

On the shape of hypersurfaces with boundary which have zero fractional mean curvature

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Abstract

We consider compact hypersurfaces with boundary in \mathbb{R}^N that are the critical points of the fractional area introduced by Paroni, Podio-Guidugli, and Seguin in [26]. In particular, we study the shape of such hypersurfaces in several simple settings. First we consider the critical points whose boundary is a smooth, orientable, closed manifold Γ of dimension $N - 2$ and lies in a hyperplane $H \subset \mathbb{R}^N$. Then we show that the critical points coincide with a smooth manifold $\mathcal{N} \subset H$ of dimension $N - 1$ with $\partial\mathcal{N} = \Gamma$. Second we consider the critical points whose boundary consists of two smooth, orientable, closed manifolds Γ_1 and Γ_2 of dimension $N - 2$ and suppose that Γ_1 lies in a hyperplane H perpendicular to the x_N -axis and that $\Gamma_2 = \Gamma_1 + d e_N$ ($d > 0$ and $e_N = (0, \dots, 0, 1) \in \mathbb{R}^N$). Then, assuming that Γ_1 has a non-negative mean curvature, we show that the critical points do not coincide with the union of two smooth manifolds $\mathcal{N}_1 \subset H$ and $\mathcal{N}_2 \subset H + d e_N$ of dimension $N - 1$ with $\partial\mathcal{N}_i = \Gamma_i$ for $i \in \{1, 2\}$. Moreover, the interior of the critical points does not intersect the boundary of the convex hull in \mathbb{R}^N of Γ_1 and Γ_2 , while this can occur in the codimension-one situation considered by Dipierro, Onoue, and Valdinoci in [14]. We also obtain a quantitative bound which may tell us how different the critical points are from $\mathcal{N}_1 \cup \mathcal{N}_2$. Finally, in the same setting as in the second case, we show that, if d is sufficiently large, then the critical points are disconnected and, if d is sufficiently small, then Γ_1 and Γ_2 are in the same connected component of the critical points when $N \geq 3$. Moreover, by computing the fractional mean curvature of a cone whose boundary is $\Gamma_1 \cup \Gamma_2$, we also obtain that the interior of the critical points does not touch the cone if the critical points are contained in either the inside or the outside of the cone.

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1 Introduction

Fractional minimal surfaces without boundary were first investigated by Caffarelli, Roquejoffre, and Savin in [6] and, since then, this topic has attracted many authors to study their geometric properties as an analogy of classical minimal surfaces. Roughly speaking, a fractional (or nonlocal) minimal surface without boundary is given as the boundary of a set which minimizes an energy functional defined by the pointwise interaction of a set and its complement. The typical interaction taken into account is scaling and translation invariant with some polynomial decay. Precisely, if $s \in (0, 1)$ and Ω is an open set with smooth boundary, one of such standard energies of a set $E \subset \mathbb{R}^N$ relative to Ω is the so-called fractional perimeter in Ω and is defined by

$$P_s(E; \Omega) := \int_{E \cap \Omega} \int_{E^c} \frac{dx dy}{|x - y|^{N+s}} + \int_{E \cap \Omega^c} \int_{E \cap \Omega} \frac{dx dy}{|x - y|^{N+s}} \quad (1.1)$$

where we denote by E^c the complement of E . With this notion, we say that a set $E \subset \mathbb{R}^N$ is a minimizer of P_s relative to Ω if it holds that

$$P_s(E; \Omega') \leq P_s(F; \Omega)$$

for any open bounded set $\Omega' \subset \Omega$ and any $F \subset \mathbb{R}^N$ with $F \setminus \Omega' = E \setminus \Omega'$. The existence and regularity of such minimizers was shown by Caffarelli, Roquejoffre, and Savin in [6]. The regularity theory of the minimizers were later strengthened in, for instance, [9, 30]. Moreover, they showed in [6] that if a set $E \subset \mathbb{R}^N$ is a minimizer of P_s , then the following Euler-Lagrange equation holds in the viscosity sense:

$$\int_{\mathbb{R}^N} \frac{\chi_{E^c}(y) - \chi_E(y)}{|y - x|^{N+s}} dy = 0 \quad (1.2)$$

for $x \in \partial E$. The integral in (1.2) is intended in the Cauchy principal value sense. This can be regarded as a nonlocal counterpart of the classical minimal surface equation and the left-hand side in (1.2) is the so-called fractional mean curvature on the boundary ∂E . Dipierro, Savin, and Valdinoci in particular have revealed many properties which classical minimal surfaces cannot possess (see, for instance, [16, 17] for the details). In addition, many authors have studied the fractional(nonlocal) minimal surfaces or minimal graphs for more than a decade since the fractional(nonlocal) minimal surfaces appear in many other topics in which a long-range interaction is involved (see [11, 29]). For further discussions about the geometric features of fractional(nonlocal) minimal surfaces without boundary, we refer to [2, 4, 5, 7, 9, 12–15, 18, 19].

Quite recently, motivated by some mathematical modelling of thin elastic structures, Paroni, Podio-Guidugli, and Seguin in [26] introduced a new notion of fractional area and fractional mean curvatures for smooth manifolds which are not necessarily closed in the following way: let $\Omega \subset \mathbb{R}^N$ be a bounded domain and let $\mathcal{M} \subset \Omega$ be any $(N - 1)$ -dimensional compact smooth manifold with or without boundary. Here we mean by “a manifold without boundary” a closed manifold. Then the fractional area of \mathcal{M} relative to Ω is defined by

$$\text{Area}_s(\mathcal{M}; \Omega) := c_N \iint_{\mathcal{X}(\mathcal{M})} \frac{\max\{\chi_\Omega(x), \chi_\Omega(y)\}}{|x - y|^{N+s}} dx dy \quad (1.3)$$

where c_N is some positive dimensional constant and $\mathcal{X}(\mathcal{M})$ is a set of all pairs $(x, y) \in \mathbb{R}^N \times \mathbb{R}^N$ such that the segment $[x; y]$ with two end points x and y has an odd number of cross intersections with \mathcal{M} and $[x; y]$ is not tangent to \mathcal{M} . Note that the presence of the term $\max\{\chi_\Omega(x), \chi_\Omega(y)\}$ in (1.3) is necessary to ensure that the integral converges whenever $\partial\mathcal{M} \neq \emptyset$.

As is explained in [26], if a $(N-1)$ -dimensional smooth manifold \mathcal{M} satisfies $\mathcal{M} = \partial E$ for some set $E \subset \mathbb{R}^N$, then the two notions (1.1) and (1.3) are equivalent, i.e., it holds that

$$\text{Area}_s(\mathcal{M}; \Omega) = P_s(E; \Omega).$$

Interestingly, Paroni, Podio-Guidugli, and Seguin also proved in [26, Theorem 3.3] that $(1-s)\text{Area}_s(\mathcal{M}; \Omega) \rightarrow \mathcal{H}^{N-1}(\mathcal{M})$ as $s \uparrow 1$ for a compact $(N-1)$ -dimensional C^1 manifold \mathcal{M} contained in a bounded domain Ω , as it happens for P_s in (1.1) (see [1, 8]). See [24, 27, 31] for further discussions on Area_s .

This manuscript is devoted to develop the theory of the fractional area Area_s for manifolds with boundary. In particular, we aim to investigate the shape and topology of critical points of Area_s . Here the critical points of Area_s are defined as a smooth manifold such that the first variation of Area_s vanishes with respect to a perturbation associated with the unit normal vector of that manifold (in the sequel, we will call these perturbations “normal variations”). The authors in [26] obtained a necessary and sufficient condition for the vanishing of the first variation for manifolds as follows: let \mathcal{M} be an orientable compact smooth manifold with or without boundary and assume that \mathcal{M} is contained in a bounded domain $\Omega \subset \mathbb{R}^N$. Then it holds that

$$\delta\text{Area}_s(\mathcal{M}; \Omega) = 0 \iff H_{\mathcal{M},s}(z) = 0 \quad \text{for any } z \in \mathcal{M}. \quad (1.4)$$

Here we denote by $\delta\text{Area}_s(\mathcal{M}; \Omega)$ the first variation of \mathcal{M} under normal variations and $H_{\mathcal{M},s}$ is the fractional mean curvature associated with Area_s which is defined by

$$H_{\mathcal{M},s}(z) := c_N \int_{\mathbb{R}^N} \frac{\chi_{\mathcal{A}_i(z)}(y) - \chi_{\mathcal{A}_e(z)}(y)}{|y-z|^{N+s}} dy$$

for any $z \in \mathcal{M}$ where c_N is as in (1.3) and the sets $\mathcal{A}_i(z)$ and $\mathcal{A}_e(z)$ are defined by

$$\begin{aligned} \mathcal{A}_i(z) := \{y \in \mathbb{R}^N \mid & \text{either } (z, y) \in \mathcal{X}(\mathcal{M}) \ \& \ (z-y) \cdot \nu_{\mathcal{M}}(z) < 0 \\ & \text{or } (z, y) \notin \mathcal{X}(\mathcal{M}) \ \& \ (z-y) \cdot \nu_{\mathcal{M}}(z) > 0\} \end{aligned} \quad (1.5)$$

$$\begin{aligned} \mathcal{A}_e(z) := \{y \in \mathbb{R}^N \mid & \text{either } (z, y) \in \mathcal{X}(\mathcal{M}) \ \& \ (z-y) \cdot \nu_{\mathcal{M}}(z) > 0 \\ & \text{or } (z, y) \notin \mathcal{X}(\mathcal{M}) \ \& \ (z-y) \cdot \nu_{\mathcal{M}}(z) < 0\} \end{aligned} \quad (1.6)$$

where $\nu_{\mathcal{M}}$ is the unit normal vector of \mathcal{M} . The sets $\mathcal{A}_i(z)$ and $\mathcal{A}_e(z)$ can be regarded as the “interior” and “exterior” of \mathcal{M} relative to the point z , respectively, and these sets are determined uniquely once the unit normal vector of \mathcal{M} at z is specified. See [26] for more discussions on the notions. Note that if a manifold is not orientable, then the unit normal vector of the manifold cannot be determined uniquely and neither can the “interior” \mathcal{A}_i and “exterior” \mathcal{A}_e . Moreover, in this paper, we require the $C^{1,\alpha}$ -regularity with $\alpha > s$ of hypersurfaces so that the fractional mean curvatures are finite everywhere.

The study of critical points or fractional minimal surfaces with boundary can be related to the classical problem on free boundary minimal surfaces in differential geometry. One of the main topics in the problem is to determine the shape of a manifold Σ (embedded or immersed) in another smooth manifold \mathcal{S} such that Σ minimizes its area in

\mathcal{S} and $\partial\Sigma \subset \partial\mathcal{S}$ with some topological constraints. The study of this classical problem was first considered by R. Courant in [10] in 1940 and, since then, a lot of authors have been intensively working on this topic. See, for instance, [21, 23, 25, 28, 32] for the details. We also refer the readers to two surveys: [20] for classical works and [22] for more recent results. The references here are obviously not exhaustive.

As an analogy of the classical free boundary minimal surfaces, it is natural to consider a fractional(nonlocal) version of free boundary minimal surfaces; however, the nonlocal version is not understood so far because, to our knowledge, suitable notions of fractional area for manifolds with boundary had not been considered until Paroni, Podio-Guidugli, and Seguin in [26] introduced the notion of Area_s in (1.3). To tackle the nonlocal version of the free boundary minimal surface problem, it is important to understand the geometric properties of critical points of Area_s .

Given the importance of critical points of Area_s from the above perspective, it is desirable to develop some intuition about their geometric features. To do this, since it is quite difficult to have explicit solutions which entirely describe critical points or minimizers of Area_s , it is often convenient to study simplified cases in which the boundary of the critical points has some special characteristics. In this paper, we basically consider three cases: the first is when the boundary of critical points in \mathbb{R}^N lies in a hyperplane and is homeomorphic ¹ to \mathbb{S}^{N-2} . The second is when the boundary is the union of two distinct parallel and co-axial manifolds each of which lies in a hyperplane, is homeomorphic to \mathbb{S}^{N-2} , and the distance between the co-axial manifolds is d . The last is when the distance d is sufficiently large or sufficiently small.

Our first goal in this paper is to determine the shape of critical points of Area_s whose boundary lies on a hyperplane. Precisely, we first define a set $\mathcal{C} \subset \mathbb{R}^N$ as

$$\mathcal{C} := \mathcal{G} \times \mathbb{R} \tag{1.7}$$

where \mathcal{G} is a non-empty bounded open subset of \mathbb{R}^{N-1} with a smooth boundary. Then we define an $(N-2)$ -dimensional smooth manifold Γ as

$$\Gamma_0 := \partial\mathcal{C} \cap \{x_N = 0\} (= \partial\mathcal{G} \times \{0\}). \tag{1.8}$$

Assume that $\mathcal{M} \subset \mathbb{R}^N$ is an orientable compact $(N-1)$ -dimensional $C^{1,\alpha}$ manifold with $\partial\mathcal{M} = \Gamma_0$ and that \mathcal{M} is a critical point of Area_s . Note that the orientability of \mathcal{M} implies the orientability of $\partial\mathcal{M} = \Gamma_0$. Then, as our first theorem, we aim to rigorously prove that \mathcal{M} must coincide with $\mathcal{C} \cap \{x_N = 0\}$, as we can intuitively expect this to be true.

Theorem 1.1. *Let $s \in (0, 1)$. Let Γ_0 be as in (1.8). Let \mathcal{M} be an orientable compact $(N-1)$ -dimensional $C^{1,\alpha}$ manifold with $\partial\mathcal{M} = \Gamma_0$. If \mathcal{M} is a critical point of Area_s under normal variations, then \mathcal{M} is a hyperplane lying on $\{x_N = 0\}$, i.e.,*

$$\mathcal{M} = \bar{\mathcal{C}} \cap \{x_N = 0\} (= \bar{\mathcal{G}} \times \{0\}).$$

Our second goal in this paper is to study the shape of critical points of Area_s whose boundary consists of two disjoint components. The problem setting in the second theorem

¹Our result (Theorem 1.1) can be also true even if the boundary of the critical point is not always homeomorphic to \mathbb{S}^{N-2} . For instance, Theorem 1.1 holds true if the boundary has more than one connected components and lies in some hyperplane. The proof is the same as the one of Theorem 1.1.

is as follows: we define two distinct compact $(N - 2)$ -dimensional smooth manifolds Γ_1 and Γ_2 by

$$\Gamma_1 := \partial\mathcal{C} \cap \{x_N = h_1\} \quad \text{and} \quad \Gamma_2 := \partial\mathcal{C} \cap \{x_N = h_2\}, \quad (1.9)$$

where \mathcal{C} is as in (1.7) and h_1 and h_2 are given constants with $h_2 < h_1$. Then a critical point exhibits a different shape from a hyperplane. Precisely we prove

Theorem 1.2. *Let $s \in (0, 1)$. Let Γ_1 and Γ_2 be as in (1.9) and let \mathcal{M} be an orientable compact $(N - 1)$ -dimensional $C^{1,\alpha}$ manifold with $\partial\mathcal{M} = \Gamma_1 \cup \Gamma_2$. Assume that \mathcal{C} is convex where \mathcal{C} is as in (1.7). If \mathcal{M} is a critical point of Area_s under normal variations, then $\mathcal{M} \subset \{h_2 \leq x_N \leq h_1\}$ and any connected component of \mathcal{M} is neither C_1 nor C_2 where we define*

$$C_1 := \bar{\mathcal{C}} \cap \{x_N = h_1\} \quad \text{and} \quad C_2 := \bar{\mathcal{C}} \cap \{x_N = h_2\}.$$

In particular, $\mathcal{M} \neq C_1 \cup C_2$. Moreover, $\mathcal{M} \setminus \partial\mathcal{M}$ does not intersect $\partial\mathcal{C} = \partial\mathcal{G} \times \mathbb{R}$.

We remark that, by using a cone whose boundary is $\Gamma_1 \cup \Gamma_2$ as in Theorem 1.2 with $h_1 = 1$ and $h_2 = -1$, we can further detect how the critical points behave. See Subsection 3.2 of Section 3 for the details.

Our third goal is to further study the shape and, in particular, the topology of critical points of Area_s in the same situation as the one in Theorem 1.2. Precisely, taking Γ_1 and Γ_2 as in Theorem 1.2 with $d := h_1 - h_2 > 0$, we will see what critical points of Area_s under normal variations look like in terms of connectedness if d is sufficiently large or sufficiently small.

To reach the third goal, we first show the following lemma which somehow tells us how different critical points are from hyperplanes.

Lemma 1.3. *Let $s \in (0, 1)$ and $d > 0$. Let Γ_1 and Γ_2 be as in (1.9) with $h_1 = 0$ and $h_2 = -d$. Assume that \mathcal{C} is convex where \mathcal{C} is as in (1.7). Then there exists a constant $\varepsilon_0 > 0$, depending only on N , s , and d , such that the following holds: let \mathcal{M} be an orientable compact $(N - 1)$ -dimensional $C^{1,\alpha}$ manifold with $\partial\mathcal{M} = \Gamma_1 \cup \Gamma_2$. If \mathcal{M} is a critical point of Area_s under normal variations, then the set enclosed by \mathcal{M} and the union of $\mathcal{C} \cap \{x_N = 0\}$ and $\mathcal{C} \cap \{x_N = -d\}$ contains two half-balls*

$$B_{\varepsilon_0}^-(0) := \{x \in B_{\varepsilon_0}(0) \mid x_N < 0\} \quad \text{and} \quad B_{\varepsilon_0}^+(p_d) := \{x \in B_{\varepsilon_0}(p_d) \mid x_N > -d\}$$

where $p_d := (0, -d) \in \mathbb{R}^{N-1} \times \mathbb{R}$.

Here, thanks to Theorem 1.2, we can precisely define the set enclosed by \mathcal{M} and $\mathcal{C} \cap (\{x_N = 0\} \cup \{x_N = -d\})$ as the collection of $x \in \mathbb{R}^N$ such that the line segment $[x, P(x)]$ where $P(x)$ realizes the minimum distance between x and $\mathcal{C} \cap (\{x_N = 0\} \cup \{x_N = -d\})$ intersects \mathcal{M} even number of times and $(x - P(x)) \cdot \nu > 0$ where $\nu = -e_N$ if $P(x) \in \mathcal{C} \cap \{x_N = 0\}$ and $\nu = e_N$ if $P(x) \in \mathcal{C} \cap \{x_N = -d\}$.

To help with the intuition, a sketch of possible critical points is given in Figure 1.

As a result of Lemma 1.3, we prove that, if the distance d between two parallel and co-axial boundaries is sufficiently small, then any critical point is connected in the sense that the two boundaries are in the same connected component when $N \geq 3$. Moreover, when $N = 2$, any critical point is disconnected and its two distinct connected components should look like the right-hand side of Figure 1 with $0 < d \ll 1$.

Precisely, our third theorem is as follows.

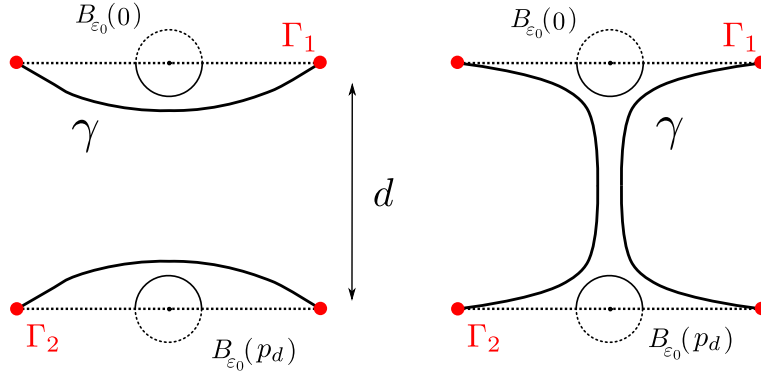


Figure 1: Two possible situations in dimension 2 in Lemma 1.3 in which the ‘interior’ or ‘exterior’ of the critical point $\mathcal{M} = \gamma$ with $\partial\gamma = \Gamma_1 \cup \Gamma_2$ contains two half-balls.

Theorem 1.4. *Let $s \in (0, 1)$. Let Γ_1 and Γ_2 be as in Lemma 1.3. Assume that \mathcal{C} is convex where \mathcal{C} is as in (1.7). Then there exists $d_0 = d_0(N, s) > 0$ such that the following holds: for any $d \in (0, d_0)$, we take any orientable compact $(N - 1)$ -dimensional $C^{1,\alpha}$ manifold $\mathcal{M} \subset \mathbb{R}^N$ with $\partial\mathcal{M} = \Gamma_1 \cup \Gamma_2$. If \mathcal{M} is a critical point of Area_s under normal variations, then Γ_1 and Γ_2 are in the same connected component of \mathcal{M} if $N \geq 3$ and \mathcal{M} is disconnected if $N = 2$.*

Moreover, when $N = 2$, there exist two distinct connected components \mathcal{M}_1 and \mathcal{M}_2 of \mathcal{M} such that $\text{dist}(\mathcal{M}_1, \mathcal{M}_2) \geq c$ with some constant $c > 0$, depending only on N and s , and $\partial\mathcal{M}_i$ intersects both Γ_1 and Γ_2 for each $i \in \{1, 2\}$.

As a counterpart of Theorem 1.4, we prove that, if the distance d between two parallel and co-axial boundaries is sufficiently large, then any critical point is disconnected in any dimension and it should look like the left-hand side of Figure 1 with $d \gg 1$.

Our last theorem is as follows.

Theorem 1.5. *Let $s \in (0, 1)$. Let Γ_1 and Γ_2 be as in Lemma 1.3. Assume that \mathcal{C} is convex where \mathcal{C} is as in (1.7). Then there exists $d_1 = d_1(N, s) > 0$ such that the following holds: for any $d > d_1$, we take any orientable compact $(N - 1)$ -dimensional $C^{1,\alpha}$ manifold $\mathcal{M} \subset \mathbb{R}^N$ with $\partial\mathcal{M} = \Gamma_1 \cup \Gamma_2$. If \mathcal{M} is a critical point of Area_s under normal variations, then \mathcal{M} is disconnected.*

Moreover, there exist two disjoint connected components \mathcal{M}_1 and \mathcal{M}_2 of \mathcal{M} such that $\partial\mathcal{M}_i = \Gamma_i$ for any $i \in \{1, 2\}$.

To help with the intuition, a sketch of possible critical points shown in Theorems 1.2, 1.4, and 1.5 is given in Figure 2

The topological properties in Theorem 1.4 and 1.5 could be expected to be true because Dipierro, Valdinoci, and the author of this paper obtained similar results in [14] on the topology of fractional minimal surfaces without boundary in similar situations. On one hand, they showed that minimizers of P_s in a given cylinder coincides with the cylinder itself for sufficiently small d where d is the distance between two disjoint parallel and co-axial external(boundary) data. On the other hand, they showed that minimizers of P_s in the cylinder are disconnected for sufficiently large d .

Interestingly, however, we show in Theorem 1.2 that the critical points (not necessarily fractional area-minimizing) cannot touch the boundary of the cylinder \mathcal{C} no matter what distance two parallel and co-axial boundaries have, while it is shown in [14] that

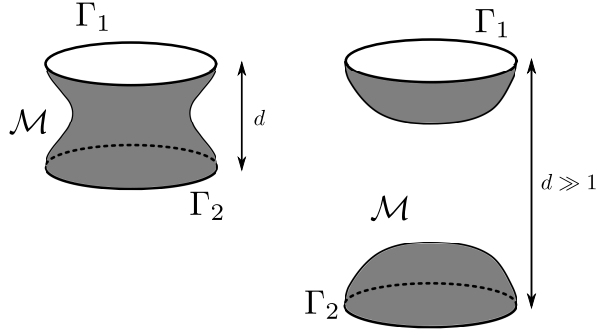


Figure 2: Possible critical points of Area_s when $N \geq 3$ in Theorem 1.2 and Theorem 1.4 on the left and in Theorem 1.5 on the right.

minimizers of P_s in a cylinder favorably stick to the boundary of the cylinder if $N = 2$ and d is large or if $N \geq 2$ and d is small. Moreover, our results together with Remark 4.1 of Section 4 possibly indicate that critical points of Area_s with two nearby parallel and co-axial compact boundaries might develop necks of catenoids, while this is not the case with fractional minimal surfaces considered in [14]. We remark that the existence of fractional minimal catenoids without boundary in \mathbb{R}^3 was shown by Dávila, Del Pino, and Wei in [13] if s is close to 1.

The organization of this paper is as follows: in Section 2, we prove Theorem 1.1 by “sliding” a hyperplane until it touches critical points (see the proof of Theorem 1.1 for the details). In Section 3, we first give the proof of Theorem 1.2 and then we study further properties of critical points of Area_s , computing the fractional mean curvature of a cone passing through the boundary of critical points. In Section 4, we first give the proof of Lemma 1.3 by constructing a suitable barrier and then, by using this lemma, we prove Theorem 1.4. Moreover, in Section 4, we also prove Theorem 1.5 by means of the “sliding method” (see Section 4 for the details).

2 Proof of Theorem 1.1

In this section, we prove Theorem 1.1. The idea of the proof is inspired by the so-called sliding method introduced by Dipierro, Savin, and Valdinoci in [16]. They developed this method in order to investigate the shape of fractional(nonlocal) minimal surfaces (see also [14, 15, 17] for further discussions).

We proceed with the proof in the following way: we slide a hyperplane, parallel to $\mathcal{C} \cap \{x_N = 0\}$, from below or above until it touches \mathcal{M} and assume by contradiction that there exists a touching point in $(\mathcal{C} \cap \{x_N = 0\})^c$. At the touching point q , we obtain the Euler-Lagrange equation (1.4). Then, taking into account all the contributions from the “interior” $\mathcal{A}_i(q)$ and the “exterior” $\mathcal{A}_e(q)$ of \mathcal{M} , we can observe that the contribution from either $\mathcal{A}_i(q)$ or $\mathcal{A}_e(q)$ turns out to be strictly larger than that from the other region. This contradicts the Euler-Lagrange equation.

Proof of Theorem 1.1. We first define a hyperplane $H_\lambda := \{(x', x_N) \mid x_N = \lambda\}$ and two half-spaces

$$H_\lambda^+ := \{(x', x_N) \mid x_N > \lambda\} \quad \text{and} \quad H_\lambda^- := \{(x', x_N) \mid x_N < \lambda\} \quad (2.1)$$

for $\lambda \in \mathbb{R}$. We set $P_\lambda : \mathbb{R}^N \rightarrow \mathbb{R}^N$ as the reflection map with respect to H_λ for $\lambda \in \mathbb{R}$ and set $x_\lambda := P_\lambda(x)$ for any $x \in \mathbb{R}^N$. Moreover, we define a (filled) cone $C_{\Gamma_0}(q)$ with vertex q by

$$\{x \in \mathbb{R}^N \mid \exists \lambda > 0 \text{ such that } q + \lambda(x - q) \in \overline{\mathcal{G}} \times \{0\}\} \cup \{q\}$$

where \mathcal{G} is as in (1.7). Note that $\partial C_{\Gamma_0}(q) \cap \{(x', x_N) \mid x_N = 0\} = \Gamma_0$. We further set $C_{\Gamma_0}^\lambda(q) := P_\lambda(C_{\Gamma_0}(q))$.

Now let $\mathcal{M} \subset \mathbb{R}^N$ be the critical point chosen in Theorem 1.1. The minimizer \mathcal{M} is bounded. Hence, we can slide the hyperplane H_λ from below until it touches the minimizer \mathcal{M} . Our result in Theorem 1.1 states that this touching does not occur in $H_0^- \cup H_0^+$ and thus, we assume by contradiction that there exist a constant $\lambda_0 < 0$ and a point $q \in \mathcal{M} \cap \Omega$ such that

$$T_q \mathcal{M} = H_{\lambda_0} \quad \text{and} \quad H_{\lambda_0}^- \cap \mathcal{M} = \emptyset$$

where $T_q \mathcal{M}$ is the tangent space of \mathcal{M} at q . Due to the symmetry of our setting, we can conduct the same argument that we will show below in the case when we slide the hyperplane from above and the touching occurs in H_0^+ . Hence, it is sufficient to show the proof in the case when the touching occurs in H_0^- . See also Figure 3 for the situation that we consider in dimension 2.

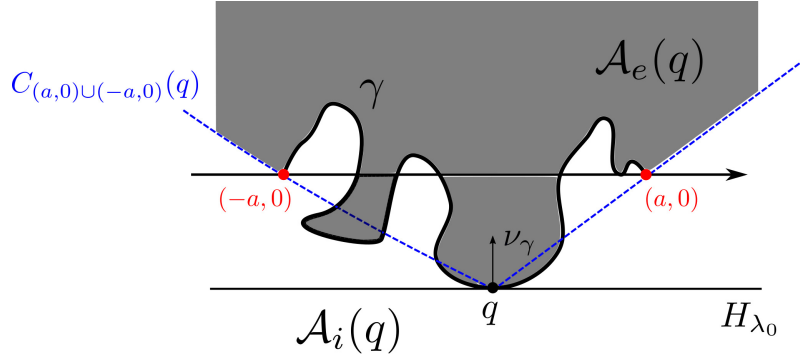


Figure 3: The situation in dimension 2 in which the critical point $\mathcal{M} = \gamma$ is a $C^{1,\alpha}$ curve with $\partial\gamma = \Gamma_0 = \{(a, 0), (-a, 0)\}$. The set $\mathcal{A}_e(q)$ is shown in dark gray, the set $\mathcal{A}_i(q)$ in white. The dashed lines represent the boundary of the cone $C_{(a,0) \cup (-a,0)}(q)$.

Since \mathcal{M} is an orientable compact critical point of Area_s , which means the vanishing of the first variation of Area_s at \mathcal{M} , and since $q \in \mathcal{M}$, we obtain, from (1.4), that

$$0 = H_{\mathcal{M},s}(q) := c_N \int_{\mathbb{R}^N} \frac{\chi_{\mathcal{A}_i(q)}(y) - \chi_{\mathcal{A}_e(q)}(y)}{|y - q|^{N+s}} dy \quad (2.2)$$

where the sets $\mathcal{A}_e(q)$ and $\mathcal{A}_i(q)$ are defined as in (1.5) and (1.6). We consider all the contributions from $\mathcal{A}_e(q)$ and $\mathcal{A}_i(q)$ in detail and show that the singular integral on the right-hand side of (2.2) is strictly positive, which is a contradiction.

Indeed, since $C_{\Gamma_0}(q) \subset H_{\lambda_0}^+$ and H_{λ_0} is tangential to \mathcal{M} , we have that $P_{\lambda_0}(\mathcal{A}_e(q)) \subset H_{\lambda_0}^- \subset \mathcal{A}_i(q)$. This implies that $\mathbb{R}^N = \mathcal{A}_e(q) \cup P_{\lambda_0}(\mathcal{A}_e(q)) \cup \mathcal{A}_i(q) \setminus P_{\lambda_0}(\mathcal{A}_e(q))$, up to negligible sets, and thus we can compute the fractional mean curvature $H_{\mathcal{M},s}$ at q as

follows:

$$\begin{aligned}
c_N^{-1} H_{\mathcal{M},s}(q) &= \int_{\mathcal{A}_e(q)} \frac{\chi_{\mathcal{A}_i(q)}(y) - \chi_{\mathcal{A}_e(q)}(y)}{|y - q|^{N+s}} dy + \int_{P_{\lambda_0}(\mathcal{A}_e(q))} \frac{\chi_{\mathcal{A}_i(q)}(y) - \chi_{\mathcal{A}_e(q)}(y)}{|y - q|^{N+s}} dy \\
&\quad + \int_{\mathcal{A}_i(q) \setminus P_{\lambda_0}(\mathcal{A}_e(q))} \frac{\chi_{\mathcal{A}_i(q)}(y) - \chi_{\mathcal{A}_e(q)}(y)}{|y - q|^{N+s}} dy \\
&= \int_{\mathcal{A}_e(q)} \frac{-1}{|y - q|^{N+s}} dy + \int_{P_{\lambda_0}(\mathcal{A}_e(q))} \frac{1}{|y - q|^{N+s}} dy \\
&\quad + \int_{\mathcal{A}_i(q) \setminus P_{\lambda_0}(\mathcal{A}_e(q))} \frac{1}{|y - q|^{N+s}} dy. \tag{2.3}
\end{aligned}$$

From the change of variables $y \mapsto P_{\lambda_0}(y)$ and the definition of P_{λ_0} , we have

$$\int_{P_{\lambda_0}(\mathcal{A}_e(q))} \frac{1}{|y - q|^{N+s}} dy = \int_{\mathcal{A}_e(q)} \frac{1}{|y - q|^{N+s}} dy. \tag{2.4}$$

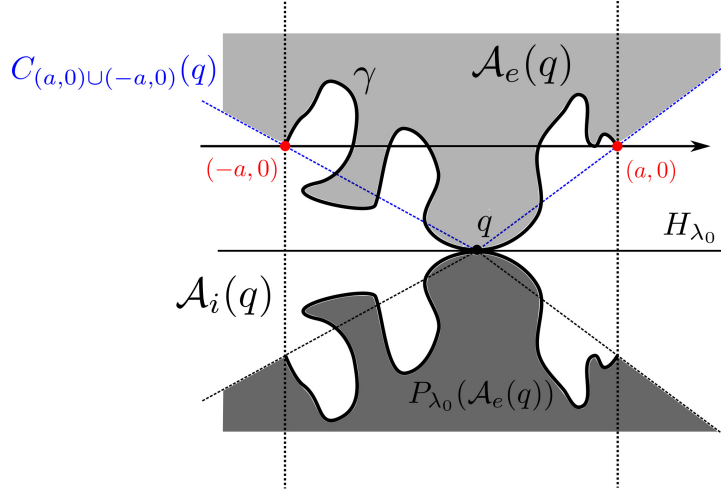


Figure 4: The same situation as in Figure 3. The reflection $P_{\lambda_0}(\mathcal{A}_e(q))$ of $\mathcal{A}_e(q)$ is shown in dark gray, the set $\mathcal{A}_e(q)$ in light gray.

Moreover, we have that the volume of the set $\mathcal{A}_i(q) \setminus P_{\lambda_0}(\mathcal{A}_e(q))$ is not zero because

$$\mathcal{A}_i(q) \setminus P_{\lambda_0}(\mathcal{A}_e(q)) \supset \{x \in \mathbb{R}^N \mid |x'| > |p'|/2\} \cap H_{\lambda_0}^+ \cap C_{\Gamma_0}(q)^c \supset B_{\frac{\lambda_0}{100}}^+(p),$$

where $p := (p', \lambda_0) \in \mathbb{R}^{N-1} \times \mathbb{R}$ for some $p' \in \mathbb{R}^{N-1}$ with $|p'| \gg |\lambda_0| + \max_{x,y \in \Gamma_0} |x - y|$ sufficiently large and $B_r^+(q) := \{x = (x', x_N) \in B_r(q) \mid x_N > q_N\}$. See also Figure 4 for illustration in dimension 2. From (2.3) and (2.4), we obtain

$$\begin{aligned}
0 &= \int_{\mathcal{A}_e(q)} \frac{-1}{|y - q|^{N+s}} dy + \int_{\mathcal{A}_e(q)} \frac{1}{|y - q|^{N+s}} dy + \int_{\mathcal{A}_i(q) \setminus P_{\lambda_0}(\mathcal{A}_e(q))} \frac{1}{|y - q|^{N+s}} dy \\
&= \int_{\mathcal{A}_i(q) \setminus P_{\lambda_0}(\mathcal{A}_e(q))} \frac{1}{|y - q|^{N+s}} dy > 0,
\end{aligned}$$

which is a contradiction. □

3 Shape of Critical Points with Two Disjoint Compact Boundaries

In this section, we first give the proof of Theorem 1.2 and then we show some properties of the critical points of Area_s and compute the fractional mean curvature of cones.

3.1 Proof of Theorem 1.2

In this subsection, we prove Theorem 1.2. The idea of the proof is basically the same as the one in the proof of Theorem 1.1. The convexity assumption on \mathcal{C} is necessary for us to use the sliding method.

Proof of Theorem 1.2. We first define

$$H_{\Gamma_i}^+ := \{(x', x_N) \mid x_N > h_i\}, \quad H_{\Gamma_i}^- := \{(x', x_N) \mid x_N < h_i\}$$

for each $i \in \{1, 2\}$. Notice that

$$\partial H_{\Gamma_i}^+ \cap \partial \mathcal{C} = \partial H_{\Gamma_i}^- \cap \partial \mathcal{C} = \Gamma_i \quad \text{and} \quad \partial H_{\Gamma_i}^+ \cap \mathcal{C} = \partial H_{\Gamma_i}^- \cap \mathcal{C} = C_i$$

for each $i \in \{1, 2\}$.

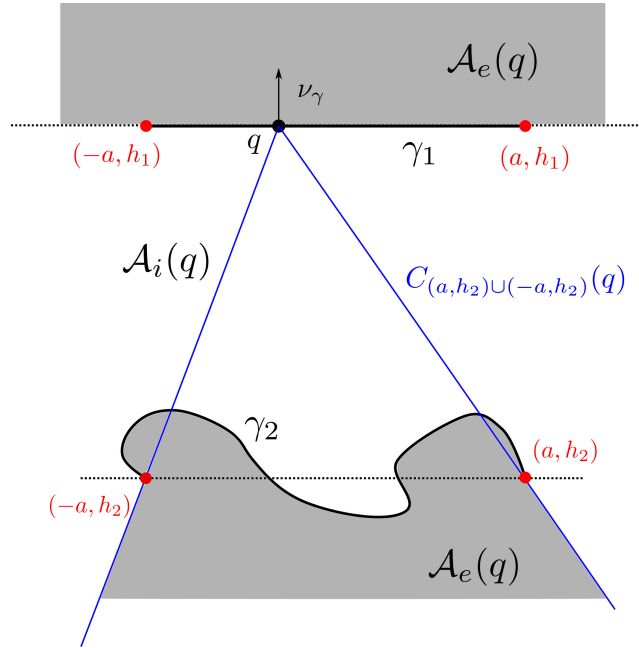


Figure 5: The situation in dimension 2 in which each component $\mathcal{M}_i = \gamma_i$ of the critical point $\mathcal{M} = \gamma$ for $i \in \{1, 2\}$ is a $C^{1,\alpha}$ curve with $\partial\gamma_i = \Gamma_i$ where $\Gamma_1 = \{(a, h_1), (-a, h_1)\}$ and $\Gamma_2 = \{(a, h_2), (-a, h_2)\}$. The set $\mathcal{A}_e(q)$ is shown in gray, the set $\mathcal{A}_i(q)$ in white.

Let $\mathcal{M} \subset \mathbb{R}^N$ be the critical point chosen in Theorem 1.2. By using the same argument as in the proof of Theorem 1.1, we obtain that \mathcal{M} cannot exist in the regions $H_{\Gamma_2}^-$ and $H_{\Gamma_1}^+$, that is, $\mathcal{M} \cap (H_{\Gamma_2}^- \cup H_{\Gamma_1}^+) = \emptyset$.

We now show that any connected component of \mathcal{M} cannot be either C_1 or C_2 . To see this, we assume by contradiction that there exists a connected component \mathcal{M}_1 of \mathcal{M}

such that \mathcal{M}_1 coincides with C_1 . Taking any $q \in \mathcal{M}_1$, we have that the cone $C_{\Gamma_2}(q)$ of vertex q whose boundary passes through Γ_2 is contained in $H_{\Gamma_1}^-$. By choosing a proper orientation of \mathcal{M} , we can have that $H_{\Gamma_1}^+ \subset \mathcal{A}_e(q)$ and $\mathcal{A}_i(q) \subset H_{\Gamma_1}^-$ where the sets $\mathcal{A}_e(q)$ and $\mathcal{A}_i(q)$ are defined as in (1.5) and (1.6), respectively. See Figure 5 for the situation in dimension 2.

Since \mathcal{M} is a critical point of Area_s , from (1.4), we have that

$$0 = H_{\mathcal{M},s}(q) = c_N \int_{\mathbb{R}^N} \frac{\chi_{\mathcal{A}_i(q)}(y) - \chi_{\mathcal{A}_e(q)}(y)}{|y - q|^{N+s}} dy. \quad (3.1)$$

Now, by employing the same argument we used in the proof of Theorem 1.1, we obtain that

$$\begin{aligned} c_N^{-1} H_{\mathcal{M},s}(q) &= \int_{\mathcal{A}_e(q) \cap H_{\Gamma_1}^+} \frac{-1}{|y - q|^{N+s}} dy + \int_{\mathcal{A}_e(q) \cap H_{\Gamma_1}^-} \frac{-1}{|y - q|^{N+s}} dy \\ &\quad + \int_{\mathcal{A}_i(q)} \frac{1}{|y - q|^{N+s}} dy \\ &\leq \int_{B_{1/2}(-\lambda e_N)} \frac{-1}{|y - q|^{N+s}} dy < 0 \end{aligned}$$

because $B_{1/2}(-\lambda e_N) \subset \mathcal{A}_e(q) \cap H_{\Gamma_2}^-$ where $\lambda > \max\{|x - z| \mid x \in C_2, z \in \mathcal{M}\} + 1$. This contradicts (3.1). Therefore, we conclude that the first claim is valid.

To prove the rest of the claim, we can argue in the same way as in the proof of the first claim. Indeed, we slide any hyperplane parallel to the x_N -axis from right to left or from left to right until it touches the boundary of the cylinder \mathcal{C} . If there is no touching point, from the convexity of \mathcal{C} , we obtain that the critical point \mathcal{M} is strictly contained in \mathcal{C} except for its boundary. Thus, we assume by contradiction that there exists a touching point q of \mathcal{M} in the complement of $\bar{\mathcal{C}}$. Then, by choosing a proper orientation of \mathcal{M} , we can show that the contribution from $\mathcal{A}_e(q)$ relative to the touching point q is strictly larger (or smaller) than that from $\mathcal{A}_i(q)$, respectively, as we see in the proof of the first claim. This contradicts that the fractional mean curvature vanishes at the touching point q . Therefore, we conclude the proof of Theorem 1.2. \square

3.2 Further Study on Critical Points and Cones

In this subsection, we further study the shape of critical points of Area_s under the same assumptions as in Theorem 1.2 with $h_1 = 1$ and $h_2 = -1$.

First, we investigate the shape of critical points in dimension 2. To begin, we divide \mathbb{R}^2 into four regions, that is, we define four regions C_0^t , C_0^b , C_0^r , and C_0^ℓ by

$$\begin{aligned} C_0^t &:= \{(x_1, x_2) \in \mathbb{R}^2 \mid x_2 > |x_1|\}, \\ C_0^b &:= \{(x_1, x_2) \in \mathbb{R}^2 \mid x_2 < -|x_1|\}, \\ C_0^r &:= \{(x_1, x_2) \in \mathbb{R}^2 \mid -|x_1| < x_2 < |x_1|, 0 < x_1\}, \\ \text{and } C_0^\ell &:= \{(x_1, x_2) \in \mathbb{R}^2 \mid -|x_1| < x_2 < |x_1|, x_1 < 0\}, \end{aligned}$$

respectively. Moreover, we set

$$C_0 := (\partial C_0^t \cup \partial C_0^b) \cap \{(x_1, x_2) \mid |x_2| \leq 1\}. \quad (3.2)$$

Notice that $\partial C_0 = \Gamma_1 \cup \Gamma_2$ where Γ_1 and Γ_2 are given in Theorem 1.2 with $h_1 = 1$ and $h_2 = -1$ in \mathbb{R}^2 . From the definition of Γ_1 and Γ_2 , we have that $\Gamma_1 = \{(\pm 1, 1)\}$ and $\Gamma_2 = \{(\pm 1, -1)\}$.

Now we prove that the fractional mean curvature of the cone C_0 vanishes at regular points, i.e.,

$$H_{C_0,s}(z) = 0 \quad (3.3)$$

for any $z \in C_0 \setminus \{0, (\pm 1, 1), (\pm 1, -1)\}$. Indeed, let $z \in C_0 \setminus \{0, (\pm 1, 1), (\pm 1, -1)\}$ and, by symmetry, we may assume that $z = (z_1, z_2)$ satisfies $-1 < z_1 < 0$ and $0 < z_2 < 1$. Then, from the definition of the ‘‘interior’’ $\mathcal{A}_i(z)$ and the ‘‘exterior’’ $\mathcal{A}_e(z)$ of the cone C_0 and by taking a suitable orientation of $C_0 \setminus \{0\}$, we may obtain that

$$\mathcal{A}_i(z) = \left(([z, (1, 1)]^- \cap [z, (1, -1)]^+) \setminus \overline{C_0^t} \right) \cup \left(([z, (-1, 1)]^- \cap [z, (-1, -1)]^+) \cup C_0^\ell \right)$$

and

$$\mathcal{A}_e(z) = \left(([z, (-1, 1)]^+ \cap [z, (1, 1)]^+) \cup C_0^t \right) \cup \left(([z, (-1, -1)]^- \cap [z, (1, -1)]^-) \setminus \overline{C_0^\ell} \right)$$

where we denote by $[p, q]$ the straight line passing through $p, q \in \mathbb{R}^2$ with $p \neq q$ and we define $[p, q]^+$ and $[p, q]^-$ by the upper part and the lower part of the region separated by the straight line $[p, q]$, respectively.

Now, because of the symmetry of the cone C_0 , we readily observe that, in dimension 2, the sets $\mathcal{A}_i(z)$ and $\mathcal{A}_e(z)$ are equivalent to each other in the sense that $T(\mathcal{A}_i(z)) = \mathcal{A}_e(z)$ where $T : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ is an isometric map such that $\frac{x+T(x)}{2} \in \{(x_1, x_2) \mid x_2 = x_1\}$ for any $x \in \mathbb{R}^2$. By definition, we notice that $T(z) = z$.

Therefore, from the change of variables $x \mapsto T(x)$, we obtain that

$$\begin{aligned} c_N^{-1} H_{C_0,s}(z) &= \int_{\mathcal{A}_i(z)} \frac{1}{|y - z|^{2+s}} dy - \int_{\mathcal{A}_e(z)} \frac{1}{|y - z|^{2+s}} dy \\ &= \int_{\mathcal{A}_i(z)} \frac{1}{|y - z|^{2+s}} dy - \int_{\mathcal{A}_i(z)} \frac{1}{|T(y) - T(z)|^{2+s}} dy \\ &= 0. \end{aligned}$$

By combining this fact with Theorem 1.2, we can prove the following proposition.

Proposition 3.1. *Let $N = 2$ and $s \in (0, 1)$. Let Γ_1 and Γ_2 be as in Theorem 1.2 with $h_1 = 1$ and $h_2 = -1$. Let $\gamma \subset \mathbb{R}^2$ be an orientable compact $C^{1,\alpha}$ curve with $\partial\gamma = \Gamma_1 \cup \Gamma_2$. Assume that $\mathcal{C} = \{(x_1, x_2) \mid |x_1| < 1\}$ where \mathcal{C} is as in (1.7). If γ is a critical point of Area_s under normal variations, then γ is not contained in either $\overline{C_0^t} \cup \overline{C_0^b}$ or $\overline{C_0^r} \cup \overline{C_0^\ell}$ whenever $(\gamma \setminus \partial\gamma) \cap (C_0 \setminus \{0\}) \neq \emptyset$.*

Remark 3.2. We may observe, by combining Proposition 3.1 with Theorem 1.2, that the possible shapes of minimizers of Area_s in dimension 2 whose boundary is $\Gamma_1 \cup \Gamma_2$ are depicted in Figure 6.

Proof. Let $\gamma \subset \mathbb{R}^2$ be as in Proposition 3.1 and assume that $(\gamma \setminus \partial\gamma) \cap (C_0 \setminus \{0\}) \neq \emptyset$. We argue by contradiction that either $\gamma \subset \overline{C_0^t} \cup \overline{C_0^b}$ or $\gamma \subset \overline{C_0^r} \cup \overline{C_0^\ell}$ holds. Due to the symmetry of C_0 , it is sufficient to consider the case that $\gamma \subset \overline{C_0^t} \cup \overline{C_0^b}$ holds. From this assumption, we can choose a point $z \in (\gamma \setminus \partial\gamma) \cap (C_0 \setminus \{0\})$.

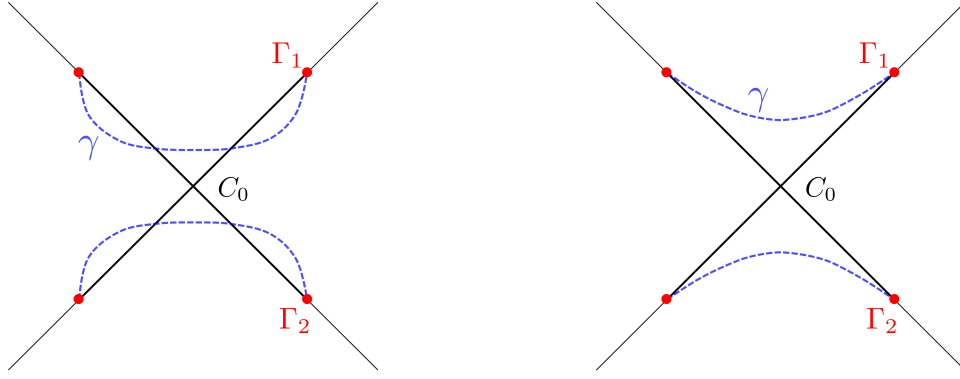


Figure 6: Possible minimizers γ of Area_s in dimension 2 with $\partial\gamma = \Gamma_1 \cup \Gamma_2$ is shown with dashed lines. On the right, γ does not intersect with C_0 except at their boundaries Γ_1 and Γ_2 .

Now, by choosing a proper orientation, we consider the “interior” and “exterior” of γ and C_0 at the touching point z . To see this, we set the interior and exterior at $q \in \eta$ of a curve $\eta \subset \mathbb{R}^2$ as $\mathcal{A}_i^\eta(z)$ and $\mathcal{A}_e^\eta(z)$, respectively. Then, from the smoothness of the critical point γ and the assumption that $\gamma \subset C_0^t \cup C_0^b$, we obtain, by taking a suitable orientation of γ and C_0 , that

$$|\mathcal{A}_e^{C_0}(z) \setminus \mathcal{A}_e^\gamma(z)| = |\mathcal{A}_i^\gamma(z) \setminus \mathcal{A}_i^{C_0}(z)| \neq 0 \quad (3.4)$$

and

$$|\mathcal{A}_e^\gamma(z) \setminus \mathcal{A}_e^{C_0}(z)| = |\mathcal{A}_i^{C_0}(z) \setminus \mathcal{A}_i^\gamma(z)| = 0. \quad (3.5)$$

Here, from Theorem 1.2, we have used the fact that all the critical points of Area_s in our situation are contained in the box $\{(x_1, x_2) \mid |x_1| < 1, |x_2| < 1\}$. See also Figure 7 for our situation.

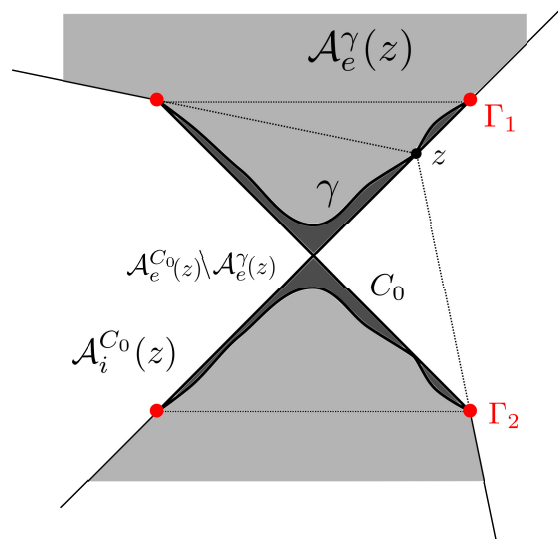


Figure 7: The situation of the critical point γ and the touching point z in which γ is included in $C_0^t \cup C_0^b$ with $\partial\gamma = \Gamma_1 \cup \Gamma_2$. The set $\mathcal{A}_e^\gamma(z)$ is shown in light gray, the set $\mathcal{A}_i^{C_0}(z)$ in white, and the set $\mathcal{A}_e^{C_0}(z) \setminus \mathcal{A}_e^\gamma(z)$ in dark gray.

Hence, since γ is a critical point of Area_s , we have that

$$H_{\gamma,s}(z) = 0.$$

From (3.3), (3.4), and (3.5), we have

$$\begin{aligned} 0 &= c_2^{-1} H_{\gamma,s}(z) = c_2^{-1} (H_{\gamma,s}(z) - H_{C_{0,s}}(z)) \\ &= \int_{\mathbb{R}^2} \frac{\chi_{\mathcal{A}_i^\gamma(z)}(y) - \chi_{\mathcal{A}_i^{C_0}(z)}(y) + \chi_{\mathcal{A}_e^{C_0}(z)}(y) - \chi_{\mathcal{A}_e^\gamma(z)}(y)}{|y-z|^{2+s}} dy \\ &= \int_{\mathcal{A}_i^\gamma(z) \setminus \mathcal{A}_i^{C_0}(z)} \frac{1}{|y-z|^{2+s}} dy - \int_{\mathcal{A}_i^{C_0}(z) \setminus \mathcal{A}_i^\gamma(z)} \frac{1}{|y-z|^{2+s}} dy \\ &\quad + \int_{\mathcal{A}_e^{C_0}(z) \setminus \mathcal{A}_e^\gamma(z)} \frac{1}{|y-z|^{2+s}} dy - \int_{\mathcal{A}_e^\gamma(z) \setminus \mathcal{A}_e^{C_0}(z)} \frac{1}{|y-z|^{2+s}} dy \\ &= \int_{\mathcal{A}_i^\gamma(z) \setminus \mathcal{A}_i^{C_0}(z)} \frac{1}{|y-z|^{2+s}} dy + \int_{\mathcal{A}_e^{C_0}(z) \setminus \mathcal{A}_e^\gamma(z)} \frac{1}{|y-z|^{2+s}} dy > 0, \end{aligned} \quad (3.6)$$

which is a contradiction. Therefore we obtain the claim. \square

Remark 3.3. We briefly consider the situation of Theorem 1.2 with $h_1 = d$ and $h_2 = -d$ for $d \neq 1$ and $d > 0$ and see what kind of shape the critical points in dimension 2 look like. Notice that we have treated the case of $d = 1$ in Proposition 3.1.

Assume that $h_1 = d$ and $h_2 = -d$ for $d > 0$. We define a cone C_d with vertex 0 by

$$C_d := \{(x_1, x_2) \in \mathbb{R}^2 \mid |x_2| = d|x_1|, |x_2| \leq d\}. \quad (3.7)$$

Notice that $\partial C_d = \Gamma_1 \cup \Gamma_2$. By slightly modifying the argument for showing that $H_{C_{0,s}} = 0$ on $C_0 \setminus (\partial C_0 \cup \{0\})$ and taking a proper orientation, we can show that the fractional mean curvature $H_{C_d,s}(z)$ of C_d is either positive or negative for any $z \in C_d \setminus \partial C_d$ with $z \neq 0$. Then, again by slightly modifying the argument in the proof of Proposition 3.1, we obtain the same result as in Proposition 3.1 even for any $d \neq 1$.

We next prove the same result as Proposition 3.1 in higher dimensions. To see this, we also show that the fractional mean curvature of a cone passing through $\Gamma_1 \cup \Gamma_2$ is either positive or negative everywhere except at its vertex in higher dimensions. The idea of the proof is the same as that in the proof of Proposition 3.1. We first introduce some notations. We define a bounded tube D_0 and a unbounded (open) cone \tilde{C}_0 by

$$\begin{aligned} D_0 &:= \{(x', x_N) \in \mathbb{R}^{N-1} \times \mathbb{R} \mid |x'| < 1, -1 < x_N < 1\} \\ \tilde{C}_0 &:= \{(x', x_N) \in \mathbb{R}^{N-1} \times \mathbb{R} \mid |x_N| > |x'|\}. \end{aligned}$$

Moreover, we set $C_0^N := \partial \tilde{C}_0 \cap \{(x', x_N) \mid |x_N| \leq 1\}$ and decompose \tilde{C}_0 into two parts \tilde{C}_0^+ and \tilde{C}_0^- which are defined by

$$\begin{aligned} \tilde{C}_0^+ &:= \{(x', x_N) \in \mathbb{R}^{N-1} \times \mathbb{R} \mid x_N > |x'|\} \\ \tilde{C}_0^- &:= \{(x', x_N) \in \mathbb{R}^{N-1} \times \mathbb{R} \mid x_N < -|x'|\}. \end{aligned}$$

Notice that C_0^N coincides with C_0 given in (3.2) if $N = 2$ and $\partial C_0^N = \Gamma_1 \cup \Gamma_2$.

Proposition 3.4. *Let $N \geq 3$ and $s \in (0, 1)$. Let Γ_1 and Γ_2 be as in Theorem 1.2 with $h_1 = 1$ and $h_2 = -1$. Let $\mathcal{M} \subset \mathbb{R}^N$ be an orientable compact $C^{1,\alpha}$ manifold with $\partial\mathcal{M} = \Gamma_1 \cup \Gamma_2$. Assume that $\mathcal{C} = \{(x', x_N) \mid |x'| < 1\}$ where \mathcal{C} is as in (1.7). If \mathcal{M} is a critical point of Area_s under normal variations, then \mathcal{M} is not contained in either $\widetilde{C}_0^+ \cup \widetilde{C}_0^-$ or $D_0 \setminus (\widetilde{C}_0^+ \cup \widetilde{C}_0^-)$ whenever $(\mathcal{M} \setminus \partial\mathcal{M}) \cap (C_0^N \setminus \{0\}) \neq \emptyset$.*

Proof. The proof is similar to that of Proposition 3.1 and we here show a rough sketch of the proof. Let \mathcal{M} be the critical point selected in Proposition 3.4. We assume that $(\mathcal{M} \setminus \partial\mathcal{M}) \cap (C_0^N \setminus \{0\})$ is not empty and we choose a point $z \in (\mathcal{M} \setminus \partial\mathcal{M}) \cap (C_0^N \setminus \{0\})$. Suppose by contradiction that either

$$\mathcal{M} \subset \overline{\widetilde{C}_0^+ \cup \widetilde{C}_0^-} \quad \text{or} \quad \mathcal{M} \subset D_0 \setminus (\overline{\widetilde{C}_0^+ \cup \widetilde{C}_0^-})$$

holds. First, by choosing an orientation, we show that

$$H_{C_0^N, s}(z) > 0. \quad (3.8)$$

Indeed, if we take the unit normal vector $\nu_{C_0^N}(z)$ of the cone C_0^N at z in such a way that the direction is towards \widetilde{C}_0 , then the ‘‘interior’’ $\mathcal{A}_i^{C_0^N}(z)$ and ‘‘exterior’’ $\mathcal{A}_e^{C_0^N}(z)$ are given by

$$\mathcal{A}_i^{C_0^N}(z) = \mathbb{R}^N \setminus \left(\mathcal{A}_e^{C_0^N}(z) \cup C_0^N \right)$$

and

$$\mathcal{A}_e^{C_0^N}(z) = (\widetilde{C}_0 \cap \{(x', x_N) \mid |x_N| \leq 1\}) \cup ((C_{\Gamma_1}(z) \cup C_{\Gamma_2}(z)) \cap \{(x', x_N) \mid |x_N| \geq 1\})$$

where $C_{\Gamma_i}(z)$ is defined by a (filled) cone of vertex z passing through Γ_i for each $i \in \{1, 2\}$. Now we take a hyperplane H_z which is tangent to $\partial\widetilde{C}_0$ and passes through z and define the reflection map T_{H_z} with respect to H_z . From the definitions of C_0^N , $\mathcal{A}_i^{C_0^N}(z)$, and $\mathcal{A}_e^{C_0^N}(z)$, we have

$$T_{H_z}(\mathcal{A}_e^{C_0^N}(z)) \subset \mathcal{A}_i^{C_0^N}(z) \quad \text{and} \quad \left| \mathcal{A}_i^{C_0^N}(z) \setminus T_{H_z}(\mathcal{A}_e^{C_0^N}(z)) \right| \neq 0.$$

Since T_{H_z} is an isometry and $T_{H_z}(z) = z$, we obtain the following:

$$\begin{aligned} c_N^{-1} H_{C_0^N, s}(z) &= \int_{\mathcal{A}_i^{C_0^N}(z) \setminus T_{H_z}(\mathcal{A}_e^{C_0^N}(z))} \frac{dx}{|x - z|^{N+s}} + \int_{T_{H_z}(\mathcal{A}_e^{C_0^N}(z))} \frac{dx}{|x - z|^{N+s}} \\ &\quad - \int_{\mathcal{A}_e^{C_0^N}(z)} \frac{dx}{|x - z|^{N+s}} \\ &= \int_{\mathcal{A}_i^{C_0^N}(z) \setminus T_{H_z}(\mathcal{A}_e^{C_0^N}(z))} \frac{dx}{|x - z|^{N+s}} + 0 > 0, \end{aligned} \quad (3.9)$$

which implies (3.8).

Now, since \mathcal{M} is a critical point of Area_s , we have the Euler-Lagrange equation

$$H_{\mathcal{M}, s}(z) = 0.$$

Thus, taking the unit normal vector $\nu_{\mathcal{M}}(z)$ of \mathcal{M} at z as $\nu_{C_0^N}(z)$, we can have the following computation:

$$\begin{aligned} 0 &= c_N^{-1} (H_{\mathcal{M}, s}(z) - H_{C_0^N, s}(z) + H_{C_0^N, s}(z)) \\ &= 2 \int_{\mathcal{A}_e^{C_0^N}(z) \setminus \mathcal{A}_e^{\mathcal{M}}(z)} \frac{1}{|x - z|^{N+s}} dx - 2 \int_{\mathcal{A}_e^{\mathcal{M}}(z) \setminus \mathcal{A}_e^{C_0^N}(z)} \frac{1}{|x - z|^{N+s}} dx + H_{C_0^N, s}(z). \end{aligned} \quad (3.10)$$

From the assumption, we can observe that

$$\left| \mathcal{A}_e^{C_0^N}(z) \setminus \mathcal{A}_e^M(z) \right| > 0 \quad \text{and} \quad \left| \mathcal{A}_e^M(z) \setminus \mathcal{A}_e^{C_0^N}(z) \right| = 0.$$

Therefore, from (3.8) and (3.10), we reach a contradiction. \square

4 Topology of Critical Points

In this section, we investigate the topology of critical points with two parallel and co-axial boundaries and prove Theorems 1.4 and 1.5.

Before proving our main theorems of this section, we establish Lemma 1.3. The idea of the proof is to construct a small barrier, whose fractional mean curvature is strictly positive or negative, and to “slide” the barrier until it touches the critical point. The construction of the barrier is inspired by the one shown in [18]. See also [14, Proof of Proposition 4.1]. In the sequel, without loss of generality, we may assume that $\mathcal{C} = \{(x', x_N) \mid |x'| < 1\}$ where \mathcal{C} is as in (1.7) for simplicity.

Proof of Lemma 1.3. We first fix $\varepsilon \in (0, 1)$ so small that $\delta = \delta(\varepsilon) := (-\log \varepsilon)^{-1/2} < \frac{1}{2}$ and we define a smooth bump function $w_\varepsilon : \mathbb{R}^{N-1} \rightarrow \mathbb{R}$ by

$$w_\varepsilon(x') := \begin{cases} -\exp\left(-\frac{1}{\delta^2 - |x'|^2}\right) & \text{for } |x'| < \delta \\ 0 & \text{otherwise.} \end{cases}$$

Notice that $w_\varepsilon \in C^\infty(\mathbb{R}^{N-1})$, $w_\varepsilon(x') = 0$ for $|x'| = \delta$, $w_\varepsilon(0) = -\varepsilon$, and

$$\lim_{\varepsilon \downarrow 0} \phi(\varepsilon) := \lim_{\varepsilon \downarrow 0} \|\nabla'^2 w_\varepsilon\|_{C^0} = 0. \quad (4.1)$$

where $\nabla' = (\partial_{x_1}, \dots, \partial_{x_{N-1}})$. If necessary, we may choose ε in such a way that $\phi(\varepsilon) < 1$. Note that, since ϕ is an increasing function in a neighborhood $I_\phi \subset [0, 1)$ of the origin, its inverse function ϕ^{-1} exists in a neighborhood $J_\phi \subset [0, 1)$ of the origin. We then set

$$r(\varepsilon) := (2(N-1)\phi(\varepsilon))^{-1} \quad \text{and} \quad d(\varepsilon) := 2r(\varepsilon). \quad (4.2)$$

Moreover, we define a positive constant ε_d as

$$\varepsilon_d := \begin{cases} \phi^{-1}((2(N-1)d)^{-1}) & \text{if } (2(N-1)d)^{-1} \in J_\phi \\ (\text{any positive constant in } J_\phi) & \text{if } (2(N-1)d)^{-1} \notin J_\phi. \end{cases}$$

By definition, we observe that $r(\varepsilon_d) \geq d$ and ε_d can be chosen independently of d if $d < (2(N-1))^{-1}$ since $J_\phi \subset [0, 1)$.

In addition, we choose a smooth function $v_\varepsilon : \mathbb{R}^{N-1} \rightarrow \mathbb{R}$ such that v_ε is radially symmetric, $0 \leq v_\varepsilon(x') \leq 1$ for $x' \in \mathbb{R}^{N-1}$, and $\text{spt } v_\varepsilon \subset B'_{1/8}(0)$ where we denote by $B'_r(0)$ an open ball centered at the origin of radius r in \mathbb{R}^{N-1} . In particular, we choose v_ε in such a way that its subgraph $\{(x', x_N) \mid 0 \leq x_N \leq v_\varepsilon(x')\}$ of v_ε contains a cylinder of height $\phi(\varepsilon)^\beta < 1$ for $\beta \in (0, s)$ with the base of radius $\frac{1}{16}$. Then we define a function $\tilde{w}_\varepsilon : \mathbb{R}^{N-1} \rightarrow \mathbb{R}$ by

$$\tilde{w}_\varepsilon(x') := \begin{cases} w_\varepsilon(x') & \text{for } |x'| < \delta \\ 0 & \text{for } \delta \leq |x'| < \frac{5}{8} \\ v_\varepsilon(x' - b') & \text{for } \frac{5}{8} \leq |x'| < \frac{7}{8} \\ 0 & \text{for } |x'| \geq \frac{7}{8} \end{cases}$$

where $b' \in \mathbb{R}^{N-1}$ is any point with $|b'| = \frac{3}{4}$. Notice that \tilde{w}_ε is smooth in \mathbb{R}^{N-1} .

Now we construct a barrier against $\widetilde{\mathcal{M}}^{\varepsilon,t}$, i.e., an orientable compact $(N-1)$ -dimensional piecewise smooth manifold $\widetilde{\mathcal{M}}^{\varepsilon,t}$ in the following way: first, taking any $t \in (0, \varepsilon]$, we define two sets

$$\begin{aligned} \mathcal{M}_1^{\varepsilon,t} &:= \{(x', x_N) \mid |x'| \leq 1, x_N = \tilde{w}_\varepsilon(x') + t\}, \\ \text{and } \mathcal{M}_2^{\varepsilon,t} &:= \{(x', x_N) \mid |x'| \leq 1, x_N = -d(\varepsilon) + t\} \end{aligned}$$

where $d(\varepsilon)$ is as in (4.2). Then we define our barrier as $\widetilde{\mathcal{M}}^{\varepsilon,t} := \mathcal{M}_1^{\varepsilon,t} \cup \mathcal{M}_2^{\varepsilon,t}$. By construction, we can easily see that $\widetilde{\mathcal{M}}^{\varepsilon,t}$ is an orientable compact $(N-1)$ -dimensional smooth manifold with $\partial\mathcal{M}_1^{\varepsilon,t} = \Gamma_1^{\varepsilon,t}$ and $\partial\mathcal{M}_2^{\varepsilon,t} = \Gamma_2^{\varepsilon,t}$ where we define

$$\Gamma_1^{\varepsilon,t} := \mathcal{C} \cap \{x_N = t\} \quad \text{and} \quad \Gamma_2^{\varepsilon,t} := \mathcal{C} \cap \{x_N = -d(\varepsilon) + t\}.$$

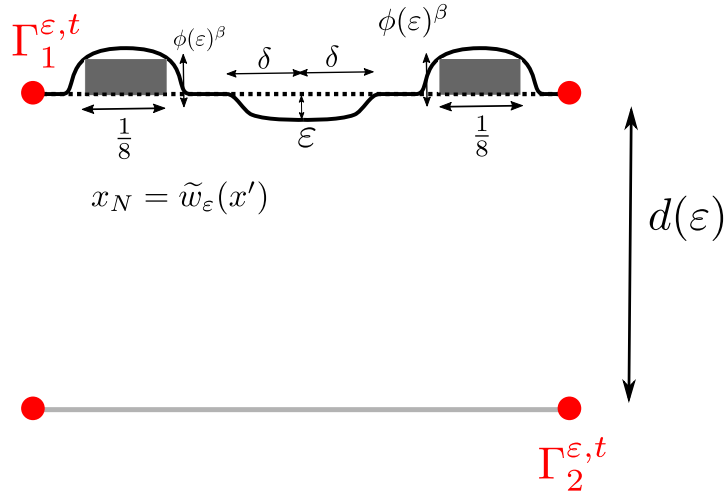


Figure 8: The barrier $\widetilde{\mathcal{M}}^{\varepsilon,t} = \mathcal{M}_1^{\varepsilon,t} \cup \mathcal{M}_2^{\varepsilon,t}$ associated with a function \tilde{w}_ε in dimension 2. The graph of \tilde{w}_ε in $\{|x'| < 1\}$ is depicted with black lines and the cylinders in dark gray.

We next construct another barrier in which the small bump associated with v_ε is removed from $\widetilde{\mathcal{M}}^{\varepsilon,t}$. First, for any $t \in (0, \varepsilon]$, we define a manifold $\mathcal{M}_3^{\varepsilon,t}$ as the graph of w_ε , i.e.,

$$\mathcal{M}_3^{\varepsilon,t} := \{(x', x_N) \mid |x'| < 1, x_N = w_\varepsilon(x') + t\}$$

and, then, define the second barrier as $\mathcal{M}^{\varepsilon,t} := \mathcal{M}_3^{\varepsilon,t} \cup \mathcal{M}_2^{\varepsilon,t}$. Notice that $\partial\mathcal{M}^{\varepsilon,t} = \Gamma_1^{\varepsilon,t} \cup \Gamma_2^{\varepsilon,t}$.

We now show, up to orientation, that the fractional mean curvature of $\widetilde{\mathcal{M}}^{\varepsilon,t}$ is negative on the graph of w_ε . Let $q \in \mathcal{M}_1^{\varepsilon,t}$ be any point such that $|q'| < \delta(\varepsilon)$ where we set $q = (q', q_N)$. We now define $C_{\Gamma_i^{\varepsilon,t}}(q)$ to be a (filled) cone of vertex q whose boundary passes through $\Gamma_i^{\varepsilon,t}$ for $i \in \{1, 2\}$. Then, up to orientation, the interior and exterior of $\widetilde{\mathcal{M}}^{\varepsilon,t}$ at q are

$$\mathcal{A}_i^{\widetilde{\mathcal{M}}^{\varepsilon,t}}(q) = \mathbb{R}^N \setminus \left(\mathcal{A}_e^{\widetilde{\mathcal{M}}^{\varepsilon,t}}(q) \cup \widetilde{\mathcal{M}}^{\varepsilon,t} \right)$$

and

$$\mathcal{A}_e^{\widetilde{\mathcal{M}}^{\varepsilon,t}}(q) = \left(C_{\Gamma_2^{\varepsilon,t}}(q) \cap \{(x', x_N) \mid x_N < -d(\varepsilon) + t\} \right) \cup \{(x', x_N) \mid x_N > \tilde{w}_\varepsilon^q(x')\},$$

respectively, where we define a function $\tilde{w}_\varepsilon^q : \mathbb{R}^{N-1} \rightarrow \mathbb{R}$ by

$$\tilde{w}_\varepsilon^q(x') := \begin{cases} \tilde{w}_\varepsilon(x') & \text{for } |x'| < 1 \\ \text{(the graph function of } \partial C_{\Gamma_1^{\varepsilon,t}}(q)) & \text{for } |x'| \geq 1. \end{cases}$$

We now compute the fractional mean curvature $H_{\widetilde{\mathcal{M}}^{\varepsilon,t}}(q)$ at q of $\widetilde{\mathcal{M}}^{\varepsilon,t}$. From the definition of the fractional mean curvature and by a change of variables, we have

$$\begin{aligned} -c_N^{-1} H_{\widetilde{\mathcal{M}}^{\varepsilon,t},s}(q) &= \int_{\mathbb{R}^N} \frac{\chi_{\mathcal{A}_e^{\widetilde{\mathcal{M}}^{\varepsilon,t}}(q)}(q-x) - \chi_{\mathcal{A}_i^{\widetilde{\mathcal{M}}^{\varepsilon,t}}(q)}(q-x)}{|x|^{N+s}} dx \\ &= \int_{B'_r(0) \times (-r,r)} \frac{\chi_{\mathcal{A}_e^{\widetilde{\mathcal{M}}^{\varepsilon,t}}(q)}(q-x) - \chi_{\mathcal{A}_i^{\widetilde{\mathcal{M}}^{\varepsilon,t}}(q)}(q-x)}{|x|^{N+s}} dx \\ &\quad + \int_{(B'_r(0) \times (-r,r))^c} \frac{\chi_{\mathcal{A}_e^{\widetilde{\mathcal{M}}^{\varepsilon,t}}(q)}(q-x) - \chi_{\mathcal{A}_i^{\widetilde{\mathcal{M}}^{\varepsilon,t}}(q)}(q-x)}{|x|^{N+s}} dx \\ &=: (I) + (II) \end{aligned} \tag{4.3}$$

where we set $r := r(\varepsilon)$ where $r(\varepsilon)$ is as in (4.2).

We first compute (I). Thanks to the choice of r and the construction of $\widetilde{\mathcal{M}}^{\varepsilon,t}$, we observe that

$$(B'_r(0) \times (-r,r)) \cap \left(C_{\Gamma_2^{\varepsilon,t}}(q) \cap \{(x', x_N) \mid x_N < -d(\varepsilon) + t\} \right) = \emptyset.$$

Thus we can represent the set $\partial \mathcal{A}_e^{\widetilde{\mathcal{M}}^{\varepsilon,t}}(q)$ in $B'_r(0) \times (-r,r)$ as the graph of \tilde{w}_ε^q . By doing a computation similar to the one in [3, Section 3], we obtain

$$\begin{aligned} (I) &= -2 \int_{B'_r(0)} F \left(\frac{\tilde{w}_\varepsilon^q(q') - \tilde{w}_\varepsilon^q(q' - x')}{|x'|} \right) \frac{dx'}{|x'|^{N-1+s}} \\ &= - \int_{B'_r(0)} F \left(\frac{\tilde{w}_\varepsilon^q(q') - \tilde{w}_\varepsilon^q(q' - x')}{|x'|} \right) \frac{dx'}{|x'|^{N-1+s}} \\ &\quad - \int_{B'_r(0)} F \left(\frac{\tilde{w}_\varepsilon^q(q') - \tilde{w}_\varepsilon^q(q' + x')}{|x'|} \right) \frac{dx'}{|x'|^{N-1+s}} \\ &= \int_{B'_r(0)} F \left(\frac{-\tilde{w}_\varepsilon^q(q') + \tilde{w}_\varepsilon^q(q' - x')}{|x'|} \right) \frac{dx'}{|x'|^{N-1+s}} \\ &\quad - \int_{B'_r(0)} F \left(\frac{\tilde{w}_\varepsilon^q(q') - \tilde{w}_\varepsilon^q(q' + x')}{|x'|} \right) \frac{dx'}{|x'|^{N-1+s}} \end{aligned} \tag{4.4}$$

where we set

$$F(t) := \int_0^t \frac{1}{(1 + \sigma^2)^{\frac{N+s}{2}}} d\sigma$$

for any $t \in \mathbb{R}$. Note that we have used the change of variables $x' \mapsto -x'$ in the second equality of (4.4) and the fact that F is odd in the last equality of (4.4). By definition, we have that $\tilde{w}_\varepsilon^q(q') = w_\varepsilon(q')$ and $\tilde{w}_\varepsilon^q \geq w_\varepsilon$ in \mathbb{R}^{N-1} . Since F is increasing, we derive from (4.4) that

$$\begin{aligned} (I) &\geq \int_{B'_r(0)} F \left(\frac{-w_\varepsilon(q') + w_\varepsilon(q' - x')}{|x'|} \right) \frac{dx'}{|x'|^{N-1+s}} \\ &\quad - \int_{B'_r(0)} F \left(\frac{w_\varepsilon(q') - w_\varepsilon(q' + x')}{|x'|} \right) \frac{dx'}{|x'|^{N-1+s}}. \end{aligned} \tag{4.5}$$

Now, by using the fundamental theorem of calculus in (4.5), we obtain

$$(I) \geq - \int_{B'_r(0)} \int_0^1 F'(a(x', q', \lambda)) d\lambda \frac{2w_\varepsilon(q') - w_\varepsilon(q' + x') - w_\varepsilon(q' - x')}{|x'|^{N+s}} dx' \quad (4.6)$$

where we set $a(x', q', \lambda)$ as

$$a(x', q', \lambda) := \lambda \frac{w_\varepsilon(q') - w_\varepsilon(q' + x')}{|x'|} + (1 - \lambda) \frac{-w_\varepsilon(q') + w_\varepsilon(q' - x')}{|x'|}$$

for $x', q' \in \mathbb{R}^{N-1}$ and $\lambda \in [0, 1]$. By using again the fundamental theorem of calculus, we have

$$\begin{aligned} & |2w_\varepsilon(q') - w_\varepsilon(q' + x') - w_\varepsilon(q' - x')| \\ &= \left| - \int_0^1 \nabla w_\varepsilon(q' + \rho x') \cdot x' d\rho + \int_0^1 \nabla w_\varepsilon(q' - \rho x') \cdot x' d\rho \right| \\ &\leq \int_0^1 |\nabla w_\varepsilon(q' + \rho x') - \nabla w_\varepsilon(q' - \rho x')| |x'| d\rho. \end{aligned} \quad (4.7)$$

Hence, combining (4.7) with (4.6), we obtain that

$$(I) \geq - \int_{B'_r(0)} \int_0^1 \frac{|\nabla' w_\varepsilon(q' + \rho x') - \nabla' w_\varepsilon(q' - \rho x')|}{|x'|^{N-1+s}} d\rho dx'.$$

Here we have used that $F'(t) = (1 + t^2)^{-\frac{N+s}{2}} \leq 1$. Since w_ε is smooth in \mathbb{R}^{N-1} , we then have

$$(I) \geq -2 \|\nabla'^2 w_\varepsilon\|_{C^0} \int_{B'_r(0)} \frac{dx'}{|x'|^{N-2+s}} = -\frac{2\omega_{N-2}}{1-s} \|\nabla'^2 w_\varepsilon\|_{C^0} r^{1-s}. \quad (4.8)$$

Now we compute (II) in the following way: since $B_r(0) \subset B'_r(0) \times (-r, r) \subset \mathbb{R}^N$, we have

$$(II) \geq - \int_{B'_r(0)} \frac{dx}{|x|^{N+s}} = -\frac{\omega_{N-1}}{s} r^{-s}. \quad (4.9)$$

Therefore, from (4.8) and (4.9), we obtain

$$-H_{\widetilde{\mathcal{M}}^{\varepsilon,t,s}}(q) \geq -(c_1 \|\nabla'^2 w_\varepsilon\|_{C^0} r^{1-s} + c_2 r^{-s}) \quad (4.10)$$

where c_1 and c_2 are defined as

$$c_1 := \frac{2\omega_{N-2}}{1-s} \quad \text{and} \quad c_2 := \frac{\omega_{N-1}}{s},$$

respectively. From (4.2), it holds that the right-hand side of (4.10) takes the maximum at $r = r(\varepsilon) \in (0, d(\varepsilon))$. Hence we finally obtain, from (4.10), that

$$-H_{\widetilde{\mathcal{M}}^{\varepsilon,t,s}}(q) \geq -c \|\nabla'^2 w_\varepsilon\|_{C^0}^s = -c \phi(\varepsilon)^s \quad (4.11)$$

where we set the constant $c = c(N, s) > 0$ as

$$c = c(N, s) := \frac{(2(N-1))^s \omega_{N-1}}{s(1-s)}.$$

Next we compute the fractional mean curvature $H_{\mathcal{M}^{\varepsilon,t},s}(q)$ at q by using Estimate (4.11) of the fractional mean curvature $H_{\widetilde{\mathcal{M}}^{\varepsilon,t},s}(q)$ at q . Indeed, from the construction of $\mathcal{M}^{\varepsilon,t}$ and $\widetilde{\mathcal{M}}^{\varepsilon,t}$, we have that, by choosing a proper orientation, $\mathcal{A}_e^{\widetilde{\mathcal{M}}^{\varepsilon,t}}(q) \subset \mathcal{A}_e^{\mathcal{M}^{\varepsilon,t}}(q)$ and thus we obtain

$$\begin{aligned}
-c_N^{-1} H_{\mathcal{M}^{\varepsilon,t},s}(q) &= -c_N^{-1} H_{\widetilde{\mathcal{M}}^{\varepsilon,t},s}(q) \\
&\quad + \int_{\mathbb{R}^N} \frac{\chi_{\mathcal{A}_e^{\mathcal{M}^{\varepsilon,t}}(q)}(x) - \chi_{\mathcal{A}_e^{\widetilde{\mathcal{M}}^{\varepsilon,t}}(q)}(x) + \chi_{\mathcal{A}_i^{\widetilde{\mathcal{M}}^{\varepsilon,t}}(q)}(x) - \chi_{\mathcal{A}_i^{\mathcal{M}^{\varepsilon,t}}(q)}(x)}{|x-q|^{N+s}} dx \\
&= -c_N^{-1} H_{\widetilde{\mathcal{M}}^{\varepsilon,t},s}(q) + \int_{\mathbb{R}^N} \frac{\chi_{\mathcal{A}_e^{\mathcal{M}^{\varepsilon,t}}(q) \setminus \mathcal{A}_e^{\widetilde{\mathcal{M}}^{\varepsilon,t}}(q)}(x) + \chi_{\mathcal{A}_i^{\widetilde{\mathcal{M}}^{\varepsilon,t}}(q) \setminus \mathcal{A}_i^{\mathcal{M}^{\varepsilon,t}}(q)}(x)}{|x-q|^{N+s}} dx \\
&= -c_N^{-1} H_{\widetilde{\mathcal{M}}^{\varepsilon,t},s}(q) + 2 \int_{\mathcal{A}_e^{\mathcal{M}^{\varepsilon,t}}(q) \setminus \mathcal{A}_e^{\widetilde{\mathcal{M}}^{\varepsilon,t}}(q)} \frac{1}{|x-q|^{N+s}} dx. \tag{4.12}
\end{aligned}$$

Recalling that $\mathcal{A}_e^{\mathcal{M}^{\varepsilon,t}}(q) \setminus \mathcal{A}_e^{\widetilde{\mathcal{M}}^{\varepsilon,t}}(q)$ contains the subgraph $\{0 \leq x_N \leq v_\varepsilon(x' - b')\}$ and the subgraph contains the cylinders of height $\phi(\varepsilon)^\beta$ with the base of radius $1/16$, we have

$$\left| \mathcal{A}_e^{\mathcal{M}^{\varepsilon,t}}(q) \setminus \mathcal{A}_e^{\widetilde{\mathcal{M}}^{\varepsilon,t}}(q) \right| \geq c' \phi(\varepsilon)^\beta.$$

where a constant $c' = c'(N) > 0$ depends only on N . Moreover, we observe that the distance between q and the cylinder is less than, at most, $2 + \phi(\varepsilon)^\beta$ and this is bounded from above by some constant depending only on N , s , and β . Hence from (4.11) and (4.12) and by recalling the choice of ϕ , we obtain

$$\begin{aligned}
-c_N^{-1} H_{\mathcal{M}^{\varepsilon,t},s}(q) &\geq -c \phi(\varepsilon)^s + \frac{2c'}{(2 + \phi(\varepsilon)^\beta)^{N+s}} \phi(\varepsilon)^\beta \\
&\geq -c \phi(\varepsilon)^s + c'' \phi(\varepsilon)^\beta \\
&= \phi(\varepsilon)^\beta (-c \phi(\varepsilon)^{s-\beta} + c'') \tag{4.13}
\end{aligned}$$

where $c'' > 0$ is a constant depending only on N , s , and β . Since $0 < \beta < s$ and $\phi(\varepsilon) \downarrow 0$ as $\varepsilon \downarrow 0$, we choose $\varepsilon_1 = \varepsilon_1(N, s, \beta) \in I_\phi \cap (0, \frac{1}{100})$ so small that the right-hand side of (4.13) is positive for any $\varepsilon \in (0, \varepsilon_1]$. Therefore, from (4.13), we obtain that $H_{\mathcal{M}^{\varepsilon,t},s}(q) < 0$ for $\varepsilon \in (0, \varepsilon_1]$.

Now we set $\varepsilon_2 := \min\{\varepsilon_1, \varepsilon_d\}$. Since $r(\varepsilon_d) \geq d$ and $\delta(\varepsilon_2) < \frac{1}{2}$, we may observe that $d(\varepsilon_2) \geq d$ and $\mathcal{M}_3^{\varepsilon_2,t} \cap \Gamma_1 = \emptyset$ for any $t \in (0, \varepsilon_2]$. For our convenience, we denote ε_2 by ε in the sequel.

We then slide the barrier $\mathcal{M}^{\varepsilon,t}$ from above, i.e., we vary the parameter t starting at ε until $\mathcal{M}^{\varepsilon,t}$ touches the critical point \mathcal{M} . To prove the claim, we assume by contradiction that there exists $t_1 \in (0, \varepsilon]$ such that $\mathcal{M} \cap \mathcal{M}_3^{\varepsilon,t_1} \neq \emptyset$ and $\mathcal{M} \cap \mathcal{M}_3^{\varepsilon,t} = \emptyset$ for any $t \in (t_1, \varepsilon]$. We pick up a point $q_{\varepsilon,t_1} \in \mathcal{M} \cap \mathcal{M}_3^{\varepsilon,t_1}$. Notice that

$$\{(x', x_N) \mid -d < x_N < 0\} \cap \mathcal{M}_2^{\varepsilon,t_1} = \emptyset$$

since $d(\varepsilon) \geq d$. See Figure 9 to favor the intuition in dimension 2.

Since \mathcal{M} is a critical point of Area_s under normal variations, we obtain

$$H_{\mathcal{M},s}(q_{\varepsilon,t_1}) = 0.$$

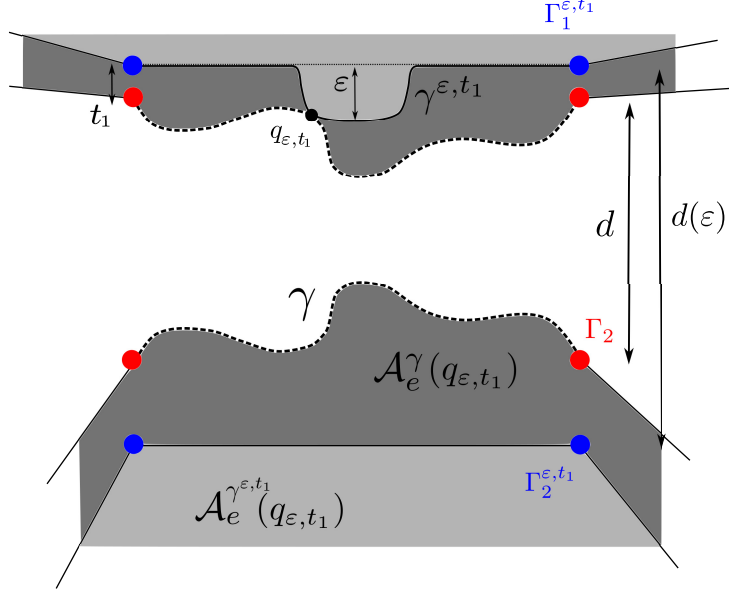


Figure 9: The critical point γ depicted with dashed lines and the barrier $\gamma^{\varepsilon, t_1}$ with black line. γ touches $\gamma^{\varepsilon, t_1}$ at q_{ε, t_1} from above. The exterior $\mathcal{A}_e^{\gamma^{\varepsilon, t_1}}(q_{\varepsilon, t_1})$ of $\gamma^{\varepsilon, t_1}$ is depicted in light gray and the exterior $\mathcal{A}_e^\gamma(q_{\varepsilon, t_1})$ of γ in both light and dark gray.

From Theorem 1.2 and the above argument, we obtain that the touching point $q_{\varepsilon, t_1} := (q'_{\varepsilon, t_1}, q_{\varepsilon, t_1}^N) \in \mathbb{R}^{N-1} \times \mathbb{R}$ satisfies $|q'_{\varepsilon, t_1}| < \delta(\varepsilon)$ and thus $H_{\mathcal{M}^{\varepsilon, t_1}, s}(q_{\varepsilon, t_1}) < 0$. Moreover, from the construction of $\mathcal{M}_1^{\varepsilon, t_1}$, we have, by choosing a proper orientation, that

$$|\mathcal{A}_e^{\mathcal{M}}(q_{\varepsilon, t_1}) \setminus \mathcal{A}_e^{\mathcal{M}^{\varepsilon, t_1}}(q_{\varepsilon, t_1})| > 0 \quad \text{and} \quad |\mathcal{A}_e^{\mathcal{M}^{\varepsilon, t_1}}(q_{\varepsilon, t_1}) \setminus \mathcal{A}_e^{\mathcal{M}}(q_{\varepsilon, t_1})| = 0.$$

Therefore, we obtain

$$\begin{aligned} 0 &= c_N^{-1} (H_{\mathcal{M}, s}(q_{\varepsilon, t_1}) - H_{\mathcal{M}^{\varepsilon, t_1}, s}(q_{\varepsilon, t_1}) + H_{\mathcal{M}^{\varepsilon, t_1}, s}(q_{\varepsilon, t_1})) \\ &< \int_{\mathbb{R}^N} \frac{\chi_{\mathcal{A}_i^{\mathcal{M}}(q_{\varepsilon, t_1})}(x) - \chi_{\mathcal{A}_i^{\mathcal{M}^{\varepsilon, t_1}}(q_{\varepsilon, t_1})}(x) + \chi_{\mathcal{A}_e^{\mathcal{M}^{\varepsilon, t_1}}(q_{\varepsilon, t_1})}(x) - \chi_{\mathcal{A}_e^{\mathcal{M}}(q_{\varepsilon, t_1})}(x)}{|x - q_{\varepsilon, t_1}|^{N+s}} dx + 0 \\ &= -2 \int_{\mathcal{A}_e^{\mathcal{M}}(q_{\varepsilon, t_1}) \setminus \mathcal{A}_e^{\mathcal{M}^{\varepsilon, t_1}}(q_{\varepsilon, t_1})} \frac{1}{|x - q_{\varepsilon, t_1}|^{N+s}} dx < 0, \end{aligned} \quad (4.14)$$

which is a contradiction. We thus conclude that we can slide the barrier $\mathcal{M}^{\varepsilon, t}$ until the boundary $\Gamma_1^{\varepsilon, t} = \partial \mathcal{M}_3^{\varepsilon, t}$ coincides with the boundary $\Gamma_1 = \partial \mathcal{M}_1$. By symmetry, we can slide the barrier from below and do the same argument. \square

Therefore we obtain that two open half-balls of radius ε_2 are contained in a set enclosed by \mathcal{M} and the union of $\mathcal{C} \cap \{x_N = 0\}$ and $\mathcal{C} \cap \{x_N = -d\}$.

As a consequence of Lemma 1.3, we now prove Theorem 1.4.

Proof of Theorem 1.4. Assume that ε_2 and $\widetilde{\mathcal{M}}^{\varepsilon, t}$ are given in the proof of Lemma 1.3 for $\varepsilon \in (0, \varepsilon_2]$ and $t \in (0, \varepsilon]$. From the definition of ε_2 , we can choose $d' > 0$ so small that $d' < (2(N-1))^{-1}$ and that ε_2 can be chosen independently of d for any $d \in (0, d')$. Moreover, if necessary, we may assume that $\varepsilon_2 \phi(\varepsilon_2) < (2(N-1))^{-1}$, which is still independent of d .

Let \mathcal{M} be the critical point chosen in Theorem 1.4. We set $d_0 := \min\{d', \varepsilon_2\}$. From the choice of ϕ and ε_2 , we have that $d_0 \phi(d_0) < (2(N-1))^{-1}$. Then we observe that $d(\varepsilon_2) - t = ((N-1)\phi(\varepsilon_2))^{-1} - t > d$ for any $t \in (0, \varepsilon_2]$ and thus we have that

$$\Gamma_2^{\varepsilon_2, t} \cap \{(x', x_N) \mid -d < x_N < 0\} = \emptyset$$

for any $t \in (0, \varepsilon_2]$ and any $d < d_0$.

Now, by Lemma 1.3, we find that we can slide the barrier $\mathcal{M}^{\varepsilon_2, t}$ until the parameter t reaches 0. Thus, by combining this with Theorem 1.2, we obtain that

$$\begin{aligned} \mathcal{M} &\subset (\bar{\mathcal{C}} \setminus \{(x', x_N) \mid |x'| < \varepsilon_2\}) \cap \{(x', x_N) \mid -d \leq x_N \leq 0\} \\ &= \{(x', x_N) \mid \varepsilon_2 \leq |x'| \leq 1, -d \leq x_N \leq 0\} \end{aligned}$$

for any $d < d_0$.

If $N = 2$, then, since Γ_i consists of two distinct points for $i \in \{1, 2\}$, by a simple geometric argument, we conclude that the critical point \mathcal{M} is disconnected for any $d \in (0, d_0)$. Moreover, from the construction of the barrier, we obtain that there exist two connected components \mathcal{M}_1 and \mathcal{M}_2 of \mathcal{M} such that $\text{dist}(\mathcal{M}_1, \mathcal{M}_2) \geq \varepsilon_2$ and \mathcal{M}_i intersects both Γ_1 and Γ_2 for each $i \in \{1, 2\}$ at its boundary (see also Remark 4.1).

If $N \geq 3$, then, by using homology theory, we conclude that Γ_1 and Γ_2 are in the same connected component of the critical point \mathcal{M} for any $d \in (0, d_0)$ (see [33]). Indeed, we assume by contradiction that there exists a connected component \mathcal{M}_0 of \mathcal{M} with $\partial\mathcal{M}_0 = \Gamma_1$. Taking the fundamental class $[\Gamma_1] \in H_{N-2}(\Gamma_1)$ where $H_k(\mathcal{S})$ is the k -th homology group of \mathcal{S} , we may have that the image in $H_{N-2}(\mathcal{M}_0)$ of $[\Gamma_1]$ by the induced map of homology from the inclusion $i : \Gamma_1 \rightarrow \mathcal{M}_0$ does not vanish because $\mathcal{M}_0 \subset \mathcal{M} \subset A_{\varepsilon_2}$ and A_{ε_2} deformation-retracts to $\Gamma_1 \simeq \mathbb{S}^{N-2}$ where

$$A_{\varepsilon_2} := \{(x', x_N) \mid \varepsilon_2 \leq |x'| \leq 1, -d \leq x_N \leq 0\}.$$

However, since $[\Gamma_1]$ is the boundary of $[\mathcal{M}_0]$ and by using an exact homology sequence of the pair $(\mathcal{M}_0, \Gamma_1)$, we obtain the contradiction. \square

Remark 4.1. Combining Remark 3.3 with Lemma 1.3 and Theorem 1.4, we may observe that two possible critical points of Area_s in dimension 2 whose boundary is $\Gamma_1 \cup \Gamma_2 = \{(\pm 1, d), (\pm 1, -d)\}$ are depicted in Figure 10.

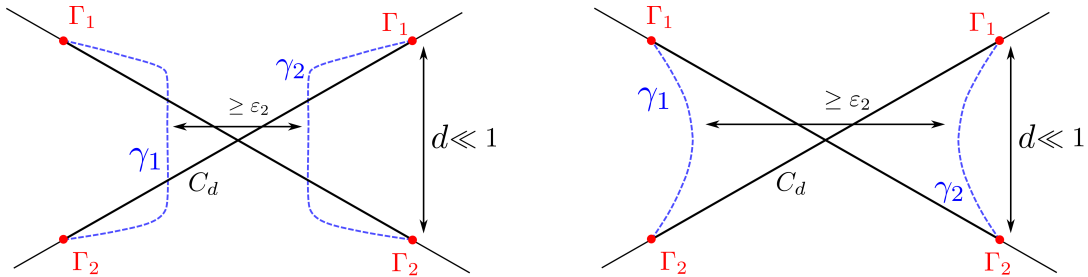


Figure 10: Two possible critical points γ of Area_s in dimension 2 with $\partial\gamma = \Gamma_1 \cup \Gamma_2$ are shown with dashed lines and the cone C_d defined in (3.7) is shown with crossed lines. On the right, γ does not intersect with C_d except at their boundaries Γ_1 and Γ_2 . In both figures, two distinct connected components γ_1 and γ_2 of γ are placed at mutually positive distance of at least $\varepsilon_2 > 0$.

Finally in this section, we prove Theorem 1.5. The idea of the proof is basically the same as the one in [14, Theorem 1.2], i.e., we use the “sliding method” that is developed by Dipierro, Savin, and Valdinoci in [15–17].

Proof of Theorem 1.5. Let \mathcal{M} be the critical point selected in Theorem 1.5 and we set $\Gamma := \Gamma_1 \cup \Gamma_2$.

Given $t \in \mathbb{R}$ and $\alpha \in (0, 1)$, we consider the open ball $B_{d^\alpha/2}(p_{t,d})$ where $p_{t,d} := (te'_1, \frac{-d}{2}) \in \mathbb{R}^{N-1} \times \mathbb{R}$ and $e'_1 := (1, 0, \dots, 0) \in \mathbb{R}^{N-1}$. Here we take d conveniently large so that $d - d^\alpha > 100$. Then we slide the ball from left to right until it touches the critical point \mathcal{M} , which means that we vary t from $t = -\infty$ to $t = +\infty$. Note that $B_{d^\alpha/2}(p_{t,d}) \subset C^c$ for $|t| > 1 + d^\alpha/2$ and $B_{d^\alpha/2}(p_{t,d}) \cap \Gamma = \emptyset$ for any t . To prove the claim, we suppose by contradiction that there exists $t_0 \in \mathbb{R}$ such that $\overline{B_{d^\alpha/2}(p_{t_0,d})} \cap \mathcal{M} = \emptyset$ for $t < t_0$ and $\partial B_{d^\alpha/2}(p_{t_0,d}) \cap \mathcal{M} \neq \emptyset$.

We choose a point $q \in \partial B_{d^\alpha/2}(p_{t_0,d}) \cap \mathcal{M}$. Note that, due to Theorem 1.2, $q \in \mathcal{C}$. By the Euler-Lagrange equation, we have that

$$H_{\mathcal{M},s}(q) = 0. \quad (4.15)$$

Moreover, by choosing a proper orientation, we can choose the interior $\mathcal{A}_i^{\mathcal{M}}(q)$ and exterior $\mathcal{A}_e^{\mathcal{M}}(q)$ at q of \mathcal{M} in such a way that

$$B_{\frac{d^\alpha}{2}}(p_{t_0,d}) \subset \mathcal{A}_i^{\mathcal{M}}(q) \quad \text{and} \quad \mathcal{A}_e^{\mathcal{M}}(q) = \mathbb{R}^N \setminus (\mathcal{A}_i^{\mathcal{M}}(q) \cup \mathcal{M}). \quad (4.16)$$

We now consider the symmetric ball of $B_{d^\alpha/2}(p_{t_0,d})$ with respect to q and we denote it by $\tilde{B} := B_{d^\alpha/2}(\tilde{p}_{t_0,d})$ where $\tilde{p}_{t_0,d} := p_{t_0,d} + 2(q - p_{t_0,d})$.

We define a cylinder S_d as

$$S_d := \{(x', x_N) \mid |x'| < 2, -d < x_N < 0\}.$$

Notice that $\mathcal{C} \cap \{(x', x_N) \mid -d < x_N < 0\} \subset S_d$ and $\mathcal{M} \subset S_d$ thanks to Theorem 1.2. From the symmetry of the balls, we have

$$\int_{S_d \cap B_{\frac{d^\alpha}{2}}(p_{t_0,d})} \frac{dx}{|x - q|^{N+s}} = \int_{S_d \cap \tilde{B}} \frac{dx}{|x - q|^{N+s}}$$

and therefore, from (4.16),

$$\begin{aligned} \int_{S_d} \frac{\chi_{\mathcal{A}_i^{\mathcal{M}}(q)}(x) - \chi_{\mathcal{A}_e^{\mathcal{M}}(q)}(x)}{|x - q|^{N+s}} dx &= \int_{S_d \cap B_{\frac{d^\alpha}{2}}(p_{t_0,d})} \frac{\chi_{\mathcal{A}_i^{\mathcal{M}}(q)}(x) - \chi_{\mathcal{A}_e^{\mathcal{M}}(q)}(x)}{|x - q|^{N+s}} dx \\ &\quad + \int_{S_d \cap \tilde{B}} \frac{\chi_{\mathcal{A}_i^{\mathcal{M}}(q)}(x) - \chi_{\mathcal{A}_e^{\mathcal{M}}(q)}(x)}{|x - q|^{N+s}} dx \\ &\quad + \int_{S_d \setminus (B_{\frac{d^\alpha}{2}}(p_{t_0,d}) \cup \tilde{B})} \frac{\chi_{\mathcal{A}_i^{\mathcal{M}}(q)}(x) - \chi_{\mathcal{A}_e^{\mathcal{M}}(q)}(x)}{|x - q|^{N+s}} dx \\ &\geq \int_{S_d \cap B_{\frac{d^\alpha}{2}}(p_{t_0,d})} \frac{dx}{|x - q|^{N+s}} - \int_{S_d \cap \tilde{B}} \frac{dx}{|x - q|^{N+s}} \\ &\quad - \int_{S_d \setminus (B_{\frac{d^\alpha}{2}}(p_{t_0,d}) \cup \tilde{B})} \frac{dx}{|x - q|^{N+s}} \\ &\geq - \int_{S_d \setminus (B_{\frac{d^\alpha}{2}}(p_{t_0,d}) \cup \tilde{B})} \frac{dx}{|x - q|^{N+s}}. \end{aligned} \quad (4.17)$$

By employing the result in [15, Lemma 3.1] with $R = d^\alpha/2$ and $\lambda = d^{-\frac{\alpha}{2}}$, we obtain

$$\int_{B_{\frac{d^\alpha}{2}}(q) \setminus \left(B_{\frac{d^\alpha}{2}}(p_{t_0, d}) \cup \tilde{B} \right)} \frac{dx}{|x - q|^{N+s}} \leq C_0 d^{-\frac{1+s}{2}\alpha}$$

where $C_0 > 0$ is a constant depending only on N and s . As a consequence, we obtain

$$\begin{aligned} \int_{S_d \setminus \left(B_{\frac{d^\alpha}{2}}(p_{t_0, d}) \cup \tilde{B} \right)} \frac{dx}{|x - q|^{N+s}} &\leq \int_{B_{\frac{d^\alpha}{2}}(q) \setminus \left(B_{\frac{d^\alpha}{2}}(p_{t_0, d}) \cup \tilde{B} \right)} \frac{dx}{|x - q|^{N+s}} \\ &\quad + \int_{S_d \setminus B_{\frac{d^\alpha}{2}}(q)} \frac{dx}{|x - q|^{N+s}} \\ &\leq C_0 d^{-\frac{1+s}{2}\alpha} + \int_{\mathbb{R}^N \setminus B_{\frac{d^\alpha}{2}}(q)} \frac{dx}{|x - q|^{N+s}} \\ &\leq C_0 d^{-\frac{1+s}{2}\alpha} + C_1 d^{-\frac{s}{2}\alpha} \leq C_2 d^{-\frac{s}{2}\alpha} \end{aligned} \quad (4.18)$$

where $C_2 := C_0 + C_1$ is a constant depending only on N and s . From (4.17) and (4.18), we have

$$\begin{aligned} \int_{\mathbb{R}^N} \frac{\chi_{\mathcal{A}_i^{\mathcal{M}}(q)}(x) - \chi_{\mathcal{A}_e^{\mathcal{M}}(q)}(x)}{|x - q|^{N+s}} dx &= \int_{S_d} \frac{\chi_{\mathcal{A}_i^{\mathcal{M}}(q)}(x) - \chi_{\mathcal{A}_e^{\mathcal{M}}(q)}(x)}{|x - q|^{N+s}} dx \\ &\quad + \int_{S_d^c} \frac{\chi_{\mathcal{A}_i^{\mathcal{M}}(q)}(x) - \chi_{\mathcal{A}_e^{\mathcal{M}}(q)}(x)}{|x - q|^{N+s}} dx \\ &\geq -C_2 d^{-\frac{s}{2}\alpha} + \int_{S_d^c} \frac{\chi_{\mathcal{A}_i^{\mathcal{M}}(q)}(x) - \chi_{\mathcal{A}_e^{\mathcal{M}}(q)}(x)}{|x - q|^{N+s}} dx. \end{aligned} \quad (4.19)$$

Now we consider the contributions from $\mathcal{A}_i^{\mathcal{M}}(q)$ and $\mathcal{A}_e^{\mathcal{M}}(q)$ in S_d^c . We now define $C_\Gamma(q)$ by a (filled) cone of vertex q whose boundary passes through Γ . Moreover we define $C_{S_d}(q)$ by a (filled) cone of vertex q whose boundary passes through

$$\partial S_d \cap \{(x', x_N) \mid x_N = 0\} \quad \text{and} \quad \partial S_d \cap \{(x', x_N) \mid x_N = -d\}.$$

From the definitions of S_d and Γ , we observe that

$$C_\Gamma(q) \subset C_{S_d}(q).$$

We now set $\widehat{C}_\Gamma(q) := C_\Gamma(q) \cap \{(x', x_N) \mid x_N > 0 \text{ or } x_N < -d\}$. We then rotate $\widehat{C}_\Gamma(q)$ by angle $\pi/2$ or $-\pi/2$ with respect to the straight line parallel to the x_1 -axis passing through q (if $N = 2$, then we just rotate $\widehat{C}_\Gamma(q)$ by angle $\pi/2$ or $-\pi/2$ with respect to q). Since we choose d so large that $d - d^\alpha > 100$, we obtain that

$$R(\widehat{C}_\Gamma(q)) \subset S_d^c \cap \mathcal{A}_i^{\mathcal{M}}(q) \cap C_{S_d}(q)^c$$

where $R(\widehat{C}_\Gamma(q))$ is an image of $\widehat{C}_\Gamma(q)$ by the rotation map $R : \mathbb{R}^N \rightarrow \mathbb{R}^N$ in the above. See Figure 11 for an intuitive understanding. Then, observing that $R(q) = q$ and

$$S_d^c \cap \mathcal{A}_e^{\mathcal{M}}(q) = \widehat{C}_\Gamma(q)$$

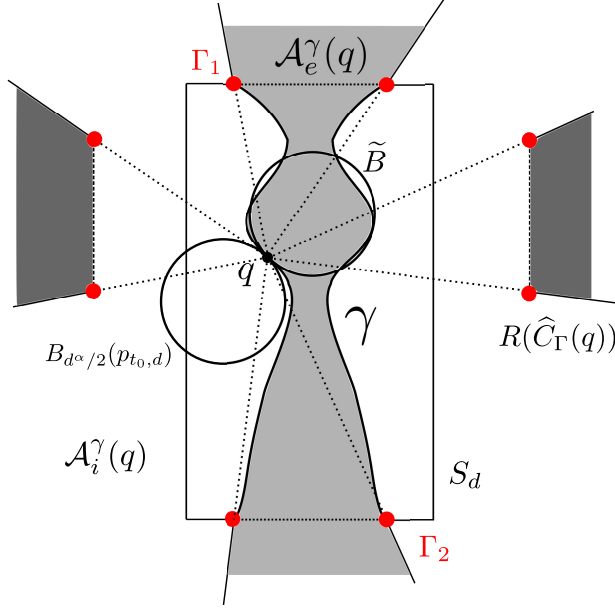


Figure 11: The touching between the ball $B_{d^\alpha/2}(p_{t_0,d})$ and its symmetric ball \tilde{B} at q . The image of the set $\widehat{C}_\Gamma(q)$ by the rotation map R is depicted in dark gray and the set $\mathcal{A}_e^\gamma(q)$ in light gray.

and by a change of variables, we have

$$\int_{R(\widehat{C}_\Gamma(q))} \frac{dx}{|x-q|^{N+s}} = \int_{\widehat{C}_\Gamma(q)} \frac{dx}{|x-q|^{N+s}} = \int_{S_d^c \cap \mathcal{A}_e^\mathcal{M}(q)} \frac{dx}{|x-q|^{N+s}}.$$

From the definitions of S_d and the rotation map R , we can choose an open ball outside S_d and $R(\widehat{C}_\Gamma(q))$ but close to q , i.e., we have

$$B_1(q + 5e_1) \subset (S_d^c \cap \mathcal{A}_i^\mathcal{M}(q)) \setminus R(\widehat{C}_\Gamma(q))$$

where we set $e_1 := (1, 0, \dots, 0) \in \mathbb{R}^N$. Thus, we obtain

$$\begin{aligned} \int_{S_d^c} \frac{\chi_{\mathcal{A}_i^\mathcal{M}(q)}(x) - \chi_{\mathcal{A}_e^\mathcal{M}(q)}(x)}{|x-q|^{N+s}} dx &= \int_{S_d^c \cap \mathcal{A}_i^\mathcal{M}(q)} \frac{dx}{|x-q|^{N+s}} - \int_{\widehat{C}_\Gamma(q)} \frac{dx}{|x-q|^{N+s}} \\ &\geq \int_{S_d^c \cap \mathcal{A}_i^\mathcal{M}(q) \cap R(\widehat{C}_\Gamma(q))} \frac{dx}{|x-q|^{N+s}} \\ &\quad + \int_{(S_d^c \cap \mathcal{A}_i^\mathcal{M}(q)) \setminus R(\widehat{C}_\Gamma(q))} \frac{dx}{|x-q|^{N+s}} \\ &\quad - \int_{\widehat{C}_\Gamma(q)} \frac{dx}{|x-q|^{N+s}} \\ &\geq \int_{B_1(q+5e_1)} \frac{dx}{|x-q|^{N+s}} \\ &= \int_{B_1(5e_1)} \frac{dx}{|x|^{N+s}} =: C_3 > 0 \end{aligned}$$

where C_3 depends only on N and s . This with (4.19) leads to

$$c_N^{-1} H_{\mathcal{M},s}(q) = \int_{\mathbb{R}^N} \frac{\chi_{\mathcal{A}_i^\mathcal{M}(q)}(x) - \chi_{\mathcal{A}_e^\mathcal{M}(q)}(x)}{|x-q|^{N+s}} dx \geq -C_2 d^{-\frac{s}{2}\alpha} + C_3.$$

Therefore, there exists $d_1 = d_1(N, s) > 0$ such that $H_{\mathcal{M},s}(q) > 0$ for any $d > d_1$ and this contradicts the Euler-Lagrange equation (4.15). \square

Remark 4.2. From Lemma 1.3 and the choice of ε_2 in the proof of Lemma 1.3, we also obtain that, for sufficiently large d , a set enclosed by \mathcal{M} and the union of $\mathcal{C} \cap \{x_N = 0\}$ and $\mathcal{C} \cap \{x_N = -d\}$ contains two half-balls of radius $\varepsilon_2 \approx \phi^{-1}(d^{-1})$ where ϕ^{-1} is as in the proof of Lemma 1.3.

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