

A Pohožaev minimization for normalized solutions: fractional sublinear equations of logarithmic type

Marco Gallo 

Università Cattolica del Sacro Cuore
Dipartimento di Matematica e Fisica
Via della Garzetta 48, 25133 Brescia, Italy
marco.gallo1@unicatt.it

Jacopo Schino 

University of Warsaw
Faculty of Mathematics, Informatics and Mechanics
ul. Banacha 2, 02-097 Warsaw, Poland
j.schino2@uw.edu.pl

Abstract

In this paper, we search for normalized solutions to a fractional, nonlinear, and possibly strongly sublinear Schrödinger equation

$$(-\Delta)^s u + \mu u = g(u) \quad \text{in } \mathbb{R}^N,$$

under the mass constraint $\int_{\mathbb{R}^N} u^2 dx = m > 0$; here, $N \geq 2$, $s \in (0, 1)$, and μ is a Lagrange multiplier. We study the case of L^2 -subcritical nonlinearities g of Berestycki–Lions type, without assuming that g is superlinear at the origin, which allows us to include examples like a logarithmic term $g(u) = u \log(u^2)$ or sublinear powers $g(u) = u^q - u^r$, $0 < r < 1 < q$. Due to the generality of g and the fact that the energy functional might be not well-defined, we implement an approximation process in combination with a Lagrangian approach and a new Pohožaev minimization in the product space, finding a solution for large values of m . In the sublinear case, we are able to find a solution for each m . Several insights on the concepts of minimality are studied as well. We highlight that some of the results are new even in the local setting $s = 1$ or for g superlinear.

Keywords: NLS equations, Fractional Laplacian, Sublinear nonlinearity, Nonsmooth analysis, L^2 -norm constraint, Prescribed-mass solutions, Least-energy solutions, Ground states.

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1 Introduction

In this paper, we study the fractional L^2 -subcritical equation

$$(-\Delta)^s u + \mu u = g(u) \quad \text{in } \mathbb{R}^N \tag{1.1}$$

under the mass constraint $\int_{\mathbb{R}^N} u^2 dx = m > 0$; here, $s \in (0, 1)$, $N > 2s$, $\mu \in \mathbb{R}$ is a suitable Lagrange multiplier (part of the unknowns), and g is a generic Berestycki–Lions [8] nonlinearity which is allowed to be (negatively) *sublinear* at the origin (see (g0)–(g4) below). In particular, $G(t) := \int_0^t g(\tau) d\tau$ is allowed to have a growth of the type

$$-|t|^{1+\theta} - |t|^{2_s^*} \lesssim G(t) \lesssim |t|^2 + |t|^{\bar{p}+1} \quad \text{for } t \in \mathbb{R},$$

where $\theta \in [0, 1)$, $2_s^* := \frac{2N}{N-2s}$ is the fractional Sobolev critical exponent, and

$$\bar{p} := 1 + \frac{4s}{N}; \tag{1.2}$$

the number $\bar{p} + 1$ is often referred to as the L^2 -critical exponent. This case includes, for example, the well-known *logarithmic* nonlinearity $g(t) = t \log(t^2)$ (here $\theta \in [0, 1)$ can be arbitrarily chosen), or *sublinear-power*-type nonlinearities $g(t) = |t|^{q-1}t - |t|^{r-1}t$, $0 < r < q < \bar{p}$ (here $\theta = r$), but also strongly sublinear nonlinearities like $g(t) \sim \frac{1}{\log(|t|)}$ as $t \rightarrow 0$ (here $\theta = 0$).

Exploiting variational methods, a key fact of these nonlinearities is that the related energy functional $K: H^s(\mathbb{R}^N) \rightarrow \mathbb{R} \cup \{+\infty\}$,

$$K(u) := \frac{1}{2} \int_{\mathbb{R}^N} |(-\Delta)^{s/2} u|^2 dx - \int_{\mathbb{R}^N} G(u) dx, \tag{1.3}$$

is only lower semicontinuous, while it is well-posed in $H^s(\mathbb{R}^N) \cap L^1(\mathbb{R}^N)$, which we highlight is not a reflexive space. Here $H^s(\mathbb{R}^N)$ denotes the usual fractional Sobolev space.

1.1 Some literature

When $s = 1$, the frequency $\mu \in \mathbb{R}$ is fixed, and $g(t) = t \log(t^2)$, the *logarithmic Schrödinger equation* [9, 10]

$$-\Delta u + \mu u = u \log(u^2) \quad \text{in } \mathbb{R}^N \quad (1.4)$$

has received much attention due to its great relevance in applied sciences, such as atomic physics, high-energy cosmic rays, Cherenkov-type shock waves, quantum hydrodynamical models, and many others; we refer to [15, 33] and the references therein. A key property of (1.4) is, indeed, the *tensorization property* (or *separability of noninteracting subsystems*): namely, if u_i are solutions of (1.4) in \mathbb{R}^{N_i} for $i = 1, 2$, then $(x_1, x_2) \mapsto u_1(x_1)u_2(x_2)$ is a solution of (1.4) in $\mathbb{R}^{N_1} \times \mathbb{R}^{N_2}$.

In order to handle the singularity at the origin, several techniques have been developed: a first approach, applied by [16], shows that the functional K in (1.3) is of class \mathcal{C}^1 on the subspace

$$X := \left\{ u \in H^1(\mathbb{R}^N) \mid \int_{\mathbb{R}^N} u^2 \log(u^2) dx < +\infty \right\};$$

see also [50, 56], and [12] for general doubling nonlinearities. We observe that $X \subsetneq H^1(\mathbb{R}^N)$, as shown by $u(x) := (|x|^{N/2} \log(|x|))^{-1}$ for $|x|$ large. Finally, in [41], a suitable norm is introduced in order to force the positive solutions far from the singularity.

A second approach has been developed in [24, 53]: here, the authors apply the *nonsmooth analysis* theory of [26, 54] and work with a suitable subdifferential of the action functional, possibly splitting it into a smooth part and a convex lower semicontinuous one.

Different approaches rely instead on some *approximation schemes*. The idea that $\frac{t^{2+\delta}-t^2}{\delta} \rightarrow t^2 \log(t)$ as $\delta \rightarrow 0^+$ was formalized in the framework of PDEs by [58] (see also [31], [15, Section 7]), and then this idea was applied in [59] to get existence of solutions. See also [33], where the regularization is of the form $\log(t + \delta)$, and [1], where the authors approximate the domain of the problem with expanding balls.

We mention that the problem of the singularity of the logarithm arises also in the study of planar Choquard equations [22, 39].

Let us move, now, to the case of sublinear powers. Although nonlinear (fractional) Schrödinger equations with combined powers have recently gained the mathematical community's attention [27, 42, 52, 55], such work usually concerns superlinear exponents. Instead, as regards sublinear-power nonlinearities (again, $s = 1$ and μ fixed), most of the work focuses on problems of the type

$$-\Delta u + V(x)u = a(x)|u|^{q-1}u + h(x) \quad \text{in } \mathbb{R}^N,$$

$q \in (0, 1)$, where suitable assumptions on V, a, h are considered in order to ensure the existence of solutions (notice, in fact, that when $V \equiv a \equiv 1$ and $h \equiv 0$, no solution exists, see e.g. [6, Lemma 2.3]). When $a \notin L^{\frac{2}{q-2}}(\mathbb{R}^N)$, the analysis of the problem is usually set in the Banach space $H^1(\mathbb{R}^N) \cap L^q(\mathbb{R}^N)$; see [5, 6, 57].

The only paper we know that deals with *general nonlinearities* in the spirit of [8] is [44]: here, the author studies

$$-\Delta u = f(u) \quad \text{in } \mathbb{R}^N,$$

where the assumptions about f include the case $\lim_{t \rightarrow 0} \frac{f(t)}{|t|^{2^*-2t}} = -\infty$. The author employs a perturbation argument at the origin and obtains a ground state for the equation.

When $s = 1$ and μ is free, very few results appear in the literature about the existence of normalized solutions $\int_{\mathbb{R}^N} u^2 dx = m > 0$. When $g(t) = t \log(t^2)$, we see that, considered a nontrivial solution u_1 of (1.4) for $\mu = 1$, then defining $u := \alpha_m u_1$ with $\alpha_m := m(\int_{\mathbb{R}^N} u_1^2 dx)^{-1}$, we have that u satisfies $\int_{\mathbb{R}^N} u^2 dx = m$ and

$$-\Delta u + (1 + \log(\alpha_m^2)) u = u \log(u^2) \quad \text{in } \mathbb{R}^N, \quad (1.5)$$

thus the existence of a normalized solution is always ensured; this was already noticed in [10]. In particular, the frequency $1 + \log(\alpha_m)$ is negative when m is sufficiently small and positive when m is sufficiently large. Additionally, an explicit L^2 -minimum can be found. Indeed, by the logarithmic Sobolev–Gross–Nelson inequality [32] there exists $C(N) \in \mathbb{R}$ such that

$$K(u) \geq -(C(N) + \log(m)) m \tag{1.6}$$

for every $u \in H^1(\mathbb{R}^N)$ with $\int_{\mathbb{R}^N} u^2 dx = m$; moreover, we know that the equality is attained by (and only by, up to a translation) the *Gausson* $u(x) = \gamma_m e^{-\frac{1}{2}|x|^2}$ with $\gamma_m := m \left(\int_{\mathbb{R}^N} e^{-|x|^2} dx \right)^{-1}$. This procedure can be found, for example, in [9, 16]. We see that the sign of the minimal energy over the L^2 -sphere depends on the size of m (and this is connected with the sign of the Lagrange multiplier, see Proposition 1.8).

When g does not enjoy scaling properties of this type, we mention the very recent papers [2, 45, 51, 60]. In [51] and [2], the authors study, respectively, logarithm-plus-power and nonautonomous problems by restricting to the subspace X . In [60], via a direct minimization, the authors study a similar equation, but with a more general g satisfying a monotonicity assumption: they find an L^2 -minimum for every mass m , which, in particular, has positive energy for m small and negative energy for m large. Finally, by exploiting the perturbation techniques developed in [44], the authors in [45] succeed in removing the monotonicity condition on g , and find an L^2 -minimum for m large, with a positive multiplier $\mu > 0$.

When $s \in (0, 1)$, the fractional logarithmic equation

$$(-\Delta)^s u + \mu u = u \log(u^2) \quad \text{in } \mathbb{R}^N \tag{1.7}$$

appears in the study of nonlinear quantized boson fields, [48]. Mathematically, several papers succeeded in generalizing the results found in the local case for sublinear nonlinearities, exploiting and adapting to the nonlocal case the aforementioned techniques, see e.g. [4, 25, 37] for the logarithmic equation, and [29] for logarithm-plus-power nonlinearities. In particular, we mention [35]: here, the author, with a perturbation argument similar to the one in [44], finds infinitely many solutions with a general nonlinear term.

Anyway, all the previous results deal with the unconstrained problem with fixed frequency. When $g(t) = t \log(t^2)$, the search for a normalized solution can be pursued as above by scaling (see (1.5)); on the other hand, the search for an L^2 -minimum through an optimal logarithmic inequality (1.6) has issues. Indeed, although such an inequality still holds in the fractional framework [17, 23], i.e., there exists $C(N, s) \in \mathbb{R}$ such that

$$K(u) \geq -(C(N, s) + \log(m)) m$$

for every $u \in H^s(\mathbb{R}^N)$ with $\int_{\mathbb{R}^N} u^2 dx = m$, it is not known if the equality is attained or not. To the authors' knowledge, the only articles dealing with normalized solutions to logarithmic fractional equations are the recent works [38, 43], where semiclassical problems are studied.

The goal of this paper is to find normalized solutions with minimality properties to fractional sublinear equations under general assumptions about the nonlinear term. When treating such general terms g , as previously commented, the splitting of the functional does not seem an option, as well as the definition of a space similar to X . To obtain the existence of an L^2 -minimum we rely, thus, on the perturbation technique employed by [44, 45]: this perturbation takes the form $g_\varepsilon(t) := g_+(t) - \varphi_\varepsilon(t)g_-(t)$, where $\varphi_\varepsilon(t) = |t|/\varepsilon$ for $|t| \leq \varepsilon$ (see Section 5 for details). Additionally, we obtain several properties that are new even for $s = 1$ and g superlinear.

1.2 Existence results

In the present paper, we study the following problem

$$\begin{cases} (-\Delta)^s u + \mu u = g(u) & \text{in } \mathbb{R}^N, \\ \int_{\mathbb{R}^N} u^2 dx = m, \\ (\mu, u) \in \mathbb{R} \times H^s(\mathbb{R}^N), \end{cases} \quad (1.8)$$

with $s \in (0, 1)$, $N \geq 2$, $m > 0$, and $g: \mathbb{R} \rightarrow \mathbb{R}$ satisfying (set $G(t) := \int_0^t g(\tau) d\tau$)

(g0) g is continuous and $g(0) = 0$,

(g1) $\limsup_{t \rightarrow 0} \frac{g(t)}{t} \leq 0$,

(g2) $\limsup_{|t| \rightarrow +\infty} \frac{|g(t)|}{|t|^{2_s^*-1}} < \infty$,

(g3) $\limsup_{|t| \rightarrow +\infty} \frac{g(t)}{|t|^{\bar{p}-1}t} \leq 0$,

(g4) there exists $t_0 \neq 0$ such that $G(t_0) > 0$.

Here and in what follows, we denote $g^\pm := \max\{\pm g, 0\}$ the standard positive and negative parts of g , and

$$g_+(t)t := (g(t)t)^+, \quad G_+(t) := \int_0^t g_+(\tau) d\tau, \quad (1.9)$$

$$g_-(t)t := (g(t)t)^-, \quad G_-(t) := \int_0^t g_-(\tau) d\tau; \quad (1.10)$$

see Section 2. With these notations, we may rewrite (g1) and (g3) as

(g1) $\lim_{t \rightarrow 0} \frac{g_+(t)}{t} = 0$,

(g3) $\lim_{|t| \rightarrow +\infty} \frac{g_+(t)}{|t|^{\bar{p}-1}t} = 0$.

We further notice that, when (g3) holds, then (g2) actually means $\limsup_{|t| \rightarrow +\infty} \frac{g_-(t)}{|t|^{2_s^*-2}t} < \infty$.

Before presenting our results, let us recall that in [18, Proposition 4.1] and [14, Proposition 1.1] (see also [21, Corollary 6.4]), when g satisfies Berestycki–Lions assumptions, the following *Pohožaev identity* holds

$$\frac{1}{2_s^*} \int_{\mathbb{R}^N} |(-\Delta)^{s/2} u|^2 dx + \frac{\mu}{2} \int_{\mathbb{R}^N} u^2 dx - \int_{\mathbb{R}^N} G(u) dx = 0 \quad (1.11)$$

whenever u is a solution of (1.8) and either $s \in (\frac{1}{2}, 1)$ or $g \in \mathcal{C}_{loc}^{0,\alpha}(\mathbb{R}^N)$ with $\alpha \in (1 - 2s, 1)$: this result can be generalized to our assumptions whenever $\int_{\mathbb{R}^N} G(u) dx$ is assumed a priori finite; however, we see that the condition on the Hölder regularity of g forces a restriction on the singularity of g_- at the origin. It is not known if this identity holds for general continuous nonlinearities g and general values of $s \in (0, 1)$. Additionally, we observe that, due to the possible

sublinearity of g , for a general solution u of equation (1.8) we cannot state a priori that it satisfies the *Nehari identity*

$$\int_{\mathbb{R}^N} |(-\Delta)^{s/2}u|^2 dx + \mu \int_{\mathbb{R}^N} u^2 dx - \int_{\mathbb{R}^N} g(u)u dx = 0. \quad (1.12)$$

To state the main theorem, we make use of a *Lagrangian formulation* of the problem (in the spirit of [20, 34]) and introduce, in addition to (1.3), a Lagrangian functional over the product space: $I^m: \mathbb{R} \times H^s(\mathbb{R}^N) \rightarrow \mathbb{R}$,

$$I^m(\mu, u) := \frac{1}{2} \int_{\mathbb{R}^N} |(-\Delta)^{s/2}u|^2 dx + \frac{\mu}{2} \left(\int_{\mathbb{R}^N} u^2 dx - m \right) - \int_{\mathbb{R}^N} G(u) dx. \quad (1.13)$$

Notice that critical points of I^m are (formally) solutions to (1.8) and $I^m(\mu, u) = K(u)$ if $\int_{\mathbb{R}^N} u^2 dx = m$.

From now on, to state our results we will use the following convention.

Convention: whenever we deal with a pair (μ, u) (or (λ, u) with $e^\lambda = \mu$, see (5.5)) and we say that such a pair is radially symmetric, nonnegative, does not change sign, etc., *we are referring to u* .

We introduce the following notations, standing for the L^2 -sphere in $H^s(\mathbb{R}^N)$ and the Pohožaev set in the product space $\mathbb{R} \times H^s(\mathbb{R}^N)$:

$$\mathcal{S}_m := \{u \in H^s(\mathbb{R}^N) \mid |u|_2^2 = m\}, \quad (1.14)$$

where $|\cdot|_q$ is the usual Lebesgue norm, $1 \leq q \leq \infty$, and

$$\mathcal{P} := \{(\mu, u) \in \mathbb{R} \times H^s(\mathbb{R}^N) \mid u \neq 0 \text{ and } P(\mu, u) = 0\}, \quad (1.15)$$

where $P: \mathbb{R} \times H^s(\mathbb{R}^N) \rightarrow \mathbb{R}$ is the Pohožaev functional¹

$$P(\mu, u) := \int_{\mathbb{R}^N} |(-\Delta)^{s/2}u|^2 dx + 2_s^* \int_{\mathbb{R}^N} \frac{\mu}{2} u^2 - G(u) dx. \quad (1.16)$$

It is clear that $(\mu, u) \in \mathcal{P}$ if and only if it satisfies the Pohožaev identity (1.11). Moreover, $\mathcal{P} \neq \emptyset$ because, fixed $u \in H_{\text{rad}}^s(\mathbb{R}^N) \setminus \{0\}$, $P(\cdot, u)$ is a continuous function such that $\lim_{\mu \rightarrow \pm\infty} P(\mu, u) = \pm\infty$. Next, we set

$$\mathcal{S}_m^{\text{rad}} := \mathcal{S}_m \cap H_{\text{rad}}^s(\mathbb{R}^N), \quad \mathcal{P}^{\text{rad}} := \mathcal{P} \cap (\mathbb{R} \times H_{\text{rad}}^s(\mathbb{R}^N)),$$

where $H^s(\mathbb{R}^N)$ is the classical fractional Sobolev space, cf. (2.1), and $H_{\text{rad}}^s(\mathbb{R}^N)$ is its subspace of radially symmetric functions, and

$$\kappa^m := \inf_{\mathcal{S}_m^{\text{rad}}} K, \quad d^m := \inf_{\mathcal{P}^{\text{rad}}} I^m; \quad (1.17)$$

here K and I^m are defined in (1.3) and (1.13). We further define

$$\bar{\mu}_0 := \sup_{t \neq 0} \frac{G(t)}{t^2/2}; \quad (1.18)$$

¹ We highlight that the terminology ‘‘Pohožaev functional’’, in the framework of normalized solutions, is often referred to a suitable combination of the Pohožaev and Nehari functionals, which allows getting rid of the Lagrange multiplier μ ; such a functional appears, in particular, in the framework of L^2 -supercritical problems. Here, instead, we focus on the proper Pohožaev identity, considering a two-variable functional.

notice that if (g4) holds, then $\bar{\mu}_0 \in (0, +\infty]$. Moreover, we introduce the L^2 -ball and the corresponding infimum:

$$\mathcal{D}_m := \{u \in H^s(\mathbb{R}^N) \mid |u|_2^2 \leq m\}, \quad \mathcal{D}_m^{\text{rad}} := \mathcal{D}_m \cap H_{\text{rad}}^s(\mathbb{R}^N), \quad \ell^m := \inf_{\mathcal{D}_m^{\text{rad}}} K; \quad (1.19)$$

we mention that the use of the L^2 -ball in normalized problems goes back to [11], while it was utilized for the first time in the L^2 -subcritical regime in [49]. We can thus state our first theorem, see Definition 2.4 for the notion of solution.

Theorem 1.1 (Existence for large masses) *If (g0)–(g4) hold, then there exists $m_0 \geq 0$ such that for every $m > m_0$ there exists $(\mu_0, u_0) \in (0, \infty) \times \mathcal{S}_m^{\text{rad}}$ with the following properties:*

- (i) (μ_0, u_0) is a solution to (1.8), and it satisfies the Nehari identity (1.12);
- (ii) u_0 is a minimum (with Lagrange multiplier μ_0) on the L^2 -sphere and the L^2 -ball, i.e., $K(u_0) = \kappa^m = \ell^m < 0$; moreover, we have the formula

$$\kappa^m = \inf_{-\infty < \mu < \bar{\mu}_0} \left(a(\mu) - \mu \frac{m}{2} \right) = a(\mu_0) - \mu_0 \frac{m}{2}, \quad (1.20)$$

where $a(\mu)$ is the least action related to the problem with fixed frequency μ , see (3.3).

- (iii) (μ_0, u_0) is a Pohožaev minimum on the product space, i.e., $(\mu_0, u_0) \in \mathcal{P}^{\text{rad}}$ and $I^m(\mu_0, u_0) = d^m$; in particular, (μ_0, u_0) satisfies the Pohožaev identity (1.11).

As a consequence of the above relations, we have $\kappa^m = d^m$.

Under suitable additional assumptions about g , we highlight that the found solution can be chosen nonnegative and radially nonincreasing, see Proposition 1.12.

Remark 1.2 *We observe the following facts.*

- The mass threshold m_0 is given by formula (5.12).
- Notice that, thanks to our method, we are able to state that the found Lagrange multiplier μ_0 is positive and the energy κ^m is negative. At the same time, the found solution (μ_0, u_0) minimizes I^m in the whole Pohožaev space over $\mathbb{R} \times H_{\text{rad}}^s(\mathbb{R}^N)$, not only over $(0, +\infty) \times H_{\text{rad}}^s(\mathbb{R}^N)$.
- In Theorem 1.1, the solution is found for m sufficiently large. In Theorem 1.3 below, we find a solution for each $m > 0$; on the other hand, in that case, we have no information on the sign of the Lagrange multiplier or that of the energy, nor do we know whether the solution minimizes the energy over $\mathcal{D}_m^{\text{rad}}$ (which is coherent with Proposition 1.5 below).
- We highlight that we cannot state that u_0 is a critical point of $I^m(\mu_0, \cdot)$ due to the lack of regularity of the functional, but only that $I^m(\mu_0, \cdot)$ is differentiable at u_0 along test functions in $H^s(\mathbb{R}^N) \cap L^1(\mathbb{R}^N)$, cf. Lemma 2.3. In particular, the validity of the Nehari identity in Theorem 1.1 is not straightforward; in this case, thanks to our method, we are able to obtain such an identity.
- We observe that (1.20) is a “Legendre transform”-type formula which relates the L^2 -minimum value to the Pohožaev minimum value (with fixed frequency); similar formulas were already obtained in [28, 30, 34].

- We notice that, fixed $\rho > 0$, we can find a threshold $m_\rho \geq m_0$ such that, for each $m > m_\rho$ there exists $(\mu_\rho, u_\rho) \in (\rho, +\infty) \times \mathcal{S}_m^{\text{rad}}$ solution of (1.8). Indeed, it is sufficient to apply Theorem 1.1 to the nonlinearity $g_\rho(t) := g(t) - \rho t$, which satisfies (g0)–(g4).

Unlike [45], in our fractional setting, two main difficulties come into play. The first is due to the nonlocality of the problem: the fact that the operator $(-\Delta)^s$ does not preserve the supports of functions creates indeed problems in a direct minimization as in [45]. To deal with this issue, we handle the perturbed problem through a Lagrangian formulation (1.13) as done in [20, 34]: this approach has additional advantages, like showing some geometrical properties of the solutions, such as the minimality of \mathcal{P}^{rad} , formula (1.20), and the form of the threshold m_0 ; this last, in turn, allows us to show the existence of a solution for every mass m (see Theorem 1.3 below).

The second difficulty is due to the absence of a Pohožaev identity: as we already pointed out, it is an open problem to determine if the Pohožaev identity holds for general continuous nonlinearities g and general values of $s \in (0, 1)$. Under our assumptions (g0)–(g4), thus, we are not able to use this identity. As already mentioned, we first study the superlinear problem (with an ε -perturbation) and pass to the limit: by construction, the solutions u_ε that we build satisfy the Pohožaev identity, which gives some information to start with. On the other hand, in the limiting process, we get $u_\varepsilon \rightharpoonup u_0$, where u_0 can be proved to be a solution, but with no a priori minimality properties: in [45], it is thus crucially used, in order to pass to the strong limit and get the existence of an L^2 -minimum, that the Pohožaev identity holds for arbitrary solutions. This obstruction is here overcome by a finer analysis on the information of u_ε , in particular, regarding the interplay between the minimization over the L^2 -sphere and the L^2 -ball.

To deal with the superlinear perturbed prescribed-mass problem, in the spirit of [20, 34], we first study the free problem (i.e., when μ is fixed), following essentially [14, 18]: here we revise and refine both the existence results contained in [14, 20]. Additionally, since we are interested in the existence of a single solution, we simplify the argument of [20] by finding the solution through a *direct minimization on the Pohožaev set* (1.15): as a matter of fact, the minimality of the solution in [20] was obtained as a by-product of a minimax approach (see Remark 5.4), which requires a nontrivial deformation theory. Here we directly study the convergence of a (precisely chosen) minimizing sequence on \mathcal{P}^{rad} : this procedure is totally unrelated with the perturbation framework (see Remark 5.3), and it seems a new approach in literature in regards of product spaces, thus it has an interest of its own.

In Theorem 1.1, we show the existence of a minimizer for $m > 0$ sufficiently large; when the mass is small, by looking at the pure logarithmic case we expect an L^2 -minimum, possibly with a Lagrange multiplier and an energy with signs different from the ones in Theorem 1.1. This is the content of the next theorem (see also the more general Theorem 7.2). This result gives a fractional counterpart of [60], as well as a generalization for $s = 1$ by dropping the monotonicity therein.

Theorem 1.3 (Existence for all masses) *If (g0)–(g4) hold and, moreover,*

$$\lim_{t \rightarrow 0} \frac{g(t)}{t} = -\infty, \quad (1.21)$$

then for every $m > 0$ there exists a solution $(\mu_0, u_0) \in \mathbb{R} \times \mathcal{S}_m^{\text{rad}}$ of (1.8) satisfying the Nehari and Pohožaev identities and such that

$$I^m(\mu_0, u_0) = K(u_0) = a(\mu_0) - \frac{\mu_0}{2}m = \kappa^m = d^m = \inf_{-\infty < \mu < \bar{\mu}_0} \left(a(\mu) - \frac{\mu}{2}m \right).$$

Finally, if $\kappa^m \leq 0$, then $\mu_0 > 0$.

We remark that in Theorem 1.3, we are not able to determine the sign of the Lagrange multiplier μ_0 and of the energy κ^m . To obtain Theorem 1.3 we make use of a shifted problem, which still falls into the assumptions of Theorem 1.1 (here the generality of (g1) is essential), and rely on the representation formula (5.12) for the mass threshold m_0 .

Remark 1.4 (On the multiplicity) *We highlight that the present approach seems suitable for generalizations in the L^2 -critical (see [19]) and supercritical cases. Moreover, by following the minimax ideas in [20], one may have the additional advantage to possibly obtain a multiplicity result: indeed, for each $\varepsilon > 0$, with the same techniques (and genus arguments, see Remark 5.4) we can get the existence of multiple solutions $(u_k^\varepsilon)_{k \in \mathbb{N}}$; on the other hand, at the moment, we are not able to exclude the possibility that these solutions collapse to the same solution of the original problem when $\varepsilon \rightarrow 0^+$; we expect that some finer topological information on the genus of these solutions is needed.*

1.3 Qualitative results and examples

We present now some outcomes about the interplay among the various concepts of minimality we introduced, as well as some symmetry results. Most of them will not require (g3).

As highlighted in Subsection 1.1, we recall that, in some cases, the Lagrange multiplier of the equation might be negative. In the following, thus, for $\rho \in [-\infty, +\infty)$ we define

$$\mathcal{P}_{(\rho)} := \{(\mu, u) \in (\rho, +\infty) \times H^s(\mathbb{R}^N) \mid u \neq 0 \text{ and } P(\mu, u) = 0\} \quad (1.22)$$

and $\mathcal{P}_{(\rho)}^{\text{rad}} := \mathcal{P}_{(\rho)} \cap H_{\text{rad}}^s(\mathbb{R}^N)$. Observe that $\mathcal{P}_{(-\infty)} = \mathcal{P}$.

In the following propositions, we are not claiming that the mentioned infima are attained or finite merely under (g0)–(g2). On the other hand, minimizers can exist under different assumptions than those in Theorems 1.1 or 1.3: for instance, with proper hypotheses on m and g , they exist when $\limsup_{|t| \rightarrow +\infty} g(t)/(|t|^{\bar{p}-1}t) < +\infty$ (which is weaker than (g3)); see, e.g., [42].

Under suitable assumptions, the several notions of minimality we introduced coincide; notice that when $m > m_0$, we have $\kappa^m < 0$, which is a particular case of such assumptions.

Proposition 1.5 (Connections among infima) *Assume (g0)–(g2). We have*

$$\inf_{\mathcal{S}_m} K = \inf_{\mathcal{P}} I^m. \quad (1.23)$$

If $\rho \in \mathbb{R}$ and $\inf_{\mathcal{S}_m} K < -\frac{\rho}{2}m$, then

$$\inf_{\mathcal{S}_m} K = \inf_{\mathcal{P}} I^m = \inf_{\mathcal{P}_{(\rho)}} I^m.$$

If $\rho \geq 0$ and $\inf_{\mathcal{D}_m} K < -\frac{\rho}{2}m$, then

$$\inf_{\mathcal{D}_m} K = \inf_{\mathcal{S}_m} K = \inf_{\mathcal{P}} I^m = \inf_{\mathcal{P}_{(\rho)}} I^m.$$

In particular, if $\inf_{\mathcal{D}_m} K < 0$ we have $\inf_{\mathcal{D}_m} K = \inf_{\mathcal{S}_m} K$.

Finally, for every $\rho \in [-\infty, +\infty)$, we have

$$\inf_{\mathcal{P}_{(\rho)}} I^m = \inf \{I^m(\mu, u) \mid (\mu, u) \in (\rho, +\infty) \times H^s(\mathbb{R}^N), u \neq 0, P(\mu, u) \leq 0\}.$$

We show now a generalization of (1.20).

Proposition 1.6 (Legendre transform formula) Assume (g0)–(g2) and let $\rho \in [-\infty, \bar{\mu}_0)$. Then,

$$\inf_{\mathcal{P}(\rho)} I^m = \inf_{\rho < \mu < \bar{\mu}_0} \left(\inf_{\mathcal{P}_\mu} J_\mu - \frac{\mu}{2}m \right).$$

We move on to show that Pohožaev minima are indeed always weak solutions.

Proposition 1.7 (Pohožaev minima are solutions) Assume (g0)–(g2) and let $\rho \in [-\infty, +\infty)$. If $(\mu, u) \in \mathcal{P}(\rho)$ satisfies $I^m(\mu, u) = \min_{\mathcal{P}(\rho)} I^m$, then it is a solution to (1.8). Moreover, if $\min_{\mathcal{P}(\rho)} I^m < -\frac{\rho}{2}m$, then $I^m(\mu, u) = K(u) = \inf_{\mathcal{S}_m} K$.

We show now that if u is an L^2 -minimum, then u satisfies the Pohožaev identity; this is obtained without assuming extra regularity. See also Proposition 4.3 for some results on the Nehari identity.

Proposition 1.8 (Pohožaev identity for L^2 -minima) Assume (g0)–(g2). Let $u \in \mathcal{S}_m$ be such that $K(u) = \inf_{\mathcal{S}_m} K$. Then, u is a solution of (1.1) for some Lagrange multiplier $\mu \in \mathbb{R}$, and (μ, u) satisfies the Pohožaev identity. Moreover, for every $\rho \in [-\infty, -\frac{2}{m} \inf_{\mathcal{S}_m} K)$, we have $K(u) = I^m(\mu, u) = \inf_{\mathcal{P}(\rho)} I^m$.

As a consequence of Propositions 1.7 and 1.8 we have the following statement.

Corollary 1.9 Under (g0)–(g2), a function $u \in H^s(\mathbb{R}^N)$ is an L^2 -minimum with Lagrange multiplier $\mu \in \mathbb{R}$ if and only if (μ, u) is a Pohožaev minimum over \mathcal{P} .

Finally, we present a fractional and possibly strongly sublinear counterpart of [36]: notice that here we deal with Pohožaev minima instead of classical ground state solutions, which are indeed equivalent notions in the case treated in [36].

Proposition 1.10 (Least energy versus least action) Assume (g0)–(g2).

(i) Let $u \in \mathcal{S}_m$ be such that $K(u) = \inf_{\mathcal{S}_m} K$ with Lagrange multiplier $\mu \in \mathbb{R}$. Then $u \in \mathcal{P}_\mu$ is such that $J_\mu(u) = \inf_{\mathcal{P}_\mu} J_\mu$ and

$$\inf_{\mathcal{S}_m} K = \inf_{\mathcal{P}_\mu} J_\mu - \frac{\mu}{2}m. \quad (1.24)$$

(ii) Let

$$\mu \in \left\{ \nu \in \mathbb{R} \mid \exists v \in \mathcal{S}_m \text{ s.t. } K(v) = \inf_{\mathcal{S}_m} K \text{ with Lagrange multiplier } \nu \right\}.$$

Let now $u \in \mathcal{P}_\mu$ be such that $J_\mu(u) = \inf_{\mathcal{P}_\mu} J_\mu$. Then $u \in \mathcal{S}_m$ and $K(u) = \inf_{\mathcal{S}_m} K$.

Remark 1.11 We highlight that with similar proofs we may show the same results of Propositions 1.5–1.8, 1.10, and Corollary 1.9 in the radial setting, extending what we found in Theorem 1.1 for $m > m_0$.

We move now to some qualitative properties of the solutions. We highlight that, dealing with general solutions of a nonregular problem, their proofs are not standard and will require some additional arguments. Since, in Theorem 1.1, we restricted to the radial framework, we investigate when this property is naturally achieved, together with positivity.

Proposition 1.12 (Symmetry and sign) Assume (g0)–(g2) and let $\rho \in [-\infty, +\infty)$.

- (i) All minimizers of I^m over $\mathcal{P}_{(\rho)}$ and all minimizers of K over \mathcal{S}_m are radial about a point.
- (ii) All nonnegative minimizers of I^m over $\mathcal{P}_{(\rho)}$ or $\mathcal{P}_{(\rho)}^{\text{rad}}$ and all nonnegative minimizers of K over \mathcal{S}_m or $\mathcal{S}_m^{\text{rad}}$ are Schwarz-symmetric up to a translation (see Definition 2.5).
- (iii) Let (μ, u) be a solution to (1.8) such that $\int_{\mathbb{R}^N} g_-(u)u \, dx < +\infty$, and assume² $g(t) - \mu t \geq 0$ for each $t \leq 0$. Then $u \geq 0$.
- (iv) Assume, in addition, (g3) and (g4). If $g|_{(-\infty, 0)} = 0$ or g is odd, then for any $m > m_0$ there exists a Schwarz-symmetric minimizer (μ_0, u_0) of I^m over \mathcal{P} , where $\mu_0 > 0$ is also a minimizer of $\mu \in (-\infty, \bar{\mu}_0) \mapsto \inf_{\mathcal{P}_\mu} J_\mu - \frac{\mu}{2}m \in \mathbb{R}$ and u_0 is also a minimizer of K over \mathcal{S}_m and \mathcal{D}_m with Lagrange multiplier μ_0 . In particular,

$$\inf_{\mathcal{P}} I^m = d^m, \quad \inf_{\mathcal{S}_m} K = \kappa^m, \quad \inf_{\mathcal{D}_m} K = \ell^m,$$

$$\inf_{-\infty < \mu < \bar{\mu}_0} \left(\inf_{\mathcal{P}_\mu} J_\mu - \mu \frac{m}{2} \right) = \inf_{-\infty < \mu < \bar{\mu}_0} \left(a(\mu) - \mu \frac{m}{2} \right),$$

and all the above quantities are indeed equal and negative. This minimizer is a solution to (1.8) and satisfies the Nehari identity.

Remark 1.13 In a way analogous to the relation between Theorems 1.1 and 1.3, if (g0)–(g4) and (1.21) hold, then the outcome of Proposition 1.12 (iv), except for the signs of μ_0 and $K(u_0)$ and the statement on ℓ_m , holds for every $m > 0$.

When searching for a normalized solution without necessarily minimality properties, we can drop (g2); moreover, we can require the solution to be Schwarz symmetric (up to the sign).

Proposition 1.14 (Existence and symmetry under relaxed assumptions) Assume (g0), (g1), (g3), and (g4). Then, for every m sufficiently large, there exists a solution $(\mu_0, u_0) \in (0, +\infty) \times \mathcal{S}_m^{\text{rad}}$ of (1.8) such that u_0 or $-u_0$ is Schwarz symmetric³. This solution satisfies the Nehari identity (1.12) and the Pohožaev identity (1.11). Finally, if (1.21) holds as well, then such a solution exists with no restriction on m .

We end this section by giving some examples of nonlinearities to which the previous results apply. In these cases, we also show nonexistence of solutions.

Corollary 1.15 (Logarithmic case) Let $m > 0$ and

$$g(t) = \alpha t \log(t^2) + \beta |t|^{q-1} t$$

with $\alpha > 0$, $\beta \in \mathbb{R}$, and $0 < q \leq 2_s^* - 1$. Then the following facts hold.

- (i) If $q > 1$ and $\beta \leq -\frac{\alpha(q+1)}{q-1} e^{-(q+1)/2}$, then there exist no solutions $(\mu, u) \in [0, +\infty) \times \mathcal{S}_m$ to (1.8) that satisfy the Pohožaev identity; in particular, no L^2 -minima with nonnegative Lagrange multipliers μ exist.
- (ii) If

$$q \in (1, \bar{p}) \quad \text{and} \quad \beta > 0$$

² Notice that if $\mu \geq 0$, this is implied by $g|_{(-\infty, 0)} \geq 0$.

³ u_0 if $t_0 > 0$, $-u_0$ if $t_0 < 0$, where t_0 is introduced in (g4).

or

$$q \in (1, 2_s^* - 1] \quad \text{and} \quad -\frac{\alpha(q+1)}{q-1}e^{-(q+1)/2} < \beta \leq 0$$

or

$$q \in (0, 1] \quad \text{and} \quad \beta \leq 0,$$

then there exists $(\mu_0, u_0) \in \mathbb{R} \times \mathcal{S}_m^{\text{rad}}$, a Schwarz-symmetric minimizer of K over \mathcal{S}_m , which is a solution of (1.8), a Pohožaev minimum, and it satisfies the Nehari identity. Moreover, the following Legendre-transform formula holds:

$$\kappa^m = \inf_{\mu \in \mathbb{R}} \left(a(\mu) - \mu \frac{m}{2} \right) = a(\mu_0) - \mu_0 \frac{m}{2}.$$

Finally, if m_0 is as in Theorem 1.1 and $m > m_0$, then $\mu_0 > 0$ and $\ell^m = \kappa^m < 0$.

As already mentioned in Subsection 1.1, we highlight that also the pure logarithmic case $\alpha = 1$ and $\beta = 0$, that is, (1.7), has an interest of its own: indeed, even if the existence of a normalized solution can be achieved by scaling arguments, in the fractional setting the existence of an L^2 -minimum is not straightforward as in the local case, due to the lack of extremals for the fractional logarithmic Sobolev inequality.

Corollary 1.16 (Low-power case) *Let*

$$g(t) = -\gamma|t|^{r-1}t + \beta|t|^{q-1}t$$

with $\gamma, \beta > 0$, $0 < r < 1 < q < \bar{p}$. Then, Corollary 1.15 (ii) holds⁴.

The results apply also to nonlinearities of the type $\alpha t \log(t^2) - \gamma|t|^{r-1}t + \beta|t|^{q-1}t$, but in this case the exact range of coefficients for existence is not explicitly available.

Remark 1.17 *We notice that all the results we obtain are valid in the local case $s = 1$, $N \geq 3$, and this leads to some improvements of [45]: in particular, in Theorem 1.1, the minimality on \mathcal{P}^{rad} of the solution, the Nehari identity, and formula (1.20) are new; moreover, Theorem 1.3 and all the results of Subsection 1.3 are new as well. With some slight adaptations, Theorems 1.1 and 1.3 and most of the results of Subsection 1.3 can be generalized also to $s = 1$ and $N = 2$ (notice that the existence of an L^2 -minimum follows from [45, Theorem 1.1]). The only result that cannot be proved via a mere adaptation of the arguments herein is Proposition 1.12 (i) because Lemma 3.3 is false, being $\int G(u) - \frac{\mu}{2}u^2 \, dx = 0$ from the Pohožaev identity, cf. [8, Remark 3.1] and [7]. We leave the details to the interested reader.*

Remark 1.18 *Some of the results contained in this paper are new even in the superlinear setting, extending or further developing [20]. First, the growth of g at infinity is slightly generalized (see (g2)-(g3)). Then, the techniques involved for the existence – that is, the direct minimization over the Pohožaev set – are new as well. Furthermore, new results appear, such as the fact that every Pohožaev minimum is a critical point (see Proposition 1.7), or the precise relation between L^2 -minima and unconstrained ground states (see Proposition 1.10), together with other results in Subsection 1.3.*

⁴The results in Corollary 1.16 hold also for $0 < r < q \leq 1$. In this case, in the Legendre transform, we have $0 < \mu < \bar{\mu}_0$, with $\bar{\mu}_0 = 2(q-r) \left(\frac{(1+r)(1-q)}{\gamma} \right)^{\frac{1-q}{q-r}} \left(\frac{\beta}{(1-r)(1+q)} \right)^{\frac{1-r}{q-r}}$ if $q < 1$ and $\bar{\mu}_0 = \beta$ if $q = 1$.

Outline of the paper. In Section 2, we prove some preliminary properties that will be used throughout the paper. In Section 3, we recall and refine some results about the free-mass problem (that is, with fixed μ). In Section 4, together with Subsection 5.2, we prove most of the results from Subsection 1.3. In Section 5, we find a solution to the perturbed problem; then, in Section 6, passing to the limit, we obtain a solution to (1.8) and prove Theorem 1.1. Finally, in Section 7, we prove Theorem 1.3, Proposition 1.14, and Corollaries 1.15 and 1.16.

Notations. The definition of (weak) solution is given in Definition 2.4, while the definition of Schwarz-symmetric function is given in Definition 2.5. The fractional Sobolev space $H^s(\mathbb{R}^N)$ is defined in (2.1), while the restrictions over the radial subspace will be denoted with “rad”. The L^2 -sphere \mathcal{S}_m and the L^2 -ball \mathcal{D}_m are given in (1.14) and (1.19), while the L^2 -critical exponent $\bar{p}+1$ is defined in (1.2). The perturbed functions g_ε (and $\varphi_\varepsilon, G_\varepsilon^-, g_\varepsilon^-, G_\varepsilon$) are given in (5.2)–(5.3); the quantities related to the perturbed problem will appear with an “ ε ”. The energy functional $K(u)$ and the L^2 -minimal value (over the sphere or the ball) κ^m, ℓ^m are given in (1.3), (1.17), (1.19). The (action, Pohožaev) functionals $J_\mu(u), P_\mu(u)$ over $H^s(\mathbb{R}^N)$ are defined in (2.3), (3.2), while the (Lagrangian, Pohožaev) functionals $I^m(\mu, u), P(\mu, u)$ over $\mathbb{R} \times H^s(\mathbb{R}^N)$ are given in (1.13), (1.16). For technical reasons, the change of variable $\mu = e^\lambda$ in the frequency will be used when $\mu > 0$, and sometimes an abuse of notation will be employed. In particular, we will write (leading to no ambiguity) $I^m(\lambda, u) \equiv I^m(e^\lambda, u)$, $P(\lambda, u) \equiv P(e^\lambda, u)$. Thus, the (action, Lagrangian, Pohožaev) functionals $J(\lambda, u), I^m(\lambda, u), P(\lambda, u)$ over $\mathbb{R} \times H^s(\mathbb{R}^N)$ are given in (5.8) and at the beginning of Subsection 5.1. The Pohožaev (set, minimal action) $\mathcal{P}_\mu, a(\mu)$ are given in (3.1), (3.3), while $\mathcal{P}, \mathcal{P}_{(\rho)}, \mathcal{P}_+, d^m, d_+^m$ can be found in (1.15), (1.22), (5.6), (1.17), and (5.7). The frequency thresholds $\bar{\mu}_0, \bar{\lambda}_0$ appear in (1.18) and (5.10). The Legendre-type minimum b^m is given in (5.11), while the mass threshold m_0 appears in (5.12).

2 Preliminaries

We first re-write here g_\pm and G_\pm , introduced in (1.9)–(1.10), in different terms:

$$G_+(t) := \begin{cases} \int_0^t g^+(\tau) d\tau & \text{if } t \geq 0, \\ \int_t^0 g^-(\tau) d\tau & \text{if } t < 0, \end{cases} \quad G_-(t) := \begin{cases} \int_0^t g^-(\tau) d\tau & \text{if } t \geq 0, \\ \int_t^0 g^+(\tau) d\tau & \text{if } t < 0, \end{cases}$$

$$g_+(t) := G'_+(t) = \begin{cases} g^+(t) & \text{if } t \geq 0, \\ -g^-(t) & \text{if } t \leq 0, \end{cases} \quad g_-(t) := G'_-(t) = \begin{cases} g^-(t) & \text{if } t \geq 0, \\ -g^+(t) & \text{if } t \leq 0. \end{cases}$$

Clearly, we have $G_\pm(t) \geq 0$ and $g_\pm(t)t \geq 0$ for each $t \in \mathbb{R}$; moreover,

$$G = G_+ - G_-, \quad g = g_+ - g_-,$$

and

$$G_+(t) \geq G(t) \geq -G_-(t), \quad g_-(t)t \leq g(t)t \leq g_+(t)t \quad \text{for } t \in \mathbb{R}.$$

Generally, G_\pm do not coincide with G^\pm , and unlike G^\pm , we have $G_\pm \in \mathcal{C}^1(\mathbb{R})$. Finally, we notice that, if g is odd, then G_+, G_- are even.

We briefly recall that the fractional Laplacian $(-\Delta)^s u := \mathcal{F}^{-1}(|\xi|^{2s} \mathcal{F}(u))$, $s \in (0, 1)$, whenever u is regular enough, can be expressed pointwise as

$$(-\Delta)^s u(x) = C_{N,s} \int_{\mathbb{R}^N} \frac{u(x) - u(y)}{|x - y|^{N+2s}} dy,$$

for some suitable $C_{N,s} > 0$; we refer to [30, Section 1.2]. We recall the fractional Sobolev space

$$H^s(\mathbb{R}^N) := \{u \in L^2(\mathbb{R}^N) \mid |\cdot|^s \mathcal{F}(u) \in L^2(\mathbb{R}^N)\} \quad (2.1)$$

and define

$$\langle u, v \rangle := \int_{\mathbb{R}^N} (-\Delta)^{s/2} u (-\Delta)^{s/2} v \, dx = C_{N,s} \int_{\mathbb{R}^N} \frac{(u(x) - u(y))(v(x) - v(y))}{|x - y|^{N+2s}} \, dx \, dy.$$

It is well-known that $H^s(\mathbb{R}^N) \hookrightarrow L^q(\mathbb{R}^N)$, $q \in [2, 2_s^*]$, while $H_{\text{rad}}^s(\mathbb{R}^N) \hookrightarrow L^q(\mathbb{R}^N)$ is compact for $q \in (2, 2_s^*)$; here $H_{\text{rad}}^s(\mathbb{R}^N)$ stands for the subspace of $H^s(\mathbb{R}^N)$ composed of radially symmetric functions. We recall, moreover, the Gagliardo-Nirenberg inequality [47]

$$|u|_{\frac{\bar{p}+1}{p+1}}^{\bar{p}+1} \leq C_{\text{GN}} |(-\Delta)^{s/2} u|_2^2 |u|_2^{\bar{p}-1}, \quad (2.2)$$

where $C_{\text{GN}} > 0$ is optimal. We also recall some basic properties (see, e.g., [30, Lemma 1.4.1]).

Lemma 2.1 *Let $u \in H^s(\mathbb{R}^N)$. Then $u^\pm, |u| \in H^s(\mathbb{R}^N)$ with $|(-\Delta)^{s/2}(u^\pm)|_2 \leq |(-\Delta)^{s/2}u|_2$ and $|(-\Delta)^{s/2}|u|_2 \leq |(-\Delta)^{s/2}u|_2$. Finally, $\langle u, u^+ \rangle \geq \langle u^+, u^+ \rangle$ and $\langle u, u^- \rangle \leq \langle u^-, u^- \rangle$.*

We show the following elementary convergence result.

Lemma 2.2 *Let $v \in H^s(\mathbb{R}^N)$. Let $v_R := v\varphi_R$, where $\varphi_R := \varphi(\cdot/R)$ and $\varphi \in \mathcal{C}_c^\infty(\mathbb{R}^N)$ is chosen $0 \leq \varphi \leq 1$, $\varphi = 1$ in $B_1(0)$, $\varphi = 0$ in $\mathbb{R}^N \setminus B_2(0)$. Then,*

$$\int_{\mathbb{R}^N} (-\Delta)^{s/2} u (-\Delta)^{s/2} v_R \, dx \rightarrow \int_{\mathbb{R}^N} (-\Delta)^{s/2} u (-\Delta)^{s/2} v \, dx \quad \text{as } R \rightarrow +\infty$$

for each $u \in H^s(\mathbb{R}^N)$.

Proof. We compute

$$\begin{aligned} & \frac{1}{C_{N,s}} \int_{\mathbb{R}^N} (-\Delta)^{s/2} u (-\Delta)^{s/2} (v_R - v) \, dx \\ &= \int_{\mathbb{R}^{2N}} \frac{(u(x) - u(y))((\varphi_R(x) - 1)v(x) - (\varphi_R(y) - 1)v(y))}{|x - y|^{N+2s}} \, dx \, dy \\ &= \int_{\mathbb{R}^{2N}} (\varphi_R(x) - 1) \frac{(u(x) - u(y))(v(x) - v(y))}{|x - y|^{N+2s}} \, dx \, dy \\ & \quad + \int_{\mathbb{R}^{2N}} \frac{(u(x) - u(y))(\varphi_R(x) - \varphi_R(y))}{|x - y|^{N+2s}} v(y) \, dx \, dy =: (I_1) + (I_2). \end{aligned}$$

As regards (I_1) , we observe that

$$\frac{|u(x) - u(y)||v(x) - v(y)|}{|x - y|^{N+2s}} = \left(\frac{|u(x) - u(y)|^2}{|x - y|^{N+2s}} \right)^{1/2} \left(\frac{|v(x) - v(y)|^2}{|x - y|^{N+2s}} \right)^{1/2} \in L^1(\mathbb{R}^{2N}),$$

thus we can apply the dominated convergence theorem. Regarding (I_2) , we have

$$\begin{aligned} & \int_{\mathbb{R}^{2N}} \frac{(u(x) - u(y))(\varphi_R(x) - \varphi_R(y))}{|x - y|^{N+2s}} v(y) \, dx \, dy \\ & \leq \left(\int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^2}{|x - y|^{N+2s}} \, dx \, dy \right)^{\frac{1}{2}} \left(\int_{\mathbb{R}^N} \frac{|\varphi_R(x) - \varphi_R(y)|^2}{|x - y|^{N+2s}} v^2(y) \, dx \, dy \right)^{\frac{1}{2}}, \end{aligned}$$

and we can apply [3, Lemma 1.4.5]. ■

Lemma 2.3 Assume (g0)–(g2). Then the functional

$$u \in H^s(\mathbb{R}^N) \mapsto \int_{\mathbb{R}^N} G(u) \, dx \in \mathbb{R}$$

is differentiable at every $u \in H^s(\mathbb{R}^N)$ along every direction $\varphi \in H^s(\mathbb{R}^N) \cap L^1(\mathbb{R}^N)$.

Proof. We study $\lim_{t \rightarrow 0} \int_{\mathbb{R}^N} \frac{G(u+t\varphi) - G(u)}{t} \, dx$ and then apply Vitali's convergence theorem. We estimate thus

$$\begin{aligned} \left| \frac{G(u+t\varphi) - G(u)}{t} \right| &= \left| \frac{1}{t} \int_u^{u+t\varphi} g(\tau) \, d\tau \right| \leq \left| \frac{1}{t} \int_u^{u+t\varphi} 1 + |\tau|^{2_s^* - 1} \, d\tau \right| \\ &\leq |\varphi| + \left| \frac{1}{t} \int_u^{u+t\varphi} |\tau|^{2_s^* - 1} \, d\tau \right| \leq |\varphi| + \left| \frac{1}{t} \int_u^{u+t\varphi} (|u| + |\varphi|)^{2_s^* - 1} \, d\tau \right| \\ &= |\varphi| + (|u| + |\varphi|)^{2_s^* - 1} |\varphi| \lesssim |\varphi| + |u|^{2_s^* - 1} |\varphi| + |\varphi|^{2_s^*}, \end{aligned}$$

where we used the convexity of $t^{2_s^* - 1}$. Then, fixed $A \subset \mathbb{R}^N$ measurable, we have

$$\int_A \left| \frac{G(u+t\varphi) - G(u)}{t} \right| \, dx \lesssim \int_A |\varphi| \, dx + |u|_{2_s^*}^{2_s^* - 1} \left(\int_A |\varphi|^{2_s^*} \, dx \right)^{\frac{1}{2_s^*}} + \int_A |\varphi|^{2_s^*} \, dx$$

and thus we get both the uniform integrability and the tightness. This allows us to pass to the limit and get the claim. \blacksquare

Definition 2.4 By weak solution (or simply solution) of (1.1) we mean $u \in H^s(\mathbb{R}^N)$ such that

$$\int_{\mathbb{R}^N} (-\Delta)^{s/2} u (-\Delta)^{s/2} \varphi \, dx + \mu \int_{\mathbb{R}^N} u \varphi \, dx = \int_{\mathbb{R}^N} g(u) \varphi \, dx$$

for each $\varphi \in H^s(\mathbb{R}^N) \cap L^1(\mathbb{R}^N)$; in particular, it holds for every $\varphi \in C_c^\infty(\mathbb{R}^N)$.

We highlight that, generally, a solution u of (1.1) is not a critical point of the action functional $J_\mu: H^s(\mathbb{R}^N) \rightarrow \mathbb{R}$,

$$J_\mu(u) := \frac{1}{2} \int_{\mathbb{R}^N} |(-\Delta)^{s/2} u|^2 \, dx + \frac{\mu}{2} \int_{\mathbb{R}^N} u^2 \, dx - \int_{\mathbb{R}^N} G(u) \, dx, \quad (2.3)$$

since, generally, J_μ is not differentiable at u along every $\varphi \in H^s(\mathbb{R}^N)$ (see Lemma 2.3); in particular, we cannot ensure that a solution satisfies the Nehari identity. On the other hand, in some circumstances, we will show that this is the case (cf. Proposition 4.3).

Finally, we will use the following definition.

Definition 2.5 A function $u: \mathbb{R}^N \rightarrow \mathbb{R}$ is called Schwarz-symmetric if it is nonnegative, radially symmetric, and radially nonincreasing.

3 The mass-free problem

The purpose of this section is to discuss the free-mass problem (i.e., μ fixed). For any $\mu \in \mathbb{R}$ we introduce a Pohožaev set

$$\mathcal{P}_\mu := \{u \in H^s(\mathbb{R}^N) \setminus \{0\} \mid P_\mu(u) = 0\}, \quad \mathcal{P}_\mu^{\text{rad}} := \mathcal{P}_\mu \cap H_{\text{rad}}^s(\mathbb{R}^N), \quad (3.1)$$

where the Pohožaev functional $P_\mu: H^s(\mathbb{R}^N) \rightarrow \mathbb{R}$ is defined by

$$P_\mu(u) := \int_{\mathbb{R}^N} |(-\Delta)^{s/2}u|^2 dx + 2_s^* \int_{\mathbb{R}^N} \frac{\mu}{2}u^2 - G(u) dx. \quad (3.2)$$

Recalled (2.3), we define the least action

$$a(\mu) := \inf_{\mathcal{P}_\mu^{\text{rad}}} J_\mu. \quad (3.3)$$

We notice that, generally, the set \mathcal{P}_μ could be empty, and in this case, $a(\mu) = +\infty$.

Here we list a series of properties that hold for $\mu \in \mathbb{R}$ under (g0)–(g2).

Lemma 3.1 *Assume (g0)–(g2) and let $\mu \in \mathbb{R}$. Then, the following statements are equivalent.*

- (a) $\mu < \bar{\mu}_0$.
- (b) There exists $t = t(\mu) \neq 0$ such that $G(t) > \frac{1}{2}\mu t^2$.
- (c) There exists $u \in \mathcal{P}_\mu^{\text{rad}}$.

Proof. (a) \iff (b) is trivial. If (c) holds, then $u \neq 0$; thus, by the Pohožaev identity,

$$\int_{\mathbb{R}^N} G(u) - \frac{\mu}{2}u^2 dx = \int_{\mathbb{R}^N} |(-\Delta)^{s/2}u|^2 dx > 0,$$

hence there exists $t(\mu) := u(x_\mu)$, $x_\mu \in \mathbb{R}^N$, such that (b) is verified. Assume (b) holds, from [8, Proof of Theorem 2] (see also [20, Lemma 3.1]) there exists $\bar{u} \in H_{\text{rad}}^1(\mathbb{R}^N) \subset H_{\text{rad}}^s(\mathbb{R}^N)$ such that $\int_{\mathbb{R}^N} (G(\bar{u}) - \frac{\mu}{2}\bar{u}^2) dx > 0$. Then, for $t > 0$, straightforward computations show that

$$P_\mu(\bar{u}(\cdot/t)) = \left(\int_{\mathbb{R}^N} |(-\Delta)^{s/2}\bar{u}|^2 dx \right) t^{N-2s} + 2_s^* \left(\int_{\mathbb{R}^N} \frac{\mu}{2}\bar{u}^2 - G(\bar{u}) dx \right) t^N$$

is positive when t is small and negative when t is large, hence there exists $\bar{t} > 0$ such that $P_\mu(\bar{u}(\cdot/\bar{t})) = 0$, whence (c). \blacksquare

Remark 3.2 *By the previous proof we notice the following fact: for every $u \in H^s(\mathbb{R}^N) \setminus \{0\}$ such that $P_\mu(u) \leq 0$, we have $\int_{\mathbb{R}^N} (G(u) - \frac{\mu}{2}u^2) dx > 0$, thus there exists $t_0 > 0$ such that $u(\cdot/t_0) \in \mathcal{P}_\mu$. Moreover, using that $P_\mu(u) \leq 0$ and that $t \in (0, \infty) \mapsto \frac{1}{2}t^{N-2s} - \frac{1}{2_s^*}t^N \in \mathbb{R}$ is maximized at $t = 1$, there holds*

$$\begin{aligned} J_\mu(u(\cdot/t_0)) &= \frac{t_0^{N-2s}}{2} \int_{\mathbb{R}^N} |(-\Delta)^s u|^2 dx - t_0^N \int_{\mathbb{R}^N} G(u) - \frac{\mu}{2}u^2 dx \\ &\leq \left(\frac{t_0^{N-2s}}{2} - \frac{t_0^N}{2_s^*} \right) \int_{\mathbb{R}^N} |(-\Delta)^s u|^2 dx \leq \frac{s}{N} \int_{\mathbb{R}^N} |(-\Delta)^s u|^2 dx \leq J_\mu(u), \end{aligned}$$

that is, $J_\mu(u(\cdot/t_0)) \leq J_\mu(u)$. Notice that this fact does not depend on the sign of μ .

Lemma 3.3 *Let $\mu \in \mathbb{R}$ and assume (g0)–(g2). If $\mathcal{P}_\mu \neq \emptyset$ and $u \in \mathcal{P}_\mu$ is such that $J_\mu(u) = \inf_{\mathcal{P}_\mu} J_\mu$, then $\omega := \int_{\mathbb{R}^N} G(u) - \frac{\mu}{2}u^2 dx > 0$ and*

$$\frac{1}{2}|(-\Delta)^{s/2}u|_2^2 = \min \left\{ \frac{1}{2}|(-\Delta)^{s/2}v|_2^2 \mid v \in H^s(\mathbb{R}^N), \int_{\mathbb{R}^N} G(v) - \frac{\mu}{2}v^2 dx = \omega \right\}. \quad (3.4)$$

Proof. First, notice that since $\mathcal{P}_\mu \neq \emptyset$, there exists a function $v \in H^s(\mathbb{R}^N)$ such that $\int_{\mathbb{R}^N} G(v) - \frac{\mu}{2}v^2 dx = 1$, thus the minimization problem

$$T := \inf \left\{ \frac{1}{2} |(-\Delta)^{s/2} v|_2^2 \mid v \in H^s(\mathbb{R}^N), \int_{\mathbb{R}^N} G(v) - \frac{\mu}{2} v^2 dx = 1 \right\}$$

is well defined. Moreover, every such function verifies $G(v) \in L^1(\mathbb{R}^N)$. We adapt now the proof of [13, Lemma 1 (ii)]. We notice that we do not need to assume that the Pohožaev identity holds for every solution or that T is a priori attained. Since $u \in \mathcal{P}_\mu$, we have that

$$J_\mu(u) \geq 2s(N - 2s)^{\frac{N}{2s} - 1} N^{-\frac{N}{2s}} T^{\frac{N}{2s}} =: \tilde{T}.$$

Now, let $(v_n)_n$ be a minimizing sequence for T , and let $t_n > 0$ be such that $v_n(\cdot/t_n) \in \mathcal{P}_\mu$. Then, since $J_\mu(u) = \inf_{\mathcal{P}_\mu} J_\mu$,

$$\tilde{T} \leq J_\mu(u) \leq J_\mu(v_n(\cdot/t_n)) = \tilde{T} + o_n(1),$$

and we conclude letting $n \rightarrow +\infty$ and via scaling arguments. ■

Proposition 3.4 *Assume (g0)–(g2) and let $\mu \in \mathbb{R}$. Then $a(\mu) \geq 0$ and the following relation holds*

$$a(\mu) = \inf \left\{ J_\mu(u) \mid u \in H_{\text{rad}}^s(\mathbb{R}^N) \setminus \{0\}, P_\mu(u) \leq 0 \right\}. \quad (3.5)$$

Moreover, $a: \mathbb{R} \rightarrow [0, +\infty]$ is nondecreasing. Finally, if $\mu > 0$ and $\mathcal{P}_\mu^{\text{rad}} \neq \emptyset$, then a is continuous at μ .

Proof. By definition of $\mathcal{P}_\mu^{\text{rad}}$, we have

$$a(\mu) = \frac{s}{N} \inf_{u \in \mathcal{P}_\mu^{\text{rad}}} |(-\Delta)^{s/2} u|_2^2 \geq 0.$$

To show relation (3.5), observe that if $\mathcal{P}_\mu^{\text{rad}} = \emptyset$, then $\{u \in H_{\text{rad}}^s(\mathbb{R}^N) \setminus \{0\} \mid P_\mu(u) \leq 0\} = \emptyset$ as well from Remark 3.2, hence both sides of (3.5) equal $+\infty$. Next, assume $P_\mu(u) \leq 0$; thus, by Remark 3.2 there exists $t_0 > 0$ such that $P_\mu(u(\cdot/t_0)) = 0$, and moreover, $J_\mu(u(\cdot/t_0)) \leq J_\mu(u)$, hence

$$a(\mu) \leq J_\mu(u(\cdot/t_0)) \leq J_\mu(u)$$

and we can pass to the infimum. The reverse inequality is obvious. We show now the monotonicity. Let $\mu_1 < \mu_2$. If $\mathcal{P}_{\mu_2}^{\text{rad}} = \emptyset$, then it clearly holds. Otherwise, let $u \in \mathcal{P}_{\mu_2}^{\text{rad}}$. Then $P_{\mu_1}(u) < P_{\mu_2}(u) = 0$ and by (3.5) we have

$$a(\mu_1) \leq J_{\mu_1}(u) < J_{\mu_2}(u); \quad (3.6)$$

passing to the infimum, we have the claim. We show now the continuity. The upper semicontinuity follows by the fact that $a(\mu)$ is the infimum of the family of continuous functions $\{\mu \mapsto J_\mu(u) \mid u \in H_{\text{rad}}^s(\mathbb{R}^N) \setminus \{0\}, P_\mu(u) = 0\}$. We focus on the lower semicontinuity and fix $\mu_n \rightarrow \mu > 0$. From (g0) and Lemma 3.1, $\mathcal{P}_{\mu_n}^{\text{rad}} \neq \emptyset$ for all sufficiently large n ; in addition, we can assume $\mu^* := \inf_n \mu_n > 0$, and by upper semicontinuity, $M := \sup_n a(\mu_n) < \infty$. For each n , let $(u_k^n)_k \subset H_{\text{rad}}^s(\mathbb{R}^N) \setminus \{0\}$ be a minimizing sequence for $a(\mu_n)$, that is, $P_{\mu_n}(u_k^n) = 0$ and $J_{\mu_n}(u_k^n) \rightarrow a(\mu_n)$ as $k \rightarrow +\infty$. Let now $v_n := u_{k_n}^n$ be such that $J_{\mu_n}(v_n) \leq a(\mu_n) + \frac{1}{n} \leq M + 1$. We show that $(v_n)_n$ is bounded. Indeed, by the Pohožaev identity we have

$$\frac{s}{N} |(-\Delta)^{s/2} v_n|_2^2 = J_{\mu_n}(v_n) \leq M + 1$$

and, again by the Pohožaev identity and the assumptions on G_+ ,

$$\frac{\mu^*}{2}|v_n|_2^2 \leq \frac{\mu_n}{2}|v_n|_2^2 \leq \int_{\mathbb{R}^N} G(v_n) dx + C \leq \frac{\delta}{2}|v_n|_2^2 + C_\delta|v_n|_{2_s^*}^{2_s^*} dx + C \leq \frac{\delta}{2}|v_n|_2^2 + C,$$

where, by choosing $\delta < \mu^*$, we obtain the boundedness of v_n in $H^s(\mathbb{R}^N)$. Thus, $v_n \rightharpoonup v$ in $H_{\text{rad}}^s(\mathbb{R}^N)$, which yields $P_\mu(v) \leq \liminf_n P_{\mu_n}(v_n) = 0$ and, in view of (3.5),

$$a(\mu) \leq J_\mu(v) \leq \liminf_n J_{\mu_n}(v_n) \leq \liminf_n a(\mu_n),$$

which is the claim. ■

Proposition 3.5 *Assume (g0)–(g2) and let $\mu \in \mathbb{R}$.*

(i) *All minimizers of J_μ over $\mathcal{P}_\mu^{\text{rad}}$ (resp. \mathcal{P}_μ) are solutions to*

$$(-\Delta)^s u + \mu u = g(u) \quad \text{in } \mathbb{R}^N. \quad (3.7)$$

(ii) *All minimizers of J_μ over \mathcal{P}_μ are radially symmetric about a point.*

(iii) *All nonnegative minimizers of J_μ over \mathcal{P}_μ or $\mathcal{P}_\mu^{\text{rad}}$ are Schwarz-symmetric up to a translation.*

(iv) *If $g(t) - \mu t \geq 0$ for all $t \leq 0$, then every solution to (3.7) such that $\int_{\mathbb{R}^N} g_-(u)u dx < +\infty$ is nonnegative.*

The proof of Proposition 3.5 is analogous to those of Propositions 1.7 and 1.12 (see Subsections 4.2, 5.2 and 6.2), hence it is omitted. We now state the main theorem of this section.

Theorem 3.6 *Let $\mu > 0$ and $g \in \mathcal{C}(\mathbb{R})$ satisfy*

- $|g(t)| \lesssim |t| + |t|^{2_s^*-1}$ for every $t \in \mathbb{R}$;
- $\limsup_{t \rightarrow 0} \frac{G(t)}{t^2} \leq 0$ and $\limsup_{|t| \rightarrow \infty} \frac{G(t)}{|t|^{2_s^*}} \leq 0$;
- there exists $t_0 \neq 0$ such that $G(t_0) > \frac{1}{2}\mu t_0^2$.

Then the following statements hold true.

(i) *There exists $v_0 \in \mathcal{P}_\mu^{\text{rad}}$ such that*

$$J_\mu(v_0) = a(\mu)$$

and thus $a(\mu) > 0$. Moreover, $a: (0, +\infty) \rightarrow (0, +\infty)$ is increasing.

(ii) *If $g|_{(-\infty, 0)} = 0$ or g is odd, then there exists a Schwarz-symmetric minimizer of J_μ over \mathcal{P}_μ ; in particular,*

$$\inf_{\mathcal{P}_\mu} J_\mu = \inf_{\mathcal{P}_\mu^{\text{rad}}} J_\mu.$$

Proof.

(i) From the assumptions, we see that for every $\delta > 0$ there exists $C_\delta > 0$ such that for every $t \in \mathbb{R}$

$$G(t) \leq \delta t^2 + C_\delta |t|^{2_s^*},$$

whence, for every $u \in H^s(\mathbb{R}^N)$,

$$\frac{1}{2_s^*} P_\mu(u) \geq \frac{1}{2_s^*} |(-\Delta)^{s/2} u|_2^2 + \frac{\mu}{2} |u|_2^2 - \delta |u|_2^2 - C_\delta |u|_{2_s^*}^{2_s^*} \geq \left(\min \left\{ \frac{1}{2_s^*}, \frac{\mu}{2} \right\} - \delta \right) \|u\|_{H^s}^2 - C_\delta \|u\|_{H^s}^{2_s^*}.$$

Taking $0 < \delta < \min\{1/2_s^*, \mu/2\}$, we see that $P_\mu(u) > 0$ if $\|u\|_{H^s}$ is sufficiently small.

Let now $(u_n)_n \subset \mathcal{P}_\mu^{\text{rad}}$ such that $\lim_n J_\mu(u_n) = \inf_{\mathcal{P}_\mu^{\text{rad}}} J_\mu$. Since $P_\mu(u_n) = 0$, we see that

$$\inf_{\mathcal{P}_\mu} J_\mu + o(1) = J_\mu(u_n) - \frac{1}{2_s^*} P_\mu(u_n) = \frac{s}{N} |(-\Delta)^{s/2} u_n|_2^2,$$

which yields that $|(-\Delta)^{s/2} u_n|_2$ is bounded, and, using the same δ -related argument as before, that

$$\begin{aligned} \inf_{\mathcal{P}_\mu} J_\mu + o(1) &= J_\mu(u_n) - \frac{1}{2} P_\mu(u_n) = -\frac{\mu}{2} \left(\frac{2_s^*}{2} - 1 \right) |u_n|_2^2 + \left(\frac{2_s^*}{2} - 1 \right) \int_{\mathbb{R}^N} G(u_n) \, dx \\ &\leq -\frac{\mu}{2} \left(\frac{2_s^*}{2} - 1 \right) |u_n|_2^2 + \left(\frac{2_s^*}{2} - 1 \right) \left(\delta |u_n|_2^2 + C_\delta |u_n|_{2_s^*}^{2_s^*} \right). \end{aligned}$$

Since $|u_n|_{2_s^*}$ is bounded from the previous step and the fractional Sobolev embedding, taking $\delta < \mu/2$ we obtain that $|u_n|_2$ is bounded as well, i.e., $\|u_n\|_{H^s}$ is bounded. We can then consider a subsequence, still denoted by u_n , and $u_0 \in H_{\text{rad}}^s(\mathbb{R}^N)$ such that $u_n \rightharpoonup u_0$ in $H^s(\mathbb{R}^N)$, $u_n \rightarrow u_0$ a.e. in \mathbb{R}^N , and $u_n \rightarrow u_0$ in $L^q(\mathbb{R}^N)$ for every $q \in (2, 2_s^*)$. In particular, using the weak lower semi-continuity of the norm and Fatou's lemma, we get that $J_\mu(u_0) \leq \inf_{\mathcal{P}_\mu^{\text{rad}}} J_\mu$ and $P_\mu(u_0) \leq 0$. In addition, recalling $P_\mu(u_n) = 0$, we have by the strong convergence in $L^q(\mathbb{R}^N)$,

$$\begin{aligned} |(-\Delta)^{s/2} u_0|_2^2 + \frac{2_s^* \mu}{2} |u_0|_2^2 + 2_s^* \int_{\mathbb{R}^N} G_-(u_0) \, dx \\ \leq \limsup_n \left(|(-\Delta)^{s/2} u_n|_2^2 + \frac{2_s^* \mu}{2} |u_n|_2^2 + 2_s^* \int_{\mathbb{R}^N} G_-(u_n) \, dx \right) \\ = \limsup_n \int_{\mathbb{R}^N} G_+(u_n) \, dx = \int_{\mathbb{R}^N} G_+(u_0) \, dx; \end{aligned}$$

if $u_0 \equiv 0$, the above relation implies $u_n \rightarrow 0$ in $H^s(\mathbb{R}^N)$, which is impossible because $P_\mu(u) > 0$ for small u . Since $u_0 \neq 0$, by Remark 3.2 we can find $t_0 > 0$ such that $P_\mu(u_0(\cdot/t_0)) = 0$ and $J_\mu(u_0(\cdot/t_0)) \leq J_\mu(u_0)$, so the sought minimizer is $v_0 := u_0(\cdot/t_0)$. By the Pohožaev identity we also have $a(\mu) = J_\mu(v_0) = \frac{s}{N} \|(-\Delta)^{s/2} v_0\|_2^2 > 0$.

To show the strict monotonicity, arguing as in (3.6), in the case $\mu_2 > 0$ we choose u as the minimum corresponding to μ_2 . Hence, we have

$$a(\mu_1) \leq J_{\mu_1}(u) < J_{\mu_2}(u) = a(\mu_2).$$

(ii) We start with the case when g is odd. Let $(u_n)_n \subset \mathcal{P}_\mu$ be such that $\lim_n J_\mu(u_n) = \inf_{\mathcal{P}_\mu} J_\mu$. From Lemma 2.1, $J_\mu(|u_n|) \leq J_\mu(u_n)$ and $P_\mu(|u_n|) \leq P_\mu(u_n) = 0$. If $t_n > 0$ is such that $P_\mu(|u_n|(\cdot/t_n)) = 0$, then we obtain $\inf_{\mathcal{P}_\mu} J_\mu \leq J_\mu(|u_n|(\cdot/t_n)) \leq J_\mu(|u_n|) \leq J_\mu(u_n)$ by Remark 3.2. This proves that, up to replacing u_n with $|u_n|(\cdot/t_n)$, we can assume that $u_n \geq 0$ for every n . In a very similar way, using the Schwarz rearrangement and the fractional Polyá–Szegő inequality [46], we can assume that every u_n is Schwarz symmetric. We can now conclude in the very same way as in point (i). If $g|_{(-\infty, 0)} = 0$, we argue as above, except we use u_n^+ instead of $|u_n|$. ■

Remark 3.7 *We observe the following facts.*

(i) Theorem 3.6 (i) is a refinement of [14, Theorem 1.2 (i)], where the main additional assumption is the existence of $q \in (2, 2_s^*)$ such that $\lim_{t \rightarrow 0} g(t)/|t|^{q-1} = 0$. At the same time, Proposition 3.5 shows that [14, Theorem 1.2 (ii)–(iii)] holds under considerably weaker assumptions (in particular, (g1)).

(ii) In Proposition 3.5 (i), we are not claiming that, for general g , there exists a minimizer over \mathcal{P}_μ ; on the other hand, from Proposition 3.5 (ii), we deduce that $\mathcal{P}_\mu^{\text{rad}}$ is a natural subset where to search for such a minimizer.

4 Some general properties

For $\rho \in [-\infty, +\infty)$, recall \mathcal{P}_ρ from (1.22) and $\bar{\mu}_0$ from (1.18). We observe that, by Lemma 3.1, if $\rho \in [-\infty, \bar{\mu}_0)$, then $\mathcal{P}_\rho \supset \mathcal{P}_\rho^{\text{rad}} \neq \emptyset$; in particular, $\mathcal{P} \neq \emptyset$ always holds.

Remark 4.1 We point out that if (μ, u) minimizes I^m over \mathcal{P}_ρ for some $\rho \in [-\infty, +\infty)$, then u minimizes J_μ over \mathcal{P}_μ .

4.1 Connection among infima

We prove now Propositions 1.5 and 1.6.

Proof of Proposition 1.5. Let $(\mu_n, u_n)_n \subset \mathcal{P}$ be a minimizing sequence for $\inf_{\mathcal{P}} I^m$, that is, $I^m(\mu_n, u_n) = \inf_{\mathcal{P}} I^m + o(1)$; recall that $u_n \neq 0$ by the definition of \mathcal{P} . Consider $t_n := (m/|u_n|_2^2)^{1/N}$ such that $v_n := u_n(\cdot/t_n) \in \mathcal{S}_m$. We notice, by the Pohožaev identity,

$$\begin{aligned} I^m(\mu_n, u_n) &= \frac{1}{2}|(-\Delta)^{s/2}u_n|_2^2 + \frac{\mu_n|u_n|_2^2}{2}(1 - t_n^N) - \int_{\mathbb{R}^N} G(u_n) \, dx \\ &= \frac{1}{2}|(-\Delta)^{s/2}u_n|_2^2 + (1 - t_n^N) \left(\int_{\mathbb{R}^N} G(u_n) \, dx - \frac{1}{2_s^*}|(-\Delta)^{s/2}u_n|_2^2 \right) - \int_{\mathbb{R}^N} G(u_n) \, dx \\ &= \left(\frac{1}{2} - \frac{1}{2_s^*}(1 - t_n^N) \right) |(-\Delta)^{s/2}u_n|_2^2 - t_n^N \int_{\mathbb{R}^N} G(u_n) \, dx \\ &= \left(\frac{1}{2} - \frac{1}{2_s^*}(1 - t_n^N) \right) \frac{1}{t_n^{N-2s}} |(-\Delta)^{s/2}v_n|_2^2 - \int_{\mathbb{R}^N} G(v_n) \, dx \\ &\geq \frac{1}{2}|(-\Delta)^{s/2}v_n|_2^2 - \int_{\mathbb{R}^N} G(v_n) \, dx = K(v_n), \end{aligned}$$

where we used that $\frac{1}{2} - \frac{1}{2_s^*}(1 - t^N) \geq \frac{1}{2}t^{N-2s}$ for all $t > 0$, which is true by a straightforward computation. Thus, we have shown

$$\inf_{\mathcal{S}_m} K \leq K(v_n) \leq I^m(\mu_n, u_n) = \inf_{\mathcal{P}} I^m + o(1).$$

Passing to the limit in n , we obtain

$$\inf_{\mathcal{S}_m} K \leq \inf_{\mathcal{P}} I^m. \quad (4.1)$$

Assume now $\rho \in [0, +\infty)$ and $\inf_{\mathcal{D}_m} K < -\frac{\rho}{2}m$. We can thus choose a minimizing sequence $(u_n)_n \subset \mathcal{D}_m$ such that $\lim_n K(u_n) = \inf_{\mathcal{D}_m} K$ and $K(u_n) < -\frac{\rho}{2}m \leq 0$ for every n (which implies $u_n \neq 0$) and define

$$\mu_n := \frac{2}{N|u_n|_2^2} \left(s|(-\Delta)^{s/2}u_n|_2^2 - NK(u_n) \right) \geq -\frac{2}{|u_n|_2^2} K(u_n) > \rho$$

so that $(\mu_n, u_n) \in \mathcal{P}(\rho)$. This, together with $\mu_n > 0$ and $|u_n|_2^2 \leq m$, yields

$$\inf_{\mathcal{D}_m} K + o_n(1) = K(u_n) \geq I^m(\mu_n, u_n) \geq \inf_{\mathcal{P}(\rho)} I^m,$$

and passing to the limit in n we get $\inf_{\mathcal{D}_m} K \geq \inf_{\mathcal{P}(\rho)} I^m$. Arguing similarly for $\rho \in \mathbb{R}$ and $\inf_{\mathcal{S}_m} K < -\frac{\rho}{2}m$ we obtain $\inf_{\mathcal{S}_m} K \geq \inf_{\mathcal{P}(\rho)} I^m$. Choosing $\rho = -\infty$ and recalling (4.1), we get (1.23). Gathering the previous inequalities we conclude the proof.

Finally, we show the last relation arguing as in the proof of (3.5). Indeed, considered $(\mu, u) \in (\rho, +\infty) \times H^s(\mathbb{R}^N)$ with $P(\mu, u) \leq 0$, there exists $t_0 > 0$ such that $P(\mu, u(\cdot/t_0)) = 0$ and $I^m(\mu, u(\cdot/t_0)) \leq I^m(\mu, u)$, thus

$$\inf_{\mathcal{P}(\rho)} I^m \leq I^m(\mu, u(\cdot/t_0)) \leq I^m(\mu, u)$$

and hence the claim passing to the infimum (the other inequality is obvious). \blacksquare

Proof of Proposition 1.6. Let $\mu \in (\rho, \bar{\mu}_0)$. From Lemma 3.1, $\mathcal{P}_\mu \neq \emptyset$ (thus $\inf_{\mathcal{P}_\mu} J_\mu$ is well defined). Then, for any $u \in \mathcal{P}_\mu$, we have

$$\inf_{\mathcal{P}(\rho)} I^m \leq I^m(\mu, u) = J_\mu(u) - \frac{\mu}{2}m.$$

Passing to the infimum over u , we have $\inf_{\mathcal{P}(\rho)} I^m \leq \inf_{\mathcal{P}_\mu} J_\mu - \frac{\mu}{2}m$. Passing to the infimum over μ , we have $\inf_{\mathcal{P}(\rho)} I^m \leq \inf_{\rho < \mu < \bar{\mu}_0} (\inf_{\mathcal{P}_\mu} J_\mu - \frac{\mu}{2}m)$.

Let now $(\nu, u) \in \mathcal{P}_\rho$, which implies $\nu \in (\rho, \bar{\mu}_0)$. As a consequence

$$I^m(\nu, u) = J_\nu(u) - \frac{\nu}{2} \geq \inf_{\mathcal{P}_\nu} J_\nu - \frac{\nu}{2} \geq \inf_{\rho < \mu < \bar{\mu}_0} \left(\inf_{\mathcal{P}_\mu} J_\mu - \frac{\mu}{2}m \right).$$

Passing to the infimum over (ν, u) we obtain $\inf_{\mathcal{P}(\rho)} I^m \geq \inf_{\rho < \mu < \bar{\mu}_0} (\inf_{\mathcal{P}_\mu} J_\mu - \frac{\mu}{2}m)$, whence the claim. \blacksquare

4.2 L^2 -minima, Pohožaev identity, and critical points

Proposition 4.2 *Assume (g0)–(g2). Let $u \in H^s(\mathbb{R}^N)$.*

- (i) *If $u \in \mathcal{S}_m$ is an L^2 -minimum for K , i.e., $K(u) = \inf_{\mathcal{S}_m} K$, then u is a solution of (1.1) for some Lagrange multiplier $\mu \in \mathbb{R}$, and (μ, u) satisfies the Pohožaev identity.*
- (ii) *If $u \in \mathcal{D}_m$ is such that $K(u) = \inf_{\mathcal{D}_m} K$ and $|u|_2^2 < m$, then u is a solution of (1.1) with Lagrange multiplier $\mu = 0$, and (μ, u) satisfies the Pohožaev identity.*
- (iii) *If $(\mu, u) \in \mathcal{P}$ and $K(u) \leq 0$, then $\mu > 0$.*
- (iv) *If $u \in \mathcal{D}_m$ is such that $K(u) = \inf_{\mathcal{D}_m} K < 0$, then $u \in \mathcal{S}_m$; in particular, $K(u) = \inf_{\mathcal{S}_m} K$.*

Proof. (i) First, we show that u is a solution. Indeed, let $\varphi \in H^s(\mathbb{R}^N) \cap L^1(\mathbb{R}^N)$; being $u \neq 0$, for $t \in \mathbb{R}$ with $|t|$ small we can assume $u + t\varphi \neq 0$. Thus, for each t we choose $\theta(t) > 0$ such that $|(u + t\varphi)(\cdot/\theta(t))|_2^2 = m$, that is, $\theta(t) := (m|u + t\varphi|_2^{-2})^{1/N}$. Let $\gamma(t) := K((u + t\varphi)(\cdot/\theta(t)))$; due to the differentiability of θ and the arguments of the proof of Lemma 2.3, we have that γ is differentiable. Being $t = 0$ a point of minimum for γ , we have $\gamma'(0) = 0$, that is (notice $\theta(0) = 1$),

$$\frac{N - 2s}{2} \theta'(0) |(-\Delta)^{s/2} u|_2^2 + \int_{\mathbb{R}^N} (-\Delta)^{s/2} u (-\Delta)^{s/2} \varphi \, dx - N \theta'(0) \int_{\mathbb{R}^N} G(u) \, dx - \int_{\mathbb{R}^N} g(u) \varphi \, dx = 0$$

and hence (being $\theta'(t) = -\frac{2}{N}m^{1/N}|u + t\varphi|_2^{-2/N-2} \int_{\mathbb{R}^N} (u + t\varphi)\varphi \, dx$)

$$\int_{\mathbb{R}^N} (-\Delta)^{s/2} u (-\Delta)^{s/2} \varphi \, dx - \frac{2}{m} \left(\frac{1}{2_s^*} |(-\Delta)^{s/2} u|_2^2 - \int_{\mathbb{R}^N} G(u) \, dx \right) \int_{\mathbb{R}^N} u \varphi \, dx - \int_{\mathbb{R}^N} g(u) \varphi \, dx = 0.$$

Set

$$\mu := -\frac{2}{m} \left(\frac{1}{2_s^*} |(-\Delta)^{s/2} u|_2^2 - \int_{\mathbb{R}^N} G(u) \, dx \right) \in \mathbb{R}, \quad (4.2)$$

we have the claim. We see that the Pohožaev identity directly comes from the definition of μ in (4.2).

(ii) The statement is obvious for $u \equiv 0$. Arguing as in (i), let $\varphi \in H^s(\mathbb{R}^N) \cap L^1(\mathbb{R}^N)$; being $u \neq 0$, for $t \in \mathbb{R}$ with $|t|$ small we can assume $u + t\varphi \neq 0$ and $|u + t\varphi|_2^2 < m$. By the minimality of $t = 0$ for $t \mapsto K(u + t\varphi)$ we obtain

$$\int_{\mathbb{R}^N} (-\Delta)^{s/2} u (-\Delta)^{s/2} \varphi \, dx - \int_{\mathbb{R}^N} g(u) \varphi \, dx = 0$$

which implies that u is a solution with zero Lagrange multiplier. To show the Pohožaev identity, consider the map $t \mapsto K(u(\cdot/t))$, where $|u(\cdot/t)|_2^2 < m$ for t close to 1. Being $t = 1$ a point of minimum, we have

$$\frac{N - 2s}{2} |(-\Delta)^{s/2} u|_2^2 - N \int_{\mathbb{R}^N} G(u) \, dx = 0$$

which is the claim.

The proof of (iii) is straightforward. Point (iv) follows by (ii) and (iii). \blacksquare

We prove now Proposition 1.7. We underline that, unlike [14], we employ a direct proof without relying on a deformation lemma.

Proof of Proposition 1.7. Let $r, t \in \mathbb{R}$ and $\varphi \in H^s(\mathbb{R}^N) \cap L^1(\mathbb{R}^N)$, and define $\sigma = \sigma(r, t) > 0$ such that $P(\mu + r, u(\cdot/\sigma) + t\varphi(\cdot/\sigma)) = 0$, that is,

$$\sigma(r, t) := \left(\frac{1}{2_s^*} |(-\Delta)^{s/2} (u + t\varphi)|_2^2 \left(\int_{\mathbb{R}^N} G(u + t\varphi) - \frac{\mu + r}{2} (u + t\varphi)^2 \, dx \right)^{-1} \right)^{\frac{1}{2s}}.$$

Indeed, since $P(\mu, u) = 0$, we know that $\int_{\mathbb{R}^N} G(u) - \frac{\mu}{2} u^2 \, dx > 0$, thus for $|r|, |t|$ small the above expression is well defined. Moreover we can assume $|r|$ small such that $\mu + r > \rho$, thus $(\mu + r, u(\cdot/\sigma) + t\varphi(\cdot/\sigma)) \in \mathcal{P}_\rho$.

Being $P(\mu, u) = 0$, we have $\sigma(0, 0) = 1$. Next, we compute

$$\begin{aligned} & I^m(\mu + r, u(\cdot/\sigma) + t\varphi(\cdot/\sigma)) \\ &= \frac{\sigma^{N-2s}}{2} |(-\Delta)^{s/2} (u + t\varphi)|_2^2 + \frac{\mu + r}{2} (\sigma^N |u + t\varphi|_2^2 - m) - \sigma^N \int_{\mathbb{R}^N} G(u + t\varphi) \, dx. \end{aligned}$$

By Lemma 2.3, the map $t \mapsto \int_{\mathbb{R}^N} G(u + t\varphi) \, dx$, and so σ as well, is differentiable at 0. Since (μ, u) is a minimizer, the gradient of $I^m(\mu + r, u(\cdot/\sigma) + t\varphi(\cdot/\sigma))$ at $(r, t) = (0, 0)$ is zero. At the same time, recalling that $\sigma(0, 0) = 1$,

$$\partial_r I^m(\mu + r, u(\cdot/\sigma))|_{(t,r)=(0,0)} = \frac{N - 2s}{2} \partial_r \sigma(0, 0) P(\mu, u) + \frac{1}{2} (|u|_2^2 - m) = \frac{1}{2} (|u|_2^2 - m)$$

and

$$\partial_t I^m(\mu, u(\cdot/\sigma) + t\varphi(\cdot/\sigma))|_{(t,r)=(0,0)} =$$

$$\begin{aligned}
&= \frac{N-2s}{2} \partial_t \sigma(0,0) P(\mu, u) + \int_{\mathbb{R}^N} (-\Delta)^{s/2} u (-\Delta)^{s/2} \varphi + \mu u \varphi - g(u) \varphi \, dx \\
&= \int_{\mathbb{R}^N} (-\Delta)^{s/2} u (-\Delta)^{s/2} \varphi + \mu u \varphi - g(u) \varphi \, dx.
\end{aligned}$$

Thus, we obtain the claim. We conclude by applying Proposition 1.5. ■

Proof of Proposition 1.8. It follows by Proposition 4.2 and Proposition 1.5. ■

Proof of Proposition 1.10.

(i) From Proposition 1.8 we have that $I^m(\mu, u) = \inf_{\mathcal{P}} I^m$, hence, by Remark 4.1, $J(\mu, u) = \inf_{\mathcal{P}_\mu} J_\mu$, and (1.24) follows from Proposition 1.5.

(ii) By (1.24) and Proposition 1.5 we have

$$J(\mu, u) = \inf_{\mathcal{P}_\mu} J_\mu = \inf_{\mathcal{S}_m} K + \frac{\mu}{2} m = \inf_{\mathcal{P}} I^m + \frac{\mu}{2} m,$$

that is, $I^m(\mu, u) = \inf_{\mathcal{P}} I^m$. From Proposition 1.7 we obtain that $u \in \mathcal{S}_m$ and we conclude. ■

4.3 Nehari identity

Proposition 4.3 (Nehari identity) *Assume (g0)–(g2) and let $\mu \in \mathbb{R}$. If $u \in H^s(\mathbb{R}^N)$ is a solution to (1.1) such that $\int_{\mathbb{R}^N} g_-(u)u \, dx < \infty$, then u satisfies the Nehari identity*

$$|(-\Delta)^{s/2} u|_2^2 + \mu |u|_2^2 = \int_{\mathbb{R}^N} g(u)u \, dx.$$

Proof. With the notation of Lemma 2.2, for every $R > 0$ there holds

$$\int_{\mathbb{R}^N} (-\Delta)^{s/2} u (-\Delta)^{s/2} (\varphi_R u) \, dx + \mu \int_{\mathbb{R}^N} u^2 \varphi_R \, dx = \int_{\mathbb{R}^N} g(u) \varphi_R u \, dx. \quad (4.3)$$

Since $|\varphi_R| \leq 1$ and $\lim_{R \rightarrow +\infty} \varphi_R = 1$ a.e. in \mathbb{R}^N , from the dominated convergence theorem we obtain

$$\lim_{R \rightarrow +\infty} \int_{\mathbb{R}^N} u^2 \varphi_R \, dx = \int_{\mathbb{R}^N} u^2 \, dx \quad \text{and} \quad \lim_{R \rightarrow +\infty} \int_{\mathbb{R}^N} g(u) \varphi_R u \, dx = \int_{\mathbb{R}^N} g(u)u \, dx;$$

using also Lemma 2.2, we can pass to the limit in (4.3) and conclude. ■

5 The superlinear (perturbed) problem

In this section, where we always assume (g0)–(g2), we study the following perturbed problem

$$\begin{cases}
(-\Delta)^s u + \mu u = g_\varepsilon(u), \\
\int_{\mathbb{R}^N} u^2 \, dx = m, \\
(\mu, u) \in (0, +\infty) \times H^s(\mathbb{R}^N),
\end{cases} \quad (5.1)$$

where, for $\varepsilon \in (0, 1]$, g_ε is defined as follows: let $\varphi_\varepsilon \in \mathcal{C}(\mathbb{R}, [0, 1])$ given by

$$\varphi_\varepsilon(t) := \begin{cases} |t|/\varepsilon & \text{for } |t| < \varepsilon, \\ 1 & \text{for } |t| \geq \varepsilon, \end{cases} \quad (5.2)$$

and set

$$g_-^\varepsilon(t) := \varphi_\varepsilon(t)g_-(t), \quad G_-^\varepsilon(t) := \int_0^t g_-^\varepsilon(\tau) d\tau.$$

Then, we define

$$g_\varepsilon(t) := g_+(t) - g_-^\varepsilon(t), \quad G_\varepsilon(t) := \int_0^t g_\varepsilon(\tau) d\tau = G^+(t) - G_-^\varepsilon(t). \quad (5.3)$$

Notice that $\varphi_\varepsilon \rightarrow 1$ pointwise, thus

$$g_\varepsilon(t) \rightarrow g(t), \quad G_\varepsilon(t) \rightarrow G(t) \quad \text{for every } t \in \mathbb{R}$$

as $\varepsilon \rightarrow 0^+$. Hence, (5.1) can be considered as an approximating family of problems; the advantage is given by the modification in zero, which suppresses the singularity in the sublinear case:

$$\lim_{t \rightarrow 0} \frac{g_\varepsilon(t)}{t} = 0. \quad (5.4)$$

We will make large use of the following monotonicity property: if $0 < \varepsilon' < \varepsilon \leq 1$, then

$$G_-^\varepsilon \leq G_-^{\varepsilon'} \leq G_-,$$

which implies, moreover,

$$G_\varepsilon \geq G_{\varepsilon'} \geq G, \quad |G_\varepsilon| \leq |G_{\varepsilon'}| \leq |G|.$$

The purpose of this section is to obtain the existence of a Pohožaev minimum for the problem (5.1), furnishing, in addition, a slight generalization of [20] by exploiting the refined Theorem 3.6; such a generalization holds for more general nonlinearities not related to the perturbation process, see Remark 5.3 below.

5.1 ε fixed: some refinement for the superlinear problem

We start by considering $\varepsilon \in (0, 1]$ fixed. Assuming $\mu > 0$, to avoid technical issues related to the boundary of $(0, +\infty)$, we exploit the identification

$$\mu \equiv e^\lambda. \quad (5.5)$$

A similar approach that does not make use of (5.5) is developed in [30, Section 4.2].

Remark 5.1 *The condition $\mu > 0$ is a classical requirement for nonlinearities with superlinear growth at the origin – cf. (5.4) – because it leads to (see [8, 18])*

$$\lim_{t \rightarrow 0} \frac{g_\varepsilon(t) - \mu t}{t} < 0.$$

Making use of (5.5), we can thus reformulate the search for solutions to (5.1) as critical points of the Lagrangian functional $I_\varepsilon^m: \mathbb{R} \times H^s(\mathbb{R}^N) \rightarrow \mathbb{R}$ defined as

$$I_\varepsilon^m(\lambda, u) := \frac{1}{2}|(-\Delta)^{s/2}u|_2^2 + \frac{e^\lambda}{2}(|u|_2^2 - m) - \int_{\mathbb{R}^N} G_\varepsilon(u) dx.$$

Let us also define $P_\varepsilon: \mathbb{R} \times H^s(\mathbb{R}^N) \rightarrow \mathbb{R}$ as

$$P_\varepsilon(\lambda, u) := |(-\Delta)^{s/2}u|_2^2 + 2_s^* \left(\frac{e^\lambda}{2}|u|_2^2 - \int_{\mathbb{R}^N} G_\varepsilon(u) dx \right),$$

$$\mathcal{P}_{\varepsilon,+} := \{(\lambda, u) \in \mathbb{R} \times H^s(\mathbb{R}^N) \mid P_\varepsilon(\lambda, u) = 0, u \neq 0\}, \quad \mathcal{P}_{\varepsilon,+}^{\text{rad}} := \mathcal{P}_{\varepsilon,+} \cap (\mathbb{R} \times H_{\text{rad}}^s(\mathbb{R}^N)),$$

and

$$d_{\varepsilon,+}^m := \inf_{\mathcal{P}_{\varepsilon,+}^{\text{rad}}} I_\varepsilon^m.$$

With an abuse of notation, we redefine consequently on $\mathbb{R} \times H^s(\mathbb{R}^N)$ also

$$I^m(\lambda, u), \quad P(\lambda, u)$$

defined in (1.13) and (1.16). We also set

$$\mathcal{P}_+ := \{(\lambda, u) \in \mathbb{R} \times H^s(\mathbb{R}^N) \mid P(\lambda, u) = 0, u \neq 0\}, \quad \mathcal{P}_+^{\text{rad}} := \mathcal{P}_+ \cap (\mathbb{R} \times H_{\text{rad}}^s(\mathbb{R}^N)). \quad (5.6)$$

Notice that, with the identification $\mu \equiv e^\lambda > 0$, we have

$$d_+^m := \inf_{(\lambda,u) \in \mathcal{P}_+^{\text{rad}}} I^m(\lambda, u) = \inf_{(\mu,u) \in \mathcal{P}_{(0)}^{\text{rad}}} I^m(\mu, u). \quad (5.7)$$

Moreover, we introduce

$$K_\varepsilon(u) := \frac{1}{2} |(-\Delta)^{s/2} u|_2^2 - \int_{\mathbb{R}^N} G_\varepsilon(u) \, dx, \quad \kappa_\varepsilon^m := \inf_{\mathcal{S}_m^{\text{rad}}} K_\varepsilon, \quad \ell_\varepsilon^m := \inf_{\mathcal{D}_m^{\text{rad}}} K_\varepsilon.$$

The following is the main theorem of this section.

Theorem 5.2 *Assume that (g0)–(g4) hold, then there exists $m_\varepsilon \geq 0$ such that for every $m > m_\varepsilon$ there exists a solution $(\mu_\varepsilon, u_\varepsilon) \in \mathcal{P}_{\varepsilon,+}^{\text{rad}}$ to (5.1) verifying $I^m(\lambda_\varepsilon, u_\varepsilon) = d_{\varepsilon,+}^m < 0$.*

Moreover, if $g|_{(-\infty,0)} = 0$ or g is odd, then for every $m > m_\varepsilon$ there exists a Schwarz-symmetric minimizer of I_ε^m over $\mathcal{P}_{\varepsilon,+}$; in particular,

$$\inf_{\mathcal{P}_{\varepsilon,+}} I_\varepsilon^m = \inf_{\mathcal{P}_{\varepsilon,+}^{\text{rad}}} I_\varepsilon^m.$$

Remark 5.3 *We highlight that the only additional property of g_ε with respect to (g0)–(g4) that is used in Theorem 5.2 is the superlinearity at zero of g_ε , i.e., (5.4) (or equivalently, from (g1), $\lim_{t \rightarrow 0} g_-^\varepsilon(t)/t = 0$). Consequently, Theorem 5.2 and all the other properties presented in this section still hold true with g instead of g_ε under the additional assumption*

$$\lim_{t \rightarrow 0} \frac{g_-(t)}{t} = 0.$$

This furnishes a small generalization of the results in [20].

Remark 5.4 *We underline that the solution found in Theorem 5.2 has also a minimax characterization. Introduced*

$$\Omega_\varepsilon := \{(\lambda, u) \in \mathbb{R} \times H^s(\mathbb{R}^N) \mid P_\varepsilon(\lambda, u) > 0\} \cup (\mathbb{R} \times \{0\}),$$

we have $\partial\Omega_\varepsilon = \mathcal{P}_{\varepsilon,+}$ (see [30, Lemma 2.3.1]), and this is treated as a “mountain” in the minimax approach: set

$$\Gamma_\varepsilon^m := \left\{ \xi \in \mathcal{C}([0, 1], \mathbb{R} \times H_{\text{rad}}^s(\mathbb{R}^N)) \mid \xi(0) \in \mathbb{R} \times \{0\}, \xi(1) \notin \Omega_\varepsilon, \max_{i=0,1} I_\varepsilon^m(\xi(i)) \leq d_{\varepsilon,+}^m - 1 \right\},$$

one can prove that [20, Section 6]

$$d_{\varepsilon,+}^m = \inf_{\xi \in \Gamma_\varepsilon^m} \max_{t \in [0,1]} I_\varepsilon^m(\xi(t)).$$

Moreover, by arguing as in [20], we also obtain (write $g_0 := g$ and $G_0 := G$ for convenience)

- if g is odd, then for each $\varepsilon \in (0, 1]$ and $k \in \mathbb{N}$ there exists $m_{\varepsilon, k} \geq 0$ such that, for $m > m_{\varepsilon, k}$, there exist k (couple of) solutions to (5.1);
- if $\lim_{t \rightarrow 0} \frac{G(t)}{|t|^{\bar{p}+1}} = +\infty$, then for each $\varepsilon \in [0, 1]$, we have $m_\varepsilon = 0$ (and $m_{\varepsilon, k} = 0$ if g is odd).
Notice that, in this case, $g = g_+$ in a neighborhood of the origin.

In [20], the existence of a normalized solution is pursued by a Mountain Pass approach, involving a suitable augmented functional and a deformation lemma; the found solution is shown a posteriori to be a Pohožaev minimum. Here we provide a different approach, by minimizing directly the functional on the Pohožaev product set: this proof is not trivial, since a precise sequence of minimizers has to be selected; on the other hand, it requires fewer tools.

To prove Theorem 5.2, we need preliminary estimates on $d_{\varepsilon, +}^m$. This will be given by the interplay of the value of the Pohožaev minimum on the product space and the one with fixed μ .

5.1.1 The Pohožaev asymptotic geometry

To estimate $d_{\varepsilon, +}^m$, we introduce now the Pohožaev geometry related to a fixed frequency λ , in the spirit of Section 3. For $\lambda \in \mathbb{R}$ we define

$$\begin{aligned}
J(\lambda, u) &:= \frac{1}{2} |(-\Delta)^{s/2} u|_2^2 + \frac{e^\lambda}{2} |u|_2^2 - \int_{\mathbb{R}^N} G(u) \, dx, \\
J_\varepsilon(\lambda, u) &:= \frac{1}{2} |(-\Delta)^{s/2} u|_2^2 + \frac{e^\lambda}{2} |u|_2^2 - \int_{\mathbb{R}^N} G_\varepsilon(u) \, dx, \\
\mathcal{P}(\lambda) &:= \{u \in H^s(\mathbb{R}^N) \mid (\lambda, u) \in \mathcal{P}_+\}, \quad \mathcal{P}^{\text{rad}}(\lambda) := \mathcal{P}(\lambda) \cap H_{\text{rad}}^s(\mathbb{R}^N), \\
\mathcal{P}_\varepsilon(\lambda) &:= \{u \in H^s(\mathbb{R}^N) \mid (\lambda, u) \in \mathcal{P}_{\varepsilon, +}\}, \quad \mathcal{P}_\varepsilon^{\text{rad}}(\lambda) := \mathcal{P}_\varepsilon(\lambda) \cap H_{\text{rad}}^s(\mathbb{R}^N), \\
a(\lambda) &:= \inf_{\mathcal{P}^{\text{rad}}(\lambda)} J(\lambda, \cdot), \quad a_\varepsilon(\lambda) := \inf_{\mathcal{P}_\varepsilon^{\text{rad}}(\lambda)} J_\varepsilon(\lambda, \cdot).
\end{aligned} \tag{5.8}$$

Notice that $J(\lambda, u) = J_{e^\lambda}(u)$ and $\mathcal{P}(\lambda) = \mathcal{P}_{e^\lambda}$ with the notations of Section 3. Moreover, we make again an abuse of notation by writing $a(\lambda) \equiv a(e^\lambda)$ – see (3.3). We observe that

$$J_\varepsilon \leq J.$$

To compare $a(\lambda)$ with $a_\varepsilon(\lambda)$, we make use of (3.5): as a matter of fact, by $P_\varepsilon \leq P$ we have $\{P = 0\} \subset \{P_\varepsilon \leq 0\}$ and thus

$$a_\varepsilon(\lambda) \leq a(\lambda). \tag{5.9}$$

We highlight that, arguing as in [14], we could give a Mountain Pass characterization of $a(\lambda)$; nonetheless, we will not make use of this characterization for the limit problem. Moreover, we introduce

$$\begin{aligned}
\bar{\mu}_\varepsilon &:= \sup_{t \neq 0} \frac{G_\varepsilon(t)}{t^2/2} & \bar{\lambda}_\varepsilon &:= \log(\bar{\mu}_\varepsilon) \\
&\geq \sup_{t \neq 0} \frac{G(t)}{t^2/2} = \bar{\mu}_0 \in (0, +\infty] & &\geq \log(\bar{\mu}_0) =: \bar{\lambda}_0 \in (-\infty, +\infty]
\end{aligned} \tag{5.10}$$

with the convention that $\log(+\infty) = +\infty$. Observe that $\bar{\mu}_0 > 0$ follows from (g4). We highlight that, in what follows, we will sometimes need to distinguish the cases $\bar{\lambda}_\varepsilon \in \mathbb{R}$ and $\bar{\lambda}_\varepsilon = +\infty$.

We define, for $m > 0$, the values

$$b_\varepsilon^m := \inf_{\lambda < \bar{\lambda}_\varepsilon} \left(a_\varepsilon(\lambda) - \frac{e^\lambda}{2} m \right) \quad \text{and} \quad b^m := \inf_{\lambda < \bar{\lambda}_0} \left(a(\lambda) - \frac{e^\lambda}{2} m \right). \tag{5.11}$$

Moreover, by Lemma 3.1 the following thresholds are well defined

$$m_\varepsilon := \inf_{\lambda < \bar{\lambda}_\varepsilon} \frac{a_\varepsilon(\lambda)}{e^{\lambda/2}} \quad \text{and} \quad m_0 := \inf_{\lambda < \bar{\lambda}_0} \frac{a(\lambda)}{e^{\lambda/2}}. \quad (5.12)$$

Observe that, from Lemma 3.1 and since $\bar{\lambda}_\varepsilon \geq \bar{\lambda}_0$ and $a_0 \geq a_\varepsilon \geq 0$ – cf. Proposition 3.4, we have

$$b^m \geq b_\varepsilon^m \quad \text{and} \quad m_0 \geq m_\varepsilon \geq 0. \quad (5.13)$$

As a straightforward consequence of Proposition 1.6, Theorem 3.6, and the definitions of $\bar{\lambda}_\varepsilon$ and m_ε , we obtain the following result.

Proposition 5.5 *Assume (g0)–(g2) and (g4). Then*

- For every $\lambda \in \mathbb{R}$ there exists a solution $u_{\lambda,\varepsilon} \in \mathcal{P}_\varepsilon^{\text{rad}}(\lambda)$ to

$$(-\Delta)^s u + e^\lambda u = g_\varepsilon(u) \quad \text{in } \mathbb{R}^N.$$

Additionally,

$$J_\varepsilon(\lambda, u_{\lambda,\varepsilon}) = \inf_{\mathcal{P}_\varepsilon^{\text{rad}}(\lambda)} J_\varepsilon(\lambda, u) = a_\varepsilon(\lambda) > 0.$$

As a consequence, $a(\lambda) > 0$.

- $d_{\varepsilon,+}^m = b_\varepsilon^m$; if $m > m_\varepsilon$, then $d_{\varepsilon,+}^m < 0$.
- If $\bar{\lambda}_\varepsilon \in \mathbb{R}$, then $\min\{P_\varepsilon(\lambda, u), J_\varepsilon(\lambda, u)\} > 0$ for all $\lambda \geq \bar{\lambda}_\varepsilon$ and $u \in H^s(\mathbb{R}^N) \setminus \{0\}$.

We focus now on properties strictly related to the L^2 -subcritical setting (g3). To start, we observe that for every $\delta > 0$ there exist $C_\delta, c_\delta > 0$ such that for every $t \in \mathbb{R}$

$$G_\varepsilon(t) \leq G_+(t) \leq \frac{\delta}{2} t^2 + \frac{C_\delta}{\bar{p}+1} |t|^{\bar{p}+1} \quad (5.14)$$

and

$$G_\varepsilon(t) \leq G_+(t) \leq \frac{c_\delta}{2} t^2 + \frac{\delta}{\bar{p}+1} |t|^{\bar{p}+1}. \quad (5.15)$$

Lemma 5.6 (Asymptotic geometry) *Assume (g0)–(g4). The following statements hold.*

- $\lim_{\lambda \rightarrow \bar{\lambda}_\varepsilon^-} a_\varepsilon(\lambda) = +\infty$. If $\bar{\lambda}_\varepsilon = +\infty$, then $\lim_{\lambda \rightarrow +\infty} \frac{a_\varepsilon(\lambda)}{e^\lambda} = +\infty$.
- $d_{\varepsilon,+}^m > -\infty$.

Proof. Recalled (5.14)–(5.15), concerning the first point we have that the first limit can be obtained with the same reasoning as in [20, Proposition 3.3] in light of the monotonicity in Proposition 3.4, while for the second one we can argue as in [20, Proposition 3.4] once the following property is proved:

$$a_\varepsilon(\lambda) \geq (e^\lambda - c_\delta) \inf_{\mathcal{P}^*(\delta)} J_\delta^* \quad \text{if } e^\lambda > c_\delta \quad \text{and} \quad \lim_{\delta \rightarrow 0^+} \inf_{\mathcal{P}^*(\delta)} J_\delta^* = +\infty, \quad (5.16)$$

where

$$\begin{aligned} J_\delta^*(u) &:= \frac{1}{2} |(-\Delta)^{s/2} u|_2^2 + \frac{1}{2} |u|_2^2 - \frac{\delta}{\bar{p}+1} |u|_{\bar{p}+1}^{\bar{p}+1}, \\ \mathcal{P}_\delta^* &:= \{u \in H_{\text{rad}}^s(\mathbb{R}^N) \mid P_\delta^*(u) = 0\} \setminus \{0\}, \end{aligned}$$

$$P_\delta^*(u) := |(-\Delta)^{s/2}u|_2^2 + 2_s^* \left(\frac{1}{2}|u|_2^2 - \frac{\delta}{\bar{p}+1}|u|_{\bar{p}+1}^{\bar{p}+1} \right)$$

(i.e., P_δ^* and \mathcal{P}_δ^* are, respectively, the Pohožaev functional and manifold related to J_δ^*). For every $\theta > 0$ and $u \in \mathcal{P}(\lambda)$, we define $u_\theta := \theta^{-N/(4s)}u(\theta^{-1/(2s)}\cdot)$. Then, using (5.15) and taking $\theta = e^\lambda - c_\delta$ (which is positive if $\lambda > \log(c_\delta)$), it is straightforward to verify that

$$J_\varepsilon(\lambda, u) \geq \theta J_\delta^*(u_\theta) \quad \text{and} \quad P_\delta^*(u_\theta) \leq 0,$$

therefore the first property in (5.16) holds thanks to (3.5). As for the second one, by scaling, one readily checks that $P_\delta^*(\delta^{-1/(\bar{p}-1)}\cdot) = \delta^{-2/(\bar{p}-1)}P_1^*$, which implies that $\mathcal{P}_\delta^* = \delta^{-1/(\bar{p}-1)}\mathcal{P}_1^*$, and $J_\delta^*(\delta^{-1/(\bar{p}-1)}\cdot) = \delta^{-2/(\bar{p}-1)}J_1^*$. From this, one obtains $\inf_{\mathcal{P}_\delta^*} J_\delta^* = \delta^{-2/(\bar{p}-1)} \inf_{\mathcal{P}_1^*} J_1^* \rightarrow +\infty$ as $\delta \rightarrow 0^+$.

Finally, recalled that $d_{\varepsilon,+}^m = b_\varepsilon^m$ by Proposition 5.5, the estimate $b_\varepsilon^m > -\infty$ is obtained as in [20, Proposition 3.7]. \blacksquare

5.1.2 Direct minimization on Pohožaev product

We are now ready to prove the existence of a Pohožaev minimum as stated in Theorem 5.2.

Proof of Theorem 5.2. By Proposition 5.5 and Lemma 5.6 we know that $d_{\varepsilon,+}^m \in (-\infty, 0)$. Let $(\lambda_n, u_n)_n \subset \mathcal{P}_{\varepsilon,+}^{\text{rad}}$ be a minimizing sequence, i.e., $I_\varepsilon^m(\lambda_n, u_n) \rightarrow d_{\varepsilon,+}^m$ and $P_\varepsilon(\lambda_n, u_n) = 0$. Notice that, being $u_n \in \mathcal{P}_\varepsilon^{\text{rad}}(\lambda_n)$, by Lemma 3.1 we have $\lambda_n < \bar{\lambda}_\varepsilon$ for each n .

Step 1. *Choice of a proper minimizing sequence.*

We build now a minimizing sequence with additional information on the energy levels and on the differential (with respect to u) of I_ε^m . For each n , by Proposition 5.5 there exists $v_n \in \mathcal{P}_\varepsilon^{\text{rad}}(\lambda_n)$ such that

$$I_\varepsilon^m(\lambda_n, v_n) = a_\varepsilon(\lambda_n) - \frac{e^{\lambda_n}}{2}m \tag{5.17}$$

and

$$\partial_u I_\varepsilon^m(\lambda_n, v_n) = 0. \tag{5.18}$$

From (5.17) we obtain

$$d_{\varepsilon,+}^m \leq I_\varepsilon^m(\lambda_n, v_n) = a_\varepsilon(\lambda_n) - \frac{e^{\lambda_n}}{2}m \leq J_\varepsilon(\lambda_n, u_n) - \frac{e^{\lambda_n}}{2}m = I_\varepsilon^m(\lambda_n, u_n) = d_{\varepsilon,+}^m + o(1).$$

As a consequence, also $(\lambda_n, v_n)_n$ is a minimizing sequence.

Step 2. $(\lambda_n)_n$ is bounded below.

We have

$$\frac{m}{2}e^{\lambda_n} = \frac{1}{2_s^*}P_\varepsilon(\lambda_n, v_n) - I_\varepsilon^m(\lambda_n, v_n) + \frac{s}{N}|(-\Delta)^{s/2}v_n|_2^2 \geq -I_\varepsilon^m(\lambda_n, v_n),$$

whence

$$\liminf_n \frac{m}{2}e^{\lambda_n} \geq -d_{\varepsilon,+}^m > 0.$$

Step 3. $(\lambda_n)_n$ is bounded above.

In this Step we use the L^2 -subcritical nature of the problem. If $\bar{\lambda}_\varepsilon < +\infty$, then we are done. Otherwise, if by contradiction $\lambda_n \rightarrow +\infty$, then by Lemma 5.6 and (5.17) we have

$$d_{\varepsilon,+}^m + o(1) = I_\varepsilon^m(\lambda_n, v_n) = e^{\lambda_n} \left(\frac{a_\varepsilon(\lambda_n)}{e^{\lambda_n}} - \frac{1}{2}m \right) \rightarrow +\infty,$$

which is impossible.

Step 4. $(|(-\Delta)^{s/2}v_n|_2)_n$ is bounded.

Indeed,

$$\frac{s}{N}|(-\Delta)^{s/2}v_n|_2^2 = I_\varepsilon^m(\lambda_n, v_n) - \frac{1}{2_s^*}P_\varepsilon(\lambda_n, v_n) + \frac{m}{2}e^{\lambda_n} = d_{\varepsilon,+}^m + \frac{m}{2}e^{\lambda_n} + o(1)$$

and thus $|(-\Delta)^{s/2}v_n|_2^2$ is bounded thanks to Step 3.

Step 5. $(|v_n|_2)_n$ is bounded.

By (5.14) we have

$$\begin{aligned} d_{\varepsilon,+}^m + \frac{m}{2}e^{\lambda_n} + o(1) &= I_\varepsilon^m(\lambda_n, v_n) - \frac{1}{2}P_\varepsilon(\lambda_n, v_n) + \frac{m}{2}e^{\lambda_n} \\ &= -\frac{e^{\lambda_n}}{2} \left(\frac{2_s^*}{2} - 1 \right) |v_n|_2^2 + \left(\frac{2_s^*}{2} - 1 \right) \int_{\mathbb{R}^N} G_\varepsilon(v_n) dx \\ &\leq \left(\frac{2_s^*}{2} - 1 \right) \left(-\frac{1}{2} (e^{\lambda_n} - \delta) |v_n|_2^2 + \frac{C_\delta}{\bar{p}+1} |v_n|_{2_s^*}^{2_s^*} \right); \end{aligned}$$

by choosing $\delta < e^{\inf\{\lambda_n\}}$ (possible thanks to Step 2) and exploiting Step 4 (which implies that $|v_n|_{2_s^*}$ is bounded), we obtain that $|v_n|_2$ is bounded as well.

Step 6. $(v_n)_n$ converges strongly in $H^s(\mathbb{R}^N)$.

In this step, we rely on the superlinearity of the perturbed problem. From Steps 2–5, there exist $\lambda \in \mathbb{R}$ and $v \in H_{\text{rad}}^s(\mathbb{R}^N)$ such that, up to a subsequence, $\lambda_n \rightarrow \lambda$, $v_n \rightharpoonup v$ in $H_{\text{rad}}^s(\mathbb{R}^N)$, $v_n \rightarrow v$ in $L^q(\mathbb{R}^N)$ for $q \in (2, 2_s^*)$, and $v_n(x) \rightarrow v(x)$ for a.e. $x \in \mathbb{R}^N$. In particular, from (g0)–(g3),

$$\lim_n \int_{\mathbb{R}^N} g_\varepsilon(v_n)v dx = \int_{\mathbb{R}^N} g_\varepsilon(v)v dx \quad \text{and} \quad \lim_n \int_{\mathbb{R}^N} g_+(v_n)v_n dx = \int_{\mathbb{R}^N} g_+(v)v dx;$$

moreover, from (5.18) we have

$$\partial_u I_\varepsilon^m(\lambda_n, v_n)[v] = 0 = \partial_u I_\varepsilon^m(\lambda_n, v_n)[v_n].$$

Thus, exploiting the weak convergence, Fatou's lemma (recall $g_-^\varepsilon(t)t \geq 0$ for all $t \in \mathbb{R}$) and that $\lambda_n \rightarrow \lambda$, we obtain

$$\begin{aligned} &|(-\Delta)^{s/2}v|_2^2 + e^\lambda |v|_2^2 - \int_{\mathbb{R}^N} g_\varepsilon(v)v dx \\ &= \lim_n \int_{\mathbb{R}^N} (-\Delta)^{s/2}v_n (-\Delta)^{s/2}v dx + e^{\lambda_n} \int_{\mathbb{R}^N} v_n v dx + \int_{\mathbb{R}^N} g_\varepsilon(v_n)v dx = 0 \\ &= \lim_n \int_{\mathbb{R}^N} |(-\Delta)^{s/2}v_n|^2 dx + e^{\lambda_n} \int_{\mathbb{R}^N} v_n^2 dx + \int_{\mathbb{R}^N} g_\varepsilon(v_n)v_n dx \\ &= \lim_n \left(\int_{\mathbb{R}^N} |(-\Delta)^{s/2}v_n|^2 dx + e^{\lambda_n} \int_{\mathbb{R}^N} v_n^2 dx + \int_{\mathbb{R}^N} g_-^\varepsilon(v_n)v_n dx \right) - \int_{\mathbb{R}^N} g_+(v)v dx \\ &\geq \liminf_n \left(\int_{\mathbb{R}^N} |(-\Delta)^{s/2}v_n|^2 dx + e^{\lambda_n} \int_{\mathbb{R}^N} v_n^2 dx \right) - \int_{\mathbb{R}^N} g_\varepsilon(v)v dx \\ &\geq |(-\Delta)^{s/2}v|_2^2 + e^\lambda |v|_2^2 - \int_{\mathbb{R}^N} g_\varepsilon(v)v dx, \end{aligned}$$

thus

$$\liminf_n \int_{\mathbb{R}^N} |(-\Delta)^{s/2}v_n|^2 dx + e^{\lambda_n} \int_{\mathbb{R}^N} v_n^2 dx = |(-\Delta)^{s/2}v|_2^2 + e^\lambda |v|_2^2;$$

this means that, up to a subsequence, $v_n \rightarrow v$ strongly in $H^s(\mathbb{R}^N)$.

Step 7. Conclusions.

By the strong convergence, we have $I_\varepsilon^m(\lambda, v) = d_{\varepsilon,+}^m$, i.e., (λ, v) is a Pohožaev minimum in the product space. The fact that $(\mu_\varepsilon, u_\varepsilon) := (e^\lambda, v)$ is a solution to (5.1) follows from Proposition 1.7. Finally, assume $g|_{(-\infty,0)} = 0$ or g odd, which imply the same properties on g_ε . Then, in Step 1, we may assume that each v_n is Schwarz-symmetric in light of Theorem 3.6 (ii), hence so is v . \blacksquare

5.2 Symmetry results

Exploiting the perturbed g_ε , we prove here Proposition 1.12 (i)–(iii). We recall that G does not map $H^s(\mathbb{R}^N)$ to $L^1(\mathbb{R}^N)$; with the intention of gaining a symmetry result arguing as in [40, Proof of Theorem 4.1], we show the following lemma.

Lemma 5.7 *Assume (g0)–(g2). For $v \in H^s(\mathbb{R}^N)$ and $x = (x_1, x') \in \mathbb{R} \times \mathbb{R}^{N-1}$, define*

$$v_1(x) := \begin{cases} v(x_1, x') & \text{if } x_1 < 0 \\ v(-x_1, x') & \text{if } x_1 \geq 0 \end{cases} \quad \text{and} \quad v_2(x) := \begin{cases} v(-x_1, x') & \text{if } x_1 < 0 \\ v(x_1, x') & \text{if } x_1 \geq 0. \end{cases}$$

Then $\int_{\mathbb{R}^N} G(v_1) dx + \int_{\mathbb{R}^N} G(v_2) dx = 2 \int_{\mathbb{R}^N} G(v) dx$.

Proof. It is clear that $\int_{\mathbb{R}^N} G_+(v_1) dx + \int_{\mathbb{R}^N} G_+(v_2) dx = 2 \int_{\mathbb{R}^N} G_+(v) dx$. Additionally, from the monotone convergence theorem,

$$\begin{aligned} \int_{\mathbb{R}^N} G_-(v_1) dx + \int_{\mathbb{R}^N} G_-(v_2) dx &= \lim_{\varepsilon \rightarrow 0^+} \int_{\mathbb{R}^N} G_-^\varepsilon(v_1) dx + \int_{\mathbb{R}^N} G_-^\varepsilon(v_2) dx \\ &= 2 \lim_{\varepsilon \rightarrow 0^+} \int_{\mathbb{R}^N} G_-^\varepsilon(v) dx = 2 \int_{\mathbb{R}^N} G_-(v) dx, \end{aligned}$$

whence the statement. \blacksquare

Proof of Proposition 1.12 (i)–(iii).

(i) Let $(\mu, u) \in \mathcal{P}(\rho)$ be a minimizer as in the statement. In view of Remark 4.1 and Lemma 3.3, u solves the minimization problem (3.4). Then, exploiting Lemma 5.7, we can argue in a similar way to [40, Proof of Theorem 4.1]⁵. Likewise if u minimizes K over \mathcal{S}_m .

(ii) Let $(\mu, u) \in \mathcal{P}(\rho)$ be a minimizer as in the statement, and let us denote by u^* the Schwarz rearrangement of u . We first observe that

$$\int_{\mathbb{R}^N} G_-(u) dx = \int_{\mathbb{R}^N} G_-(u^*) dx; \tag{5.19}$$

indeed, (5.19) holds with G_-^ε instead of G_- (see [14, Proposition B.1]), thus from the monotone convergence theorem

$$\int_{\mathbb{R}^N} G_-(u) dx = \lim_{\varepsilon \rightarrow 0^+} \int_{\mathbb{R}^N} G_-^\varepsilon(u) dx = \lim_{\varepsilon \rightarrow 0^+} \int_{\mathbb{R}^N} G_-^\varepsilon(u^*) dx = \int_{\mathbb{R}^N} G_-(u^*) dx,$$

and we get the claim. Now, assume by contradiction that u does not equal its Schwarz rearrangement $u^* \neq u$. Then, from [14, Propositions B.1 and B.3],

$$|(-\Delta)^{s/2} u^*|_2^2 < |(-\Delta)^{s/2} u|_2^2 \quad \text{and} \quad |u^*|_2 = |u|_2,$$

⁵ We observe that we formally deal with *Case A* therein but minimize over $H^s(\mathbb{R}^N)$.

which imply

$$J_\mu(u^*) < J_\mu(u) \quad \text{and} \quad P_\mu(u^*) < P_\mu(u) = 0.$$

From Remark 3.2, we find $t^* > 0$ such that $P_\mu(u^*(\cdot/t^*)) = 0$ and

$$\inf_{\mathcal{P}_\mu} J_\mu \leq J_\mu(u^*(\cdot/t^*)) \leq J_\mu(u^*) < J_\mu(u) = \inf_{\mathcal{P}_\mu} J_\mu,$$

a contradiction. Similarly for L^2 -minima.

(iii) We claim that $\int_{\mathbb{R}^N} |g_-(u)u^-| dx < +\infty$, where, we recall, $u^\pm = \max\{\pm u, 0\}$. As a matter of fact, since $g_-(u)u^- = g_-(-u^-)u^- \leq 0$, there holds

$$|g_-(u)u^-| = -g_-(-u^-)u^- \leq g_-(u^+)u^+ - g_-(-u^-)u^- = g_-(u)u^+ - g_-(u)u^- = g_-(u)u \in L^1(\mathbb{R}^N),$$

which holds by the assumptions. As a consequence, with the notation of Lemma 2.2, for every $R > 0$ there holds

$$\int_{\mathbb{R}^N} (-\Delta)^{s/2} u (-\Delta)^{s/2} (\varphi_R u^-) dx + \mu \int_{\mathbb{R}^N} u \varphi_R u^- dx = \int_{\mathbb{R}^N} g(u) \varphi_R u^- dx,$$

and passing to the limit we obtain

$$\int_{\mathbb{R}^N} (-\Delta)^s u (-\Delta)^s u^- dx + \mu \int_{\mathbb{R}^N} u u^- dx = \int_{\mathbb{R}^N} g(u) u^- dx. \quad (5.20)$$

Finally, using (5.20) and Lemma 2.1,

$$\begin{aligned} 0 &\leq \int_{\mathbb{R}^N} (g(-u^-) - \mu(-u^-)) u^- dx = \int_{\mathbb{R}^N} g(u) u^- - \mu u u^- dx = \int_{\mathbb{R}^N} (-\Delta)^{s/2} u (-\Delta)^{s/2} u^- dx \\ &\leq -|(-\Delta)^{s/2} u^-|_2^2 \leq 0, \end{aligned}$$

and we conclude. ■

6 The sublinear problem

6.1 ε -independent estimates

Let $\mu_\varepsilon = e^{\lambda_\varepsilon}$, u_ε , m_ε , and $d_{\varepsilon,+}^m$ be the quantities found in Theorem 5.2 and m_0 be as in (5.12). From now on, we are interested in the dependence on ε ; in particular, we will ensure that some objects related to problem (5.1) enjoy some uniformity in $\varepsilon > 0$. We start by showing that $(d_{\varepsilon,+}^m)_{\varepsilon \in (0,1]}$ is bounded and far from zero.

Lemma 6.1 *Let $m > m_0$. There exist $\zeta_m, \beta_m \in \mathbb{R}$ such that*

$$-\infty < \zeta_m \leq d_{\varepsilon,+}^m \leq \beta_m < 0 \quad \text{for every } \varepsilon \in (0, 1].$$

Proof. By Lemma 5.6 and the monotonicity of $\varepsilon \mapsto a_\varepsilon(\lambda)$ and $\varepsilon \mapsto \bar{\lambda}_\varepsilon$ - cf. (5.9) and (5.10), we have $d_{\varepsilon,+}^m = b_\varepsilon^m \geq b_1^m =: \zeta_m$. Again by Lemma 5.6 we have, for every $\lambda < \bar{\lambda}_0 \leq \bar{\lambda}_\varepsilon$,

$$d_{\varepsilon,+}^m \leq \frac{e^\lambda}{2} \left(\frac{a_\varepsilon(\lambda)}{e^{\lambda/2}} - m \right) \leq \frac{e^\lambda}{2} \left(\frac{a(\lambda)}{e^{\lambda/2}} - m \right).$$

Let now $\delta := \frac{m-m_0}{2} > 0$. By the definition of infimum $m_0 = \inf_{\lambda < \bar{\lambda}_0} \frac{a(\lambda)}{e^{\lambda/2}}$ there exists $\lambda_\delta < \bar{\lambda}_0$ such that

$$d_{\varepsilon,+}^m \leq \frac{e^{\lambda_\delta}}{2} \left(\frac{a(\lambda_\delta)}{e^{\lambda_\delta/2}} - m \right) \leq \frac{e^{\lambda_\delta}}{2} (m_0 + \delta - m) = -\frac{e^{\lambda_\delta}}{4} (m - m_0) =: \beta_m < 0. \quad \blacksquare$$

We show now some estimates on the solutions $(\lambda_\varepsilon, u_\varepsilon)$. Here we use the L^2 -subcritical growth (g3).

Lemma 6.2 *Let $m > m_0$. Then $(\lambda_\varepsilon, u_\varepsilon)_{\varepsilon \in (0,1]}$ is bounded in $\mathbb{R} \times H_{\text{rad}}^s(\mathbb{R}^N)$.*

Proof. Observe first that, since $(\lambda_\varepsilon, u_\varepsilon)$ satisfies the Pohožaev identity and by Lemma 6.1, we have

$$\frac{m}{2}e^{\lambda_\varepsilon} = \frac{1}{2_s^*}P_\varepsilon(\lambda_\varepsilon, u_\varepsilon) - I_\varepsilon^m(\lambda_\varepsilon, u_\varepsilon) + \frac{s}{N}|(-\Delta)^{s/2}u_\varepsilon|_2^2 \geq -d_{\varepsilon,+}^m \geq -\beta_m$$

which gives a bound below on $(\lambda_\varepsilon)_{\varepsilon \in (0,1]}$. Next, from (g0), (g1), and (g3), in a similar way to (5.15), for every $\delta > 0$ there exists $c_\delta > 0$ such that for all $t \in \mathbb{R}$

$$g_+(t)t \leq c_\delta t^2 + \delta |t|^{\bar{p}+1}.$$

Then, from (2.2) and the fact that $(\lambda_\varepsilon, u_\varepsilon)$ is a critical point of I_ε^m , we have

$$\begin{aligned} 0 &= |(-\Delta)^{s/2}u_\varepsilon|_2^2 + e^{\lambda_\varepsilon}|u_\varepsilon|_2^2 + \int_{\mathbb{R}^N} g_-^\varepsilon(u_\varepsilon)u_\varepsilon \, dx - \int_{\mathbb{R}^N} g_+(u_\varepsilon)u_\varepsilon \, dx \\ &\geq |(-\Delta)^{s/2}u_\varepsilon|_2^2 + (e^{\lambda_\varepsilon} - c_\delta)|u_\varepsilon|_2^2 + \int_{\mathbb{R}^N} g_-^\varepsilon(u_\varepsilon)u_\varepsilon \, dx - \delta |u_\varepsilon|_{\bar{p}+1}^{\bar{p}+1} \\ &\geq |(-\Delta)^{s/2}u_\varepsilon|_2^2 + (e^{\lambda_\varepsilon} - c_\delta)|u_\varepsilon|_2^2 - \delta C_{\text{GN}}|(-\Delta)^{s/2}u_\varepsilon|_2^2 |u_\varepsilon|_2^{4s/N} \\ &= \left(1 - \delta C_{\text{GN}}m^{2s/N}\right)|(-\Delta)^{s/2}u_\varepsilon|_2^2 + (e^{\lambda_\varepsilon} - c_\delta)m, \end{aligned}$$

and we conclude by taking $\delta < (C_{\text{GN}}m^{2s/N})^{-1}$. ■

6.2 Passage to the limit

We want to pass to the limit with the solutions $(\lambda_\varepsilon, u_\varepsilon)$ found in Theorem 5.2. Recall that, by Lemma 6.1 and Proposition 1.5, we have

$$-\infty < \zeta_m \leq d_{\varepsilon,+}^m = I_\varepsilon^m(\lambda_\varepsilon, u_\varepsilon) = \kappa_\varepsilon^m = \ell_\varepsilon^m = K_\varepsilon(u_\varepsilon) \leq \beta_m < 0;$$

moreover, by Lemma 6.2, $(\lambda_\varepsilon, u_\varepsilon)_{\varepsilon \in (0,1]}$ is bounded in $\mathbb{R} \times H_{\text{rad}}^s(\mathbb{R}^N)$. Therefore, up to a subsequence,

$$\begin{cases} u_\varepsilon \rightharpoonup u_0 & \text{in } H_{\text{rad}}^s(\mathbb{R}^N), \\ \lambda_\varepsilon \rightarrow \lambda_0 & \text{in } \mathbb{R}. \end{cases} \quad (6.1)$$

In addition, we have

$$|u_0|_2^2 \leq m$$

by the weak convergence. We write $\mu_\varepsilon := e^{\lambda_\varepsilon}$, $\mu_0 := e^{\lambda_0}$.

We want to show that $(\mu_0, u_0) \in (0, +\infty) \times H_{\text{rad}}^s(\mathbb{R}^N)$ is the desired solution of Theorem 1.1. To this aim, the first step is to prove the pointwise convergence of $G_\varepsilon(u_\varepsilon)$ to $G(u_0)$.

Lemma 6.3 *Let $u_0 \in H_{\text{rad}}^s(\mathbb{R}^N)$ as in (6.1). Then, $g_\varepsilon(u_\varepsilon(x)) \rightarrow g(u_0(x))$ and $G_\varepsilon(u_\varepsilon(x)) \rightarrow G(u_0(x))$ as $\varepsilon \rightarrow 0^+$ for a.e. $x \in \mathbb{R}^N$.*

Proof of Lemma 6.3. Let $x \in \mathbb{R}^N$ such that $u_\varepsilon(x) \rightarrow u_0(x) \in \mathbb{R}$. We prove first the claim on g_ε . We write

$$(0, 1] = \{\varepsilon \in (0, 1] \mid |u_\varepsilon(x)| > \varepsilon\} \cup \{\varepsilon \in (0, 1] \mid |u_\varepsilon(x)| \leq \varepsilon\} =: I_1^x \cup I_2^x.$$

Clearly, $0 \in \overline{I_1^x} \cup \overline{I_2^x}$. If $0 \in \overline{I_1^x}$ then, observed that $g_\varepsilon(u_\varepsilon(x)) = g(u_\varepsilon(x))$ for every $\varepsilon \in I_1^x$, we get

$$\lim_{\varepsilon \rightarrow 0^+, \varepsilon \in I_1^x} g_\varepsilon(u_\varepsilon(x)) = \lim_{\varepsilon \rightarrow 0^+, \varepsilon \in I_1^x} g(u_\varepsilon(x)) = g(u_0(x)).$$

If $0 \in \overline{I_2^x}$, let $\delta > 0$. By the assumptions, we know that there exists $\omega_\delta > 0$ such that $|g(t)| \leq \delta$ for every $|t| \leq \omega_\delta$. Observed that

$$\begin{aligned} |g_\varepsilon(u_\varepsilon(x))| &\leq |g_+(u_\varepsilon(x))| + |g_-^\varepsilon(u_\varepsilon(x))| \leq |g_+(u_\varepsilon(x))| + |g_-(u_\varepsilon(x))| \\ &= |g^+(u_\varepsilon(x)) + g^-(u_\varepsilon(x))| = |g(u_\varepsilon(x))| \leq \delta \end{aligned}$$

for every $\varepsilon \in I_2^x \cap (0, \omega_\delta]$, we obtain

$$\limsup_{\varepsilon \rightarrow 0^+, \varepsilon \in I_2^x} |g_\varepsilon(u_\varepsilon(x))| \leq \delta;$$

by the arbitrariness of δ , we achieve

$$\lim_{\varepsilon \rightarrow 0^+, \varepsilon \in I_2^x} g_\varepsilon(u_\varepsilon(x)) = 0;$$

on the other hand, by definition of I_2^x , we have

$$\lim_{\varepsilon \rightarrow 0^+, \varepsilon \in I_2^x} |u_\varepsilon(x)| \leq \lim_{\varepsilon \rightarrow 0^+, \varepsilon \in I_2^x} \varepsilon = 0$$

and thus $u_0(x) = \lim_{\varepsilon \rightarrow 0^+, \varepsilon \in I_2^x} u_\varepsilon(x) = 0$; since $g(u_0(x)) = g(0) = 0$, we have reached

$$\lim_{\varepsilon \rightarrow 0^+, \varepsilon \in I_2^x} g_\varepsilon(u_\varepsilon(x)) = g(u_0(x)),$$

which is the claim. We prove now the claim on G_ε . Obviously, $G_+(u_\varepsilon(x)) \rightarrow G_+(u_0(x))$; noticed that, in a similar way as before, we can show $g_-^\varepsilon(u_0(x)t) \rightarrow g_-^\varepsilon(u_0(x)t)$ for each $t \in [0, 1]$, and observed that

$$|g_-^\varepsilon(u_\varepsilon(x)t)| \leq |g_-(u_\varepsilon(x)t)| \leq 1 + |u_\varepsilon(x)|^{2^*} \leq 2 + |u_0(x)|^{2^*}$$

for $0 < \varepsilon \ll 1$ and any $t \in [0, 1]$, we can use the dominated convergence theorem and obtain

$$G_-^\varepsilon(u_\varepsilon(x)) = u_\varepsilon(x) \int_0^1 g_-^\varepsilon(u_\varepsilon(x)t) dt \rightarrow u_0(x) \int_0^1 g_-(u_0(x)t) dt = G_-(u_0(x)). \quad \blacksquare$$

We observe that the limit function obtained in (6.1), at the moment, could be trivial. In what follows, we show that this is not the case and, moreover, u_0 is a solution enjoying minimality properties.

Proposition 6.4 *Let $(\mu_0, u_0) \in (0, +\infty) \times H_{\text{rad}}^s(\mathbb{R}^N)$ as in (6.1). Then $u_0 \in \mathcal{S}_m$, $(\mu_0, u_0) \in \mathcal{P}$, it is a solution to (1.8), and*

$$K(u_0) = I^m(\mu_0, u_0) = \ell^m = \kappa^m = d_+^m = d^m < 0.$$

Moreover, $u_\varepsilon \rightarrow u_0$ in $H^s(\mathbb{R}^N)$, $\int_{\mathbb{R}^N} G_-^\varepsilon(u_\varepsilon) dx \rightarrow \int_{\mathbb{R}^N} G_-(u_0) dx < \infty$, and $d_{\varepsilon,+}^m \rightarrow d_+^m$ as $\varepsilon \rightarrow 0^+$.

Proof. We show first that u_0 minimizes K over the L^2 -ball. Since $G_\varepsilon \geq G$, we have

$$\ell_\varepsilon^m = \inf_{\mathcal{D}_m^{\text{rad}}} K_\varepsilon \leq \inf_{\mathcal{D}_m^{\text{rad}}} K \leq K(u_0).$$

On the other hand, by exploiting that $(-\Delta)^{s/2}u_\varepsilon \rightharpoonup (-\Delta)^{s/2}u_0$ in $L^2(\mathbb{R}^N)$, the strong convergence in $L^1(\mathbb{R}^N)$ of $G_+(u_\varepsilon)$, and Fatou's lemma for $G_-^\varepsilon(u_\varepsilon)$ (here we use Lemma 6.3),

$$\begin{aligned} K(u_0) &= \frac{1}{2}|(-\Delta)^{s/2}u_0|_2^2 - \int_{\mathbb{R}^N} G_+(u_0) dx + \int_{\mathbb{R}^N} G_-(u_0) dx \\ &\leq \liminf_{\varepsilon \rightarrow 0^+} \left(\frac{1}{2}|(-\Delta)^{s/2}u_\varepsilon|_2^2 \right) - \lim_{\varepsilon \rightarrow 0^+} \left(\int_{\mathbb{R}^N} G_+(u_\varepsilon) dx \right) + \liminf_{\varepsilon \rightarrow 0^+} \left(\int_{\mathbb{R}^N} G_-^\varepsilon(u_\varepsilon) dx \right) \\ &\leq \liminf_{\varepsilon \rightarrow 0^+} K_\varepsilon(u_\varepsilon) = \liminf_{\varepsilon \rightarrow 0^+} \ell_\varepsilon^m \leq \limsup_{\varepsilon \rightarrow 0^+} \ell_\varepsilon^m \leq \inf_{\mathcal{D}_m^{\text{rad}}} K \leq K(u_0). \end{aligned}$$

This yields that u_0 is a minimum over the L^2 -ball and the following holds as $\varepsilon \rightarrow 0^+$: $\ell_\varepsilon^m \rightarrow \ell^m$, $\int_{\mathbb{R}^N} G_-^\varepsilon(u_\varepsilon) dx \rightarrow \int_{\mathbb{R}^N} G_-(u_0) dx < \infty$, and $(-\Delta)^{s/2}u_\varepsilon \rightarrow (-\Delta)^{s/2}u_0$ in $L^2(\mathbb{R}^N)$. By the fact that $\ell_\varepsilon^m \leq \beta_m < 0$, we obtain $K(u_0) \leq \beta_m < 0$; in particular, $u_0 \neq 0$. Moreover, by Proposition 4.2 and Proposition 1.5 we know that (μ_0, u_0) is a solution of (1.1), it satisfies the Pohožaev identity, $u_0 \in \mathcal{S}_m$ (whence $I^m(\mu, u_0) = K(u_0)$), and $\kappa^m = \ell^m = d_+^m = d^m$, thus $u_\varepsilon \rightarrow u_0$ in $L^2(\mathbb{R}^N)$ and hence in $H^s(\mathbb{R}^N)$. \blacksquare

To show the Nehari identity, we need the following lemma.

Lemma 6.5 *The solution u_0 given in (6.1) satisfies $\int_{\mathbb{R}^N} g_-(u_0)u_0 dx < \infty$.*

Proof. Let us recall that, as $\varepsilon \rightarrow 0^+$, $\mu_\varepsilon \rightarrow \mu_0$, $u_\varepsilon \rightarrow u_0$ in $H^s(\mathbb{R}^N)$ and a.e. in \mathbb{R}^N , and $g_\varepsilon(u_\varepsilon) \rightarrow g(u_0)$ a.e. in \mathbb{R}^N . Additionally, $(\mu_\varepsilon, u_\varepsilon)$ satisfies the Nehari identity

$$|(-\Delta)^{s/2}u_\varepsilon|_2^2 + \mu_\varepsilon|u_0|_2^2 = \int_{\mathbb{R}^N} g_+(u_\varepsilon) - g_-^\varepsilon(u_\varepsilon) dx.$$

Since $g_-(t)t \geq 0$ and $\varphi_\varepsilon(t) \geq 0$ for all $t \in \mathbb{R}$, we also have that $g_-^\varepsilon(t)t \geq 0$ for all $t \in \mathbb{R}$, hence u_0 satisfies the following ‘‘Nehari inequality’’

$$\begin{aligned} |(-\Delta)^{s/2}u_0|_2^2 + \mu_0|u_0|_2^2 + \int_{\mathbb{R}^N} g_-(u_0)u_0 dx &\leq \lim_{\varepsilon \rightarrow 0^+} |(-\Delta)^{s/2}u_\varepsilon|_2^2 + \mu_\varepsilon|u_\varepsilon|_2^2 + \int_{\mathbb{R}^N} g_-^\varepsilon(u_\varepsilon)u_\varepsilon dx \\ &= \lim_{\varepsilon \rightarrow 0^+} \int_{\mathbb{R}^N} g_+(u_\varepsilon)u_\varepsilon dx = \int_{\mathbb{R}^N} g_+(u_0)u_0 dx \end{aligned}$$

thanks to (g1), (g3), and Fatou's lemma. In particular, this gives the claim. \blacksquare

Proof of Theorem 1.1. The theorem is a consequence of Lemma 6.5, Proposition 4.3, Proposition 6.4, Proposition 1.6 and Proposition 1.10. \blacksquare

Proof of Proposition 1.12 (iv). From Theorem 5.2, for every $\varepsilon \in (0, 1]$ there exists $(\lambda_\varepsilon, u_\varepsilon) \in \mathcal{P}_{\varepsilon,+}$ Schwarz-symmetric such that $I_\varepsilon^m(\lambda_\varepsilon, u_\varepsilon) = \inf_{\mathcal{P}_{\varepsilon,+}} I_\varepsilon^m = \inf_{\mathcal{P}_{\varepsilon,+}^{\text{rad}}} I_\varepsilon^m$. Then, repeating all the steps in Section 6, not only do we re-obtain Theorem 1.1, but we also get the same outcome where κ^m , ℓ^m , $a(\mu)$, and d^m are replaced, respectively, with $\inf_{\mathcal{S}_m} K$, $\inf_{\mathcal{D}_m} K$, $\inf_{\mathcal{P}_\mu} J_\mu$, and $\inf_{\mathcal{P}} I^m$, but the pair (μ_0, u_0) is the same in both cases. In particular, $\kappa^m = \inf_{\mathcal{S}_m} K$, $\ell^m = \inf_{\mathcal{D}_m} K$, $a(\mu) = \inf_{\mathcal{P}_\mu} J_\mu$, and $d^m = \inf_{\mathcal{P}} I^m$. \blacksquare

7 Further existence results

7.1 Existence for all masses

In Theorem 1.1, we obtained the existence of a solution for large masses. When we rule out the possibility of g to be superlinear at the origin (or positive nearby), which clearly mimics the logarithmic case, we can achieve existence also for all masses. Obviously, we already have existence for all masses if $m_0 = 0$, which is why we assume $m_0 > 0$ in Theorem 7.2 below.

Lemma 7.1 *Under assumptions (g0)–(g4), for every $\omega \geq 0$ there exists $\tilde{\mu}_\omega \in (0, +\infty)$ such that*

$$\inf_{0 < \mu < \bar{\mu}_0} \frac{a(\mu)}{\mu + \omega'} = \inf_{0 < \mu < \tilde{\mu}_\omega} \frac{a(\mu)}{\mu + \omega'} \quad (7.1)$$

for every $\omega' \in [0, \omega]$. In particular, $m_0 = \inf_{0 < \mu < \tilde{\mu}_0} \frac{a(\mu)}{\mu/2}$. Moreover,

$$\frac{\tilde{\mu}_\omega}{\omega} \rightarrow 0 \quad \text{as } \omega \rightarrow +\infty.$$

Proof. If $\bar{\mu}_0 < +\infty$, we take $\tilde{\mu}_\omega := \bar{\mu}_0$. Otherwise, we set $M := \inf_{\mu \in (0, +\infty)} \frac{a(\mu)}{\mu}$ and

$$\tilde{\mu}_\omega := \inf \left\{ \tilde{\mu} > 0 \mid \frac{a(\mu)}{\mu + \omega} \geq M + 1 \text{ for each } \mu \geq \tilde{\mu} \right\}.$$

We observe that $\tilde{\mu}_\omega < +\infty$ from Lemma 5.6, while $\tilde{\mu}_\omega > 0$ from its definition and that of M . Additionally, from the continuity of $a: (0, +\infty) \rightarrow [0, +\infty)$ (see Proposition 3.4), we have

$$\frac{a(\tilde{\mu}_\omega)}{\tilde{\mu}_\omega + \omega} = M + 1. \quad (7.2)$$

Consequently, if $\nu \geq \tilde{\mu}_\omega$ and $\omega' \in [0, \omega]$, then

$$\frac{a(\nu)}{\nu + \omega'} \geq M + 1 > \inf_{0 < \mu < +\infty} \frac{a(\mu)}{\mu + \omega'},$$

whence (7.1). We move to the last claim; we can assume $\tilde{\mu}_\omega \rightarrow +\infty$ as $\omega \rightarrow +\infty$ (otherwise, the claim is obvious). Since (7.2) can be rewritten as

$$\frac{\omega}{\tilde{\mu}_\omega} = \frac{1}{M + 1} \frac{a(\tilde{\mu}_\omega)}{\tilde{\mu}_\omega} - 1,$$

the claim follows from Lemma 5.6. ■

Theorem 7.2 *Assume (g0)–(g4), $m_0 > 0$, and*

$$\limsup_{t \rightarrow 0} \frac{g(t)}{t} =: -\bar{\omega} \in [-\infty, 0).$$

Then, there exists $m_1 = m_1(\bar{\omega}) \in [0, m_0)$ such that, for every $m > m_1$, there exists a solution $(\mu_0, u_0) \in \mathbb{R} \times \mathcal{S}_m^{\text{rad}}$ to (1.8), which satisfies the Nehari and Pohožaev identities, and such that

$$I^m(\mu, u_0) = K(u_0) = a(\mu_0) - \frac{\mu_0}{2}m = \kappa^m = d^m = \inf_{-\infty < \nu < \bar{\mu}_0} \left(a(\nu) - \frac{\nu}{2}m \right);$$

if $\kappa^m \leq 0$ then $\mu_0 > 0$. Finally, if $\bar{\omega} = \infty$, then $m_1 = 0$.

Proof. Let $\omega \in (0, \bar{\omega})$ and define, for $t \in \mathbb{R}$,

$$h_\omega(t) := g(t) + \omega t, \quad H_\omega(t) := G(t) + \frac{\omega}{2}t^2.$$

We observe that

- h_ω is continuous and $h_\omega(0) = 0$,
- $\limsup_{|t| \rightarrow +\infty} \frac{|h_\omega(t)|}{|t|^{2_s^* - 1}} < \infty$,
- there exists $t_0 \neq 0$ such that $H_\omega(t_0) > 0$.

Moreover, it is straightforward to check that $\limsup_{t \rightarrow 0} \frac{h_\omega(t)}{t} < 0$ and $\limsup_{|t| \rightarrow +\infty} \frac{h_\omega(t)}{|t|^{p-1}t} \leq 0$, thus h_ω satisfies the assumptions of Theorem 1.1. In particular, for

$$m > m_\omega := \inf_{0 < \mu < \bar{\mu}_\omega} \frac{a_\omega(\mu)}{\mu/2}$$

there exists an L^2 -minimum $u_\omega \in \mathcal{S}_m^{\text{rad}}$ with Lagrange multiplier $\mu_\omega > 0$, i.e.,

$$K_\omega(u_\omega) = \inf_{\mathcal{D}_m^{\text{rad}}} K_\omega = \inf_{\mathcal{S}_m^{\text{rad}}} K_\omega =: \kappa_\omega^m < 0; \quad (7.3)$$

here

$$\bar{\mu}_\omega := \sup_{t \neq 0} \frac{H_\omega(t)}{t^2/2} = \bar{\mu}_0 + \omega, \quad K_\omega(u) := \frac{1}{2}|(-\Delta)^{s/2}u|_2^2 - \int_{\mathbb{R}^N} H_\omega(u) \, dx,$$

$$a_\omega(\mu) := \inf_{\mathcal{P}_{\mu,\omega}^{\text{rad}}} J_{\mu,\omega}, \quad J_{\mu,\omega}(u) := \frac{1}{2}|(-\Delta)^{s/2}u|_2^2 + \frac{\mu}{2}|u|_2^2 - \int_{\mathbb{R}^N} H_\omega(u) \, dx,$$

$$P_{\mu,\omega}(u) := \int_{\mathbb{R}^N} |(-\Delta)^{s/2}u|^2 \, dx + 2_s^* \int_{\mathbb{R}^N} \frac{\mu}{2}u^2 - H_\omega(u) \, dx,$$

$$\mathcal{P}_{\mu,\omega} := \{u \in H^s(\mathbb{R}^N) \setminus \{0\} \mid P_{\mu,\omega}(u) = 0\}, \quad \mathcal{P}_{\mu,\omega}^{\text{rad}} := \mathcal{P}_{\mu,\omega} \cap H_{\text{rad}}^s(\mathbb{R}^N).$$

We notice that (μ_ω, u_ω) satisfies the Nehari and Pohožaev identities with respect to H_ω . Since $K_\omega(u) = K(u) - \frac{\omega}{2}|u|_2^2$, we obtain

$$K(u_\omega) - \frac{\omega}{2}m = K(u_\omega) - \frac{\omega}{2}|u_\omega|_2^2 = K_\omega(u_\omega) = \inf_{\mathcal{S}_m^{\text{rad}}} K_\omega = \kappa_\omega^m - \frac{\omega}{2}m,$$

whence $K(u_\omega) = \kappa_\omega^m$, which means that u_ω is an L^2 -minimum for the original problem. Moreover,

$$\kappa^m = \kappa_\omega^m + \frac{\omega}{2}m. \quad (7.4)$$

In addition, u_ω is a solution to the Euler–Lagrange equation corresponding to K_ω , i.e., for some Lagrange multiplier $\mu_\omega > 0$,

$$(-\Delta)^{s/2}u_\omega + \mu_\omega u_\omega = h_\omega(u_\omega) \quad \text{in } \mathbb{R}^N,$$

which means that $u_0 := u_\omega$ is a solution to (1.8) with Lagrange multiplier

$$\mu_0 := \mu_\omega - \omega \quad (7.5)$$

and (μ_0, u_0) satisfies the Nehari and Pohožaev identities with respect to G . Additionally, for any $\nu > \omega$, we have $J_{\nu, \omega}(u) = J_{\nu - \omega}(u)$ and $P_{\nu, \omega}(u) = P_{\nu - \omega}(u)$, therefore

$$\frac{1}{2}m_\omega \leq \inf_{\omega < \nu < \tilde{\mu}_0 + \omega} \frac{a_\omega(\nu)}{\nu} = \inf_{0 < \nu - \omega < \tilde{\mu}_0} \frac{a(\nu - \omega)}{\nu} = \inf_{0 < \nu < \tilde{\mu}_0} \frac{a(\nu)}{\nu + \omega} = \inf_{0 < \nu < \tilde{\mu}_0} \frac{a(\nu)}{\nu} \frac{\nu}{\nu + \omega}. \quad (7.6)$$

Let $\tilde{\mu} = \tilde{\mu}_\omega \in (0, +\infty)$ be the value given in Lemma 7.1. Since the function $\nu \in (0, +\infty) \mapsto \frac{\nu}{\nu + \omega} \in (0, +\infty)$ is increasing, from (7.6) we obtain

$$\frac{1}{2}m_\omega \leq \inf_{0 < \nu < \tilde{\mu}} \frac{a(\nu)}{\nu} \frac{\nu}{\nu + \omega} \leq \frac{\tilde{\mu}}{\tilde{\mu} + \omega} \inf_{0 < \nu < \tilde{\mu}} \frac{a(\nu)}{\nu} = \frac{\tilde{\mu}}{\tilde{\mu} + \omega} \inf_{0 < \nu < \tilde{\mu}_0} \frac{a(\nu)}{\nu} = \frac{\tilde{\mu}}{\tilde{\mu} + \omega} \frac{m_0}{2} < \frac{1}{2}m_0.$$

In addition, if $\bar{\omega} = \infty$, taking ω arbitrarily large we obtain

$$m_\omega \rightarrow 0 \quad \text{as } \omega \rightarrow +\infty$$

again from Lemma 7.1. As a consequence, for any $m > 0$, there exists $m_\omega \in [0, m)$ and thus an L^2 -minimum u_ω . We conclude setting

$$m_1 := \inf_{0 < \omega < \bar{\omega}} m_\omega;$$

noticed that $\omega \mapsto m_\omega$ is nonincreasing (this is a direct consequence of $J_{\mu, \omega} \leq J_\mu$, $P_{\mu, \omega} \leq P_\mu$, $\mu_\omega \geq \mu_0$, and an argument similar to (5.9)–(5.13)), we further observe that $m_1 = \lim_{\omega \rightarrow \bar{\omega}^-} m_\omega$. The sign of μ_0 is a simple consequence of the Pohožaev identity. Finally, the fact that

$$K(u_0) = I^m(\mu_0, u_0) = d^m = \inf_{-\infty < \nu < \tilde{\mu}_0} \left(a(\nu) - \frac{\nu}{2} m \right) = a(\mu_0) - \frac{\mu_0}{2} m$$

follows from Propositions 1.5, 1.6, and 1.10. ■

Remark 7.3 *By definition, m_1 is nonincreasing with respect to $\bar{\omega}$. Moreover, in the case of small masses, Theorem 7.2 gives no information on the sign of the Lagrange multiplier and of the minimal energy – see indeed (7.4) and (7.5).*

Proof of Theorem 1.3. It is a consequence of Theorem 7.2. ■

7.2 Existence under relaxed assumptions

We deal now with Proposition 1.14. We thus search for a normalized solution which is nonnegative, possibly considering also the case

$$\limsup_{|t| \rightarrow +\infty} \frac{g_-(t)}{|t|^{2^*_s - 1}} = +\infty.$$

The idea is to modify the nonlinearity g . We begin with the following result.

Lemma 7.4 *Let (g0)–(g2) hold. Assume $g(t) = 0$ for $|t| \geq t_1$, and let u be a nonnegative solution of (1.1) for some $\mu \geq 0$. Then $u \in L^\infty(\mathbb{R}^N)$ and $|u|_\infty \leq t_1$.*

Proof. Let $v := (u - t_1)^+$. It is easy to check, via explicit computations, that $v \in H^s(\mathbb{R}^N)$. Let φ_R be as in Lemma 2.2 and set $v_R := v\varphi_R$. Then $v_R \in H^s(\mathbb{R}^N) \cap L^1(\mathbb{R}^N)$ and hence

$$\frac{1}{2} \int_{\mathbb{R}^N} (-\Delta)^{s/2} u (-\Delta)^{s/2} v_R dx + \frac{\mu}{2} \int_{\mathbb{R}^N} u v_R dx = \int_{\mathbb{R}^N} g(u) v_R dx.$$

We notice that $g(u)v_R \equiv 0$ by the assumptions. Since $u \geq 0$ and $\mu \geq 0$, we have $\mu uv_R \geq 0$ and thus

$$\int_{\mathbb{R}^N} (-\Delta)^{s/2} u (-\Delta)^{s/2} v_R \, dx \leq 0.$$

By Lemma 2.2 we have

$$\int_{\mathbb{R}^N} (-\Delta)^{s/2} u (-\Delta)^{s/2} (u - t_1)^+ \, dx \leq 0.$$

On the other hand,

$$\begin{aligned} & \int_{\mathbb{R}^N} (-\Delta)^{s/2} u (-\Delta)^{s/2} (u - t_1)^+ \, dx \\ &= C_{N,s} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{((u(x) - t_1) - (u(y) - t_1))((u(x) - t_1)^+ - (u(y) - t_1)^+)}{|x - y|^{N+2s}} \\ &\geq C_{N,s} \int_{\{u(x) \geq t_1\}} \int_{\{u(y) \geq t_1\}} \frac{|(u(x) - t_1)^+ - (u(y) - t_1)^+|^2}{|x - y|^{N+2s}} \\ &= C_{N,s} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|(u(x) - t_1)^+ - (u(y) - t_1)^+|^2}{|x - y|^{N+2s}} = |(-\Delta)^{s/2} (u - t_1)^+|_2^2. \end{aligned}$$

Joining the two relations found we obtain $(u - t_1)^+ = 0$, i.e., $u \leq t_1$. Finally, using $v_R := (u + t_1)^- \varphi_R$ and $g(t) = 0$ for $t \leq -t_1$, we similarly obtain $u \geq -t_1$. \blacksquare

Proof of Proposition 1.14. Let us assume that t_0 , defined in (g4), is positive (if it is negative, then we argue likewise). This implies that g is positive somewhere on $(0, +\infty)$. Let t_1 be the first positive zero of g such that g is positive somewhere on $(0, t_1)$ (if $g(t) > 0$ for all $t > 0$, we set $t_1 = +\infty$). Then, we set

$$\bar{g}(t) := g(t) \chi_{(0, t_1)}(t).$$

Clearly, \bar{g} satisfies (g0)–(g4), thus Proposition 1.12 applies (or Remark 1.13 if (1.21) holds), and we have a Schwarz-symmetric solution (μ_0, u_0) , which satisfies the Nehari and Pohožaev identities, to

$$(-\Delta)^s u + \mu u = \bar{g}(u).$$

If $t_1 = +\infty$, then $\bar{g}(u_0) = g(u_0)$, so (μ_0, u_0) is a solution to (1.8); otherwise, from Lemma 7.4, $|u_0|_\infty \leq t_1$, hence $\bar{g}(u_0) = g(u_0)$ once again, and we conclude as before. \blacksquare

We end this section with the examples mentioned in the introduction.

Proof of Corollary 1.15. As in [45, Proof of Theorem 1.6], we prove that when $\beta < 0$ and $q \in (1, 2_s^* - 1]$, $\max G \leq 0$ if and only if $\beta \leq -\frac{\alpha(q+1)}{q-1} e^{-(q+1)/2}$, hence the second part (existence) follows from Theorem 1.3 and Proposition 1.12 (iv) (see also Remark 1.13). If $m > m_0$, then the claim follows by Theorem 1.1. The first part (nonexistence) is a direct consequence of the Pohožaev identity, together with Proposition 1.8. \blacksquare

Proof of Corollary 1.16. We argue as in Corollary 1.15. \blacksquare

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