$. Then we have

$$\begin{cases} p < \gamma'_0 \leq p^* & \text{if } 1 < p < N, \\ p < \gamma'_0 < \infty, & \text{if } p = N, \\ p < \gamma'_0 \leq \infty, & \text{if } p > N. \end{cases}$$

Observe that for the limit case $\gamma'_0 = p^*$ we always have

$$\lambda_{p,p^*}(\Omega) > 0.$$

Indeed, this constant coincides with the sharp Sobolev constant, which is independent of the open set Ω (see [37, Chapter I, Section 4.5]). For the remaining case, we recall that [32, Theorem 15.4.1] shows

$$\lambda_p(\Omega) > 0 \quad \Longleftrightarrow \quad \lambda_{p,\gamma'_0}(\Omega) > 0.$$

This concludes the proof. $\square$

The following simple result is concerned with the functions G considered in the Main Theorem. This should be well-known, we include it for completeness.

Lemma 2.2. *Let $1 < p < \infty$ and let $G : \mathbb{R}^N \rightarrow \mathbb{R}$ be a C^1 convex function, such that*

$$g_1 |z|^p \leq G(z) \leq g_2 |z|^p,$$

for two constants $g_2 \geq g_1 > 0$. Then we have

$$|\nabla G(z)| \leq (2^p g_2 - g_1) |z|^{p-1}, \quad \text{for every } z \in \mathbb{R}^N,$$

and

$$\langle \nabla G(z), z \rangle \geq g_1 |z|^p, \quad \text{for every } z \in \mathbb{R}^N.$$

Proof. We take $z \in \mathbb{R}^N$ and $\omega \in \mathbb{S}^{N-1}$. By the “above tangent” property of convex functions, we have

$$G(z + |z|\omega) \geq G(z) + \langle \nabla G(z), \omega \rangle |z|.$$

In particular, by using the growth conditions on G , we obtain

$$g_2 |z + |z|\omega|^p \geq g_1 |z|^p + \langle \nabla G(z), \omega \rangle |z|.$$

By the triangle inequality, we can further estimate the left-hand side as follows

$$|z + |z|\omega|^p \leq 2^p |z|^p.$$

Thus, in conclusion we obtain

$$\langle \nabla G(z), \omega \rangle |z| \leq (2^p g_2 - g_1) |z|^p.$$

By taking the maximum over $\omega \in \mathbb{S}^{N-1}$, we get the first estimate.

For the second one, observe that for $z = 0$ there is nothing to prove. We take $z \neq 0$ and we observe that for every $t \in [0, 1]$ we have

$$\langle \nabla G(z) - \nabla G(tz), z - tz \rangle \geq 0.$$

This follows from the monotonicity of ∇G , which in turn is a consequence of the convexity of G . In particular, the previous inequality implies that

$$(2.2) \quad \langle \nabla G(z), z \rangle \geq \langle \nabla G(tz), z \rangle, \quad \text{for every } t \in [0, 1].$$

We can now write

$$g_1 |z|^p \leq G(z) = G(z) - G(0) = \int_0^1 \langle \nabla G(tz), z \rangle dt \leq \langle \nabla G(z), z \rangle,$$

thanks to (2.2). This gives the second estimate, as well. $\square$

As in [26], we will make a crucial use of the following classical result, the so-called *Gehring Lemma* (see [18]) in the version of Giaquinta and Modica [20] and Stredulinsky [36]. Differently from [26], we prefer to work here with cubes, instead of balls. The statement is taken from [21, Theorem 6.6], to which we refer for the proof.

Lemma 2.3. *Let $h \in L^1(Q_L(x_0))$ and $k \in L^s(Q_L(x_0))$, for some $s > 1$. Let us suppose that there exists $0 < \varepsilon < 1$ and $C > 0$ such that for every $Q_r(y_0) \subseteq Q_L(x_0)$ we have*

$$(2.3) \quad \int_{Q_{r/2}(y_0)} |h| dx \leq C \left[\left(\int_{Q_r(y_0)} |h|^{(1-\varepsilon)} dx \right)^{\frac{1}{1-\varepsilon}} + \int_{Q_r(y_0)} |k| dx \right].$$

Then, there exists an exponent $\delta = \delta(N, \varepsilon, C, s) > 0$ such that $\delta \leq \min\{1, s-1\}$ and a constant $\tilde{C} > 0$ such that

$$h \in L^{1+\delta}(Q_{L/2}(x_0)),$$

with the following estimate

$$(2.4) \quad \int_{Q_{L/2}(x_0)} |h|^{1+\delta} dx \leq \tilde{C} \left[\left(\int_{Q_L(x_0)} |h| dx \right)^{1+\delta} + \int_{Q_L(x_0)} |k|^{1+\delta} dx \right].$$

2.3. Open sets with thick complement. In the proof of the Main Theorem we will exploit some properties of *variational p -capacity*. We recall at first the definition employed in this paper.

Definition 2.4. Let $1 \leq p < \infty$, for every $E \subseteq \mathbb{R}^N$ open set and every $K \subseteq E$ compact set, the *p -capacity of K relative to E* is defined as

$$\text{cap}_p(K; E) = \inf_{\varphi \in C_0^\infty(E)} \left\{ \int_E |\nabla \varphi|^p dx : \varphi \geq 1 \text{ on } K \right\}.$$

We refer to [32] for a comprehensive study on the properties of the p -capacity.

We observe that the notion of uniform p -thickness is usually stated in terms of balls. It is not difficult to show that this is equivalent to our definition in terms of cubes. This is a straightforward consequence of the next result, that we enclose for completeness.

Lemma 2.5. *Let $1 \leq p < \infty$ and let $K \subseteq \mathbb{R}^N$ be a non-empty compact set. There exist two constants $0 < \beta_1 < \beta_2$, both depending on N and p only, such that for every $x \in K$ and every $r > 0$ we have that*

$$\beta_2 \text{cap}_p(\overline{Q_r(x)} \cap K, Q_{2r}(x)) \geq \text{cap}_p(\overline{B_r(x)} \cap K, B_{2r}(x)) \geq \beta_1 \text{cap}_p(\overline{Q_r(x)} \cap K, Q_{2r}(x)).$$

Proof. We start by observing that $B_r(x) \subseteq Q_r(x) \subseteq B_{r\sqrt{N}}(x)$. By the monotonicity properties of the relative capacity, we infer that

$$\text{cap}_p(\overline{Q_r(x)} \cap K, Q_{2r}(x)) \geq \text{cap}_p(\overline{B_r(x)} \cap K, Q_{2r}(x)) \geq \text{cap}_p(\overline{B_r(x)} \cap K, B_{2r\sqrt{N}}(x)).$$

We can appeal again to [32, Proposition 13.1.1.1], in order to get

$$\text{cap}_p(\overline{B_r(x)} \cap K, B_{2r\sqrt{N}}(x)) \geq c \text{cap}_p(\overline{B_r(x)} \cap K, B_{2r}(x)),$$

for a constant $c = c(N, p) > 0$. Thus, we get

$$\text{cap}_p(\overline{Q_r(x)} \cap K, Q_{2r}(x)) \geq c \text{cap}_p(\overline{B_r(x)} \cap K, B_{2r}(x)),$$

i.e., the leftmost estimate. The proof of the other estimate is similar, it is sufficient to notice that $Q_{r/\sqrt{N}}(x) \subseteq B_r(x) \subseteq Q_r(x)$. $\square$

The next basic result is quite straightforward (see also [2, Remark 4]), we give the details in order to keep track of the constants.

Lemma 2.6. *Let $1 \leq p_0 < \infty$ and let $\Omega \in \text{UT}_{p_0}(c_0, r_0)$, for some $c_0, r_0 > 0$. For every $q > p_0$ there exists a constant $c_1 = c_1(N, p_0, q) > 0$ such that*

$$\Omega \in \text{UT}_q(c_1, r_0),$$

as well.

Proof. We take $x_0 \in \mathbb{R}^N \setminus \Omega$ and $0 < r \leq r_0$. For every $\varphi \in C_0^\infty(Q_{2r}(x_0))$ such that $\varphi \geq 1$ on $\overline{Q_r(x_0)} \setminus \Omega$, we have

$$\begin{aligned} c_0 \text{cap}_{p_0}(\overline{Q_r(x_0)}; Q_{2r}(x_0)) &\leq \text{cap}_{p_0}(\overline{Q_r(x_0)} \setminus \Omega; Q_{2r}(x_0)) \\ &\leq \int_{Q_{2r}(x_0)} |\nabla \varphi|^{p_0} dx \leq |Q_{2r}(x_0)|^{1-\frac{p_0}{q}} \left(\int_{Q_{2r}(x_0)} |\nabla \varphi|^q dx \right)^{\frac{q}{p_0}}. \end{aligned}$$

By taking the infimum over φ , we get the desired conclusion. $\square$

Remark 2.7. By inspecting the previous proof, we see that we can choose c_1 to be

$$c_1 = 4^N \left(\frac{c_0}{4^N} \right)^{\frac{q}{p_0}} \frac{(\text{cap}_{p_0}(\overline{Q_1}; Q_2))^{\frac{q}{p_0}}}{\text{cap}_q(\overline{Q_1}; Q_2)}.$$

The following simple result expresses the scaling properties of the thickness condition. This will be useful in order to get the a priori estimate of the main result.

Lemma 2.8. *Let $1 \leq p_0 < \infty$ and let $\Omega \subsetneq \mathbb{R}^N$ be an open set such that $\Omega \in \text{UT}_{p_0}(c_0, r_0)$. For every $t > 0$, the scaled set $t\Omega$ satisfies*

$$t\Omega \in \text{UT}_{p_0}(c_0, tr_0).$$

Proof. Let $x_0 \in \mathbb{R}^N \setminus (t\Omega)$, then $x_0/t \in \mathbb{R}^N \setminus \Omega$. Accordingly, for every $0 < r \leq tr_0$ we have

$$\begin{aligned} (2.5) \quad \text{cap}_{p_0}(\overline{Q_{r/t}(x_0/t)} \setminus \Omega; Q_{2r/t}(x_0/t)) &\geq c_0 \text{cap}_{p_0}(\overline{Q_{r/t}}; Q_{2r/t}) \\ &= c_0 \text{cap}_{p_0}(\overline{Q_r}; Q_{2r}) t^{p_0-N}. \end{aligned}$$

We used both the definition of p_0 -thickness and the scaling properties of the relative p_0 -capacity. On the other hand, we have that

$$\begin{aligned} t^{p_0-N} \text{cap}_{p_0}(\overline{Q_r(x_0)} \setminus (t\Omega); Q_{2r}(x_0)) &= \text{cap}_{p_0}(t^{-1}(\overline{Q_r(x_0)} \setminus (t\Omega)); t^{-1}Q_{2r}(x_0)) \\ &= \text{cap}_{p_0}(\overline{Q_{r/t}(x_0/t)} \setminus \Omega; Q_{2r/t}(x_0/t)). \end{aligned}$$

By using (2.5) to estimate the right-hand side, we conclude. $\square$

The following technical lemma will be needed in order to control Lebesgue norms on certain cubes lying “sufficiently close” to the boundary. This is a slight generalization of [34, Lemma 3.8], with cubes in place of balls. It will be needed in order to properly fix the argument used in² [26].

Lemma 2.9. *Let $1 \leq p_0 < \infty$ and let $\Omega \subsetneq \mathbb{R}^N$ be an open set such that $\Omega \in \text{UT}_{p_0}(c_0, r_0)$. We fix a point $x_0 \in \Omega$ such that*

$$\mathfrak{d}(x_0) := \inf_{y \in \partial\Omega} \|x_0 - y\|_{\ell^\infty}.$$

There exists a constant $C_2 > 1$, depending only on N, p_0 , and c_0 , such that

$$\begin{aligned} \text{cap}_{p_0} \left(\overline{Q_r(x_0)} \setminus \Omega; Q_{2r}(x_0) \right) &\geq \frac{1}{C_2} \left(\frac{r - \mathfrak{d}(x_0)}{2r + \mathfrak{d}(x_0)} \right)^{p_0} \\ &\quad \times (r - \mathfrak{d}(x_0))^{N-p_0}, \quad \text{for every } \mathfrak{d}(x_0) < r \leq r_0 + \mathfrak{d}(x_0). \end{aligned}$$

Proof. Since Ω is a proper open subset of $\mathbb{R}^N$, its boundary $\partial\Omega$ is closed and non-empty. Therefore, we can find a point $y_0 \in \partial\Omega$ such that

$$\|y_0 - x_0\|_\infty = \mathfrak{d}(x_0).$$

For simplicity, in what follows we write $\mathfrak{d}$ in place of $\mathfrak{d}(x_0)$. For every $r > \mathfrak{d}$, we observe that

$$Q_{r-\mathfrak{d}}(y_0) \subseteq Q_r(x_0) \quad \text{and} \quad Q_{2r}(x_0) \subseteq Q_{2r+\mathfrak{d}}(y_0).$$

By the monotonicity properties of the capacity, we can estimate from below

$$\begin{aligned} \text{cap}_{p_0} \left(\overline{Q_r(x_0)} \setminus \Omega; Q_{2r}(x_0) \right) &\geq \text{cap}_{p_0} \left(\overline{Q_{r-\mathfrak{d}}(y_0)} \setminus \Omega; Q_{2r}(x_0) \right) \\ &\geq \text{cap}_{p_0} \left(\overline{Q_{r-\mathfrak{d}}(y_0)} \setminus \Omega; Q_{2r+\mathfrak{d}}(y_0) \right). \end{aligned}$$

We can now apply [32, Proposition 13.1.1.1], in order to further estimate from below the last term as follows

$$\text{cap}_{p_0} \left(\overline{Q_{r-\mathfrak{d}}(y_0)} \setminus \Omega; Q_{2r+\mathfrak{d}}(y_0) \right) \geq c \left(\frac{r - \mathfrak{d}}{2r + \mathfrak{d}} \right)^{p_0} \text{cap}_{p_0} \left(\overline{Q_{r-\mathfrak{d}}(y_0)} \setminus \Omega; Q_{2(r-\mathfrak{d})}(y_0) \right),$$

where $c = c(N, p_0) > 0$. It is now sufficient to use the property of p_0 -thickness, by observing that $0 < r - \mathfrak{d} \leq r_0$, for every r as in the statement. $\square$

The next result will be crucial in order to achieve the desired gain of regularity. This is analogous to [26, Lemma 3.1], but we will avoid here to use the fine properties of Sobolev functions. It can be obtained by a density argument, as a straightforward consequence of a classical result by Maz'ya.

Lemma 2.10 (Maz'ya-Poincaré-Sobolev inequality). *Let $1 \leq p_0 \leq p < \infty$ and let $\theta \geq 1$ be such that*

$$\begin{cases} \theta \leq p_0^*, & \text{if } 1 \leq p_0 < N, \\ \theta < \infty, & \text{if } p_0 \geq N. \end{cases} \quad \text{where } p_0^* = \frac{N p_0}{N - p_0}.$$

Let $\Omega \subsetneq \mathbb{R}^N$ be an open set such that $\Omega \in \text{UT}_{p_0}(c_0, r_0)$. There exists a constant $C_3 = C_3(N, p_0, \theta, c_0) > 0$ such that for every $u \in W_0^{1,p}(\Omega)$ we have

$$\left(\int_{Q_r(x_0)} |u|^\theta dx \right)^{\frac{p_0}{\theta}} \leq C_3 r^{p_0} \int_{Q_r(x_0)} |\nabla u|^{p_0} dx,$$

²In [26], the proof of the main result is slightly incorrect, since it uses the thickness property of $\mathbb{R}^N \setminus \Omega$ also in the case of balls lying partly outside Ω , but whose center still belongs to Ω . The paper by Parviainen [34] fixed this issue in order to prove the parabolic counterpart of [26, Theorem 1.1].

where $Q_r(x_0)$ is such that $0 < r \leq r_0$ and

- either $x_0 \in \mathbb{R}^N \setminus \Omega$;
- or $x_0 \in \Omega$ with

$$(2.6) \quad \mathfrak{d}(x_0) := \min_{y \in \partial\Omega} \|x_0 - y\|_{\ell^\infty} \leq \frac{2}{3} r.$$

Proof. By definition of $W_0^{1,p}(\Omega)$, there exists a sequence $\{u_n\}_{n \in \mathbb{N}} \subseteq C_0^\infty(\Omega)$ such that

$$(2.7) \quad \lim_{n \rightarrow \infty} \|u_n - u\|_{W^{1,p}(\Omega)} = 0.$$

By extending the functions to be 0 on the complement $\mathbb{R}^N \setminus \Omega$, we can thus in particular infer that

$$\lim_{n \rightarrow \infty} \int_{Q_r(x_0)} |\nabla u_n - \nabla u|^{p_0} dx = 0,$$

for $1 \leq p_0 \leq p$ and for every cube as in the statement. As for the L^θ norm of $u_n - u$, we observe that for $1 \leq \theta \leq p$ we directly get

$$(2.8) \quad \lim_{n \rightarrow \infty} \int_{Q_r(x_0)} |u_n - u|^\theta dx = 0,$$

by (2.7) and Jensen's inequality. On the other hand, for an exponent $\theta > p$ satisfying the restrictions of the statement, it is sufficient to use the following standard Gagliardo-Nirenberg interpolation inequality

$$\|u_n - u\|_{L^\theta(\mathbb{R}^N)} \leq G_{N,p,\theta} \|\nabla u_n - \nabla u\|_{L^p(\mathbb{R}^N)}^\vartheta \|u_n - u\|_{L^p(\mathbb{R}^N)}^{1-\vartheta}, \quad \text{for every } n \in \mathbb{N},$$

see for example [7, Theorem 3.8.1] or [29, Theorem 12.83]. Here the exponent $0 < \vartheta \leq 1$ is dictated by scale invariance, its precise expression does not matter. From (2.7), the previous estimate implies that

$$\lim_{n \rightarrow \infty} \|u_n - u\|_{L^\theta(\mathbb{R}^N)} = 0,$$

and thus we get (2.8), as well.

The previous discussion shows that it is sufficient to prove the claimed inequality for a function $u \in C_0^\infty(\Omega)$. As before, we consider u to be defined on the whole $\mathbb{R}^N$, by extending it to 0 in $\mathbb{R}^N \setminus \Omega$. In this case, we can directly appeal to [32, Theorem 14.1.2] and obtain

$$\text{cap}_{p_0} \left(\overline{Q_r(x_0)} \setminus \Omega; Q_{2r}(x_0) \right) \left(\int_{Q_r(x_0)} |u|^\theta dx \right)^{\frac{p_0}{\theta}} \leq C r^N \int_{Q_r(x_0)} |\nabla u|^{p_0} dx,$$

for every $x_0 \in \mathbb{R}^N$ and every $r > 0$. Here the constant $C = C(N, p_0, \theta) > 0$. We need to estimate from below the capacity of $\overline{Q_r(x_0)} \setminus \Omega$. If $x_0 \in \mathbb{R}^N \setminus \Omega$, we can use the p_0 -thickness assumption on $\mathbb{R}^N \setminus \Omega$ and get for every $0 < r \leq r_0$

$$\begin{aligned} \text{cap}_{p_0} \left(\overline{Q_r(x_0)} \setminus \Omega; Q_{2r}(x_0) \right) &\geq c_0 \text{cap}_{p_0} \left(\overline{Q_r(x_0)}; Q_{2r}(x_0) \right) \\ &= c_0 \text{cap}_{p_0} (\overline{Q_1}; Q_2) r^{N-p_0}. \end{aligned}$$

Observe that we used the scaling properties of the relative capacity. On the other hand, if $x_0 \in \Omega$ and (2.6) holds, we can apply Lemma 2.9 and obtain

$$\begin{aligned} \text{cap}_{p_0} \left(\overline{Q_r(x_0)} \setminus \Omega; Q_{2r}(x_0) \right) &\geq \frac{1}{C_2} \left(\frac{r - \mathfrak{d}(x_0)}{2r + \mathfrak{d}(x_0)} \right)^{p_0} (r - \mathfrak{d}(x_0))^{N-p_0} \\ &\geq \frac{1}{C} \left(\frac{1}{8} \right)^{p_0} \left(\frac{1}{3} \right)^{N-p_0} r^{N-p_0}. \end{aligned}$$

The desired conclusion now follows. $\square$

3. AN APPROXIMATING PROBLEM

3.1. Construction of the approximating sequence. As explained in the Introduction, in order to achieve the desired regularity result we will need to construct a suitable sequence $\{u_n\}_{n \in \mathbb{N}} \subseteq W_0^{1,p(1+\delta)}(\Omega)$ weakly converging to the solution u of the Main Theorem.

To this aim, it is useful to remark that u is the unique minimizer of the following strictly convex functional

$$\mathfrak{F}(\varphi) = \int_{\Omega} G(\nabla \varphi) dx - \int_{\Omega} f \varphi dx, \quad \text{for every } \varphi \in W_0^{1,p}(\Omega).$$

We are going to construct such a sequence through a penalized minimization problem. This is the content of the next result.

Proposition 3.1. *Under the assumptions of the Main Theorem, we set*

$$(3.1) \quad q := 2p,$$

and we consider the family of energy functionals

$$\mathfrak{F}_n(\varphi) = \int_{\Omega} G(\nabla \varphi) dx + \frac{\alpha_n}{q} \int_{\Omega} |\nabla \varphi|^q dx - \int_{\Omega} f \varphi dx, \quad \text{for every } \varphi \in W_0^{1,p}(\Omega) \cap W_0^{1,q}(\Omega),$$

where $\{\alpha_n\}_{n \in \mathbb{N}}$ is a sequence of positive numbers, such that

$$\lim_{n \rightarrow \infty} \alpha_n = 0.$$

Then, for every $n \in \mathbb{N}$ there exists a unique minimizer $u_n \in W_0^{1,p}(\Omega) \cap W_0^{1,q}(\Omega)$ for $\mathfrak{F}_n$. Moreover, u_n is a weak solution of the equation

$$(3.2) \quad -\operatorname{div} \nabla G(\nabla u_n) - \alpha_n \operatorname{div} (|\nabla u_n|^{q-2} \nabla u_n) = f, \quad \text{in } \Omega.$$

Proof. Thanks to the choice of γ_0 and the assumption $\lambda_p(\Omega) > 0$, by Proposition 2.1 we have that the embedding

$$W_0^{1,p}(\Omega) \hookrightarrow L^{\gamma'_0}(\Omega),$$

is continuous. Thus, for every $\varphi \in W_0^{1,p}(\Omega)$ we have

$$(3.3) \quad \left| \int_{\Omega} f \varphi dx \right| \leq \|f\|_{L^{\gamma_0}(\Omega)} \|\varphi\|_{L^{\gamma'_0}(\Omega)} \leq \|f\|_{L^{\gamma_0}(\Omega)} \frac{\|\nabla \varphi\|_{L^p(\Omega)}}{\left(\lambda_{p,\gamma'_0}(\Omega) \right)^{\frac{1}{p}}}.$$

We recall that $\lambda_{p,\gamma'_0}(\Omega)$ is defined in (2.1). This shows that $\mathfrak{F}_n$ is well-defined on $W_0^{1,p}(\Omega) \cap W_0^{1,q}(\Omega)$. Moreover, since $\lambda_p(\Omega) > 0$, by Proposition 2.1 we have that

$$\varphi \mapsto \|\nabla \varphi\|_{L^p(\Omega)},$$

is an equivalent norm on $W_0^{1,p}(\Omega)$. The fact that $\lambda_p(\Omega) > 0$ implies $\lambda_q(\Omega) > 0$, as well (see for example [12, Proposition 2.1]). Thus, also the space $W_0^{1,q}(\Omega)$ can be equivalently normed by

$$\varphi \mapsto \|\nabla \varphi\|_{L^q(\Omega)}.$$

In force of all these preliminary observations, existence of a minimizer u_n follows by a standard application of the Direct Methods in the Calculus of Variations. Uniqueness is a consequence of the strict convexity of the functionals. Finally, the fact that u_n weakly solves (3.2) simply follows from the minimality, once noticed that the functional $\mathfrak{F}_n$ is Gateaux differentiable. $\square$

We now prove that the sequence of minimizers $\{u_n\}_{n \in \mathbb{N}}$ converges to the original solution u .

Proposition 3.2. *With the notation of Proposition 3.1, we have*

$$(3.4) \quad \lim_{n \rightarrow \infty} \left(\min_{\varphi \in W_0^{1,p}(\Omega) \cap W_0^{1,q}(\Omega)} \mathfrak{F}_n(\varphi) \right) = \min_{\varphi \in W_0^{1,p}(\Omega)} \mathfrak{F}(\varphi).$$

Moreover, we have

$$u_n \xrightarrow{n \rightarrow \infty} u, \quad \text{in } W^{1,p}(\Omega),$$

where $u \in W_0^{1,p}(\Omega)$ is the unique minimizer of $\mathfrak{F}$. Finally, we have

$$(3.5) \quad \limsup_{n \rightarrow \infty} \int_{\Omega} \left[|\nabla u_n|^p + \alpha_n |\nabla u_n|^q \right] dx \leq C_{g_1, g_2, q} \int_{\Omega} |\nabla u|^p dx.$$

Proof. In order to prove (3.4), observe that for every $\varphi \in C_0^\infty(\Omega)$ we have

$$\begin{aligned} \mathfrak{F}(\varphi) &= \int_{\Omega} G(\nabla \varphi) dx - \int_{\Omega} f \varphi dx = \lim_{n \rightarrow \infty} \mathfrak{F}_n(\varphi) \geq \limsup_{n \rightarrow \infty} \mathfrak{F}_n(u_n) \\ &= \limsup_{n \rightarrow \infty} \left(\min_{\psi \in W_0^{1,p}(\Omega) \cap W_0^{1,q}(\Omega)} \mathfrak{F}_n(\psi) \right). \end{aligned}$$

By the arbitrariness of $\varphi \in C_0^\infty(\Omega)$, this shows that

$$\limsup_{n \rightarrow \infty} \left(\min_{\varphi \in W_0^{1,p}(\Omega) \cap W_0^{1,q}(\Omega)} \mathfrak{F}_n(\varphi) \right) \leq \inf_{\varphi \in C_0^\infty(\Omega)} \mathfrak{F}(\varphi) = \min_{\varphi \in W_0^{1,p}(\Omega)} \mathfrak{F}(\varphi) < +\infty.$$

To get the reverse inequality, we observe that the previous estimate implies in particular that $\mathfrak{F}_n(u_n) \leq M < +\infty$, for every $n \in \mathbb{N}$. Thus, by recalling that $\alpha_n > 0$ we get

$$\begin{aligned} M &\geq \int_{\Omega} G(\nabla u_n) dx - \int_{\Omega} f u_n dx \\ &\geq g_1 \int_{\Omega} |\nabla u_n|^p dx - \|f\|_{L^{\gamma_0}(\Omega)} \frac{\|\nabla u_n\|_{L^p(\Omega)}^{\frac{1}{p}}}{\left(\lambda_{p, \gamma'_0}(\Omega)\right)^{\frac{1}{p}}}. \end{aligned}$$

In the last inequality, we used (3.3). By using Young's inequality, for every $\varepsilon > 0$ we get

$$M \geq g_1 \int_{\Omega} |\nabla u_n|^p dx - \frac{1}{\left(\lambda_{p, \gamma'_0}(\Omega)\right)^{\frac{1}{p-1}}} \frac{(p-1) \varepsilon^{-\frac{1}{p-1}}}{p} \|f\|_{L^{\gamma_0}(\Omega)}^{p'} - \frac{\varepsilon}{p} \int_{\Omega} |\nabla u_n|^p dx.$$

It is now sufficient to choose $\varepsilon = (g_1 p)/2$, so to finally get

$$\int_{\Omega} |\nabla u_n|^p dx \leq C = C(N, p, \lambda_{p, \gamma'_0}(\Omega), g_1, \|f\|_{L^{\gamma_0}(\Omega)}), \quad \text{for every } n \in \mathbb{N}.$$

This proves that $\{u_n\}_{n \in \mathbb{N}}$ is a bounded sequence in $W_0^{1,p}(\Omega)$. Thus, up to a subsequence, we have that u_n weakly converges to v in $W^{1,p}(\Omega)$. Moreover, we still have $v \in W_0^{1,p}(\Omega)$, since the latter is weakly closed, as well.

In light of all these informations, we get

$$\begin{aligned} \min_{\varphi \in W_0^{1,p}(\Omega)} \mathfrak{F}(\varphi) &\leq \mathfrak{F}(v) \leq \liminf_{n \rightarrow \infty} \mathfrak{F}(u_n) \leq \liminf_{n \rightarrow \infty} \mathfrak{F}_n(u_n) \\ &= \liminf_{n \rightarrow \infty} \left(\min_{\varphi \in W_0^{1,p}(\Omega) \cap W_0^{1,q}(\Omega)} \mathfrak{F}_n(\varphi) \right). \end{aligned}$$

In the second inequality, we used the lower semicontinuity of the functional $\mathfrak{F}$, with respect to the weak convergence of the gradients. This eventually establishes (3.4).

We can now ensure that all the inequalities in the previous estimates must hold as equalities. Thus, the limit function $v = u$ is the unique minimizer of $\mathfrak{F}$. This also implies that

$$\lim_{n \rightarrow \infty} \left[\int_{\Omega} G(\nabla u_n) dx + \frac{\alpha_n}{q} \int_{\Omega} |\nabla u_n|^q dx - \int_{\Omega} f u_n dx \right] = \int_{\Omega} G(\nabla u) dx - \int_{\Omega} f u dx.$$

By using the continuous embedding $W_0^{1,p}(\Omega) \hookrightarrow L^{\gamma'_0}(\Omega)$, we have that $\{u_n\}_{n \in \mathbb{N}}$ weakly converges³ also in $L^{\gamma'_0}(\Omega)$. Thus, we get

$$\lim_{n \rightarrow \infty} \int_{\Omega} f u_n dx = \int_{\Omega} f u dx.$$

By comparing the last two equations, we then obtain

$$\lim_{n \rightarrow \infty} \left[\int_{\Omega} G(\nabla u_n) dx + \frac{\alpha_n}{q} \int_{\Omega} |\nabla u_n|^q dx \right] = \int_{\Omega} G(\nabla u) dx.$$

In particular, by recalling the assumption on G , we get

$$\limsup_{n \rightarrow \infty} \int_{\Omega} \left[|\nabla u_n|^p + \alpha_n |\nabla u_n|^q \right] dx \leq \frac{g_2 q}{\min\{g_1 q, 1\}} \int_{\Omega} |\nabla u|^p dx.$$

The proof is now concluded. $\square$

3.2. Local energy estimates. By using the equation (3.2), we can in particular obtain some suitable Caccioppoli inequalities for the minimizers u_n .

Lemma 3.3 (Caccioppoli inequality – internal cubes). *Under the assumption of Proposition 3.1, let $Q_r(x_0) \subseteq \Omega$. Then, for every $0 < \sigma < 1$ and every $\mathbf{m} \in \mathbb{R}$ we have*

$$\begin{aligned} \int_{Q_{\sigma r}(x_0)} \left[|\nabla u_n|^p + \alpha_n |\nabla u_n|^q \right] dx &\leq \frac{C_4}{\sigma} \int_{Q_r(x_0)} \left[\frac{|u_n - \mathbf{m}|^p}{(1 - \sigma)^p r^p} + \alpha_n \frac{|u_n - \mathbf{m}|^q}{(1 - \sigma)^q r^q} \right] dx \\ &\quad + \frac{C_4}{\sigma} r^{p'} \left(\int_{Q_r(x_0)} |f|^{\gamma_0} dx \right)^{\frac{p'}{\gamma_0}}. \end{aligned}$$

The constant C depends only on the dimension N , the exponents p, q, γ_0 and the structural constants g_1, g_2 , but not on n .

³To be more precise, for $p > N$ and $\gamma'_0 = \infty$, the sequence $*$ -weakly converges.

Proof. Let us fix $0 < \sigma < 1$ and take a Lipschitz cut-off function η such that

$$0 \leq \eta \leq 1, \quad \eta = 1 \text{ on } Q_{\sigma r}(x_0), \quad \eta = 0 \text{ on } \partial Q_r(x_0), \quad |\nabla \eta| \leq \frac{1}{(1-\sigma)r}.$$

We then test the weak formulation of (3.2) with $\varphi = \eta^q (u_n - \mathbf{m}) \in W_0^{1,p}(\Omega) \cap W_0^{1,q}(\Omega)$. This yields

$$\begin{aligned} \int_{Q_r(x_0)} \left[\langle \nabla G(\nabla u_n), \nabla u_n \rangle + \alpha_n |\nabla u_n|^q \right] \eta^q dx &+ q \int_{Q_r(x_0)} \langle \nabla G(\nabla u_n), \nabla \eta \rangle \eta^{q-1} (u_n - \mathbf{m}) dx \\ &+ q \alpha_n \int_{Q_r(x_0)} \langle |\nabla u_n|^{q-2} \nabla u_n, \nabla \eta \rangle \eta^{q-1} (u_n - \mathbf{m}) dx \\ &= \int_{Q_r(x_0)} f \eta^q (u_n - \mathbf{m}) dx. \end{aligned}$$

By using the assumptions on G , Lemma 2.2 and Cauchy-Schwarz inequality, we get in particular

$$\begin{aligned} \int_{Q_r(x_0)} \left[g_1 |\nabla u_n|^p + \alpha_n |\nabla u_n|^q \right] \eta^q dx &\leq (2^p g_2 - g_1) q \int_{Q_r(x_0)} |\nabla u_n|^{p-1} |\nabla \eta| \eta^{q-1} |u_n - \mathbf{m}| dx \\ &+ q \alpha_n \int_{Q_r(x_0)} |\nabla u_n|^{q-1} |\nabla \eta| \eta^{q-1} |u_n - \mathbf{m}| dx \\ &+ \int_{Q_r(x_0)} f \eta^q (u_n - \mathbf{m}) dx. \end{aligned}$$

By using Young's inequality, we can infer that for every $\varepsilon > 0$

$$\begin{aligned} \int_{Q_r(x_0)} \left[|\nabla u_n|^p + \alpha_n |\nabla u_n|^q \right] \eta^q dx &\leq C \varepsilon \int_{Q_r(x_0)} |\nabla u_n|^p \eta^{(q-1)p'} dx \\ &+ C \varepsilon^{1-p} \int_{Q_r(x_0)} |\nabla \eta|^p |u_n - \mathbf{m}|^p dx \\ &+ C \alpha_n \varepsilon \int_{Q_r(x_0)} |\nabla u_n|^q \eta^q dx \\ &+ C \alpha_n \varepsilon^{1-q} \int_{Q_r(x_0)} |\nabla \eta|^q |u_n - \mathbf{m}|^q dx \\ &+ C \int_{Q_r(x_0)} f \eta^q (u_n - \mathbf{m}) dx, \end{aligned}$$

for a constant $C = C(p, q, g_1, g_2) > 0$. Observe that, since $0 \leq \eta \leq 1$ and $p < q$, we have

$$\eta^{(q-1)p'} \leq \eta^q.$$

By using this fact, choosing $\varepsilon = 1/(2C)$ and then reabsorbing the two terms containing ∇u_n in the right-hand side, we get after some simple manipulations

$$\begin{aligned} \int_{Q_r(x_0)} \left[g_1 |\nabla u_n|^p + \alpha_n |\nabla u_n|^q \right] \eta^q dx &\leq C \int_{Q_r(x_0)} |\nabla \eta|^p |u_n - \mathbf{m}|^p dx \\ &+ C \alpha_n \int_{Q_r(x_0)} |\nabla \eta|^q |u_n - \mathbf{m}|^q dx \\ &+ C \int_{Q_r(x_0)} f \eta^q (u_n - \mathbf{m}) dx, \end{aligned} \tag{3.6}$$

possibly for a different constant $C = C(p, q, g_1, g_2) > 0$. We now estimate the last term, thanks to Hölder and Poincaré-Sobolev inequalities. Indeed, thanks to the choice of the exponent γ_0 , we have

$$\begin{cases} \gamma'_0 < p^*, & \text{if } 1 < p < N, \\ 1 < \gamma'_0 < \infty, & \text{if } p = N, \\ \gamma'_0 \leq \infty, & \text{if } p > N. \end{cases}$$

Thus, we obtain

$$\begin{aligned} \int_{Q_r(x_0)} f \eta^q (u_n - \mathbf{m}) dx &\leq \left\| f \eta^{q-\frac{q}{p}} \right\|_{L^{\gamma_0}(Q_r(x_0))} \left\| \eta^{\frac{q}{p}} (u_n - \mathbf{m}) \right\|_{L^{\gamma'_0}(Q_r(x_0))} \\ &\leq \frac{\left\| f \eta^{q-\frac{q}{p}} \right\|_{L^{\gamma_0}(Q_r(x_0))}}{(\lambda_{p,\gamma'_0}(Q_r(x_0)))^{\frac{1}{p}}} \left(\int_{Q_r(x_0)} \left| \eta^{\frac{q}{p}} \nabla u_n + \frac{q}{p} (u_n - \mathbf{m}) \eta^{\frac{q}{p}-1} \nabla \eta \right|^p dx \right)^{\frac{1}{p}} \\ &\leq \frac{r^{\frac{N}{\gamma'_0}+1-\frac{N}{p}}}{(\lambda_{p,\gamma'_0}(Q_1(x_0)))^{\frac{1}{p}}} \left\| f \eta^{q-\frac{q}{p}} \right\|_{L^{\gamma_0}(Q_r(x_0))} \left(\int_{Q_r(x_0)} |\nabla u_n|^p \eta^q dx \right)^{\frac{1}{p}} \\ &\quad + \frac{q}{p} \frac{r^{\frac{N}{\gamma'_0}+1-\frac{N}{p}}}{(\lambda_{p,\gamma'_0}(Q_1(x_0)))^{\frac{1}{p}}} \left\| f \eta^{q-\frac{q}{p}} \right\|_{L^{\gamma_0}(Q_r(x_0))} \\ &\quad \times \left(\int_{Q_r(x_0)} |u_n - \mathbf{m}|^p \eta^{q-p} |\nabla \eta|^p dx \right)^{\frac{1}{p}}. \end{aligned}$$

Observe that we used the scaling properties of the Sobolev-Poincaré constant λ_{p,γ'_0} , to infer that

$$\lambda_{p,\gamma'_0}(Q_r(x_0)) = r^{N-p-\frac{p}{\gamma'_0}N} \lambda_{p,\gamma'_0}(Q_1(x_0)).$$

By using Young's inequality on both terms, we obtain for every $\varepsilon > 0$

$$\begin{aligned} \int_{Q_r(x_0)} f \eta^q (u_n - \mathbf{m}) dx &\leq \left(\frac{r^{\frac{N}{\gamma'_0}+1-\frac{N}{p}}}{(\lambda_{p,\gamma'_0}(Q_1(x_0)))^{\frac{1}{p}}} \right)^{p'} \frac{p-1}{p} \varepsilon^{-\frac{1}{p-1}} \left\| f \eta^{q-\frac{q}{p}} \right\|_{L^{\gamma_0}(Q_r(x_0))}^{p'} \\ &\quad + \frac{\varepsilon}{p} \int_{Q_r(x_0)} |\nabla u_n|^p \eta^q dx \\ &\quad + \left(\frac{r^{\frac{N}{\gamma'_0}+1-\frac{N}{p}}}{(\lambda_{p,\gamma'_0}(Q_1(x_0)))^{\frac{1}{p}}} \right)^{p'} \frac{p-1}{p} \left(\frac{q}{p} \right)^{p'} \left\| f \eta^{q-\frac{q}{p}} \right\|_{L^{\gamma_0}(Q_r(x_0))}^{p'} \\ &\quad + \frac{1}{p} \int_{Q_r(x_0)} |u_n - \mathbf{m}|^p |\nabla \eta|^p dx. \end{aligned}$$

By using this estimate in (3.6) and choosing $\varepsilon = p/4$, we can again absorb the term containing ∇u_n and obtain

$$\begin{aligned} \int_{Q_r(x_0)} \left[|\nabla u_n|^p + \alpha_n |\nabla u_n|^q \right] \eta^q dx &\leq C \int_{Q_r(x_0)} |\nabla \eta|^p |u_n - \mathbf{m}|^p dx \\ &\quad + C \alpha_n \int_{Q_r(x_0)} |\nabla \eta|^q |u_n - \mathbf{m}|^q dx \\ &\quad + C \left(r^{\frac{N}{\gamma_0} + 1 - \frac{N}{p}} \left\| f \eta^{q - \frac{q}{p}} \right\|_{L^{\gamma_0}(Q_r(x_0))} \right)^{p'}, \end{aligned}$$

for a constant $C = C(N, p, q, g_1, g_2, \gamma_0) > 0$. By using the properties of η , we can finally eliminate it and obtain the claimed estimate. $\square$

In order to get estimates near the boundary, we will also need the following variant of the previous result. This time, we do not require the cubes to be contained in Ω .

Lemma 3.4 (Caccioppoli inequality – general cubes). *Let $Q_r(x_0)$ be such that $Q_r(x_0) \cap \Omega \neq \emptyset$. Then, for every $0 < \sigma < 1$ we have*

$$\begin{aligned} \int_{Q_{\sigma r}(x_0)} \left[|\nabla u_n|^p + \alpha_n |\nabla u_n|^q \right] dx &\leq \frac{C_4}{\sigma} \int_{Q_r(x_0)} \left[\frac{|u_n|^p}{(1-\sigma)^p r^p} + \alpha_n \frac{|u_n|^q}{(1-\sigma)^q r^q} \right] dx \\ &\quad + \frac{C_4}{\sigma} r^{p'} \left(\int_{Q_r(x_0)} |f|^{\gamma_0} dx \right)^{\frac{p'}{\gamma_0}}, \end{aligned}$$

with the same constant C_4 of Lemma 3.3. Here the functions u_n and f are extended by 0 in the complement $\mathbb{R}^N \setminus \Omega$.

Proof. As above, we fix $0 < \sigma < 1$ and take a Lipschitz cut-off function η such that

$$0 \leq \eta \leq 1, \quad \eta = 1 \text{ on } Q_{\sigma r}(x_0), \quad \eta = 0 \text{ on } \partial Q_r(x_0), \quad |\nabla \eta| \leq \frac{1}{(1-\sigma)r}.$$

This time, we test the weak formulation of (3.2) with $\varphi = \eta^q u_n \in W_0^{1,p}(\Omega) \cap W_0^{1,q}(\Omega)$. We can then repeat verbatim the estimates of Lemma 3.3 and get the result. $\square$

4. PROOF OF THE MAIN THEOREM

We will divide the proof in two parts: in the first part, we will prove the claimed result under the further assumption

$$(4.1) \quad \|f\|_{L^{\gamma_0}(\Omega)} = 1.$$

In the second one, we will show how to remove this assumption by means of a scaling argument and get the desired a priori estimate.

4.1. Part 1: estimate at scale 1. We thus assume (4.1) and, for ease of readability, we further divide this part in various steps.

Step 1: set-up. For every $n \in \mathbb{N}$, we take u_n the minimizer constructed in Proposition 3.1. From now on, we set for brevity

$$H_n(t) := |t|^p + \alpha_n |t|^q, \quad \text{for every } t \in \mathbb{R}.$$

We take a cube $Q_1(x_0)$ such that $Q_1(x_0) \cap \Omega \neq \emptyset$. We wish to prove that for every $Q_r(y_0) \subseteq Q_1(x_0)$ we have

$$(4.2) \quad \int_{Q_{r/2}(y_0)} H_n(|\nabla u_n|) dx \leq C \left[\left(\int_{Q_r(y_0)} H_n(|\nabla u_n|)^{1-\varepsilon} dx \right)^{\frac{1}{1-\varepsilon}} + \int_{Q_r(y_0)} |f|^{\gamma_0} dx \right],$$

for an exponent $\varepsilon = \varepsilon(N, p_0, p) \in (0, 1)$ and a constant $C = C(N, p, q, p_0, c_0) > 0$, both independent of n . Here, the exponent γ_0 is the same of the Main Theorem.

In order to get (4.2), we must distinguish between different cases, depending on the position of the cube, or its center, with respect to Ω .

Step 2: proof of (4.2) for internal cubes. Let us suppose that

$$Q_r(y_0) \subseteq \Omega.$$

For every $n \in \mathbb{N}$, we apply the Caccioppoli inequality of Lemma 3.3, with the choices

$$\mathbf{m} = \mathbf{m}_n := \overline{u_n}_{Q_r(y_0)},$$

and $\sigma = 1/2$. This gives

$$(4.3) \quad \int_{Q_{r/2}(y_0)} H_n(|\nabla u_n|) dx \leq C \int_{Q_r(y_0)} H_n\left(\frac{u_n - \mathbf{m}_n}{r}\right) dx + C r^{p'} \left(\int_{Q_r(y_0)} |f|^{\gamma_0} dx \right)^{\frac{p'}{\gamma_0}}.$$

We set

$$(4.4) \quad p_1 = \max\left\{\frac{Np}{N+p}, 1\right\} \quad \text{and} \quad q_1 = \max\left\{\frac{Nq}{N+q}, 1\right\}.$$

Observe that we have

$$p_1 < p \leq p_1^*, \quad q_1 < q \leq q_1^* \quad \text{and} \quad \frac{q_1}{q} < \frac{p_1}{p},$$

hence by Sobolev-Poincaré inequality on cubes (see for example [21, Theorem 3.17]) we can infer

$$\int_{Q_r(y_0)} \frac{|u_n - \mathbf{m}_n|^p}{r^p} \leq \frac{1}{\mu_{p_1, p}} \left(\int_{Q_r(y_0)} |\nabla u_n|^{p_1} dx \right)^{\frac{p}{p_1}},$$

and

$$\int_{Q_r(y_0)} \frac{|u_n - \mathbf{m}_n|^q}{r^q} \leq \frac{1}{\mu_{q_1, q}} \left(\int_{Q_r(y_0)} |\nabla u_n|^{q_1} dx \right)^{\frac{q}{q_1}}.$$

By recalling the definition of H_n , these estimates entail that

$$(4.5) \quad \begin{aligned} \int_{Q_r(y_0)} H_n\left(\frac{u_n - \mathbf{m}_n}{r}\right) dx &\leq \frac{1}{\mu_{p_1, p}} \left(\int_{Q_r(y_0)} |\nabla u_n|^{p_1} dx \right)^{\frac{p}{p_1}} \\ &\quad + \frac{\alpha_n}{\mu_{q_1, q}} \left(\int_{Q_r(y_0)} |\nabla u_n|^{q_1} dx \right)^{\frac{q}{q_1}}. \end{aligned}$$

By using Hölder's inequality, the second integral can be estimated as follows

$$\left(\int_{Q_r(y_0)} |\nabla u_n|^{q_1} dx \right)^{\frac{q}{q_1}} = \left(\int_{Q_r(y_0)} (|\nabla u_n|^q)^{\frac{q_1}{q}} dx \right)^{\frac{q}{q_1}} \leq \left(\int_{Q_r(y_0)} (|\nabla u_n|^q)^{\frac{p_1}{p}} dx \right)^{\frac{p}{p_1}},$$

thanks to the fact that $p_1/p > q_1/q$. By using this fact in (4.5), we get

$$\begin{aligned} \int_{Q_r(y_0)} H_n \left(\frac{u_n - \mathbf{m}_n}{r} \right) dx &\leq \frac{1}{\mu_{p_1,p}} \left(\int_{Q_r(y_0)} (|\nabla u_n|^p)^{\frac{p_1}{p}} dx \right)^{\frac{p}{p_1}} \\ &\quad + \frac{1}{\mu_{q_1,q}} \left(\int_{Q_r(y_0)} (\alpha_n |\nabla u_n|^q)^{\frac{p_1}{p}} dx \right)^{\frac{p}{p_1}}. \end{aligned}$$

We can apply the following two elementary inequalities (here $a, b \geq 0$)

$$a^{\frac{p}{p_1}} + b^{\frac{p}{p_1}} \leq (a+b)^{\frac{p}{p_1}} \quad \text{and} \quad a^{\frac{p_1}{p}} + b^{\frac{p_1}{p}} \leq 2^{\frac{p-p_1}{p}} (a+b)^{\frac{p_1}{p}},$$

so to obtain

$$\begin{aligned} \int_{Q_r(y_0)} H_n \left(\frac{u_n - \mathbf{m}_n}{r} \right) dx &\leq 2^{\frac{p-p_1}{p_1}} \max \left\{ \frac{1}{\mu_{p_1,p}}, \frac{1}{\mu_{q_1,q}} \right\} \\ (4.6) \quad &\quad \times \left(\int_{Q_r(y_0)} [|\nabla u_n|^p + \alpha_n |\nabla u_n|^q]^{\frac{p_1}{p}} dx \right)^{\frac{p}{p_1}}. \end{aligned}$$

By using the definition of H_n and inserting the previous estimate in (4.3), we obtain

$$\int_{Q_{r/2}(y_0)} H_n(|\nabla u_n|) dx \leq C \left(\int_{Q_r(y_0)} H_n(|\nabla u_n|)^{\frac{p_1}{p}} dx \right)^{\frac{p}{p_1}} + C r^{p'} \left(\int_{Q_r(y_0)} |f|^{\gamma_0} dx \right)^{\frac{p'}{\gamma_0}}.$$

The last integral can be estimated by

$$\begin{aligned} r^{p'} \left(\int_{Q_r(y_0)} |f|^{\gamma_0} dx \right)^{\frac{p'}{\gamma_0}} &= \frac{r^{p'}}{|Q_r(y_0)|^{\frac{p'}{\gamma_0}-1}} \left(\int_{Q_r(y_0)} |f|^{\gamma_0} dx \right)^{\frac{p'}{\gamma_0}-1} \int_{Q_r(y_0)} |f|^{\gamma_0} dx \\ &\leq \frac{r^{p'}}{|Q_r(y_0)|^{\frac{p'}{\gamma_0}-1}} \int_{Q_r(y_0)} |f|^{\gamma_0} dx. \end{aligned}$$

Observe that we used (4.1), so that

$$\int_{Q_r(y_0)} |f|^{\gamma_0} dx \leq \int_{\Omega} |f|^{\gamma_0} dx = 1,$$

and the fact that $\gamma_0 \leq p'$, by Assumption 1. Moreover, by recalling that $Q_r(y_0) \subseteq Q_1(x_0)$, we have $r \leq 1$ and

$$\frac{r^{p'}}{|Q_r(y_0)|^{\frac{p'}{\gamma_0}-1}} = 2^{N-\frac{N}{\gamma_0} p'} r^{p'+N-\frac{N}{\gamma_0} p'} \leq 2^{N-\frac{N}{\gamma_0} p'},$$

thanks to the fact that the exponent of r is non-negative. We thus get (4.2), with

$$\varepsilon_1 = 1 - \frac{p_1}{p}.$$

Step 3: proof of (4.2) for “peripheral” cubes. Let us suppose that

$$Q_r(y_0) \setminus \Omega \neq \emptyset \quad \text{and} \quad y_0 \in (\mathbb{R}^N \setminus \Omega) \cup \left\{ x \in \Omega : \min_{y \in \partial\Omega} \|x - y\|_{\ell^\infty} \leq \frac{2}{3} r \right\}.$$

This means that the center y_0 of the cube falls either outside Ω or “not too far” from the boundary $\partial\Omega$. For this cube, we will apply the Caccioppoli inequality of Lemma 3.4, again with $\sigma = 1/2$. We get

$$(4.7) \quad \int_{Q_{r/2}(y_0)} H_n(|\nabla u_n|) dx \leq C \int_{Q_r(y_0)} H_n\left(\frac{u_n}{r}\right) dx + C r^{p'} \left(\int_{Q_r(x_0)} |f|^{\gamma_0} dx \right)^{\frac{p'}{\gamma_0}}.$$

We proceed similarly as before, but we now have to use the vanishing of u_n in $\mathbb{R}^N \setminus \Omega$, in order to get the desired reverse estimate. More precisely, we set

$$p_2 = \max\{p_1, p_0\} \quad \text{and} \quad q_2 = \max\{q_1, p_0\},$$

where p_1 and q_1 are the same as in (4.4). By definition, we still have

$$p_2 < p \begin{cases} \leq p_2^*, & \text{if } 1 \leq p_0 < N, \\ < \infty, & \text{if } p_0 \geq N. \end{cases}, \quad q_2 < q \begin{cases} \leq q_2^*, & \text{if } 1 \leq p_0 < N, \\ < \infty, & \text{if } p_0 \geq N. \end{cases} \quad \text{and} \quad \frac{q_2}{q} < \frac{p_2}{p}.$$

Observe in particular that $p_2 \geq p_0$ and $q_2 \geq p_0$, thus from the assumption $\Omega \in \text{UT}_{p_0}(c_0)$ and Lemma 2.6 we get that

$$\Omega \in \text{UT}_{p_2}(c_1) \cap \text{UT}_{q_2}(c_2),$$

as well, for two constants $c_1 = c_1(N, p_0, p)$ and $c_2 = c_2(N, p_0, q)$. By using these thickness conditions and Lemma 2.10, we get

$$\int_{Q_r(y_0)} \frac{|u_n|^p}{r^p} dx \leq C'_3 \left(\int_{Q_r(y_0)} |\nabla u_n|^{p_2} dx \right)^{\frac{p}{p_2}},$$

and

$$\int_{Q_r(y_0)} \frac{|u_n|^q}{r^q} \leq C''_3 \left(\int_{Q_r(y_0)} |\nabla u_n|^{q_2} dx \right)^{\frac{q}{q_2}},$$

for two constants $C'_3 = C'_3(N, p_0, p, c_0) > 0$ and $C''_3 = C''_3(N, p_0, q, c_0) > 0$. We can now repeat the arguments used in *Step 2* in order to get (4.6). This time we get

$$\int_{Q_r(y_0)} H_n\left(\frac{u_n}{r}\right) dx \leq 2^{\frac{p-p_1}{p_1}} \max\{C'_3, C''_3\} \left(\int_{Q_r(y_0)} H_n(|\nabla u_n|)^{\frac{p_2}{p}} dx \right)^{\frac{p}{p_2}}.$$

We insert this estimate in (4.7) and proceed exactly as in *Step 2*. We get again (4.2), this time with

$$\varepsilon_2 = 1 - \frac{p_2}{p}.$$

Step 4: proof of (4.2) for “suburban” cubes. Finally, we consider the case

$$Q_r(y_0) \setminus \Omega \neq \emptyset \quad \text{and} \quad y_0 \in \left\{ x \in \Omega : \min_{y \in \partial\Omega} \|x - y\|_{\ell^\infty} > \frac{2}{3} r \right\}.$$

In this case, observe that we have $Q_{(2r)/3}(y_0) \subseteq \Omega$. Thus, we can appeal again to the Caccioppoli inequality for internal cubes, i.e. Lemma 3.3, as in *Step 2*. We use it on the cube $Q_{(2r)/3}(y_0)$, by choosing $\sigma = 3/4$, so that

$$Q_{\sigma \frac{2r}{3}}(y_0) = Q_{\frac{r}{2}}(y_0).$$

This gives

$$(4.8) \quad \int_{Q_{\frac{r}{2}}(y_0)} H_n(|\nabla u_n|) dx \leq C \int_{Q_{\frac{2r}{3}}(y_0)} H_n\left(\frac{u_n - \mathbf{m}_n}{r}\right) dx + C r^{p'} \left(\int_{Q_{\frac{2r}{3}}(y_0)} |f|^{\gamma_0} dx \right)^{\frac{p'}{\gamma_0}},$$

where again $\mathbf{m}_n := \overline{u_n}_{Q_r(y_0)}$. It is easily seen that

$$\int_{Q_{\frac{2r}{3}}(y_0)} H_n\left(\frac{u_n - \mathbf{m}_n}{r}\right) dx \leq \left(\frac{3}{2}\right)^N \int_{Q_r(y_0)} H_n\left(\frac{u_n - \mathbf{m}_n}{r}\right) dx,$$

and

$$\int_{Q_{\frac{2r}{3}}(y_0)} |f|^{\gamma_0} dx \leq \left(\frac{3}{2}\right)^N \int_{Q_r(y_0)} |f|^{\gamma_0} dx.$$

Thus, from (4.8) we can obtain (4.3), possibly with a different constant C (still independent of n). The desired conclusion now follows by proceeding exactly as in *Step 2*. We then get (4.2), again with

$$\varepsilon_1 = 1 - \frac{p_1}{p}.$$

Step 5: local gain of integrability. By setting

$$\varepsilon = \max\{\varepsilon_1, \varepsilon_2\} = \max\left\{1 - \frac{p_1}{p}, 1 - \frac{p_2}{p}\right\},$$

and using Hölder inequality, we thus have established (4.2) for every cube $Q_r(y_0) \subseteq Q_1(x_0)$. We can now appeal to Lemma 2.3 and infer that $H_n(\nabla u_n) \in L^{1+\delta}(Q_{1/2}(x_0))$, for some $\delta = \delta(N, p_0, p, c_0, \gamma, g_1, g_2) > 0$ such that

$$\delta \leq \min\left\{1, \frac{\gamma - \gamma_0}{\gamma_0}\right\}.$$

Moreover, we have the estimate

$$\int_{Q_{1/2}(x_0)} H_n(\nabla u_n)^{1+\delta} dx \leq C \left[\left(\int_{Q_1(x_0)} H_n(\nabla u_n) dx \right)^{1+\delta} + \int_{Q_1(x_0)} |f|^{\gamma_0(1+\delta)} dx \right].$$

The constant C does not depend on $n \in \mathbb{N}$. In particular, we also get

$$(4.9) \quad \int_{Q_{1/2}(x_0)} |\nabla u_n|^{p(1+\delta)} dx \leq C \left[\left(\int_{Q_1(x_0)} H_n(|\nabla u_n|) dx \right)^{1+\delta} + \int_{Q_1(x_0)} |f|^{\gamma_0(1+\delta)} dx \right],$$

possibly for a different constant C , still independent of n .

Step 6: global uniform gain. We can now get a global uniform estimate on Ω for ∇u_n , by means of a covering argument. We consider the countable tiling of $\mathbb{R}^N$ given by the cubes

$$\{Q_{1/2}(\mathbf{i})\}_{\mathbf{i} \in \mathbb{Z}^N}$$

with sides parallel to the coordinate axis and with disjoint interiors. We set

$$\mathbb{Z}_\Omega^N = \{\mathbf{i} \in \mathbb{Z}^N : Q_{1/2}(\mathbf{i}) \cap \Omega \neq \emptyset\},$$

so that

$$\sum_{\mathbf{i} \in \mathbb{Z}_\Omega^N} \int_{Q_{1/2}(\mathbf{i}) \cap \Omega} |\nabla u_n|^{p(1+\delta)} dx = \int_{\Omega} |\nabla u_n|^{p(1+\delta)} dx.$$

We then double the side length of the cubes, by considering $\{Q_1(\mathbf{i})\}_{\mathbf{i} \in \mathbb{Z}_\Omega^N}$: observe that this gives a locally finite covering of Ω . More precisely, for a fixed $\mathbf{i} \in \mathbb{Z}_\Omega^N$ we have

$$\#\{\mathbf{j} \in \mathbb{Z}_\Omega^N : Q_1(\mathbf{i}) \cap Q_1(\mathbf{j})\} \leq \#\{\mathbf{j} \in \mathbb{Z}^N : Q_1(\mathbf{i}) \cap Q_1(\mathbf{j})\} = 3^N - 1.$$

In other words, when we double the length of the side of the tiling, we get a covering such that each cube intersects only a (universally) finite number of other cubes. This is needed in order to trigger the covering argument. In particular, we can rearrange the family $\{Q_1(\mathbf{i})\}_{\mathbf{i} \in \mathbb{Z}_\Omega^N}$ into 2^N subfamilies $\mathfrak{Q}_1, \dots, \mathfrak{Q}_{2^N}$ satisfying

$$Q_1(\mathbf{i}) \cap Q_1(\mathbf{j}) = \emptyset, \quad \text{if } Q_1(\mathbf{i}) \neq Q_1(\mathbf{j}) \in \mathfrak{Q}_k, \quad \text{with } 1 \leq k \leq 2^N.$$

From this observation and estimate (4.9), we can infer

$$\begin{aligned} \int_{\Omega} |\nabla u_n|^{p(1+\delta)} dx &= \sum_{\mathbf{i} \in \mathbb{Z}_\Omega^N} \int_{Q_{1/2}(\mathbf{i}) \cap \Omega} |\nabla u_n|^{p(1+\delta)} dx \\ (4.10) \quad &\leq C \sum_{\mathbf{i} \in \mathbb{Z}_\Omega^N} \left(\int_{Q_1(\mathbf{i}) \cap \Omega} H_n(|\nabla u_n|) dx \right)^{1+\delta} \\ &\quad + C \sum_{\mathbf{i} \in \mathbb{Z}_\Omega^N} \int_{Q_1(\mathbf{i}) \cap \Omega} |f|^{\gamma_0(1+\delta)} dx. \end{aligned}$$

For the last term, we simply observe that

$$\begin{aligned} \sum_{\mathbf{i} \in \mathbb{Z}_\Omega^N} \int_{Q_1(\mathbf{i}) \cap \Omega} |f|^{\gamma_0(1+\delta)} dx &= \sum_{j=1}^{2^N} \left(\sum_{\{\mathbf{i} \in \mathbb{Z}_\Omega^N : Q_1(\mathbf{i}) \in \mathfrak{Q}_j\}} \int_{Q_1(\mathbf{i}) \cap \Omega} |f|^{\gamma_0(1+\delta)} dx \right) \\ &\leq 2^N \int_{\Omega} |f|^{\gamma_0(1+\delta)} dx, \end{aligned}$$

where we used the fact that the cubes in each $\mathfrak{Q}_j$ have disjoint interiors. We can further estimate the last integral, by using (4.1) and interpolation in Lebesgue spaces: since $\gamma_0 < \gamma_0(1+\delta) \leq \gamma$, we get

$$\int_{\Omega} |f|^{\gamma_0(1+\delta)} dx \leq \left(\int_{\Omega} |f|^{\gamma} dx \right)^{\frac{\delta \gamma_0}{\gamma - \gamma_0}}.$$

In a similar way, by using the super-additivity of convex powers, we get

$$\begin{aligned} \sum_{\mathbf{i} \in \mathbb{Z}_\Omega^N} \left(\int_{Q_1(\mathbf{i}) \cap \Omega} H_n(|\nabla u_n|) dx \right)^{1+\delta} &= \sum_{j=1}^{2^N} \sum_{\{\mathbf{i} \in \mathbb{Z}_\Omega^N : Q_1(\mathbf{i}) \in \mathfrak{Q}_j\}} \left(\int_{Q_1(\mathbf{i})} H_n(|\nabla u_n|) dx \right)^{1+\delta} \\ &\leq 2^N \left(\int_{\Omega} H_n(|\nabla u_n|) dx \right)^{1+\delta}. \end{aligned}$$

Thus, from (4.10) we get

$$(4.11) \quad \int_{\Omega} |\nabla u_n|^{p(1+\delta)} dx \leq C \left(\int_{\Omega} H_n(|\nabla u_n|) dx \right)^{1+\delta} + C \left(\int_{\Omega} |f|^{\gamma} dx \right)^{\frac{\delta \gamma_0}{\gamma - \gamma_0}}.$$

Step 7: conclusion. We now wish to pass to the limit as n goes to ∞ . Observe that $p < p(1+\delta) \leq q$: the second inequality follows by recalling that $q = 2p$ (see the definition (3.1)) and using that

$$\delta \leq \min \left\{ 1, \frac{\gamma - \gamma_0}{\gamma_0} \right\}.$$

Thus, from $u_n \in W_0^{1,p}(\Omega) \cap W_0^{1,q}(\Omega)$ we get that $u_n \in W_0^{1,p(1+\delta)}(\Omega)$, as well. Moreover, by recalling (3.5), we have that the right-hand side of (4.11) is uniformly bounded in n . Thus, we get that

$$\{u_n\}_{n \in \mathbb{N}} \subseteq W_0^{1,p(1+\delta)}(\Omega) \quad \text{with} \quad \sup_{n \in \mathbb{N}} \|u_n\|_{W^{1,p(1+\delta)}(\Omega)} < +\infty.$$

By reflexivity, we know that $\{u_n\}_{n \in \mathbb{N}}$ weakly converges in $W^{1,p(1+\delta)}(\Omega)$ to some function v , up to a subsequence. Moreover, since $W_0^{1,p(1+\delta)}(\Omega)$ is weakly closed, we get that $v \in W_0^{1,p(1+\delta)}(\Omega)$. By uniqueness of the limit, from Proposition 3.2 we can thus conclude

$$u = v \in W_0^{1,p(1+\delta)}(\Omega),$$

which establishes the desired regularity result on u , under the condition (4.1).

This result comes with the relevant a priori estimate: by taking the limit as n goes to ∞ in (4.11), using the lower semicontinuity of the left-hand side and (3.5) in the right-hand side, we finally get

$$(4.12) \quad \int_{\Omega} |\nabla u|^{p(1+\delta)} dx \leq C \left(\int_{\Omega} |\nabla u|^p dx \right)^{1+\delta} + C \left(\int_{\Omega} |f|^{\gamma} dx \right)^{\frac{\delta \gamma_0}{\gamma - \gamma_0}},$$

possibly for a different constant $C = C(N, p, \gamma_0, \gamma, \delta, g_1, g_2) > 0$.

4.2. Part 2: estimate at a general scale. We now remove the restriction (4.1) and make the estimate scale invariant.

Step 1: vertical scalings. Given the solution $u \in W_0^{1,p}(\Omega)$, we set

$$u_{\lambda}(x) = \lambda u(x), \quad \text{with } \lambda = \|f\|_{L^{\gamma_0}(\Omega)}^{\frac{1}{1-p}}.$$

It is not difficult to see that $u_{\lambda} \in W_0^{1,p}(\Omega)$ is the weak solution of

$$-\operatorname{div}(\nabla G_{\lambda}(\nabla u_{\lambda})) = f_{\lambda}, \quad \text{in } \Omega,$$

where

$$(4.13) \quad G_{\lambda}(z) = \lambda^p G\left(\frac{z}{\lambda}\right) \quad \text{and} \quad f_{\lambda} = \lambda^{p-1} f.$$

Observe that G_{λ} is still C^1 strictly convex and satisfies

$$g_1 |z|^p \leq G_{\lambda}(z) \leq g_2 |z|^p,$$

i.e. the constants entering in the growth conditions are still the same g_1 and g_2 . Moreover, we have

$$\|f_{\lambda}\|_{L^{\gamma_0}(\Omega)} = 1.$$

We can thus apply the result of **Part 1** to u_λ and obtain $u_\lambda \in W_0^{1,p(1+\delta)}(\Omega)$, together with the estimate (4.12). By recalling the definition of u_λ and f_λ , we get $u \in W_0^{1,p(1+\delta)}(\Omega)$ with the estimate

$$\int_{\Omega} |\nabla u|^{p(1+\delta)} dx \leq C \left[\left(\int_{\Omega} |\nabla u|^p dx \right)^{1+\delta} + \lambda^{\gamma(p-1) \frac{\delta \gamma_0}{\gamma - \gamma_0} - p(1+\delta)} (\|f\|_{L^\gamma(\Omega)})^{\frac{\delta \gamma_0 \gamma}{\gamma - \gamma_0}} \right].$$

Step 2: horizontal scalings. We finally adjust the a priori estimate, by making it scale invariant with respect to horizontal scalings, as well. We fix $r > 0$, given the solution $u \in W_0^{1,p}(\Omega)$, we set

$$u_r(x) = u(rx).$$

As before, we see that $u_r \in W_0^{1,p}(\Omega/r)$ is the weak solution of

$$-\operatorname{div}(\nabla G_r(\nabla u_r)) = \tilde{f}_r, \quad \text{in } \frac{1}{r} \Omega,$$

where G_r is still given by (4.13), with r in place of λ , and

$$\tilde{f}_r(x) = r^p f(rx).$$

We notice that Ω/r is still uniformly p_0 -thick, with the same constant c_0 , thanks to Lemma 2.8 (used with $r_0 = +\infty$). Thus, we can appeal to the previous step and get $u_r \in W_0^{1,p(1+\delta)}(\Omega/r)$, with the estimate

$$\begin{aligned} \left(\int_{\Omega/r} |\nabla u_r|^{p(1+\delta)} dx \right)^{\frac{1}{1+\delta}} &\leq C \int_{\Omega/r} |\nabla u_r|^p dx \\ &\quad + C \left(\|\tilde{f}_r\|_{L^{\gamma_0}(\Omega/r)} \right)^{p' - \frac{\delta}{1+\delta} \frac{\gamma_0 \gamma}{\gamma - \gamma_0}} \left(\|\tilde{f}_r\|_{L^\gamma(\Omega/r)} \right)^{\frac{\delta}{1+\delta} \frac{\gamma_0 \gamma}{\gamma - \gamma_0}}. \end{aligned}$$

By recalling the definitions of u_r and $\tilde{f}_r$, we can scale back this estimate and get the estimate for every $r > 0$

$$\begin{aligned} \left(\frac{1}{r^{N-p(1+\delta)}} \int_{\Omega} |\nabla u|^{p(1+\delta)} dx \right)^{\frac{1}{1+\delta}} &\leq \frac{C}{r^{N-p}} \int_{\Omega} |\nabla u|^p dx \\ (4.14) \quad &\quad + C \left(\frac{1}{r^{N-p\gamma_0}} \|f\|_{L^{\gamma_0}(\Omega)}^{\gamma_0} \right)^{\frac{p'}{\gamma_0} - \frac{\delta}{1+\delta} \frac{\gamma}{\gamma - \gamma_0}} \\ &\quad \times \left(\frac{1}{r^{N-p\gamma}} \|f\|_{L^\gamma(\Omega)}^\gamma \right)^{\frac{\delta}{1+\delta} \frac{\gamma_0}{\gamma - \gamma_0}}. \end{aligned}$$

We multiply both sides of (4.14) by

$$r^{\frac{N}{1+\delta} - p},$$

so to get (after some tedious, yet elementary, algebraic computations)

$$\begin{aligned} (4.15) \quad \left(\int_{\Omega} |\nabla u|^{p(1+\delta)} dx \right)^{\frac{1}{1+\delta}} &\leq \frac{C_1}{r^{N \frac{\delta}{1+\delta}}} \int_{\Omega} |\nabla u|^p dx \\ &\quad + C_1 r^{p' - \frac{N}{\gamma_0} (p' - \gamma_0)} (\|f\|_{L^{\gamma_0}(\Omega)})^{p' - \frac{\delta}{1+\delta} \frac{\gamma \gamma_0}{\gamma - \gamma_0}} (\|f\|_{L^\gamma(\Omega)})^{\frac{\delta}{1+\delta} \frac{\gamma \gamma_0}{\gamma - \gamma_0}}, \end{aligned}$$

which holds for every $r > 0$, where $C_1 = C_1(N, p, p_0, g_1, g_2, \gamma_0, \gamma, \delta) > 0$. We now want to optimize this estimate with respect to the scale parameter $r > 0$. At first, we notice that we can rewrite the second exponent on r as follows

$$\begin{aligned} p' - \frac{N}{\gamma_0} (p' - \gamma_0) &= p' N \left(\frac{1}{N} - \frac{1}{\gamma_0} + \frac{1}{p'} \right) \\ &= p' N \left(\frac{1}{N} - \frac{1}{\gamma_0} + 1 - \frac{1}{p} \right) = p' N \left(\frac{1}{\gamma_0'} - \frac{N-p}{Np} \right). \end{aligned}$$

By recalling the assumption on γ_0 , it is not difficult to see that the latter is positive. Thus, the right-hand side of (4.15) has the form

$$\frac{A}{r^\alpha} + B r^\beta, \quad \text{with } A, B, \alpha, \beta > 0.$$

The previous quantity is minimal for the choice

$$r = \left(\frac{\alpha}{\beta} \frac{A}{B} \right)^{\frac{1}{\alpha+\beta}}.$$

By making this choice in (4.15), we can get the desired scale invariant estimate and conclude the proof.

5. EXAMPLE: CONTRACTIBLE SETS

We want to complement the Main Theorem by discussing a sufficient topological condition ensuring that $\mathbb{R}^N \setminus \Omega$ is uniformly p -thick. Namely, we show that contractible sets satisfies this property, provided $p > N - 1$. This fact should be already known, but it is not always easy to find a reference containing a proof. For this reason, for completeness we provide it: in this way, one can also get a quantitative information about the relevant constant c_0 of p -thickness. We will also show that for contractible open sets, the analytic condition (E_Ω) admits a simple geometric equivalent reformulation.

5.1. Thickness of the complement. The first step is the following condition implying p -thickness of the complement. This is a consequence of [28, Theorem 3.3] and Lemma 2.5. We give the details needed in order to keep track of the relevant constant.

Lemma 5.1. *Let $1 \leq p < \infty$ and let $\Omega \subsetneq \mathbb{R}^N$ be an open set. Let us suppose that there exists a constant $\mathcal{C} > 0$ and a radius $r_0 > 0$ such that*

$$(5.1) \quad \int_{B_r(x_0)} |\varphi|^p dx \leq \mathcal{C} r^p \int_{B_r(x_0)} |\nabla \varphi|^p dx, \quad \text{for every } x_0 \in \mathbb{R}^N \setminus \Omega, \ 0 < r \leq r_0, \ \varphi \in C_0^\infty(\Omega).$$

Then $\Omega \in \text{UT}_p(\tilde{c}_0, r_0)$, with $\tilde{c}_0$ given by

$$\tilde{c}_0 = \frac{1}{\beta_2} \frac{\omega_N}{2^{p+N}} \min \left\{ \lambda_p(B_2), \frac{2^p}{\mathcal{C}} \right\} \frac{1}{\text{cap}_p(\overline{Q_1}; Q_2)},$$

and $\beta_2 = \beta_2(N, p) > 0$ is the same constant as in Lemma 2.5.

Proof. According to [28, Theorem 3.3], the assumption on Ω is *equivalent* to the fact that there exists $c_0 > 0$ such that

$$(5.2) \quad \text{cap}_p \left(\overline{B_r(x_0)} \setminus \Omega; B_{2r}(x_0) \right) \geq c_0 r^{N-p}, \quad \text{for every } x_0 \in \mathbb{R}^N \setminus \Omega, \ 0 < r \leq r_0.$$

An inspection of the proof of [28, Theorem 3.3] informs us that we can take

$$c_0 = \frac{\omega_N}{2^{p+N}} \min \left\{ \lambda_p(B_2), \frac{2^p}{C} \right\}.$$

If we now set

$$\tilde{c}_0 = \frac{1}{\beta_2} \frac{c_0}{\text{cap}_p(\overline{Q_1}; Q_2)},$$

use the leftmost inequality in Lemma 2.5 and then observe that

$$c_0 r^{N-p} = \beta_2 \tilde{c}_0 r^{N-p} \text{cap}_p(\overline{Q_1}; Q_2) = \beta_2 \tilde{c}_0 \text{cap}_p(\overline{Q_r(x_0)}; Q_{2r}(x_0)),$$

from the previous estimate (5.2) we get the desired conclusion. $\square$

We now prove the claimed result for contractible sets. We start with the two-dimensional case $N = 2$, which is somehow exceptional. In this case, we can even include the limit case $p = N - 1 = 1$. This is due to the fact that the following one-dimensional sharp Poincaré constant for the interval $I = (0, 1)$

$$\pi_p = \inf_{\varphi \in C_0^\infty(I)} \{ \|\varphi'\|_{L^p(I)} : \|\varphi\|_{L^p(I)} = 1 \},$$

is positive for every $1 \leq p \leq \infty$. We have the following

Proposition 5.2. *Let $\Omega \subsetneq \mathbb{R}^2$ be an open simply connected set. Then, for every $1 \leq p < \infty$ there exists a constant $c_p > 0$ such that*

$$\Omega \in \text{UT}_p(c_p).$$

Proof. In light of Lemma 5.1, we just have to prove the validity of the boundary Poincaré inequality. We will use the argument by Hayman contained in [23, Lemma 1]. Let $\varphi \in C_0^\infty(\Omega)$ and let $x_0 \in \mathbb{R}^2 \setminus \Omega$. We observe that for every $\varrho > 0$, we must have

$$\partial B_\varrho(x_0) \cap (\mathbb{R}^2 \setminus \Omega) \neq \emptyset.$$

Indeed, if $B_\varrho(x_0) \subseteq \Omega$ this would violate the topological property of Ω , since the loop $\gamma(t) = x_0 + (\varrho \cos t, \varrho \sin t)$ could not be null-homotopic. By using polar coordinates and a slight abuse of notation, we thus get that for every $\varrho > 0$ there exists $\vartheta_\varrho \in [0, 2\pi)$ such that

$$\varphi(\varrho, \vartheta_\varrho) = 0.$$

We can thus apply the one-dimensional Poincaré inequality for the function $\vartheta \mapsto \varphi(\varrho, \vartheta)$, so to get

$$\left(\frac{\pi_p}{2\pi}\right)^p \int_0^{2\pi} |\varphi(\varrho, \vartheta)|^p d\vartheta \leq \int_0^{2\pi} |\partial_\vartheta \varphi(\varrho, \vartheta)|^p d\vartheta.$$

We multiply both sides by ϱ and then integrate over the interval $[0, r]$: this yields

$$\begin{aligned} \left(\frac{\pi_p}{2\pi}\right)^p \iint_{[0, r] \times [0, 2\pi]} |\varphi(\varrho, \vartheta)|^p \varrho d\vartheta d\varrho &\leq \iint_{[0, r] \times [0, 2\pi]} |\partial_\vartheta \varphi(\varrho, \vartheta)|^p \varrho d\vartheta d\varrho \\ &\leq \iint_{[0, r] \times [0, 2\pi]} \left[|\partial_\varrho \varphi|^2 + \frac{1}{\varrho^2} |\partial_\vartheta \varphi(\varrho, \vartheta)|^2 \right]^{\frac{p}{2}} \varrho^{p+1} d\vartheta d\varrho \\ &\leq r^p \iint_{[0, r] \times [0, 2\pi]} \left[|\partial_\varrho \varphi|^2 + \frac{1}{\varrho^2} |\partial_\vartheta \varphi(\varrho, \vartheta)|^2 \right]^{\frac{p}{2}} \varrho d\vartheta d\varrho \end{aligned}$$

This proves that

$$\left(\frac{\pi_p}{2\pi}\right)^p \frac{1}{r^p} \int_{B_r(x_0)} |\varphi|^p dx \leq \int_{B_r(x_0)} |\nabla \varphi|^p dx, \quad \text{for every } x_0 \in \mathbb{R}^2 \setminus \Omega, \varphi \in C_0^\infty(\Omega),$$

as desired. $\square$

The previous result can be generalized to contractible sets in every dimension, as announced above. However, in this case we need to have the strict inequality $p > N - 1$.

Proposition 5.3. *Let $N \geq 3$ and let $\Omega \subsetneq \mathbb{R}^N$ be a contractible open set. Then, for every $N - 1 < p < \infty$ there exists a constant $c_{N,p} > 0$ such that*

$$\Omega \in \text{UT}_p(c_{N,p}).$$

Proof. As before, the goal is to show the validity of the boundary Poincaré inequality (5.1). We still use Hayman's idea, as in the two-dimensional case. Fix $x_0 \in \mathbb{R}^N \setminus \Omega$ and $\varphi \in C_0^\infty(\Omega)$. For every $\varrho > 0$ we have

$$\partial B_\varrho(x_0) \cap (\mathbb{R}^N \setminus \Omega) \neq \emptyset.$$

Indeed, if $\partial B_\varrho(x_0) \subseteq \Omega$ one could define the continuous map

$$\begin{aligned} f : \mathbb{S}^{N-1} &\rightarrow \Omega \\ \omega &\mapsto x_0 + \varrho \omega \end{aligned}$$

and this could not be null-homotopic, since $x_0 \notin \Omega$. Thus, the $(N - 1)$ -th homotopy group would be non-trivial, by violating the fact that Ω is contractible.

The previous observation guarantees that for every $\varrho > 0$ there exists $\omega_\varrho \in \mathbb{S}^{N-1}$ such that $x_0 + \varrho \omega_\varrho \notin \Omega$. Thus, we have $\varphi(x_0 + \varrho \omega_\varrho) = 0$ and we can apply the Poincaré inequality of Lemma 5.4 below, to the function $\omega \mapsto \varphi(x_0 + \varrho \omega)$. We then obtain

$$\int_{\mathbb{S}^{N-1}} |\varphi(x_0 + \varrho \omega)|^p d\sigma(\omega) \leq C \int_{\mathbb{S}^{N-1}} |\nabla_\tau \varphi(x_0 + \varrho \omega)|^p d\sigma(\omega).$$

Here, ∇_τ is the tangential gradient and $d\sigma$ is the standard surface measure on $\mathbb{S}^{N-1}$. If we multiply both sides by ϱ^{N-1} and integrate with respect to $\varrho \in [0, r]$, we get

$$\begin{aligned} \int_0^r \int_{\mathbb{S}^{N-1}} |\varphi(x_0 + \varrho \omega)|^p \varrho^{N-1} d\sigma(\omega) d\varrho &\leq C \int_0^r \int_{\mathbb{S}^{N-1}} |\nabla_\tau \varphi(x_0 + \varrho \omega)|^p \varrho^{N-1} d\sigma(\omega) d\varrho \\ &\leq C \int_0^r \int_{\mathbb{S}^{N-1}} \left[|\partial_\varrho \varphi|^2 + \frac{1}{\varrho^2} |\nabla_\tau \varphi|^2 \right]^{\frac{p}{2}} \varrho^{N+p-1} d\varrho d\sigma(\omega) \\ &\leq C r^p \int_0^r \int_{\mathbb{S}^{N-1}} \left[|\partial_\varrho \varphi|^2 + \frac{1}{\varrho^2} |\nabla_\tau \varphi|^2 \right]^{\frac{p}{2}} \varrho^{N-1} d\varrho d\sigma(\omega). \end{aligned}$$

This finally proves that

$$\frac{1}{C r^p} \int_{B_r(x_0)} |\varphi|^p dx \leq \int_{B_r(x_0)} |\nabla \varphi|^p dx, \quad \text{for every } x_0 \in \mathbb{R}^N \setminus \Omega, \varphi \in C_0^\infty(\Omega),$$

as desired. $\square$

We complement this section with the following Poincaré inequality on the sphere, for functions vanishing at a point: it has been an essential ingredient of the previous proof.

Lemma 5.4. *Let $N \geq 3$ and let $N - 1 < p < \infty$. There exists a constant $C = C(N, p) > 0$ such that for every $x_0 \in \mathbb{S}^{N-1}$ and every $\varphi \in C^\infty(\mathbb{S}^{N-1})$ with $\varphi(x_0) = 0$, we have*

$$\int_{\mathbb{S}^{N-1}} |\varphi|^p d\sigma \leq C \int_{\mathbb{S}^{N-1}} |\nabla_\tau \varphi|^p d\sigma.$$

Proof. Thanks to the fact that $p > N - 1$, we have the Sobolev-Morrey continuous embedding

$$W^{1,p}(\mathbb{S}^{N-1}) \hookrightarrow C^{0,1-\frac{N-1}{p}}(\mathbb{S}^{N-1}),$$

see [24, Theorem 2.8]. Moreover, by the same result, for every $\varphi \in C^\infty(\mathbb{S}^{N-1})$ we have

$$[\varphi]_{C^{0,1-\frac{N-1}{p}}(\mathbb{S}^{N-1})} := \frac{|\varphi(x) - \varphi(y)|}{d_{\mathbb{S}^{N-1}}(x, y)^{1-\frac{N-1}{p}}} \leq C \left(\int_{\mathbb{S}^{N-1}} |\nabla_\tau \varphi|^p d\sigma + \int_{\mathbb{S}^{N-1}} |\varphi|^p d\sigma \right)^{\frac{1}{p}},$$

for a constant $C = C(N, p) > 0$. Here $d_{\mathbb{S}^{N-1}}$ is the geodesic distance on $\mathbb{S}^{N-1}$, computed with respect to the standard round metric, i.e. the Riemannian metric induced by the standard Euclidean one. If we set

$$\bar{\varphi} = \int_{\mathbb{S}^{N-1}} \varphi d\sigma,$$

we can apply the previous inequality to the smooth function $\varphi - \bar{\varphi}$ and obtain

$$[\varphi]_{C^{0,1-\frac{N-1}{p}}(\mathbb{S}^{N-1})} \leq C \left(\int_{\mathbb{S}^{N-1}} |\nabla_\tau \varphi|^p d\sigma + \int_{\mathbb{S}^{N-1}} |\varphi - \bar{\varphi}|^p d\sigma \right)^{\frac{1}{p}}.$$

We can estimate the rightmost term by means of the Poincaré-Wirtinger inequality on $\mathbb{S}^{N-1}$ (see [24, Theorem 2.10]), so to get

$$[\varphi]_{C^{0,1-\frac{N-1}{p}}(\mathbb{S}^{N-1})} \leq C \left(\int_{\mathbb{S}^{N-1}} |\nabla_\tau \varphi|^p d\sigma \right)^{\frac{1}{p}},$$

possibly for a different constant $C = C(N, p) > 0$. The previous Morrey-type inequality is valid for every $\varphi \in C^\infty(\mathbb{S}^{N-1})$. We now fix a point $x_0 \in \mathbb{S}^{N-1}$ and use it for a smooth function vanishing at x_0 . We get in particular

$$|\varphi(x)| = |\varphi(x) - \varphi(x_0)| \leq C d_{\mathbb{S}^{N-1}}(x, x_0)^{1-\frac{N-1}{p}} \left(\int_{\mathbb{S}^{N-1}} |\nabla_\tau \varphi|^p d\sigma \right)^{\frac{1}{p}}.$$

By using that the geodesic distance on $\mathbb{S}^{N-1}$ is bounded, we can raise the previous estimate to the power p and integrate over $\mathbb{S}^{N-1}$. This gives the desired conclusion. $\square$

5.2. Poincaré inequality Vs. inradius. Finally, we conclude this section by giving an equivalent condition for (E_Ω) to hold, in the case of contractible sets. We recall the definition of inradius for an open set

$$r_\Omega = \sup \left\{ r > 0 : \exists B_r(x_0) \subseteq \Omega \right\}.$$

We have that contractible open subsets of $\mathbb{R}^N$ supports the L^p Poincaré inequality if and only if they have finite inradius, for p large enough. This characterization is interesting in itself and will follow from the results of Appendix A.

Proposition 5.5. *Let $N - 1 < p < \infty$ and let $\Omega \subsetneq \mathbb{R}^N$ be an open contractible set. Then*

$$\lambda_p(\Omega) > 0 \quad \Longleftrightarrow \quad r_\Omega < +\infty.$$

Proof. The implication $\implies$ is valid for every open set, as previously recalled in Proposition 2.1.

For the converse implication, we first observe that the case $p > N$ is contained in [11, Theorem 5.4]. In this case the result actually holds *for every open set*, without any further assumption.

We thus consider the case $N - 1 < p \leq N$. By Propositions 5.2 and 5.3, we have that $\Omega \in \text{UT}_p(c_{N,p})$. Then the desired conclusion follows directly from Theorem A.2 below. $\square$

Remark 5.6. The case $N = 2$ of the previous result was already contained in [6, Theorem 1.2], by means of a different proof. Actually, the latter provides the stronger result

$$c_p \left(\frac{1}{r_\Omega} \right)^p \leq \lambda_p(\Omega),$$

which gives a *quantitative* version of the implication $\Leftarrow$. Observe that in the two-dimensional case, we can even allow the choice $p = N - 1 = 1$. Finally, we recall that for $N \geq 3$ and $p \leq N - 1$, the equivalence stated in Proposition 5.5 cease to be true (see for example the Introduction of [6] for more details).

6. REGULARITY FOR THE LANE-EMDEN EQUATION

We specialize the main result of the paper to the case of weak solutions of the equation

$$(6.1) \quad -\Delta_p u = \lambda |u|^{m-2} u, \quad \text{in } \Omega,$$

for some $\lambda > 0$ and $m \geq p$. We will not take any further assumption on the open sets, apart for those of the Main Theorem. The following result will be useful, in order to get the desired regularity estimates in scale invariant form.

Lemma 6.1. *Let $1 < p, m < \infty$, let $\Omega \subseteq \mathbb{R}^N$ be an open set and let $u \in W_0^{1,p}(\Omega) \cap L^m(\Omega)$ be a non-trivial weak solution of (6.1), for some $\lambda > 0$. There exists $R > 0$ such that the scaled function*

$$v(x) := \frac{R^{-\frac{N}{m}}}{\|u\|_{L^m(\Omega)}} u\left(\frac{x}{R}\right), \quad \text{for } x \in R\Omega,$$

is a weak solution of (6.1) in $R\Omega$, with $\lambda = 1$. Moreover, we have

$$\int_{R\Omega} |v|^m dx = 1.$$

Proof. It is sufficient to define

$$(6.2) \quad R = \left(\lambda^m \left(\int_{\Omega} |u|^m dx \right)^{m-p} \right)^{\frac{1}{p m + N(p-m)}}.$$

Then the proof is by direct verification and thus omitted. $\square$

We now come to the higher integrability result for solutions of (6.1). For the sake of clarify, it will be convenient to distinguish the two cases: $m > p$ and $p = m$. We stress that we are here particularly interested in the precise form of the global a priori estimates.

6.1. Super-homogeneous case. We start with the case $m > p$.

Corollary 6.2. *Let $1 < p < m$, where m verifies*

$$(6.3) \quad \begin{cases} m < p^* & \text{if } 1 < p < N, \\ m < \infty, & \text{if } p \geq N. \end{cases}$$

Let $\Omega \subseteq \mathbb{R}^N$ be an open set satisfying Assumptions 1 and 2. There exists an exponent $\delta = \delta(N, p, p_0, c_0, m) > 0$ such that every weak solution $u \in W_0^{1,p}(\Omega)$ of (6.1) satisfies

$$u \in W_0^{1,p(1+\delta)}(\Omega).$$

Moreover, by denoting with

$$\sigma := \frac{pm + N(p-m)}{m}, \quad \xi := 1 + \frac{N}{m} - \frac{N}{p(1+\delta)},$$

the following estimate holds

$$\|\nabla u\|_{L^{p(1+\delta)}(\Omega)} \leq C \lambda^{\frac{\xi}{\sigma}} \left(\|u\|_{L^m(\Omega)} \right)^{1+(m-p)\frac{\xi}{\sigma}},$$

for a constant $C = C(N, p, p_0, c_0, m, \delta) > 0$.

Proof. Let $u \in W_0^{1,p}(\Omega)$ be a solution of (6.1) and consider the scaled function

$$v(x) = \frac{R^{-\frac{N}{m}}}{\|u\|_{L^m(\Omega)}} u\left(\frac{x}{R}\right), \quad \text{for } x \in R\Omega,$$

of Lemma 6.1. We now want to apply the Main Theorem to the function v , which is such that

$$-\Delta_p v = |v|^{m-2} v, \quad \text{in } R\Omega, \quad \|v\|_{L^m(R\Omega)} = 1.$$

We thus apply it to the scaled open set $R\Omega$, with

$$f = |v|^{m-2} v, \quad \gamma_0 = m' \quad \text{and} \quad \gamma = +\infty.$$

Indeed, observe that

$$\|f\|_{L^{m'}(R\Omega)}^{m'} = \int_{R\Omega} |v|^m dx = 1 \quad \text{and} \quad \|f\|_{L^\infty(R\Omega)} = \|v\|_{L^\infty(R\Omega)}^{m-1} \leq C,$$

for a constant $C = C(N, p, m) > 0$. The second fact follows from the L^∞ estimate for solutions of the super-homogeneous Lane-Emden equation (see for example [9, Proposition 2.4] and [8, Lemma 2.3]) and the normalization conditions on v .

As for the needed assumptions on the open set $R\Omega$, we observe that, in view of Lemma 2.8, its complement is still uniformly p_0 -thick with the same constant c_0 as Ω . Moreover, the assumption $\lambda_p(\Omega) > 0$ implies that $\lambda_p(R\Omega) > 0$, as well. Therefore, we can apply the Main Theorem and obtain existence of $\delta = \delta(N, p, p_0, c_0, m) > 0$ such that $v \in W_0^{1,p(1+\delta)}(R\Omega)$, with the estimate

$$\|\nabla v\|_{L^{p(1+\delta)}(R\Omega)} \leq C_1, \quad \text{where } C_1 = C_1(N, p, p_0, c_0, \delta) > 0.$$

By recalling the definition of v and scaling back we obtain $u \in W_0^{1,p(1+\delta)}(\Omega)$, with the estimate

$$\|\nabla u\|_{L^{p(1+\delta)}(\Omega)} \leq C_1 R^{1+\frac{N}{m}-\frac{N}{p(1+\delta)}} \|u\|_{L^m(\Omega)}.$$

Finally, if we use the definition (6.2) of R , we get the desired conclusion. $\square$

6.2. Eigenfunctions of the p -Laplacian. In the homogeneous case $m = p$, the equation (6.1) turns out to be

$$(6.4) \quad -\Delta_p u = \lambda |u|^{p-2} u, \quad \text{in } \Omega,$$

i.e. the so-called eigenvalue equation for the p -Laplacian (see for example [31]). We have the following result, which just follows by repeating the proof of Corollary 6.2 with $m = p$ and using the classical L^∞ estimate for eigenfunctions

$$\|u\|_{L^\infty(\Omega)} \leq C_{N,p} \lambda^{\frac{N}{p^2}} \|u\|_{L^p(\Omega)}.$$

Corollary 6.3. *Let $1 < p < \infty$ and let $\Omega \subseteq \mathbb{R}^N$ be an open set satisfying Assumptions 1 and 2. There exists an exponent $\delta = \delta(N, p, p_0, c_0) > 0$ such that every weak solution $u \in W_0^{1,p}(\Omega)$ of (6.4) satisfies*

$$u \in W_0^{1,p(1+\delta)}(\Omega).$$

Moreover, the following estimate holds

$$\|\nabla u\|_{L^{p(1+\delta)}(\Omega)} \leq C \lambda^{\frac{p+\delta(N+p)}{p^2(1+\delta)}} \|u\|_{L^p(\Omega)},$$

for a constant $C = C(N, p, p_0, c_0, \delta) > 0$.

At this point, the Theorem for the eigenfunctions of the Dirichler-Laplacian, stated in the Introduction, simply follows from the previous results. More generally, we can obtain

Corollary 6.4 (Eigenfunctions in the conformal case). *There exists a universal exponent $\delta = \delta(N) > 0$ such that for every contractible open set $\Omega \subseteq \mathbb{R}^N$ with $r_\Omega < +\infty$ and every Dirichlet eigenfunction $u \in W_0^{1,N}(\Omega) \setminus \{0\}$ of the N -Laplacian with eigenvalue λ , we have*

$$u \in W_0^{1,N(1+\delta)}(\Omega) \quad \text{and} \quad u \in C_0^{0, \frac{\delta}{1+\delta}}(\overline{\Omega}).$$

Moreover, it holds

$$\|\nabla u\|_{L^{N(1+\delta)}(\Omega)} \leq C \left(\sqrt[N]{\lambda} \right)^{\frac{1+2\delta}{1+\delta}} \|u\|_{L^N(\Omega)},$$

and

$$|u|_{C_0^{0, \frac{\delta}{1+\delta}}(\overline{\Omega})} \leq C \left(\sqrt[N]{\lambda} \right)^{\frac{1+2\delta}{1+\delta}} \|u\|_{L^N(\Omega)},$$

for a universal constant $C = C(N) > 0$.

Proof. By using Propositions 5.2, 5.3 and 5.5, we know that Ω satisfies Assumptions 1 and 2. We can choose for example

$$p_0 = N - \frac{1}{2}.$$

Accordingly, we have $\Omega \in \text{UT}_{p_0}(c_0)$, with c_0 depending on the dimension N , only. We can then apply Corollary 6.3 with $N = p$, so to get $u \in W_0^{1,N(1+\delta)}(\Omega)$ with the associated estimate. The global Hölder continuity follows from the Sobolev-Morrey-embedding. Observe that the relevant embedding constant only depends on N and δ (and the latter in turn is only determined by N). $\square$

APPENDIX A. THICK COMPLEMENT AND POINCARÉ INEQUALITY

In this appendix, we will prove some properties of open sets with uniform thick complement, which are interesting in themselves.

We start with a simple result, assuring that we can always enlarge the radius r_0 in the thickness condition, by possibly spoiling the constant c_0 .

Lemma A.1. *Let $1 \leq p_0 < \infty$ and let $\Omega \subsetneq \mathbb{R}^N$ be an open set such that $\Omega \in \text{UT}_{p_0}(c_0, r_0)$, for some $c_0 > 0$ and $r_0 > 0$. There exists a constant $c = c(N, p_0) \in (0, 1)$ such that for every $R_0 > r_0$ we also have*

$$\Omega \in \text{UT}_{p_0} \left(c \cdot c_0 \left(\frac{r_0}{R_0} \right)^N, R_0 \right).$$

In other words, for every $R_0 > r_0$ we have

$$\text{cap}_{p_0} \left(\overline{Q_r(x_0)} \setminus \Omega; Q_{2r}(x_0) \right) \geq c \cdot c_0 \left(\frac{r_0}{R_0} \right)^N \text{cap}_{p_0} \left(\overline{Q_r(x_0)}; Q_{2r}(x_0) \right), \quad \text{for every } 0 < r \leq R_0,$$

as well.

Proof. We fix $R_0 > r_0$ and consider a cube $Q_r(x_0)$ with $x_0 \in \mathbb{R}^N \setminus \Omega$ and $r_0 < r \leq R_0$. By using the monotonicity properties of p_0 -capacity, we have

$$\begin{aligned} \text{cap}_{p_0} \left(\overline{Q_r(x_0)} \setminus \Omega; Q_{2r}(x_0) \right) &\geq \text{cap}_{p_0} \left(\overline{Q_{r_0}(x_0)} \setminus \Omega; Q_{2R_0}(x_0) \right) \\ &= \frac{\text{cap}_{p_0} \left(\overline{Q_{r_0}(x_0)} \setminus \Omega; Q_{2R_0}(x_0) \right)}{\text{cap}_{p_0} \left(\overline{Q_r(x_0)}; Q_{2r}(x_0) \right)} \text{cap}_{p_0} \left(\overline{Q_r(x_0)}; Q_{2r}(x_0) \right) \\ &= \frac{\text{cap}_{p_0} \left(\overline{Q_{r_0}(x_0)} \setminus \Omega; Q_{2R_0}(x_0) \right)}{r^{N-p_0} \text{cap}_{p_0} (\overline{Q_1}; Q_2)} \text{cap}_{p_0} \left(\overline{Q_r(x_0)}; Q_{2r}(x_0) \right) \\ &\geq \frac{\text{cap}_{p_0} \left(\overline{Q_{r_0}(x_0)} \setminus \Omega; Q_{2R_0}(x_0) \right)}{R_0^{N-p_0} \text{cap}_{p_0} (\overline{Q_1}; Q_2)} \text{cap}_{p_0} \left(\overline{Q_r(x_0)}; Q_{2r}(x_0) \right). \end{aligned}$$

On the other hand, by [32, Proposition 13.1.1.1], we can estimate

$$\text{cap}_{p_0} \left(\overline{Q_{r_0}(x_0)} \setminus \Omega; Q_{2R_0}(x_0) \right) \geq c \left(\frac{r_0}{R_0} \right)^{p_0} \text{cap}_{p_0} \left(\overline{Q_{r_0}(x_0)} \setminus \Omega; Q_{2r_0}(x_0) \right).$$

A further application of the definition of p_0 -thickness permits to bound from below the right-hand side. Thus, we get

$$\text{cap}_{p_0} \left(\overline{Q_r(x_0)} \setminus \Omega; Q_{2r}(x_0) \right) \geq c \cdot c_0 \left(\frac{r_0}{R_0} \right)^N \text{cap}_{p_0} \left(\overline{Q_r(x_0)}; Q_{2r}(x_0) \right).$$

This concludes the proof. $\square$

The following result assures that the finiteness of the inradius, in conjunction with the thickness of the complement, is a sufficient condition for Poincaré inequality to hold. This is the main result of this appendix. Apart for its applications to our problem, it is interesting in itself.

Theorem A.2. *Let $1 \leq p_0 < \infty$ and let $\Omega \subseteq \mathbb{R}^N$ be an open set such that:*

- $r_\Omega < +\infty$;

- $\Omega \in \text{UT}_{p_0}(c_0, r_0)$ for some $c_0 > 0$ and $r_0 > 0$.

Then

$$\lambda_{p_0}(\Omega) > 0 \quad \text{and} \quad \Omega \in \text{UT}_{p_0}(\tilde{c}_0),$$

for some $\tilde{c}_0 = \tilde{c}_0(N, p, \Omega) > 0$.

Proof. We divide the proof in two parts: we first prove that Ω supports the L^{p_0} Poincaré inequality and then we show that the p_0 -thickness condition holds uniformly for every $r > 0$.

Part 1: validity of the Poincaré inequality. We claim that there exists a constant $C = C(N, p_0, r_0/r_\Omega) > 0$ such that

$$(A.1) \quad \text{cap}_{p_0} \left(\overline{Q_{2r_\Omega}(x_0)} \setminus \Omega; Q_{4r_\Omega}(x_0) \right) \geq \frac{1}{C} r_\Omega^{N-p_0}, \quad \text{for every } x_0 \in \mathbb{R}^N.$$

In light of [32, Theorem 15.4.2.1], this would be enough to get the desired conclusion. In order to establish (A.1), we distinguish two cases: either $x_0 \in \mathbb{R}^N \setminus \Omega$ or $x_0 \in \Omega$.

Case $x_0 \in \mathbb{R}^N \setminus \Omega$. Our assumptions guarantee that there exists $\tilde{c}_0 > 0$ such that for every $x_0 \in \mathbb{R}^N \setminus \Omega$ we have

$$(A.2) \quad \text{cap}_{p_0} \left(\overline{Q_r(x_0)} \setminus \Omega; Q_{2r}(x_0) \right) \geq \tilde{c}_0 \text{cap}_{p_0} \left(\overline{Q_r(x_0)}; Q_{2r}(x_0) \right), \quad \text{for every } 0 < r \leq 2r_\Omega.$$

Indeed, if $r_0 \geq 2r_\Omega$ this is immediate from the definition of p_0 -thickness (and we can take $\tilde{c}_0 = c_0$). If $r_0 < 2r_\Omega$, we could apply Lemma A.1 with the choice $R_0 = 2r_\Omega$ and get (A.2) with

$$\tilde{c}_0 = c \cdot c_0 \left(\frac{r_0}{2r_\Omega} \right)^N.$$

In particular, from (A.2) with $r = 2r_\Omega$ we easily get (A.1) for $x_0 \in \mathbb{R}^N \setminus \Omega$.

Case $x_0 \in \Omega$. We recall that

$$r_\Omega = \sup_{x \in \Omega} d_\Omega(x), \quad \text{where} \quad d_\Omega(x) = \min_{y \in \partial\Omega} |x - y|, \quad \text{for every } x \in \Omega.$$

By observing that

$$\|x - y\|_{\ell^\infty} \leq |x - y|, \quad \text{for every } x, y \in \mathbb{R}^N,$$

we thus obtain that

$$(A.3) \quad \mathfrak{d}(x) := \inf_{y \in \partial\Omega} \|x - y\|_{\ell^\infty} \leq r_\Omega, \quad \text{for every } x \in \Omega.$$

Then, for every $x_0 \in \Omega$ and every $\mathfrak{d}(x_0) < r \leq 2r_\Omega + \mathfrak{d}(x_0)$, we have from Lemma 2.9

$$\text{cap}_{p_0} \left(\overline{Q_r(x_0)} \setminus \Omega; Q_{2r}(x_0) \right) \geq \frac{1}{C_2} \left(\frac{r - \mathfrak{d}(x_0)}{2r + \mathfrak{d}(x_0)} \right)^{p_0} (r - \mathfrak{d}(x_0))^{N-p_0}.$$

By choosing $r = 2r_\Omega$ and using (A.3), we obtain in particular (A.1) for $x_0 \in \Omega$, as well.

Part 2: universality of the radius. We now need to show that there exists a constant $\tilde{c}_0$ such that

$$\text{cap}_{p_0} \left(\overline{Q_r(x_0)} \setminus \Omega; Q_{2r}(x_0) \right) \geq \tilde{c}_0 \text{cap}_{p_0} \left(\overline{Q_r(x_0)}; Q_{2r}(x_0) \right), \quad \text{for every } x_0 \in \mathbb{R}^N \setminus \Omega, \quad r > 0.$$

More precisely, we are only left with verifying this property for $r > r_0$, since for $0 < r \leq r_0$ this is true by assumption.

The case $p_0 > N$ is easy, it is sufficient to observe that for every $x_0 \in \mathbb{R}^N \setminus \Omega$ and every $r > 0$, we have

$$\begin{aligned} \text{cap}_{p_0} \left(\overline{Q_r(x_0)} \setminus \Omega; Q_{2r}(x_0) \right) &\geq \text{cap}_{p_0} (\{x_0\}; Q_{2r}(x_0)) \\ &= r^{N-p_0} \text{cap}_{p_0} (\{x_0\}; Q_2(x_0)) \\ &= \frac{\text{cap}_{p_0} (\{x_0\}; Q_2(x_0))}{\text{cap}_{p_0} (\overline{Q_1(x_0)}; Q_2(x_0))} r^{N-p_0} \text{cap}_{p_0} (\overline{Q_1(x_0)}; Q_2(x_0)) \\ &=: \frac{1}{C} \text{cap}_{p_0} (\overline{Q_r(x_0)}; Q_{2r}(x_0)). \end{aligned}$$

Observe that we used the scale property of the relative p_0 -capacity and the fact that points have positive p_0 -capacity, if $p_0 > N$.

Thus, let us consider the case $1 \leq p_0 \leq N$. We need to recall at first the definition of *capacitary inradius* of Ω , taken from [5]. For every $0 < \gamma < 1$, this is defined by

$$R_{p_0, \gamma}(\Omega) := \sup \left\{ r > 0 : \exists x_0 \in \mathbb{R}^N \text{ s. t. } \text{cap}_{p_0} (\overline{B_r(x_0)} \setminus \Omega; B_{2r}(x_0)) \leq \gamma \text{cap}_{p_0} (\overline{B_r(x_0)}; B_{2r}(x_0)) \right\}.$$

Then we have the following upper bound (see [5, Main Theorem])

$$\lambda_{p_0}(\Omega) \leq C_{N, p_0, \gamma} \left(\frac{1}{R_{p_0, \gamma}(\Omega)} \right)^{p_0},$$

with the constant $C_{N, p_0, \gamma}$ which diverges to $+\infty$, as γ goes to 1. Observe that $\lambda_{p_0}(\Omega) > 0$, by the first part of the proof.

Moreover, we can also consider the following variant of the capacitary inradius, with cubes in place of balls

$$R_{p_0, \gamma}(\Omega; Q_1) := \sup \left\{ r > 0 : \exists x_0 \in \mathbb{R}^N \text{ s. t. } \text{cap}_{p_0} (\overline{Q_r(x_0)} \setminus \Omega; Q_{2r}(x_0)) \leq \gamma \text{cap}_{p_0} (\overline{Q_r(x_0)}; Q_{2r}(x_0)) \right\}.$$

This quantity has been introduced by Maz'ya, under the name of *inner cubic diameter* of Ω , see [33, Definition 10.2.2]. By [4, Proposition 3.6], we have

$$(A.4) \quad R_{p_0, c \cdot \gamma}(\Omega; Q_1) \leq R_{p_0, \gamma}(\Omega),$$

for a constant $0 < c < 1$ depending only on N and p_0 . We now choose for simplicity $\gamma = 1/2$, from (A.4) we get

$$R_{p_0, c/2}(\Omega; Q_1) \leq R_{p_0, 1/2}(\Omega) =: \mathfrak{r}_0.$$

Thus, by the very definition of $R_{p_0, c \cdot \gamma}(\Omega; Q_1)$, we get that for every $r > \mathfrak{r}_0$ and for every $x_0 \in \mathbb{R}^N$, we must have

$$\text{cap}_{p_0} (\overline{Q_r(x_0)} \setminus \Omega; Q_{2r}(x_0)) > \frac{c}{2} \text{cap}_{p_0} (\overline{Q_r(x_0)}; Q_{2r}(x_0)).$$

If $\mathfrak{r}_0 \leq r_0$, the proof is over. On the contrary, if $\mathfrak{r}_0 > r_0$ we can simply apply Lemma A.1 with $R_0 = \mathfrak{r}_0$. $\square$

Remark A.3 (Sharpness of the assumptions). In the previous result, the assumption $r_\Omega < +\infty$ is crucial, i.e. in general its conclusion does not hold if the inradius is infinite. For example, by taking $1 \leq p_0 < N$ and

$$\Omega = \mathbb{R}^N \setminus \overline{Q_1},$$

we see that its complement $\mathbb{R}^N \setminus \Omega = \overline{Q_1}$ is uniformly p_0 -thick, for suitable $c_0 > 0$ and $0 < r_0 < +\infty$. Indeed, for every $x_0 \in \overline{Q_1}$, we have

$$\text{cap}_{p_0}(\overline{Q_r(x_0)} \setminus \Omega; Q_{2r}(x_0)) = \text{cap}_{p_0}(\overline{Q_r(x_0)} \cap \overline{Q_1}; Q_{2r}(x_0)).$$

By the monotonicity properties of the p_0 -capacity and the lower bound [32, (2.2.12)], we can estimate from below the right-hand side as follows

$$\begin{aligned} \text{cap}_{p_0}(\overline{Q_r(x_0)} \cap \overline{Q_1}; Q_{2r}(x_0)) &\geq \text{cap}_{p_0}(\overline{Q_r(x_0)} \cap \overline{Q_1}; \mathbb{R}^N) \\ &\geq (N \omega_N)^{\frac{p_0}{N}} N^{\frac{N-p_0}{N}} \left(\frac{N-p_0}{p_0-1} \right)^{p_0-1} \left| \overline{Q_r(x_0)} \cap \overline{Q_1} \right|^{\frac{N-p_0}{N}}. \end{aligned}$$

We now observe that

$$\left| \overline{Q_r(x_0)} \cap \overline{Q_1} \right| \geq \frac{1}{2^N} |Q_r(x_0)| = r^N, \quad \text{for every } x_0 \in \overline{Q_1}, \quad 0 < r \leq 1.$$

Thus, we conclude that $\Omega \in \text{UT}_{p_0}(c_0, r_0)$, with

$$c_0 = \frac{(N \omega_N)^{\frac{p_0}{N}} N^{\frac{N-p_0}{N}} \left(\frac{N-p_0}{p_0-1} \right)^{p_0-1}}{\text{cap}_{p_0}(\overline{Q_1}; Q_2)} \quad \text{and} \quad r_0 = 1.$$

On the other hand, for this set Ω we have

$$\lambda_{p_0}(\Omega) = 0 \quad \text{and} \quad \Omega \notin \bigcup_{c>0} \text{UT}_{p_0}(c; +\infty).$$

The first fact follows since Ω contains arbitrarily large balls. The second one follows since

$$(A.5) \quad \lim_{r \rightarrow +\infty} \frac{\text{cap}_{p_0}(\overline{Q_r} \setminus \Omega; Q_{2r})}{\text{cap}_{p_0}(\overline{Q_r}; Q_{2r})} = 0.$$

It is sufficient to observe that

$$\overline{Q_r} \setminus \Omega = \overline{Q_1}, \quad \text{for every } r \geq 1.$$

Then, by taking $\eta \in C_0^\infty(Q_2)$ such that $\eta \geq 1$ on $\overline{Q_1}$, we get for every $r \geq 1$

$$\frac{\text{cap}_{p_0}(\overline{Q_r} \setminus \Omega; Q_{2r})}{\text{cap}_{p_0}(\overline{Q_r}; Q_{2r})} = \frac{\text{cap}_{p_0}(\overline{Q_1}; Q_{2r})}{\text{cap}_{p_0}(\overline{Q_r}; Q_{2r})} \leq \frac{\int_{Q_2} |\nabla \eta|^{p_0} dx}{r^{N-p_0}} \frac{1}{\text{cap}_{p_0}(\overline{Q_1}; Q_2)}.$$

By taking the limit as r goes to ∞ and using that $p_0 < N$, we get (A.5).

From Theorem A.2, we immediately get the following consequence.

Corollary A.4. *Let $1 \leq p_0 < \infty$ and let $\Omega \subseteq \mathbb{R}^N$ be an open set such that:*

- $r_\Omega < +\infty$;
- $\Omega \in \text{UT}_{p_0}(c_0, r_0)$ for some $c_0 > 0$ and $r_0 > 0$.

Then $\lambda_p(\Omega) > 0$, for every $p \geq p_0$.

Proof. Once we have obtained $\lambda_{p_0}(\Omega) > 0$ from Theorem A.2, the desired conclusion follows from the inequality

$$\left(\frac{p_0}{p} \right)^p (\lambda_{p_0}(\Omega))^{\frac{p}{p_0}} \leq \lambda_p(\Omega),$$

see for example [12, Proposition 2.1]. □

APPENDIX B. A NOTE ON THE CASE $1 < p < N$

In the case $1 < p < N$, if we allow $\gamma_0 = (p^*)'$ in (E_f) of Assumptions 1, one could be dispensed with the assumption (E_Ω) . However, the natural functional space where the solution lies is now the homogeneous Sobolev space $\mathcal{D}_0^{1,p}(\Omega)$. Accordingly, the statement of the Main Theorem should be amended as follows. The proof would be just a minor variant of that of the Main Theorem: therefore, it is omitted.

Theorem B.1 (Exceptional case). *Let $1 < p < N$ and let $\Omega \subseteq \mathbb{R}^N$ be an open set such that*

$$\Omega \in \text{UT}_{p_0}(c_0, r_0), \quad \text{for some } 1 \leq p_0 < p \quad \text{and} \quad c_0, r_0 > 0.$$

Let $f \in L^{(p^)'}(\Omega) \cap L^\gamma(\Omega)$, for some $(p^*)' < \gamma \leq \infty$. Finally, let $u \in \mathcal{D}_0^{1,p}(\Omega)$ be the weak solution to the equation (1.1), where G still satisfies the assumption (E_G) .*

Then, there exists an exponent $\delta = \delta(N, p, p_0, \gamma, c_0, g_1, g_2) > 0$ with $\delta \leq (\gamma - (p^)')/(p^*)'$, such that*

$$\nabla u \in L^{p(1+\delta)}(\Omega).$$

Moreover, there exists a sequence $\{\varphi_n\}_{n \in \mathbb{N}} \subseteq C_0^\infty(\Omega)$ such that

$$\lim_{n \rightarrow \infty} \|\nabla u - \nabla \varphi_n\|_{L^{p(1+\delta)}(\Omega)} = 0.$$

Finally, the following scale invariant estimate holds, for $C_1 = C_1(N, p, p_0, \gamma, c_0, g_1, g_2) > 0$

$$\begin{aligned} \left(\int_{\Omega} |\nabla u|^{p(1+\delta)} dx \right)^{\frac{1}{1+\delta}} &\leq \frac{C_1}{r_0^{\frac{\delta}{N \frac{1}{1+\delta}}}} \int_{\Omega} |\nabla u|^p dx \\ &\quad + C_1 \left(\|f\|_{L^{\gamma_0}(\Omega)}^{\gamma_0} \right)^{\frac{p'}{\gamma_0} - \frac{\delta}{1+\delta} \frac{\gamma}{\gamma-\gamma_0}} \left(\|f\|_{L^\gamma(\Omega)}^\gamma \right)^{\frac{\delta}{1+\delta} \frac{\gamma_0}{\gamma-\gamma_0}}. \end{aligned}$$

It is intended that the first term on the right-hand side has to be considered as zero, in the case $r_0 = +\infty$.

Remark B.2. We remark that, under the previous assumptions, in general we could have $r_0 < +\infty$. It is sufficient to keep in mind the example of Remark A.3.

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