

# A PROOF OF FINITE CRYSTALLIZATION VIA STRATIFICATION

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ABSTRACT. We devise a new technique to prove two-dimensional crystallization results in the square lattice for finite particle systems. We apply this strategy to energy minimizers of configurational energies featuring two-body short-ranged particle interactions and three-body angular potentials favoring bond-angles of the square lattice. To each configuration, we associate its bond graph which is then suitably modified by identifying chains of successive atoms. This method, called *stratification*, reduces the crystallization problem to a simple minimization that corresponds to a proof via slicing of the isoperimetric inequality in  $\ell^1$ . As a byproduct, we also prove a fluctuation estimate for minimizers of the configurational energy, known as the  $n^{3/4}$ -law.

## 1. INTRODUCTION

At low temperature, atoms and molecules typically arrange themselves into crystalline order. Tackling this phenomenon by using mathematical models consists in proving or disproving that ground states of particle systems for certain configurational energies with interatomic interactions exhibit crystalline order. This issue, referred to as the crystallization problem [5], has attracted a great deal of attention in the physics and mathematics community. By now, various mathematically rigorous crystallization results are available both for systems with a fixed, *finite* number of atoms, and in the so-called *thermodynamic limit* dealing with the infinite particle limit. The reader is referred to [5, 24] for a general overview and also to [33] for a detailed account of available results. The goal of this paper is to revisit the problem of finite crystallization in dimension two, and to present a novel and substantially different proof strategy.

We consider a model where configurations are identified with the respective positions of atoms  $\{x_1, \dots, x_n\}$  in the plane with an associated configurational energy  $\mathcal{E}(\{x_1, \dots, x_n\})$  comprising classical interaction potentials. More specifically,  $\mathcal{E} = \mathcal{E}_2 + \mathcal{E}_3$  decomposes into  $\mathcal{E}_2$  and  $\mathcal{E}_3$  describing two- and three-body interactions, respectively. The two-body interaction potential  $\mathcal{E}_2$  is short-ranged and attractive-repulsive favoring atoms sitting at some specific reference distance. For  $\mathcal{E}_3 \equiv 0$  and for a specific choice of  $\mathcal{E}_2$ , namely the so-called *sticky disc potential*, crystallization in the triangular lattice has been proved by HEITMANN & RADIN [29] (see also [36, 42] for generalizations) and recently revisited in [13], via an approach from discrete differential geometry. If instead  $\mathcal{E}_3$  is active satisfying specific quantitative assumptions, optimal geometries can be identified as the square or the hexagonal lattice [33, 35], depending on whether  $\mathcal{E}_3$  favors triples of particles forming angles which are multiples of  $\frac{\pi}{2}$  or  $\frac{2\pi}{3}$ , respectively. Besides crystallization, fine characterizations of ground-state geometries are available by proving the emergence of hexagonal or square macroscopic Wulff shapes for growing particle numbers [2, 9, 10, 22]. We also refer to related rigorous crystallization results for particle systems involving different types of atoms [4, 7, 38, 20, 21, 23] and to [25, 26, 37, 41] for a nonexhaustive list of results in dimension one.

Although the exact realization of the proof of each result is different depending on the used potentials and the underlying optimal geometry, all proofs follow the very same strategy, originally devised in [28, 29]. One associates a planar graph to the configuration where vertices and edges

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*Key words and phrases.* Crystallization, square lattice, atomic interaction potentials, stratification, edge isoperimetric inequality.

correspond to particles and bonds, respectively. The graph is then separated into the boundary and bulk atoms. The *boundary energy* (roughly, the number of bonds at the boundary of the configuration) is carefully estimated by geometric arguments involving the angles between atoms and relying on the sum of interior angles in planar polygons. Moreover, by means of Euler's formula for planar graphs a connection between the number of bonds and atoms in the configuration is derived. Then, the essential idea of the proof lies in an induction argument over the number of particles: one removes a *bond graph layer*, i.e., the boundary atoms of the configuration, and by induction hypothesis one uses information of the remaining configuration consisting of less atoms. The approach in [13] is different in the sense that it endows the bond graph with a suitable notion of discrete combinatorial curvature and uses a discrete version of the Gauss-Bonnet theorem from differential geometry. However, it still vitally hinges on specific geometric arguments and the induction method over bond graph layers.

It appears to be challenging to generalize this strategy to problems beyond the setting described above. On the one hand, it is hardly conceivable to extend the delicate estimates on the boundary energy to particle systems in three dimensions where surfaces have a much richer structure. On the other hand, the induction method over bond graph layers is often not flexible enough to handle more general situations such as particles systems with two types of atoms with prescribed ratio since this ratio might not be preserved by removing a bond graph layer.

In this work, we propose a new strategy to tackle finite crystallization problems which does not use the induction method over bond graph layers and comes along without arguments from the theory of planar graphs and discrete differential geometry such as Euler's formula or Gauss-Bonnet. It relies on an idea that we call *stratification*. In this paper, we present our technique for the model by MAININI, PIOVANO, & STEFANELLI [33] and reprove finite crystallization in the square lattice, see Theorem 2.1. We are confident, however, that the strategy carries over to other lattices as well, such as the triangular [13, 29] and the hexagonal [35] lattice.

As observed in [33], ground states correspond to configurations minimizing a specific edge perimeter of the configuration, essentially counting the number of missing bonds of atoms having less than four bonds. For ground-state competitors, the bond graph can be locally interpreted as a deformed version of  $\mathbb{Z}^2$ , apart from possible defects in the lattice, see Definition 3.1. Therefore, we can identify chains of atoms in the bond graph where the angle between three successive atoms is near to  $\pi$ , called *strata*. Strata can be *open*, where the first and the last atom of the stratum lie at the boundary of the configuration, or they can be *closed* forming a closed cycle. In contrast to open strata, closed strata do not contribute to the edge perimeter. Therefore, for a correct estimate, we aim at excluding the existence of closed strata. To this end, we observe that, due to the cycle structure of closed strata, there need to exist angles deviating from  $\pi$  and thus contributing to the three-body energy  $\mathcal{E}_3$ . Given specific quantitative assumptions on the potentials similar to the ones in [33], the contribution of  $\mathcal{E}_3$  is large enough to allow us to erase a bond from the stratum to turn it into an open stratum. This procedure is made precise in Lemma 3.7 and referred to as *stratification*. Once all strata are open, the graph satisfies specific properties (see Lemmas 3.3 and 3.6) which reduce our crystallization problem to a simple argument related to an edge isoperimetric inequality on the square lattice. (Compare to [33] for a problem on  $\mathbb{Z}^2$ , and see also [6, 27] for some related classical issues in Discrete Mathematics.)

In contrast to uniqueness of Wulff shapes for continuum crystalline isoperimetric problems, minimizers for a finite number of particles  $n$  are in general not unique. For different lattices in 2D, it has been shown that there are arbitrarily large  $n$  with ground-state configurations deviating from the hexagonal or square macroscopic Wulff shape by a number of  $n^{3/4}$ -particles [9, 10, 33, 39]. Later, this analysis has been extended to the cubic lattice in higher dimensions [32, 34]. The proof of such maximal asymptotic deviation, also known as *maximal fluctuation estimate*, relies on

careful rearrangement techniques for atoms at the boundary and edge-isoperimetric inequalities. In our setting of the square lattice, we can immediately reobtain this so-called  $n^{3/4}$ -law as a mere byproduct of our crystallization proof, see Theorem 2.2. Our argument is similar to [33] with the interesting difference however that our strategy can be applied even if configurations are *not* subset of  $\mathbb{Z}^2$ . We also mention the complementary approach [8], yet restricted to subsets of periodic lattices, where maximal fluctuation estimates are derived via a quantitative version of the edge isoperimetric inequality, based on the quantitative version of the anisotropic isoperimetric inequality proved in [16].

One goal of our work is to revisit finite crystallization results and to suggest a substantially different proof strategy which does not use the induction method over bond graph layers and comes along without arguments from the theory of planar graphs and discrete differential geometry. Besides providing, to our view, a simpler and more direct proof of known results, our main motivation is that our techniques seem promising to tackle more challenging crystallization problems. For example, we expect that our approach can contribute to understand finite crystallization in three dimensions or crystallization for double-bubble problems [11, 12, 19] (configuration with two types of atoms).

Let us highlight that our proof strategy is tailor-made for the problem of finite crystallization. Concerning the thermodynamic limit, i.e., as the number of particles tends to infinity, other techniques are used and allow to prove results under less restrictive assumptions on the potentials. We refer to [3, 14, 15, 40] for results in the plane and to some few available rigorous results [17, 18] in three dimensions.

The article is organized as follows. In Section 2 we introduce our setting and state the main results. Section 3 is devoted to the concept of stratification and in Section 4 we prove our main results. We close the introduction with basic notation. The Euclidian distance between a point  $x \in \mathbb{R}^d$  and a set  $A \subset \mathbb{R}^d$  is denoted by  $\text{dist}(A, x)$ . By  $\#A$  we denote the cardinality of a set  $A$ . By  $B_r(x)$  we indicate the open ball with center  $x \in \mathbb{R}^d$  and radius  $r > 0$ , and simply write  $B_r$  if  $x = 0$ . We define the ceil function by  $\lceil t \rceil := \min\{z \in \mathbb{Z} : z \geq t\}$  for  $t \in \mathbb{R}$ .

## 2. SETTING AND MAIN RESULTS

We consider particle systems in two dimensions, and model their interaction by classical potentials in the frame of Molecular Mechanics [1, 30]. Indicating the configuration of particles by  $C_n = \{x_1, \dots, x_n\} \subset \mathbb{R}^2$ , we define its energy by

$$\mathcal{F}(C_n) = \frac{1}{2} \sum_{i \neq j} v_2(|x_i - x_j|) + \frac{1}{2} \sum_{i,j,k} v_3(\theta_{i,j,k}), \quad (2.1)$$

where the second sum runs over triples  $(i, j, k)$  and  $\theta_{i,j,k}$  denotes the angle formed by the vectors  $x_j - x_i$  and  $x_k - x_i$  (counted clockwise). The factor  $\frac{1}{2}$  accounts for double counting of bonds and angles. In the following, for simplicity we denote the angle formed by the vectors  $x - y$  and  $z - y$  by  $\theta_{x,y,z}$ . We fix  $0 < \varepsilon < \varepsilon_0$  for  $\varepsilon_0 < \frac{\pi}{6}$  specified in Lemma 3.2. The two-body potential  $v_2: [0, +\infty) \rightarrow \mathbb{R} \cup \{+\infty\}$  satisfies

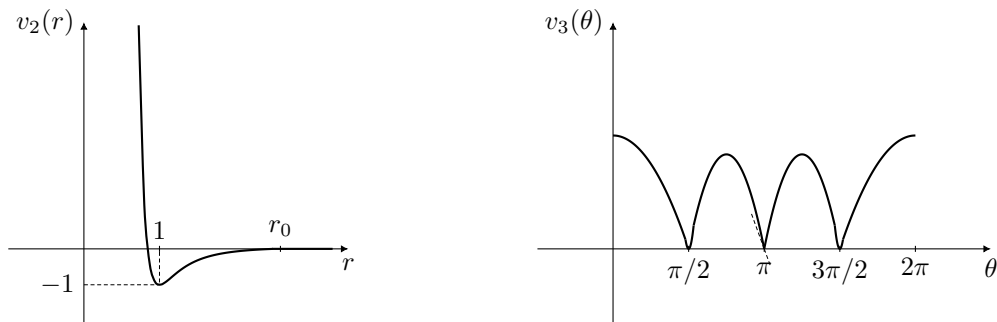
- (i<sub>2</sub>)  $\min_{r \geq 0} v_2(r) = v_2(1) = -1$  and  $v_2(r) > -1$  if  $r \neq 1$ ;
- (ii<sub>2</sub>) There exists  $1 < r_0 < \sqrt{2}$  such that  $v_2(r) = 0$  for all  $r \geq r_0$ ;
- (iii<sub>2</sub>) For all  $r \in [0, 1 - \varepsilon]$  it holds that  $v_2(r) > \varepsilon^{-1}$ .

The three-body potential  $v_3: [0, 2\pi] \rightarrow \mathbb{R}$  satisfies

- (i<sub>3</sub>)  $v_3(\theta) = v_3(2\pi - \theta)$  for all  $\theta \in [0, 2\pi]$ ;
- (ii<sub>3</sub>)  $v_3(k\pi/2) = 0$  for  $k = 1, 2, 3$  and  $v_3(\theta) > 0$  if  $\theta \notin \{\pi/2, \pi, 3\pi/2\}$ ;
- (iii<sub>3</sub>)  $v_3(\theta) \geq 4(\pi/6 - \varepsilon)^{-1} |\theta - \pi|$  for all  $\theta \in [\pi - \varepsilon, \pi + \varepsilon]$  with equality only if  $\theta = \pi$ ;

$$(iv_3) \theta \notin [\pi/2 - \varepsilon, \pi/2 + \varepsilon] \cup [\pi - \varepsilon, \pi + \varepsilon] \cup [3\pi/2 - \varepsilon, 3\pi/2 + \varepsilon] \implies v_3(\theta) > \frac{4}{(1-\varepsilon)^2} (\sqrt{2} + \frac{1}{2})^2.$$

We briefly comment on the assumptions. Condition (i<sub>2</sub>) on a unique minimum (here normalized to 1) is natural, e.g., it is valid for Lennard-Jones-type potentials. Assumption (ii<sub>2</sub>) states that  $v_2$  has compact support. In particular, it ensures that for configurations  $C_n \subset \mathbb{Z}^2$  only atoms at distance 1 interact. These atoms are usually referred to as *nearest neighbors* in the literature. Eventually, (iii<sub>2</sub>) prevents clustering of points. In fact, along with (ii<sub>2</sub>) it shows Definition 3.1(i) in the proof of Lemma 3.2 below. Condition (i<sub>3</sub>) ensures that the potential  $v_3$  does not depend on how (clockwise or counter-clockwise) bond angles are measured, and (ii<sub>3</sub>) guarantees that for  $C_n \subset \mathbb{Z}^2$  there is no contribution stemming from the three-body interaction. Slope conditions similar to (iii<sub>3</sub>) have been used in [20, 21, 33, 35] in order to obtain crystallization on the square or hexagonal lattice. Let us mention that in the other works the condition is needed at *all* minimum points of  $v_3$ , whereas here only at  $\pi$ . As a consequence, the potential is necessarily non-smooth at  $\pi$ . We also point out that in this work the focus lies on a new proof strategy and all appearing specific numerical constants are chosen for computational simplicity rather than optimality. The two potentials are illustrated in Figure 1.



**Figure 1.** The potentials  $v_2$  and  $v_3$ .

We now state the main theorems of the paper.

**Theorem 2.1** (Crystallization). *For each  $C_n \in (\mathbb{R}^2)^n$ , it holds that*

$$\mathcal{F}(C_n) \geq -2n + \lceil 2\sqrt{n} \rceil \tag{2.2}$$

*with equality only if  $C_n \subset \mathbb{Z}^2$  (up to a rigid motion).*

Some configurations of minimal energy are depicted in Figure 2.

**Theorem 2.2** ( $n^{3/4}$ -law). *There exists  $c > 0$  such that for all  $n \in \mathbb{N}$  it holds that each ground state  $C_n$ , up to a rigid motion, satisfies*

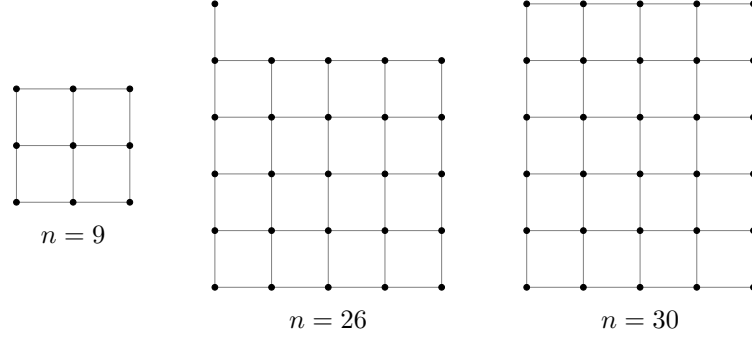
$$\#(C_n \triangle S_n) \leq cn^{3/4},$$

*where  $S_n := [1, \lceil \sqrt{n} \rceil]^2 \cap \mathbb{Z}^2$ .*

Let us note that the scaling is sharp: the construction in [21, Section 3.2] shows that there exists a sequence  $(n_k)_{k \in \mathbb{N}}$  with  $n_k \rightarrow +\infty$  and corresponding ground states  $C_{n_k}$  such that, up to applying any rigid motion to  $C_{n_k}$ , it holds that

$$\#(C_{n_k} \triangle S_{n_k}) \geq cn_k^{3/4}$$

for some  $c > 0$ .



**Figure 2.** Configurations of minimal energy for different cardinality.

### 3. STRATIFICATION

After a short preliminary on graph theory, this section is devoted to the main technique of this paper: modification of bond graphs, called stratification.

**3.1. Bond graph.** We denote by  $G = (V, E)$  a graph, where  $V \subset \mathbb{R}^2$  indicates the set of vertices and  $E \subset \{\{x, y\}: x, y \in V \text{ and } x \neq y\}$  is the set of edges. For  $x \in V$ , we denote the neighborhood with respect to  $G$  by

$$\mathcal{N}(x, E) := \{y \in V : \{x, y\} \in E\}.$$

Given  $G = (V, E)$  we define

$$F(G) = F_{\text{bond}}(G) + F_{\text{ex}}(G),$$

where

$$F_{\text{bond}}(G) = \sum_{x \in V} (4 - \#\mathcal{N}(x, E))$$

is the *bond energy* and

$$F_{\text{ex}}(G) = \sum_{\{x, y\} \in E} (v_2(|x - y|) + 1) + \sum_{\{x, y\}, \{y, z\} \in E} v_3(\theta_{x, y, z})$$

the *excess energy*. For  $V' \subset V$ , we also define the localized elastic energy by

$$F_{\text{ex}}(V') = F_{\text{ex}}(G[V']), \tag{3.1}$$

where  $G[V']$  is the (vertex) induced subgraph of  $V'$  in  $G$ , that is  $G[V'] = (V', E')$  with  $E' = \{\{x, y\} \in E : x, y \in V'\}$ .

We will identify each  $C_n \subset \mathbb{R}^2$  with its *natural bond graph*  $G_{\text{nat}} = (V, E_{\text{nat}})$ , where  $V = C_n$  and the *natural edges* are given by

$$E_{\text{nat}} = \{\{x, y\} : x, y \in C_n, |x - y| \leq r_0\}, \tag{3.2}$$

for  $r_0 > 0$  as given in (ii<sub>2</sub>). This definition is motivated by the relation to (2.1), namely

$$2\mathcal{F}(C_n) = -4n + F(G_{\text{nat}}). \tag{3.3}$$

In Subsection 3.2 below, we will successively modify  $E_{\text{nat}}$  to a smaller set of edges  $E \subset E_{\text{nat}}$ .

**Definition 3.1.** We say that  $G = (V, E)$  is  $\varepsilon$ -regular if:

- (i) If  $\{x, y\} \in E$ , then

$$|x - y| \geq 1 - \varepsilon;$$

(ii) If  $\theta$  is a bond angle, then

$$\theta \in [\pi/2 - \varepsilon, \pi/2 + \varepsilon] \cup [\pi - \varepsilon, \pi + \varepsilon] \cup [3\pi/2 - \varepsilon, 3\pi/2 + \varepsilon].$$

Note that, if  $G_{\text{nat}} = (V, E_{\text{nat}})$  is  $\varepsilon$ -regular, then it is easy to see that  $G = (V, E)$  is  $\varepsilon$ -regular for all  $E \subset E_{\text{nat}}$ .

**Lemma 3.2.** *There exists  $\varepsilon_0 > 0$  such that the following holds true: if  $v_2, v_3$  satisfy (i<sub>2</sub>)–(iii<sub>2</sub>) and (i<sub>3</sub>)–(iv<sub>3</sub>) for some  $0 < \varepsilon < \varepsilon_0$  and if  $C_n$  is a minimizer of (2.1), then its natural bond graph  $G_{\text{nat}} = (V, E_{\text{nat}})$  is  $\varepsilon$ -regular. Moreover, it holds that  $\#\mathcal{N}(x, E_{\text{nat}}) \leq 4$  for all  $x \in V$ .*

Analogous properties have been derived in [33, Propostion 2.1] and [40, Lemma 2.2]. However, as our assumptions on the potentials are slightly different, we include a sketch of the proof for the reader's convenience in Appendix A. For the remainder of this paper, we assume that  $\varepsilon_0 > 0$  is chosen small enough such that Lemma 3.2 holds true and that  $v_2, v_3$  satisfy (i<sub>2</sub>)–(iii<sub>2</sub>) and (i<sub>3</sub>)–(iv<sub>3</sub>) for some  $0 < \varepsilon < \varepsilon_0$ .

**3.2. Stratified bond graph.** Given  $G = (V, E)$ , we say that  $\gamma = (x_1, \dots, x_N)$  with  $x_i \in V$  for all  $i = 1, \dots, N$  is a *straight path* if  $N \geq 2$  and the following holds:

- (i)  $\{x_i, x_{i+1}\} \in E$  for all  $i = 1, \dots, N - 1$ ;
- (ii)  $\theta_i \in [\pi - \varepsilon, \pi + \varepsilon]$  for all  $i \in 2, \dots, N - 1$ , where  $\theta_i = \theta_{x_{i+1}, x_i, x_{i-1}}$ ;
- (iii)  $\{x_i, x_{i+1}\} \neq \{x_j, x_{j+1}\}$  for all  $i, j = 1, \dots, N - 1, j \neq i$ .

(If  $N = 2$ , (ii) and (iii) are empty.) The set of straight paths is denoted by

$$\Gamma(G) := \{\gamma \text{ straight path}\}.$$

We drop  $G$  and write  $\Gamma$  if no confusion arises. If  $\gamma \in \Gamma$  and  $x_1 = x_N$ , we say that  $\gamma$  is *closed* and otherwise that  $\gamma$  is *open*. In the following, we add some strata for degenerate points which will be convenient for Lemma 3.3. Specifically, we define

$$\begin{aligned} V_i &:= \{x \in V : \#\mathcal{N}(x, E) = i\} \text{ for } i = 0, \dots, 4, \\ V_2^\pi &:= \{x \in V_2 : \theta_{x_1, x, x_2} \in [\pi - \varepsilon, \pi + \varepsilon] \text{ where } \mathcal{N}(x, E) = \{x_1, x_2\}\}. \end{aligned} \quad (3.4)$$

Note that in the second definition, one could equally use the angle  $\theta_{x_2, x, x_1}$  as  $\theta_{x_2, x, x_1} = 2\pi - \theta_{x_1, x, x_2}$ . If  $x \in V_0$  we set  $s(x) = \{(x), (x)\}$ , if  $x \in V_1 \cup V_2^\pi$  we set  $s(x) = \{(x)\}$ , and we define the *set of strata* by

$$\mathcal{S}(G) := \mathcal{S}_\Gamma \cup \bigcup_{x \in V_0 \cup V_1 \cup V_2^\pi} s(x), \quad \text{where } \mathcal{S}_\Gamma := \{\gamma \in \Gamma : \gamma \text{ is a maximal element w.r.t. } \subseteq\}. \quad (3.5)$$

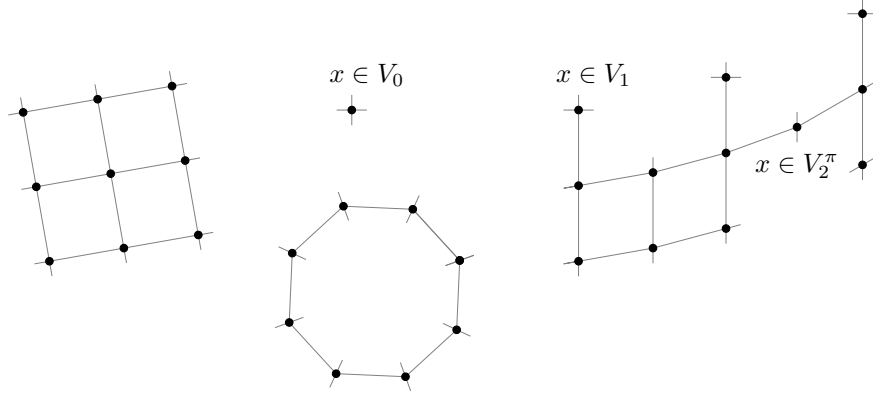
We drop  $G$  and write  $\mathcal{S}$  if no confusion arises. Some closed, open, and degenerate strata are illustrated in Figure 3. Adding the degenerate stratum  $(x)$  with one element twice for  $V_0$  and once for  $V_1 \cup V_2^\pi$  has no geometrical interpretation but is merely convenient to relate the overall number of strata to  $F_{\text{bond}}$ . More precisely, denoting by  $l(s) := \#s$  the length of  $s \in \mathcal{S}$ , we have the following.

**Lemma 3.3.** *(Properties of graphs only containing open paths) Let  $G = (V, E)$  be an  $\varepsilon$ -regular graph. Assume that all  $\gamma \in \Gamma$  are open. Then, the following holds:*

- (i)  $\sum_{s \in \mathcal{S}} l(s) = 2n$ ;
- (ii)  $F_{\text{bond}}(G) = \sum_{x \in V} (4 - \#\mathcal{N}(x, E)) = 2\#\mathcal{S}$ .

*Proof.* We prove the two statements in separate steps.

(i) Since all  $\gamma \in \Gamma$  are open, by definitions (3.4) and (3.5), each  $x \in V$  belongs to exactly two  $s \in \mathcal{S}$ . Indeed, each  $x \in V \setminus (V_0 \cup V_1 \cup V_2^\pi)$  lies in exactly two elements of  $\mathcal{S}_\Gamma$ . (It cannot be contained only in one as this would contradict the openness of straight paths.) Each  $x \in V_1 \cup V_2^\pi$  lies in exactly



**Figure 3.** Illustration of some strata in the bond graph of  $G$ .

one element of  $\mathcal{S}_\Gamma$  and  $x \in V_0$  is not contained in any element of  $\mathcal{S}_\Gamma$ . The definition of  $s(x)$  for  $x \in V_0 \cup V_1 \cup V_2^\pi$  implies that each  $x \in V$  belongs to exactly two  $s \in \mathcal{S}$ . Hence, (i) follows.

(ii) We prove the statement by induction over  $m = \#E$ . It is clearly true for  $m = 0$  since, by definition of (3.5),  $x \in V \implies x \in V_0$  and thus

$$\#\mathcal{S} = 2\#V = \frac{1}{2} \sum_{x \in V} (4 - \#\mathcal{N}(x, \emptyset)) = \frac{1}{2} \sum_{x \in V} (4 - \#\mathcal{N}(x, E)).$$

Let now  $\#E = m \geq 1$  and let  $s = \{x_1, \dots, x_N\} \in \mathcal{S}$  be arbitrary. Consider  $\hat{E} := E \setminus \{x_1, x_2\}$  and the corresponding graph  $\hat{G} = (V, \hat{E})$ . Then,  $\#\hat{E} = m - 1$  and thus, by the induction hypothesis,

$$\sum_{x \in V} (4 - \#\mathcal{N}(x, \hat{E})) = 2\#\mathcal{S}(\hat{G}),$$

where  $\mathcal{S}(\hat{G})$  is the set of strata of  $\hat{G}$ , defined in (3.5). Note that  $\mathcal{S}(\hat{G}) = (\mathcal{S} \cup \{(x_1)\} \cup \{(x_2, \dots, x_N)\}) \setminus s$  and thus  $\#\mathcal{S}(\hat{G}) = \#\mathcal{S} + 1$ . As  $\#\mathcal{N}(x_i, E) = \#\mathcal{N}(x_i, \hat{E}) + 1$  for  $i = 1, 2$ , we have

$$\sum_{x \in V} (4 - \#\mathcal{N}(x, E)) = -2 + \sum_{x \in V} (4 - \#\mathcal{N}(x, \hat{E})) = -2 + 2\#\mathcal{S}(\hat{G}) = 2\#\mathcal{S}.$$

This concludes the proof.  $\square$

We proceed with two definitions and a lemma on graphs with small angle excess.

**Definition 3.4** (Angle excess). Given  $\gamma = \{x_1, \dots, x_N\} \in \Gamma$ , we define the *angle excess* by

$$\theta_{\text{ex}}(\gamma) := \sum_{i=2}^{N-1} |\theta_i - \pi|, \quad \text{where } \theta_i = \theta_{x_{i+1}, x_i, x_{i-1}}.$$

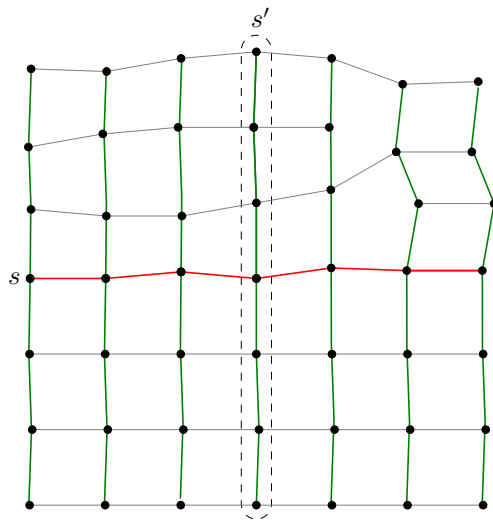
**Definition 3.5** (Orthogonal strata). Let  $s \in \mathcal{S}$ . We define the *set of orthogonal strata* to  $s$  by

$$\mathcal{S}^\perp(s) = \{s' \in \mathcal{S} \setminus \{s\} : s \cap s' \neq \emptyset\}.$$

A stratum  $s \in \mathcal{S}$  and its orthogonal strata are illustrated in Figure 4.

**Lemma 3.6.** (Small angle excess) Let  $G = (V, E)$  be an  $\varepsilon$ -regular graph. The following implications hold true:

- (i) If  $\max_{\gamma \in \Gamma} \theta_{\text{ex}}(\gamma) < \frac{3\pi}{2} - \varepsilon$ , then all  $s \in \Gamma$  are open;



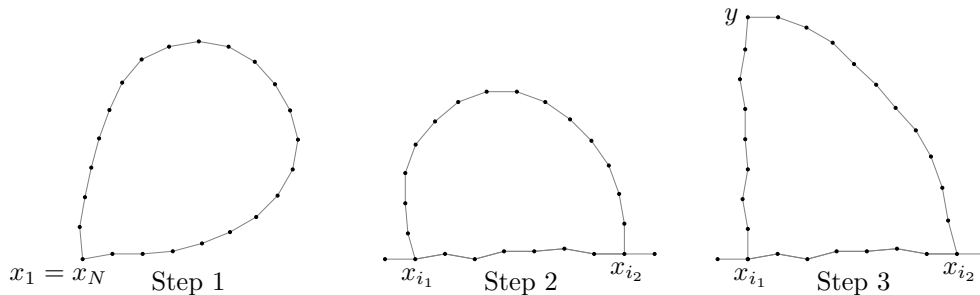
**Figure 4.** The stratum  $s$ , in red, and its orthogonal strata  $\mathcal{S}^\perp(s)$  in green. One  $s' \in \mathcal{S}^\perp(s)$  is encircled.

- (ii) If  $\max_{\gamma \in \Gamma} \theta_{\text{ex}}(\gamma) < \frac{\pi}{2} - \varepsilon$ , then  $\#\mathcal{S}^\perp(s) = l(s)$  for all  $s \in \mathcal{S}$ ;
- (iii) If  $\max_{\gamma \in \Gamma} \theta_{\text{ex}}(\gamma) < \frac{\pi}{6} - \varepsilon$ , then  $s_1 \cap s_2 = \emptyset$  for all  $s_1, s_2 \in \mathcal{S}^\perp(s)$  and for all  $s \in \mathcal{S}$ .

*Proof.* We first introduce some notation that will be used throughout the proof. Let  $p = \{x_1, \dots, x_N\}$  be such that the edges  $e_i = \{x_i, x_{i+1}\}$ ,  $i = 1, \dots, N-1$ , form a closed polygon. We denote by  $\theta(e_i, e_{i+1})$  the interior angle formed by the edges  $e_i$  and  $e_{i+1}$ ,  $i = 1, \dots, N-1$ , with the convention  $e_1 = e_N$ . By the interior angle sum of polygons it holds

$$\sum_{i=2}^N (\theta(e_{i-1}, e_i) - \pi) = -2\pi. \quad (3.6)$$

For the reader's convenience, the proof of the three different statements is aided by Figure 5.



**Figure 5.** The three cases discussed in Steps 1-3.

Step 1.(Proof of (i)) Assume by contradiction that  $\theta_{\text{ex}}(\gamma) < \frac{3\pi}{2} - \varepsilon$  for all  $\gamma \in \Gamma$  and that there exists  $\gamma \in \Gamma$  closed. Let  $\gamma = \{x_1, \dots, x_N\} \in \Gamma$  be a minimal (w.r.t. set inclusion) closed path. The edges  $e_i = \{x_i, x_{i+1}\}$ ,  $i = 1, \dots, N-1$ , form a closed polygon. Therefore, by (3.6) and the triangle inequality, we have  $\theta_{\text{ex}}(\gamma) + |\theta(e_{N-1}, e_1) - \pi| \geq 2\pi$ . Since  $|\theta(e_{N-1}, e_1) - \pi| \leq \frac{\pi}{2} + \varepsilon$  by



Definition 3.1(ii), this yields a contradiction and concludes Step 1.

Step 2.(Proof of (ii)) Assume by contradiction that  $\theta_{\text{ex}}(\gamma) < \frac{\pi}{2} - \varepsilon$  for all  $\gamma \in \Gamma$  and that there exists  $s = \{x_1, \dots, x_N\} \in \mathcal{S}$  with  $\#\mathcal{S}^\perp(s) < l(s)$ . This implies that  $N \geq 2$ . Moreover, there exists  $s' \in \mathcal{S}^\perp(s)$  and  $1 \leq i_1 < i_2 \leq N$  such that  $\{x_{i_1}, x_{i_2}\} \subset s' \cap s$ . Let us consider  $\gamma \subset s'$  connecting  $x_{i_1}$  and  $x_{i_2}$  such that  $\gamma \cap s = \{x_{i_1}, x_{i_2}\}$ . We now consider  $p = \{y_1, \dots, y_M\} = \phi \cup \gamma$ , where  $\phi := \{x_{i_1}, \dots, x_{i_2}\} \subset s$ , and observe that its edges  $e_i$ ,  $i = 1, \dots, M-1$ , form a closed polygon. Let  $y_j = x_{i_1}$  and  $y_k = x_{i_2}$ . Note that  $|\theta(e_{j-1}, e_j) - \pi|, |\theta(e_{k-1}, e_k) - \pi| \leq \frac{\pi}{2} + \varepsilon$  by Definition 3.1(ii). Identity (3.6) applied to  $p$  implies  $\sum_{i=2}^M (\theta(e_{i-1}, e_i) - \pi) = -2\pi$ . Furthermore,  $\phi, \gamma \in \Gamma$  and therefore we obtain

$$\begin{aligned} 2\pi &= \left| \sum_{i=2}^M (\theta(e_{i-1}, e_i) - \pi) \right| \leq |\theta(e_{j-1}, e_j) - \pi| + |\theta(e_{k-1}, e_k) - \pi| + \sum_{\substack{i=2 \\ i \neq \{j,k\}}}^M |\theta(e_{i-1}, e_i) - \pi| \\ &\leq \pi + 2\varepsilon + \sum_{\substack{i=2 \\ i \neq \{j,k\}}}^M |\theta(e_{i-1}, e_i) - \pi| = \pi + 2\varepsilon + \theta_{\text{ex}}(\gamma) + \theta_{\text{ex}}(\phi). \end{aligned}$$

This implies that  $\theta_{\text{ex}}(\gamma) \geq \pi/2 - \varepsilon$  or  $\theta_{\text{ex}}(\phi) \geq \pi/2 - \varepsilon$  and yields therefore a contradiction.

Step 3.(Proof of (iii)) Assume by contradiction that  $\theta_{\text{ex}}(\gamma) < \pi/6 - \varepsilon$  for all  $\gamma \in \Gamma$  and that there exists  $s \in \mathcal{S}$  and  $s_1, s_2 \in \mathcal{S}^\perp(s)$  such that  $s_1 \cap s_2 \neq \emptyset$ . Writing  $s = \{x_1, \dots, x_N\}$  there exists  $i_1 < i_2$  such that  $s_1 \cap s = \{x_{i_1}\}$  and  $s_2 \cap s = \{x_{i_2}\}$ . (Due to Step 2, there is only one point of intersection between  $s_i$  and  $s$ .) Again by Step 2, there holds  $\{y\} = s_1 \cap s_2$  for some  $y \in V$ . Denote by  $\gamma_1 = \{y_1, \dots, y_{l_1}\} \subset s$  the path connecting  $x_{i_1}$  with  $x_{i_2}$ ,  $\gamma_2 = \{y_{l_1}, \dots, y_{l_2}\} \subset s_2$  the path connecting  $x_{i_2}$  with  $y$ , and by  $\gamma_3 = \{y_{l_2}, \dots, y_{l_3}\} \subset s_1$  the path connecting  $y$  with  $x_{i_1}$ . Note that  $\gamma_i \in \Gamma$  for  $i = 1, 2, 3$ . We set  $p = \gamma_1 \cup \gamma_2 \cup \gamma_3$  and observe that the edges of  $p$  form a closed polygon. Equation (3.6) applied for  $p$  implies  $\sum_{i=2}^{l_3} (\theta(e_{i-1}, e_i) - \pi) = -2\pi$ . Note that  $|\theta(e_{l_k-1}, e_{l_k}) - \pi| \leq \frac{\pi}{2} + \varepsilon$  for all  $k = 1, 2, 3$  by Definition 3.1(ii). Therefore, we obtain

$$\begin{aligned} 2\pi &= \left| \sum_{i=2}^{l_3} (\theta(e_{i-1}, e_i) - \pi) \right| \leq \sum_{k=1}^3 |\theta(e_{l_k-1}, e_{l_k}) - \pi| + \sum_{\substack{i=2 \\ i \neq \{l_1, l_2, l_3\}}}^{l_3} |\theta(e_{i-1}, e_i) - \pi| \\ &\leq \frac{3\pi}{2} + 3\varepsilon + \sum_{k=1}^3 \theta_{\text{ex}}(\gamma_k). \end{aligned}$$

This implies that there exists  $k \in \{1, 2, 3\}$  such that  $\theta_{\text{ex}}(\gamma_k) \geq \pi/6 - \varepsilon$ , a contradiction.  $\square$

We now come to the *stratification* of bond graphs. The following lemma allows to reduce the problem of crystallization to a purely geometric problem of minimizing the number of strata in graphs containing only open strata with small angle excess.

**Lemma 3.7.** (Construction of a graph with small angle excess) *Let  $G = (V, E)$  be  $\varepsilon$ -regular. Then, there exists  $G_o = (V, E_o)$  with  $E_o \subset E$  such that*

- (i)  $\max_{\gamma \in \Gamma(G_o)} \theta_{\text{ex}}(\gamma) < \frac{\pi}{6} - \varepsilon$ ;
- (ii)  $G_o$  satisfies

$$F(G) \geq F_{\text{bond}}(G_o)$$

with equality only if  $E = E_o$ ,  $|x - y| = 1$  for all  $x \in V$ ,  $y \in \mathcal{N}(x, E)$ , and  $\theta \in \{\pi/2, \pi, 3\pi/2\}$  for all bond angles  $\theta$ .

*Proof.* We construct  $G_o = (V, E_o)$  by iteratively erasing edges. We start by setting  $G^0 = (V, E)$  and we suppose that  $G^k = (V, E^k)$  is already given. We construct  $G^{k+1} = (V, E^{k+1})$  by suitably modifying the set of edges  $E^k$ . If (i) is satisfied, we may stop. Thus, we assume that there exists  $\gamma \in \Gamma(G_k)$  such that  $\theta_{\text{ex}}(\gamma) \geq \pi/6 - \varepsilon$ . Let  $\gamma_k \in \Gamma(G_k)$  be minimal (w.r.t. set inclusion) such that  $\theta_{\text{ex}}(\gamma_k) \geq \pi/6 - \varepsilon$ , i.e.,  $\gamma_k = \{x_1, x_2, \dots, x_{N-1}, x_N\}$  and  $\hat{\gamma}_k := \{x_2, \dots, x_{N-1}\}$  satisfies  $\theta_{\text{ex}}(\hat{\gamma}_k) < \pi/6 - \varepsilon$ . We define  $E^{k+1} := E^k \setminus (\{x_1, x_2\} \cup \{x_{N-1}, x_N\})$  and  $G^{k+1} = (V, E^{k+1})$ . Then,

$$\sum_{x \in V} (4 - \#\mathcal{N}(x, E^{k+1})) = 4 + \sum_{x \in V} (4 - \#\mathcal{N}(x, E^k)). \quad (3.7)$$

Additionally, due to (iii<sub>3</sub>), with  $L := 4(\pi/6 - \varepsilon)^{-1} > 0$  we have

$$\sum_{j=2}^{N-1} v_3(\theta_j) > L \sum_{j=2}^{N-1} |\theta_j - \pi| \geq L(\pi/6 - \varepsilon) = 4. \quad (3.8)$$

Therefore, due to (3.7) and (3.8), we have that

$$F_{\text{ex}}(\gamma_k) + F_{\text{bond}}(G^k) > F_{\text{bond}}(G^{k+1}), \quad (3.9)$$

where  $F_{\text{ex}}(\gamma_k)$  is defined in (3.1). Since  $G = (V, E)$  is finite, the procedure terminates for some  $K \in \mathbb{N}$  and we set  $G_o := (V, E^K)$ . By construction,  $G_o$  satisfies (i). It remains to show (ii). Note that, due to the minimal selection of  $\gamma_k \in \Gamma$ , once  $\gamma_k$  is selected this way, we will not select any  $\gamma' \subset \gamma_k$  in any successive step  $j > k$ . Thus, using (3.9) and the previous observation, we have

$$\begin{aligned} F_{\text{ex}}(G) + F_{\text{bond}}(G) &= F_{\text{ex}}(G^0) + E_{\text{bond}}(G^0) \geq \sum_{k=0}^{K-1} F_{\text{ex}}(\gamma_k) + E_{\text{bond}}(G^0) \\ &\geq F_{\text{bond}}(G^K) = F_{\text{bond}}(G_o) \end{aligned} \quad (3.10)$$

with strict inequality whenever  $K \geq 1$ . In particular, if equality holds in (3.10), we have that  $G_o = G$ . This necessarily gives  $F_{\text{ex}}(G) = 0$  which implies that  $|x - y| = 1$  for all  $x \in V, y \in \mathcal{N}(x, E)$ , and  $\theta \in \{\pi/2, \pi, 3\pi/2\}$  for all bond angles by (i<sub>2</sub>) and (i<sub>3</sub>). This concludes the proof.  $\square$

#### 4. PROOF OF THE MAIN RESULTS

This section is devoted to the proofs of Theorems 2.1–2.2.

**4.1. Crystallization.** We will show that the minimum of  $F$  is given by  $2\lceil 2\sqrt{n} \rceil$ , and that it is attained by subsets of  $\mathbb{Z}^2$ . In view of (3.3), this shows Theorem 2.1. Recall the definition of  $G_{\text{nat}}$  in (3.2). We first state the following upper bound.

**Lemma 4.1.** (*Upper bound*) *Let  $C_n$  be a minimizer of (2.1). Then,  $G_{\text{nat}}$  satisfies*

$$F(G_{\text{nat}}) \leq 2\lceil 2\sqrt{n} \rceil.$$

*Proof.* The statement is obtained by direct construction of configurations  $C_n$  with  $C_n \subset \mathbb{Z}^2$  satisfying the energy bound. We refer to [33, Section 4] for details, see also Figure 2.  $\square$

The core of the proof now consists in proving a lower bound.

*Proof of Theorem 2.1.* Let  $C_n$  be a minimizer of (2.1). Then  $G_{\text{nat}}$  is  $\varepsilon$ -regular by Lemma 3.2. We denote by  $G_o = (V, E_o)$  the graph obtained in Lemma 3.7. The graph  $G_o$  is also  $\varepsilon$ -regular and satisfies

$$\max_{\gamma \in \Gamma(G_o)} \theta_{\text{ex}}(\gamma) < \pi/6 - \varepsilon. \quad (4.1)$$

The main part of the proof consists in verifying

$$F_{\text{bond}}(G_o) \geq 2\lceil 2\sqrt{n} \rceil. \quad (4.2)$$

Once (4.2) is proven, we conclude as follows. First, (2.2) holds due to Lemma 3.7(ii) and (3.3). To characterize the equality case, we get from Lemma 3.7 that  $G = G_o$ , that all bond lengths are 1, and all bond angles lie in  $\{\pi/2, \pi, 3/2\pi\}$ . This shows that each connected component (in the sense of graph theory) of  $G$  lies in a rotated and shifted version of  $\mathbb{Z}^2$ . If there existed more than one connected component, one could obtain a modified configuration with an additional bond. This contradicts minimality, and we therefore obtain  $V \subset \mathbb{Z}^2$  after a rigid motion.

We now show (4.2). In the following, we write  $\mathcal{S}$  in place of  $\mathcal{S}(G_o)$  for simplicity. By Lemma 4.1, Lemma 3.3(ii), and Lemma 3.7 we have that

$$2\#\mathcal{S} = F_{\text{bond}}(G_o) \leq 2\lceil 2\sqrt{n} \rceil. \quad (4.3)$$

Furthermore, by Lemma 3.3(i) we have

$$\sum_{s \in \mathcal{S}} l(s) = 2n.$$

Hence, there exists  $s_0 \in \mathcal{S}$  such that

$$l(s_0) = \max_{s \in \mathcal{S}} l(s) \geq \frac{1}{\#\mathcal{S}} \sum_{s \in \mathcal{S}} l(s) \geq \frac{2n}{\lceil 2\sqrt{n} \rceil}. \quad (4.4)$$

Recall Definition 3.5 and define  $l^v := \max_{s \in \mathcal{S}^\perp(s_0)} l(s)$  and  $s^v \in \operatorname{argmax}_{s \in \mathcal{S}^\perp(s_0)} l(s)$ . We claim that

$$\#(\mathcal{S}^\perp(s_0) \cup \mathcal{S}^\perp(s^v)) = \#\mathcal{S}^\perp(s_0) + \#\mathcal{S}^\perp(s^v) = l(s_0) + l^v. \quad (4.5)$$

In fact, by (4.1) and Lemma 3.6(ii) we have that  $\#\mathcal{S}^\perp(s_0) = l(s_0)$ ,  $\#\mathcal{S}^\perp(s^v) = l^v$  and, by Lemma 3.6(iii), if  $s \in \mathcal{S}^\perp(s_0)$ , then  $s \notin \mathcal{S}^\perp(s^v)$  (and vice versa). This yields (4.5). We set  $\operatorname{span}(s_0) = \bigcup_{s' \in \mathcal{S}^\perp(s_0)} s' \subseteq V$  and we consider two cases:

- (a)  $\operatorname{span}(s_0) = V$ ;
- (b)  $\operatorname{span}(s_0) \subsetneq V$ ;

*Proof in case (a):* Due to (4.5), we get

$$\#\mathcal{S} \geq \#(\mathcal{S}^\perp(s_0) \cup \mathcal{S}^\perp(s^v)) = l^v + l(s_0).$$

Now, since  $\operatorname{span}(s_0) = V$ , we have in particular that  $l(s_0) \cdot l^v \geq n$  and therefore, noting that  $\lceil t \rceil < t + 1$  we obtain by Lemma 3.3(ii)

$$F_{\text{bond}}(G_o) = 2\#\mathcal{S} \geq 2(l(s_0) + l^v) \geq 2 \left( l(s_0) + \frac{n}{l(s_0)} \right) \geq 4\sqrt{n} > 2\lceil 2\sqrt{n} \rceil - 2. \quad (4.6)$$

Since  $2\#\mathcal{S} \in 2\mathbb{N}$ , the previous estimate yields the claim (4.2) in case (a).

*Proof in case (b):* We claim that in this case we have that

$$\#\mathcal{S} \geq l^v + l(s_0) + 1. \quad (4.7)$$

In fact, by definition, there exists  $x \in V \setminus \operatorname{span}(s_0)$ . Due to Lemma 3.6(iii), for  $s, s' \in \mathcal{S}^\perp(s^v)$  we have that  $s \cap s' = \emptyset$ , and thus there exists at most one  $s \in \mathcal{S}^\perp(s^v)$  such that  $s \cap \{x\} \neq \emptyset$ . We also note that  $s' \cap \{x\} = \emptyset$  for all  $s' \in \mathcal{S}^\perp(s_0)$ . Since for all  $x \in V$  there exist two strata  $s, s'$  such that  $x \in s, s'$  (see proof of Lemma 3.3(i)), there exists at least one stratum  $s \notin \mathcal{S}^\perp(s_0) \cup \mathcal{S}^\perp(s^v)$ . Therefore (4.7) follows.

We denote by  $\mathcal{S}^a := \mathcal{S} \setminus (\mathcal{S}^\perp(s_0) \cup \mathcal{S}^\perp(s^v))$  and observe that by (4.3) and (4.5) it holds that

$$\#\mathcal{S}^a = \#\mathcal{S} - l^v - l(s_0) \leq \lceil 2\sqrt{n} \rceil - l^v - l(s_0).$$

Now, by Lemma 3.6(ii) and the choice of  $s_0$ , see (4.4), and  $s^v$  respectively, we have  $\#\mathcal{S}^\perp(s_0) = l(s_0)$ ,  $\#\mathcal{S}^\perp(s^v) = l^v$ ,  $l(s) \leq l(s_0)$  for all  $s \in \mathcal{S}$ , and  $l(s) \leq l^v$  for all  $s \in \mathcal{S}^\perp(s_0)$ . Due to Lemma 3.3(i) and  $\mathcal{S}^\perp(s_0) \cap \mathcal{S}^\perp(s^v) = \emptyset$ , we get

$$2l(s_0) \cdot l^v + l(s_0)(\lceil 2\sqrt{n} \rceil - l^v - l(s_0)) \geq \sum_{s \in \mathcal{S}^\perp(s_0)} l(s) + \sum_{s \in \mathcal{S}^\perp(s^v)} l(s) + \sum_{s \in \mathcal{S}^a} l(s) = \sum_{s \in \mathcal{S}} l(s) = 2n,$$

and thus

$$l^v \geq \frac{2n}{l(s_0)} + l(s_0) - \lceil 2\sqrt{n} \rceil.$$

This together with Lemma 3.3(ii), (4.5), and  $\lceil t \rceil < t + 1$  implies

$$\begin{aligned} F_{\text{bond}}(G_o) &\geq 2(l^v + l(s_0) + \#\mathcal{S}^a) \geq 2 \left( \frac{2n}{l(s_0)} + 2l(s_0) - \lceil 2\sqrt{n} \rceil + \#\mathcal{S}^a \right) \geq 2(4\sqrt{n} - \lceil 2\sqrt{n} \rceil + \#\mathcal{S}^a) \\ &> 2(2\lceil 2\sqrt{n} \rceil - \lceil 2\sqrt{n} \rceil + \#\mathcal{S}^a - 2) = 2\lceil 2\sqrt{n} \rceil + 2(\#\mathcal{S}^a - 2). \end{aligned}$$

Again, since  $2(l^v + l(s_0) + \#\mathcal{S}^a) \in 2\mathbb{N}$  and,  $\#\mathcal{S}^a \geq 1$  by (4.5)–(4.7), the claim (4.2) follows also in case (b). This finishes the proof.

Finally, we make the following observation: the argument along with Lemma 3.7(ii) also shows that  $\#\mathcal{S}^a \geq 2$  would induce that  $G$  is not a ground state. From this and (4.7) we deduce

$$\#\mathcal{S} = l^v + l(s_0) + 1 \tag{4.8}$$

for the number of strata of a ground state  $G$  with  $\text{span}(s_0) \subsetneq V$ .  $\square$

Estimate (4.6) is related to proving an isoperimetric inequality with respect to the  $l^1$ -perimeter via slicing. We present a corresponding argument in the continuum setting in Appendix B.

**4.2. The  $n^{3/4}$ -law.** We close with a fluctuation estimate for minimizers.

*Proof of Theorem 2.2.* Clearly, it is enough to prove the statement for  $n \geq n_0$  for some  $n_0 \in \mathbb{N}$ . We use the notation of the previous proof. In particular, we choose  $s_0$  and  $l^v$  as done before (4.5). As each  $x \in V$  belongs to exactly two strata, by (4.3) we find that  $l(s_0) = \max_{s \in \mathcal{S}} l(s) \leq \lceil 2\sqrt{n} \rceil$ . We start by noting that

$$l(s_0) \cdot l^v \geq n - \lceil 2\sqrt{n} \rceil. \tag{4.9}$$

Indeed, if  $\text{span}(s_0) = V$ , we have  $l(s_0) \cdot l^v \geq n$ . Otherwise, in view of (4.8), the span missed exactly one stratum, consisting of at most  $\lceil 2\sqrt{n} \rceil$  points.

Now, by (4.4), (4.5), and (4.9) we compute for  $n$  sufficiently large

$$\begin{aligned} 4\sqrt{n} + 2 &\geq 2\lceil 2\sqrt{n} \rceil = F(G) = 2\#\mathcal{S} \geq 2(l(s_0) + l^v) \geq 2l(s_0) + 2\frac{n - \lceil 2\sqrt{n} \rceil}{l(s_0)} \\ &\geq 2l(s_0) + 2\frac{n}{l(s_0)} - \frac{(2\sqrt{n} + 1)^2}{n} \geq 2l(s_0) + 2\frac{n}{l(s_0)} - 5. \end{aligned}$$

This yields

$$2l(s_0)^2 - (4\sqrt{n} + 7)l(s_0) + 2n \leq 0.$$

Therefore

$$x_- \leq l(s_0) \leq x_+,$$

where

$$x_\pm = \frac{4\sqrt{n} + 7 \pm \sqrt{(4\sqrt{n} + 7)^2 - 16n}}{4}.$$

Then, a short computation yields

$$\sqrt{n} - cn^{1/4} \leq l(s_0) \leq \sqrt{n} + cn^{1/4} \quad (4.10)$$

for some universal  $c > 0$  large enough. Using again (4.9) and  $l^v \leq l(s_0)$ , we also find

$$\sqrt{n} - cn^{1/4} \leq l^v \leq \sqrt{n} + cn^{1/4} \quad (4.11)$$

for a larger  $c > 0$ . In view of (4.8), we get that, up to a translation and up to one stratum,  $C_n$  is contained in the rectangular subset of  $\mathbb{Z}^2$  defined by

$$R_n := \{(k_1, k_2): k_1 \in \{1, \dots, l(s_0)\}, k_2 \in \{1, \dots, l^v\}\}.$$

As each stratum consists of at most  $\lceil 2\sqrt{n} \rceil$  points, we get

$$\#(C_n \setminus R_n) \leq \lceil 2\sqrt{n} \rceil. \quad (4.12)$$

Note that (4.10)–(4.11) imply

$$\#(R_n \triangle S_n) \leq cn^{3/4}.$$

Thus, recalling (4.12), to conclude it now suffices to prove that

$$l(s_0) \cdot l^v - n \leq 5\sqrt{n}$$

for some  $c > 0$ . Assume by contradiction that  $l(s_0) \cdot l^v - n > 5\sqrt{n}$ . Then, since  $l(s_0) \leq \lceil 2\sqrt{n} \rceil$ , we get

$$4\sqrt{n} + 2 \geq 2\lceil 2\sqrt{n} \rceil = 2\#\mathcal{S} \geq 2(l(s_0) + l^v) > 2l(s_0) + 2\frac{n}{l(s_0)} + 2\frac{5\sqrt{n}}{l(s_0)} \geq 2l(s_0) + 2\frac{n}{l(s_0)} + \frac{5}{2}$$

for  $n$  large enough. Since  $2l(s_0) + 2\frac{n}{l(s_0)} \geq 4\sqrt{n}$ , we obtain a contradiction. This concludes the proof.  $\square$

#### ACKNOWLEDGEMENTS

This work was supported by the DFG project FR 4083/3-1 and by the Deutsche Forschungsgemeinschaft (DFG, German Research Foundation) under Germany's Excellence Strategy EXC 2044 -390685587, Mathematics Münster: Dynamics–Geometry–Structure.

#### APPENDIX A. PROOF OF LEMMA 3.2

*Proof of Lemma 3.2.* Let  $C_n$  be a minimizer of (2.1). For simplicity, we write  $G = (V, E)$  instead of  $G_{\text{nat}} = (V, E_{\text{nat}})$  for the associated natural bond graph.

Step 1. In this step, we show Definition 3.1(i), that is

$$|x - y| \geq 1 - \varepsilon \text{ for all } \{x, y\} \in E. \quad (A.1)$$

Define

$$M := \max_{x \in \mathbb{R}^2} \#(V \cap B_{\frac{1}{2}(1-\varepsilon)}(x)). \quad (A.2)$$

It suffices to show  $M = 1$ . Let  $x_0 \in \mathbb{R}^2$  be a maximizer. After translation of  $V$ , it is not restrictive to assume that  $x_0 = 0$ . By assumption (iii)<sub>2</sub> we have

$$\sum_{\substack{x, y \in B_{\frac{1}{2}(1-\varepsilon)} \\ \{x, y\} \in E}} v_2(|x - y|) \geq \frac{1}{2\varepsilon} M(M - 1). \quad (A.3)$$

Consider the annulus  $A_\varepsilon := B_{\frac{1}{2}(1-\varepsilon)+\sqrt{2}} \setminus B_{\frac{1}{2}(1-\varepsilon)} \subset B_{\frac{1}{2}+\sqrt{2}}$ . There exists  $N \in \mathbb{N}$  and  $\{z_i\}_{i=1}^N \subset \mathbb{R}^2$  such that for all  $0 < \varepsilon \leq \frac{1}{2}$

$$A_\varepsilon \subset B_{\frac{1}{2}+\sqrt{2}} \subset \bigcup_{i=1}^N B_{\frac{1}{4}}(z_i) \subset \bigcup_{i=1}^N B_{\frac{1}{2}(1-\varepsilon)}(z_i).$$

Thus, recalling (A.2), we have

$$\#(V \cap A_\varepsilon) \leq \# \left( V \cap \bigcup_{i=1}^N B_{\frac{1}{2}(1-\varepsilon)}(z_i) \right) \leq \sum_{i=1}^N \# \left( V \cap B_{\frac{1}{2}(1-\varepsilon)}(z_i) \right) \leq NM. \quad (\text{A.4})$$

By (A.4), the definition of  $M$ , and (i<sub>2</sub>), (ii<sub>2</sub>) we have

$$\sum_{\substack{x \in B_{\frac{1}{2}(1-\varepsilon)}, y \in A_\varepsilon \\ \{x, y\} \in E}} v_2(|x - y|) \geq -1 \cdot \#\{(x, y) : x \in V \cap B_{\frac{1}{2}(1-\varepsilon)}, y \in V \cap A_\varepsilon\} \geq -NM^2. \quad (\text{A.5})$$

We write  $V \cap B_{\frac{1}{2}(1-\varepsilon)} = \{x_i\}_{i=1}^M$  and consider a competitor  $\hat{V}$  (with associated natural bond graph  $\hat{G}$ ) given by

$$\hat{V} = (V \setminus B_{\frac{1}{2}(1-\varepsilon)}) \cup \bigcup_{i=1}^M \{x_i + \tau_i\},$$

where  $\tau_i \in \mathbb{R}^2$  are chosen such that

$$\text{dist}(x_i + \tau_i, \hat{V} \setminus \{x_i + \tau_i\}) \geq \sqrt{2} \text{ for all } i = 1, \dots, M. \quad (\text{A.6})$$

By (A.6), (ii<sub>2</sub>), and the optimality of  $G$  we have

$$F(G) \leq F(\hat{G}) \leq F(G) - \sum_{\substack{x, y \in B_{\frac{1}{2}(1-\varepsilon)} \\ \{x, y\} \in E}} v_2(|x - y|) - 2 \sum_{\substack{x \in B_{\frac{1}{2}(1-\varepsilon)}, y \in A_\varepsilon \\ \{x, y\} \in E}} v_2(|x - y|). \quad (\text{A.7})$$

Now, using (A.7), (A.3), and (A.5), we obtain

$$\frac{1}{2\varepsilon} M(M-1) \leq \sum_{\substack{x, y \in B_{\frac{1}{2}(1-\varepsilon)} \\ \{x, y\} \in E}} v_2(|x - y|) \leq -2 \sum_{\substack{x \in B_{\frac{1}{2}(1-\varepsilon)}, y \in A_\varepsilon \\ \{x, y\} \in E}} v_2(|x - y|) \leq 2NM^2.$$

For  $\varepsilon > 0$  small enough ( $\varepsilon < \frac{1}{8N}$  suffices), this inequality can only be true for  $M = 1$ . This yields (A.1) and concludes Step 1.

Step 2. In this step we prove Definition 3.1(ii), i.e., all bond angles satisfy

$$\theta \in [\pi/2 - \varepsilon, \pi/2 + \varepsilon] \cup [\pi - \varepsilon, \pi + \varepsilon] \cup [3\pi/2 - \varepsilon, 3\pi/2 + \varepsilon]. \quad (\text{A.8})$$

In particular, for  $\varepsilon < \frac{1}{10}\pi$ , we then also have  $\#\mathcal{N}(x, E) \leq 4$  for all  $x \in V$  since all bond angles at  $x \in V$  sum up to  $2\pi$ . To see (A.8), we first of all claim that

$$\#\mathcal{N}(x, E) \leq 4 \frac{(\sqrt{2} + \frac{1}{2})^2}{(1-\varepsilon)^2} \text{ for all } x \in V. \quad (\text{A.9})$$

This follows by Step 1 and, due to (ii<sub>2</sub>), by the fact that  $B_{\frac{1}{2}(1-\varepsilon)}(y) \subset B_{\sqrt{2}+\frac{1}{2}}(x)$  for all  $y \in \mathcal{N}(x, E)$ . More precisely,

$$(\sqrt{2} + \frac{1}{2})^2 \pi = |B_{\sqrt{2}+\frac{1}{2}}(x)| \geq \sum_{y \in \mathcal{N}(x, E)} |B_{\frac{1}{2}(1-\varepsilon)}(y)| \geq \frac{1}{4}(1-\varepsilon)^2 \pi \#\mathcal{N}(x, E),$$

i.e., (A.9) holds. Now, (A.8) follows. In fact, if  $x$  has a bond angle that does not satisfy (A.8) we could define  $\hat{V} = (V \setminus \{x\}) \cup \{x + \tau\}$  for some  $\tau \in \mathbb{R}^2$  such that  $\text{dist}(x + \tau, \hat{V} \setminus \{x + \tau\}) \geq \sqrt{2}$ . Then, by (i<sub>2</sub>), (iv<sub>3</sub>), and (A.9) we obtain a contradiction to the minimality of  $G$ . Summarizing, with the choice  $\varepsilon_0 := \min\{\frac{1}{10}\pi, \frac{1}{8N}\}$ , the statement holds.  $\square$

#### APPENDIX B. PROOF OF ISOPERIMETRIC INEQUALITIES IN $l^1$ VIA SLICING

In this short excursion, we show how isoperimetric inequalities with respect to the  $l^1$ -perimeter can be obtained by a slicing argument similar to the one used in case (a) of the proof of Theorem 2.1. Indeed, our main proof was inspired by such an argument. We present it directly in any space dimension  $d \geq 1$ . First, given  $m \in \mathbb{N}$  and  $x \in \mathbb{R}^m$  we denote by  $|x|_1 = \sum_{k=1}^m |x_k|$  its  $l^1$ -norm. Now, for a set of finite perimeter  $E \subset \mathbb{R}^d$ , see [31], we introduce the  $l^1$ -perimeter by

$$P_{l^1}^d(E) = \int_{\partial^* E} |\nu_E|_1 \, d\mathcal{H}^{d-1},$$

where  $\nu_E$  denotes the measure theoretical outer normal of  $\partial^* E$ .

**Theorem B.1.** *For each set of finite perimeter  $E \subset \mathbb{R}^d$ ,  $d \geq 1$ , it holds that*

$$P_{l^1}^d(E) \geq 2d(\mathcal{L}^d(E))^{1-1/d}.$$

A proof of this result via slicing hinges on the following lemma, for which we use the notation

$$E_t = E \cap \{(x', t) : x' \in \mathbb{R}^{d-1}\} \quad \text{for all } t \in \mathbb{R}. \quad (\text{B.1})$$

**Lemma B.2.** *Suppose that  $E \subset \mathbb{R}^d$  is a bounded set of finite perimeter. Then,*

$$P_{l^1}^d(E) \geq \int_{\mathbb{R}} P_{l^1}^{d-1}(E_t) \, dt + 2 \sup_{t \in \mathbb{R}} \mathcal{H}^{d-1}(E_t). \quad (\text{B.2})$$

We postpone the proof of the lemma to the end.

*Proof of Theorem B.1.* We prove the statement by induction. The case  $d = 1$  is clear as each set  $E \subset \mathbb{R}$  with finite volume satisfies  $P_{l^1}^1(E) \geq 2$ . Suppose that the statement holds for  $d - 1$  and consider  $E \subset \mathbb{R}^d$ . Then, by Lemma B.2 and the induction hypothesis we have

$$\begin{aligned} P_{l^1}^d(E) &\geq \int_{\mathbb{R}} P_{l^1}^{d-1}(E_t) \, dt + 2 \sup_{t \in \mathbb{R}} \mathcal{H}^{d-1}(E_t) \\ &\geq \int_{\mathbb{R}} 2(d-1)(\mathcal{H}^{d-1}(E_t))^{1-1/(d-1)} \, dt + 2 \sup_{t \in \mathbb{R}} \mathcal{H}^{d-1}(E_t). \end{aligned}$$

Using the shorthand  $M := \sup_{t \in \mathbb{R}} \mathcal{H}^{d-1}(E_t)$  and integrating over the slices  $E_t$  we get

$$P_{l^1}^d(E) \geq \int_{\mathbb{R}} 2(d-1)M^{-1/(d-1)} \mathcal{H}^{d-1}(E_t) \, dt + 2M = 2(d-1)M^{-1/(d-1)} \mathcal{L}^d(E) + 2M.$$

By optimizing with respect to  $M$  we get  $M = (\mathcal{L}^d(E))^{1-1/d}$ , and thus we conclude

$$P_{l^1}^d(E) \geq 2d(\mathcal{L}^d(E))^{1-1/d}.$$

$\square$

*Proof of Lemma B.2.* We start by splitting the  $l^1$ -perimeter into

$$P_{l^1}^d(E) = \int_{\partial^* E} |\nu_E|_1 \, d\mathcal{H}^{d-1} = \int_{\partial^* E} (|\nu'_E|_1 + |(\nu_E)_d|) \, d\mathcal{H}^{d-1}, \quad (\text{B.3})$$

where  $\nu'_E = ((\nu_E)_1, \dots, (\nu_E)_{d-1}) \in \mathbb{R}^{d-1}$ . Introducing the function

$$\bar{g} = \frac{|\nu'_E|_1}{\sqrt{1 - |(\nu_E)_d|^2}} \quad \text{on } \partial^* E,$$

the coarea formula, see [31, (18.25)], implies

$$\int_{\partial^* E} |\nu'_E|_1 \, d\mathcal{H}^{d-1} = \int_{\partial^* E} \bar{g} \sqrt{1 - (\nu_E)_d^2} \, d\mathcal{H}^{d-1} = \int_{\mathbb{R}} \int_{(\partial^* E)_t} \bar{g} \, d\mathcal{H}^{d-2} \, dt = \int_{\mathbb{R}} P_{t^1}^{d-1}(E_t) \, dt, \quad (\text{B.4})$$

where in the last step we used the fact that  $\frac{1}{\sqrt{1 - (\nu_E)_d^2}} \nu'_E \in \mathbb{R}^{d-1}$  is a unit normal to  $E_t$ . On the other hand, using the notation  $(\partial^* E)^{x'} := \partial^* E \cap \{(x', t) : t \in \mathbb{R}\}$  for  $x' \in \mathbb{R}^{d-1}$ , by slicing properties of  $BV$ -functions, we obtain

$$\int_{\partial^* E} |(\nu_E)_d| \, d\mathcal{H}^{d-1} = \int_{\mathbb{R}^{d-1}} \mathcal{H}^0((\partial^* E)^{x'}) \, d\mathcal{H}^{d-1}(x') \geq 2 \sup_{t \in \mathbb{R}} \mathcal{H}^{d-1}(E_t), \quad (\text{B.5})$$

where we used that for  $\mathcal{H}^1$ -a.e.  $t \in \mathbb{R}$  and  $\mathcal{H}^{d-1}$ -a.e.  $x'$  with  $(x', t) \in E_t$  we have  $\mathcal{H}^0((\partial^* E)^{x'}) \geq 2$ . Combining the two estimates (B.4)–(B.5), the desired results follows from (B.3).  $\square$

#### REFERENCES

- [1] N.L. ALLINGER. *Molecular structure: understanding steric and electronic effects from molecular mechanics*. John Wiley & Sons (2010).
- [2] Y. AU YEUNG, G. FRIESECKE, B. SCHMIDT. *Minimizing atomic configurations of short range pair potentials in two dimensions: crystallization in the Wulff-shape*. Calc. Var. Partial Differential Equations **44** (2012), 81–100.
- [3] L. BÉTERMIN, L. DE LUCA, M. PETRACHE. *Crystallization to the square lattice for a two-body potential*. Arch. Ration. Mech. Anal. **240** (2021), 987–1053.
- [4] L. BÉTERMIN, H. KNÜPFER, F. NOLTE. *Crystallization of one-dimensional alternating two-component systems*. J. Stat. Phys. **2** (2020), 803–815.
- [5] X. BLANC, M. LEWIN. *The crystallization conjecture: a review*. EMS Surv. Math. Sci. **2** (2015), 255–306.
- [6] B. BOLLOBAS, I. LEADER. *Edge-isoperimetric inequalities in the grid*. Combinatorica **11** (1991), 299–314.
- [7] D. C. BRYDGES, P. A. MARTIN. *Coulomb systems at low density: A review*. J. Stat. Phys. **96** (1999), 1163–1330.
- [8] M. CICALESE, G. P. LEONARDI. *Maximal fluctuations on periodic lattices: an approach via quantitative Wulff inequalities*. Comm. Math. Phys. **375** (2020), 1931–1944.
- [9] E. DAVOLI, P. PIOVANO, U. STEFANELLI. *Wulff shape emergence in graphene*. Math. Models Methods Appl. Sci. (M3AS) **26** (2016), 12:2277–2310.
- [10] E. DAVOLI, P. PIOVANO, U. STEFANELLI. *Sharp  $n^{3/4}$  law for the minimizers of the edge-isoperimetric problem in the triangular lattice*. J. Nonlin. Sci. **27** (2017), 627–660.
- [11] P. DUNCAN, R. O'DWYER, E. B. PROCACCIA. *An elementary proof for the double bubble problem in  $\ell^1$  norm*. Preprint 2020, arXiv:2008.07767.
- [12] P. DUNCAN, R. O'DWYER, E. B. PROCACCIA. *Discrete  $\ell^1$  Double Bubble solution is at most ceiling  $+2$  of the continuous solution*. Preprint 2021, arXiv:2109.11879.
- [13] L. DE LUCA, G. FRIESECKE. *Crystallization in two dimensions and a discrete GaussBonnet Theorem*. J. Nonlinear Sci. **28** (2017), 69–90.
- [14] W. E, D. LI. *On the crystallization of 2D hexagonal lattices*. Comm. Math. Phys. **286** (2009), 1099–1140.
- [15] B. FARMER, S. ESEDOĞLU, P. SMERKA. *Crystallization for a Brenner-like potential*. Comm. Math. Phys. **349** (2017), 1029–1061.
- [16] A. FIGALLI, F. MAGGI, A. PRATELLI. *A mass transportation approach to quantitative isoperimetric inequalities*. Inventiones mathematicae **182** (2010), 167–211.
- [17] L. FLATLEY, M. TAYLOR, A. TARASOV, F. THEIL. *Packing twelve spherical caps to maximize tangencies*. J. Comput. Appl. Math. **254** (2013), 220–225.
- [18] L. FLATLEY, F. THEIL. *Face-centered cubic crystallization of atomistic configurations*. Arch. Ration. Mech. Anal. **218** (2015), 363–416.
- [19] M. FRIEDRICH, W. GÓRNY, U. STEFANELLI. *The double-bubble problem on the square lattice*. Preprint 2021, arXiv:2109.01697.
- [20] M. FRIEDRICH, L. KREUTZ. *Crystallization in the hexagonal lattice for ionic dimers*. Math. Models Methods Appl. Sci. (M3AS) **29** (2019), 1853–1900.



- [21] M. FRIEDRICH, L. KREUTZ. *Finite crystallization and Wulff shape emergence for ionic compounds in the square lattice*. Nonlinearity **33** (2020), 1240–1296.
- [22] M. FRIEDRICH, L. KREUTZ, B. SCHMIDT. *Emergence of rigid polycrystals from atomistic systems with Heitmann-Radin sticky disk energy*. Arch. Ration. Mech. Anal. **240** (2021), 627–698.
- [23] M. FRIEDRICH, U. STEFANELLI. *Crystallization in a one-dimensional periodic landscape*. J. Stat. Phys. (2020), **179** (2020), 485–501.
- [24] G. FRIESECKE, F. THEIL. *Molecular geometry optimization, models*. In the Encyclopedia of Applied and Computational Mathematics, B. Engquist (Ed.), Springer, 2015.
- [25] C. S. GARDNER, C. RADIN. *The infinite-volume ground state of the Lennard-Jones potential*. J. Stat. Phys. **20** (1979), 719–724.
- [26] G. C. HAMRICK, C. RADIN. *The symmetry of ground states under perturbation*. J. Stat. Phys. **21** (1979), 601–607.
- [27] F. HARARY, H. HARBORTH. *Extremal animals*. J. Combin. Inform. and System Sci. **1** (1976), 1–8.
- [28] H. HARBORTH. *Solution to Problem 664A*. Elem. Math. **29** (1974), 14–15.
- [29] R. HEITMAN, C. RADIN. *Ground states for sticky disks*. J. Stat. Phys. **22** (1980), 281–287.
- [30] E. G. LEWARS. *Computational Chemistry*. 2nd edition, Springer, 2011.
- [31] F. MAGGI. *Sets of finite perimeter and geometric variational problems: an introduction to Geometric Measure Theory 135*, Cambridge University Press, 2012.
- [32] E. MAININI, P. PIOVANO, B. SCHMIDT, U. STEFANELLI.  *$N^{3/4}$  law in the cubic lattice*. J. Stat. Phys. **176** (2019), 1480–1499.
- [33] E. MAININI, P. PIOVANO, U. STEFANELLI. *Finite crystallization in the square lattice*. Nonlinearity **27** (2014), 717–737.
- [34] E. MAININI, B. SCHMIDT. *Maximal Fluctuations Around the Wulff Shape for Edge-Isoperimetric Sets in  $\mathbb{Z}^d$ : A Sharp Scaling Law*. Comm. Math. Phys. **380** (2020), 947–971.
- [35] E. MAININI, U. STEFANELLI. *Crystallization in carbon nanostructures*. Comm. Math. Phys. **328** (2014), 545–571.
- [36] C. RADIN. *The ground state for soft disks*. J. Stat. Phys. **26** (1981), 365–373.
- [37] C. RADIN. *Classical ground states in one dimension*. J. Stat. Phys. **35** (1983), 109–117.
- [38] C. RADIN. *Crystals and quasicrystals: a continuum model*. Comm. Math. Phys. **105** (1986), 385–390.
- [39] B. SCHMIDT. *Ground states of the 2D sticky disc model: fine properties and  $N^{3/4}$  law for the deviation from the asymptotic Wulff-shape*. J. Stat. Phys. **153** (2013), 727–738.
- [40] F. THEIL. *A proof of crystallization in two dimensions*. Comm. Math. Phys. **262** (2006), 209–236.
- [41] W. J. VENTEVOGEL, B. R. A. NIJBOER. *On the configuration of systems of interacting atom with minimum potential energy per particle*. Phys. A **98** (1979), 274–288.
- [42] H. J. WAGNER. *Crystallinity in two dimensions: a note on a paper of C. Radin*. J. Stat. Phys. **33**, (1983), 523–526.

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