A review of Lorentzian synthetic theory of timelike Ricci curvature bounds

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April 28, 2022

Abstract

The scope of this survey is to give a self-contained introduction to synthetic timelike Ricci curvature bounds for (possibly non-smooth) Lorentzian spaces via optimal transport & entropy tools, including a synthetic version of Hawking's singularity theorem and a synthetic characterisation of Einstein's vacuum equations. We will also discuss some motivations arising from the smooth world and some possible directions for future research.

Contents

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Introduction

Optimal transport revealed to be a very effective tool to study spaces with extremely low regularity in Riemannian signature. More precisely, one can use optimal transport to define and study metric measure spaces with Ricci $\geq K \in \mathbb{R}$ and dimension $\leq N \in [1, \infty)$ in a synthetic sense; this is indeed the celebrated theory of $\mathsf{CD}(K, N)$ spaces pioneered by Sturm [86, 87] and Lott-Villani [65] independently.

Motivated by the success of optimal transport techniques in the Riemannian signature, during the last few years the development of optimal transport tools in the Lorentzian setting has become an increasingly popular topic of research, let us give a brief account.

The optimal transport problem for relativistic costs was proposed by Brenier [18], further investigated in [13, 14], and pushed to a geometric Lorentzian setting in [30, 89, 54]. An intriguing physical motivation for studying optimal transport in a Lorentzian setting is the *early universe reconstruction problem*. Studied in [19] and in the Nature paper [35] with methods of optimal transportation, such a problem amounts to reconstructing the trajectories of masses from the big bang to their present day positions in Robertson-Walker spacetimes. A mathematical formulation for general globally hyperbolic spacetimes would read as follows: given two probability measures, one concentrated on a Cauchy hypersurface, the other on the past cone of a point, what can be said about the trajectories of the minimizers in a dynamical optimal coupling of the two measure? In [35] it is explained why the problem can be attacked with methods of optimal transportation.

Due to the central role played by the Ricci curvature in general relativity (e.g. the Einstein's equations can be formulated just in terms of the Ricci tensor, the energy momentum tensor, and the cosmological constant) and due to the success of optimal transport tools to study Ricci curvature bounds in the Riemannian signature, there has been a growing interest in the interplay of optimal transport and Ricci curvature bounds in a Lorentzian setting. Indeed, for a smooth Lorentzian manifold (see Section 1.5), McCann [71] characterised non-negative timelike Ricci curvature lower bounds in terms of dispacement convexity, and the second author joint with Suhr [75] characterised any lower and upper bound on the timelike Ricci curvature (and thus the full Einstein equations) in terms of optimal transport and entropy.

While the non-smooth (or "synthetic") counterpart of a Riemannian manifold is a metric (measure) space, and thus the latter gives a natural framework for Lott-Sturm-Villani theory, the interplay of causal and metric structures in a Lorentzian manifold is more subtle and forces to modify the setting in order to develop a synthetic theory. Such a program was pioneered in seminal work by Kronheimer-Penrose [57] and has been recently formalised by Kunzinger-Sämann [58] into the concept of Lorentzian (pre-)length spaces (see Section 1.1). Lorentzian (pre)-length spaces give a natural framework to develop a theory of optimal transport and of synthetic curvature bounds.

Synthetic *timelike sectional curvature bounds* for Lorentzian (pre-)length spaces have been investigated in [58] (see also the more recent [10]). Synthetic timelike Ricci curvature lower bounds in Lorentzian (pre-)lenght spaces, which constitute the subject of the present survey, have been developed by the authors in [24] (inspired by the aforementioned smooth characterisations obtained in [71, 75]). The aim of [24] has been to initiate a theory of Lorentzian CD spaces, denoted by TCD for "timelike CD" (see Section 2), and to obtain some first geometric and functional insights; for instance the following inequalities are among the consequences of the TCD condition: volume comparison, diameter bound and Poincaré inequality. Moreover, a far reaching synthetic version of Hawking singularity theorem holds (in a sharp form, and in a formulation which includes Lorentzian metrics of very low regularity: $C^{0,1}$ or, more generally, C^0 plus causally plain). We refer to Section 3 for an overview of (some of) the applications.

In Section 4 we survey synthetic timelike Ricci upper bounds and a synthetic formulation of Ein-

stein's vacuum equations for Lorentzian (pre)-length spaces, obtained in [75, App. B] after [88, 24]. We will end the review by proposing some possible directions for future research, see Section 5.

Let us conclude the introduction by stressing that, apart from the specific synthetic framework and the aforementioned results, we expect optimal transport tools to be useful in a wide range of applications. Indeed, typically, classical arguments make use of Jacobi-fields computations (e.g. Raychaudhuri equation) heavily relying on the C^2-regularity of the metric (with some effort, one can lower to $C^{1,1}$ by approximation). The main advantage of the optimal transport tools surveyed here (see [24] for more results and details) is that they often allow to carry over arguments in a setting of very low regularity (including $C^{0,1}$ or, more generally, C^0 and causally plain) where the classical terms are even not well defined.

Motivations for studying Lorentzian metrics of low regularity (or lowering the regularity even to a Lorentzian synthetic framework) come both from the PDE point of view in general relativity (i.e. the Cauchy initial value problem for the Einstein equations) and from physically relevant models.

From the PDE point of view, the standard local existence results for the vacuum Einstein equations assume the metric to be of Sobolev regularity H^s_{loc}, with $s > \frac{5}{2}$ (see for instance [81]). The Sobolev regularity of the metric has been lowered even further (e.g. [56]). Related to the initial value problem for the Einstein equations, one of the main open problems in the field is the so called (weak/strong) censorship conjecture proposed originally by Penrose and later refined in [26] (see also [28] for a more updated overview). The strong form of such a conjecture asserts (roughly) that the maximal globally hyperbolic development of generic initial data for the Einstein equations is inextendible as a suitably regular Lorentzian manifold. Formulating a precise statement of the conjecture is itself non-trivial since one needs to give a precise meaning to "generic initial data" and "suitably regular Lorentzian manifold". Understanding the latter is where Lorentzian metrics of low regularity and related inextendibility results play a pivotal role. The strongest form of the conjecture would prove inextendibility for a C^0 metric. As pointed out by Chrusciel-Grant [27] (see also [64] for a survey), causality theory for C^0 metrics departs significantly from classical theory (e.g. the light-like curves emanating from a point may span a set with non-empty interior, a phenomenon called "bubbling"). Nevertheless, Sbierski [84] gave a clever proof of C^0-inextendibility of Schwarzschild, Minguzzi-Suhr [74] showed C^0-inextendibility for timelike geodesically complete spacetimes, and Grant-Kunzinger-Sämann [41] pushed the inextendibility to Lorentzian length spaces.

As already envisaged in seminal work by Lichnerowitz in the 50'ies [63], from the point of view of physically relevant models, several types of matter in a spacetime may give a discontinuous energy-momentum tensor and thus, via the Einstein's equations, lead to a Lorentzian metric of regularity lower than C^2. Examples of such a phenomenon are spacetimes that model the inside and outside of a star, matched spacetimes [68], self-gravitating compressible fluids [21] or shock waves. Some physically relevant models require even lower regularity, for instance: spacetimes with conical singularities [92], cosmic strings [91] and (impulsive) gravitational waves (see for instance the pioneering work of Penrose [80], and the more recent [44, Chapt. 20]), or models for cyclic spacetimes [62].

Finally, a long term motivation for studying non-regular Lorentzian spaces is the desire of understanding the ultimate nature of spacetime. The rough picture is that at the quantum level (and thus in extreme physical conditions such as gravitational collapse or the origin of the universe), the spacetime may be very singular and possibly not approximable by smooth structures (see the end of Section 1.1).

In case of a metric of low regularity, the approach to curvature used so far is distributional, exploiting the smoothness of the underline spacetime. This allows [37] (see also [85]) to define distributional curvature tensors for $W^{1,2}_{loc}$-Lorentzian metrics satisfying a suitable non-degeneracy condition (fulfilled for instance when the metric is C^1, see [40]). One of the goals of the approach surveyed here is to address the question of (timelike Ricci) curvature when *not only the the metric tensor, but the space itself is singular.*

1 Optimal Transport in Lorentzian spaces

1.1 Classical setting and Lorentzian length spaces

We start by introducing some basic terminology.

Given a Lorentzian manifold (M^n, g), we also fix an auxiliary complete Riemannian metric h on M and we denote by d^h the associated distance. We say that $v \in T_x M \setminus \{0\}$ is *timelike* if $g(v, v) < 0$. (M, g) is *time-oriented* if it admits a continuous timelike vector field Y. Such a Y induces a partition on the set of timelike vectors: the equivalence class of *future pointing* vectors $\{v \in T_x M \mid g(v, Y) < 0\}$ and the class of *past pointing* vectors $\{v \in T_x M \mid g(v, Y) > 0\}$.

The closure of the set of future pointing timelike vectors is called *future causal cone* and denoted by $\mathcal{C} \subset TM$. Let $I = [0, 1]$ denote the unit interval. A locally Lipshitz curve $\gamma : I \to M$ is a *causal* if $\dot\gamma \in \mathcal{C}$ for a.e. t; *timelike* if $\dot\gamma \in \mathrm{Int}(\mathcal{C})$ for a.e. t.

These definitions induce two causal relations for $x, y \in M$:

- we say that $x \leq y$ if $\exists\, \gamma$ causal curve with $\gamma_0 = x$ and $\gamma_1 = y$;

- we say that $x \ll y$ if $\exists\, \gamma$ timelike curve with $\gamma_0 = x$ and $\gamma_1 = y$.

The *length* of a causal curve $\gamma : I \to M$ is defined as

$$\mathrm{L}_g(\gamma) := \int_I \sqrt{|g(\dot\gamma, \dot\gamma)|}\,\mathrm{d}t.$$

The *time separation* function $\tau : M \times M \to [0, \infty]$ is defined as

$$\tau(x, y) := \begin{cases} \sup\{\mathrm{L}_g(\gamma) : \gamma \text{ causal},\ \gamma_0 = x,\ \gamma_1 = y\}, & \text{if } x \leq y, \\ 0, & \text{otherwise.} \end{cases} \qquad (1.1)$$

The following reverse triangular inequality holds true for τ:

$$\tau(x, y) + \tau(y, z) \leq \tau(x, z), \quad \text{if } x \leq y \leq z.$$

If g is $C^{0,1}$-regular (C^0 is not enough, but C^0 and causally plain suffices [58] after [27]): τ is lower-semicontinuous and $\tau(x, y) > 0$ if and only if $x \ll y$.

The idea to formulate a synthetic theory for Lorentzian spaces is to turn the previous properties into definitions. Such a point of view has roots in the seminal work of Kronheimer-Penrose [57] and has been formalised in recent work by Kunzinger-Sämann [58].

Definition 1.1 (Causal space $(X, \ll, \leq)$). A *causal space* $(X, \ll, \leq)$ is a set X endowed with a preorder $\leq$ and a transitive relation $\ll$ contained in $\leq$.

One says that x and y are *timelike* (resp. *causally*) related if $x \ll y$ (resp. $x \leq y$). Let $A \subset X$ be an arbitrary subset of X. We define the *chronological* (resp. *causal*) future of A the set

$$I^+(A) := \{y \in X \mid \exists x \in A,\ x \ll y\}, \quad J^+(A) := \{y \in X \mid \exists x \in A,\ x \leq y\}$$

respectively. Analogously, one can define the *chronological* (resp. *causal*) past of A. In order to keep the notation short, it is also useful to set

$$X_{\leq}^2 := \{(x, y) \in X^2 \mid x \leq y\}, \quad X_{\ll}^2 := \{(x, y) \in X^2 \mid x \ll y\}.$$

Additionally requiring the existence of a time separation function produces the following notion that has to be understood like the Lorentzian analog of metric spaces.

Definition 1.2 (Lorentzian pre-length space $(X, \mathsf{d}, \ll, \leq, \tau)$). A *Lorentzian pre-length space* $(X, \mathsf{d}, \ll, \leq, \tau)$ is a casual space $(X, \ll, \leq)$ additionally equipped with a proper metric d (i.e. closed and bounded subsets are compact) and a lower semicontinuous function $\tau : X \times X \to [0, \infty]$, called *time-separation function*, satisfying

$$\tau(x, y) + \tau(y, z) \leq \tau(x, z) \quad \forall x \leq y \leq z \quad \text{reverse triangle inequality}$$
$$\tau(x, y) = 0, \text{ if } x \not\leq y, \quad \tau(x, y) > 0 \Leftrightarrow x \ll y. \tag{1.2}$$

Notice that X is endowed with the metric topology induced by d and all the topological concepts on X (like l.s.c. of τ) are formulated in terms of such metric topology. Note that the lower semicontinuity of τ implies that $I^{\pm}(x)$ is open, for any $x \in X$.

Accordingly to the causal relations, $\gamma : I \to X$ is said *timelike* (resp. *causal*) if it is locally Lipschitz (w.r.t. d) and $\gamma_{t_1} \ll \gamma_{t_2}$ (resp. $\gamma_{t_1} \leq \gamma_{t_2}$) for all $t_1 < t_2$. The length induced by τ for a causal curve γ is then defined as:

$$\mathrm{L}_{\tau}(\gamma) := \inf \sum_i \tau(\gamma_{t_i}, \gamma_{t_{i+1}}),$$

where the inf is taken over all finite partitions $0 = t_0 < t_1 < \ldots < t_m = 1$, $m \in \mathbb{N}$, of $I = [0, 1]$.
A causal curve $\gamma : [a, b] \to X$ is a *geodesic* if $\mathrm{L}_{\tau}(\gamma) = \tau(\gamma_a, \gamma_b)$.
We denote the set of causal (resp. timelike) geodesics as:

$$\mathrm{Geo}(X) := \{\gamma : [0, 1] \to X \, : \, \tau(\gamma_s, \gamma_t) = (t - s)\tau(\gamma_0, \gamma_1) \, \forall s < t\}, \tag{1.3}$$
$$\mathrm{TGeo}(X) := \{\gamma \in \mathrm{Geo}(X) \, : \, \tau(\gamma_0, \gamma_1) > 0\}. \tag{1.4}$$

Note that, by definition, geodesics are always *maximising, future oriented,* and *parametrised at constant speed on* $[0, 1]$.

Moreover a Lorentzian pre-length space $(X, \mathsf{d}, \ll, \leq, \tau)$ is called:

- *Causally closed:* $\{x \leq y\} \subset X \times X$ is closed

- *Non-totally imprisoning* if $\forall K \subset X$ compact, $\exists C > 0$ s.t. $\mathrm{L}_{\mathsf{d}}(\gamma) \leq C$ for all γ causal curve in K.

- *Globally hyperbolic* if it is non-totally imprisoning and $\forall x, y \in X$, $J^+(x) \cap J^-(y)$ is compact.

- $\mathcal{K}$-*Globally hyperbolic:* if it is non-totally imprisoning and $\forall K_1, K_2 \subset X$ compact, $J^+(K_1) \cap J^-(K_2)$ is compact.

- *Geodesic*: if $\forall x, y \in X$ with $x \leq y$ there exists a geodesic γ from x to y.

It was proved in [58, Theorem 3.28] that for a globally hyperbolic Lorentzian geodesic (actually length would suffice) space $(X, \mathsf{d}, \ll, \leq, \tau)$, the time-separation function τ is finite and continuous. Moreover, any globally hyperbolic Lorentzian length space (for the definition of Lorenzian length space see [58, Definition 3.22], we skip it for brevity since we will not use it) is geodesic [58, Theorem 3.30].

Examples entering the class of Lorentzian synthetic spaces

Spacetimes with a continuous Lorentzian metric. Let M be a smooth manifold endowed with a continuous Lorentzian metric g. Assume that (M, g) is time-oriented. Observe that, for C^0-metrics, the natural class of differentiability of the underlying manifolds is C^1; now, C^1 manifolds always admit a C^∞ subatlas, and one can pick such sub-atlas whenever convenient. We endow such a spacetime (M, g) with the time separation function τ defined in (1.1).

For a spacetime with a Lorentzian C^0-metric:

- Global hyperbolicity implies causal closedness and $\mathcal{K}$-global hyperbolicity [83, Prop. 3.3 and Cor. 3.4].

- Recall that a Cauchy hypersurface is a subset which intersects exactly once every inextendible causal curve. It was proved in [83] that every Cauchy hypersurface is a closed acasual topological hypersurface and that global hyperbolicity is equivalent to the existence of a Cauchy hypersurface.

- If g is a causally plain (or, more strongly, locally Lipschitz) Lorentzian C^0-metric on M then the associated synthetic structure is a pre-length Lorentzian space, see [58, Prop. 5.8]. More strongly, if g is a globally hyperbolic and causally plain Lorentzian C^0-metric on M then the associated synthetic structure is a causally closed, $\mathcal{K}$-globally hyperbolic Lorentzian geodesic space (see [58, Thm. 3.30 and Thm. 5.12]).

Summarising, if M is a smooth manifold endowed with a C^0-Lorentzian metric g making (M, g) time-oriented then $(M, \ll, \leq)$ is a causal space but it is not necessarily a Lorentzian pre-length space. If g is "causally plain" [27] (for instance $g \in C^{0,1}$), then $(M, \mathsf{d}^h, \ll, \leq, \tau)$ is a Lorentzian pre-length space. If in addition g is globally hyperbolic, then $(M, \mathsf{d}^h, \ll, \leq, \tau)$ is causally closed, $\mathcal{K}$-globally hyperbolic and geodesic Lorentzian space.

Thus, the framework of causally closed, $\mathcal{K}$-globally hyperbolic and geodesic Lorentzian spaces is rather natural and will provide the setting of our work.

Closed cone structures. Closed cone structures can be seen as the synthetic-Lorentzian analogue of Finsler manifolds. They provide a rich source of examples of Lorentzian pre-length and length spaces (see [58, Sec. 5.2] for more details). We refer to Minguzzi's comprehensive review paper [73] for a thorough analysis of causality theory in the framework of closed cone structures, including embedding and singularity theorems.

Some examples towards quantum gravity. The general setting of Lorentzian pre-length spaces allows to consider more general structures than Lorentz(-Finsler) metrics on smooth manifolds. A remarkable motivation for such a general framework is given by certain approaches to quantum gravity. For instance let us mention [69] where it is shown that, in a purely order theoretic manner, one can reconstruct a globally hyperbolic spacetime and the causality relation from a countable dense set of events.

Two approaches to quantum gravity, particularly close in spirit to Lorentzian pre-length spaces, are the theory of *causal Fermion systems* [33, 34] and the *theory of causal sets* [17]. The basic idea in both the approaches is that the structure of spacetime needs to be relaxed on a microscopic scale to include quantum effects. This leads to non-smoothness of the underlying geometry, and the classical structure of a Lorentzian spacetime emerges only in the macroscopic regime. For the connection to the theory of Lorentzian (pre-)length spaces, the reader is referred to [58, Sec. 5.3] and [33, Sec. 5.1].

1.2 Optimal transport in Lorentzian spaces

The space of Borel probability measures (resp. with compact support) over a metric space (X, d) is denoted by $\mathcal{P}(X)$ (resp. $\mathcal{P}_c(X)$). For $\mu, \nu \in \mathcal{P}(X)$, we denote by $\mu \otimes \nu \in \mathcal{P}(X^2)$ the unique probability measure such that

$$\mu \otimes \nu(A \times B) := \mu(A) \cdot \nu(B), \quad \text{for all } A, B \subset X \text{ Borel subsets.} \tag{1.5}$$

Note that (1.5) uniquely defines a probability measure on X^2, as the σ-algebra of Borel sets of X^2 is generated by products of Borel sets of X.

Let $P_i : X \times X \to X$, $i = 1, 2$, denote the projection maps on the factors. For a Borel map $f : X \to Y$, the associated push-forward map $f_\sharp : \mathcal{P}(X) \to \mathcal{P}(Y)$ is defined by

$$f_\sharp \mu(B) := \mu(f^{-1}(B)), \quad \text{for all } \mu \in \mathcal{P}(X) \text{ and } B \subset Y \text{ Borel subset.}$$

Notice that $(P_1)_\sharp(\mu \otimes \nu) = \mu$ and $(P_2)_\sharp(\mu \otimes \nu) = \nu$.

Definition 1.3 (Transport plans)**.** If $(X, \mathsf{d}, \ll, \leq, \tau)$ is a Lorentzian pre-length space and $\mu, \nu \in \mathcal{P}(X)$, we can distinguish different families of transport plans:

- Classic: $\Pi(\mu, \nu) := \{\pi \in \mathcal{P}(X^2) \mid (P_1)_\sharp \pi = \mu, (P_2)_\sharp \pi = \nu\}$.

- Causal: $\Pi_\leq(\mu, \nu) := \{\pi \in \Pi(\mu, \nu) \mid \pi(X_\leq^2) = 1\}$, where $X_\leq^2 := X \times X \cap \leq$,

- Timelike: $\Pi_\ll(\mu, \nu) := \{\pi \in \Pi(\mu, \nu) \mid \pi(X_\ll^2) = 1\}$, where $X_\ll^2 := X \times X \cap \ll$.

Some intuition behind the notion of transport plan. Since the survey is meant for an audience possibly non specialised in optimal transport, let us discuss some basic heuristics.

The rough idea of a classic transport plan $\pi \in \Pi(\mu, \nu)$ is that the "mass $\mu(\mathrm{d}x)$ at x" is possibly split and transported following the law $\pi(\mathrm{d}x\mathrm{d}y)$. Note that if $f : X \to X$ is a Borel map, then $\pi := (\mathrm{Id}, f)_\sharp \mu \in \Pi(\mu, f_\sharp \mu)$ is a classic transport plan from μ to $f_\sharp \mu$. However this is a very special case of a transport plan, namely a transport plan induced by a map (called transport map). A (trivial) example of a transport plan *not* induced by a map is $\pi := \mu \otimes \nu$ obtained by taking the product of μ and ν as in (1.5). It is also instructive to consider the case of $\mu = \delta_{x_0}, \nu = \frac{1}{2}(\delta_{x_1} + \delta_{x_2})$, where δ_x denotes the Dirac mass at $x \in X$: in this case it is not possible to find a trasport map, and a transport plan is given by the product $\mu \otimes \nu$, which corresponds to "split" the Dirac mass at x_0 in half and transport each half into x_1 and x_2 respectively.

The rough idea of a causal (resp. timelike) transport plan $\pi \in \Pi_\leq(\mu, \nu)$ is that the "mass $\mu(\mathrm{d}x)$ at x" is possibly split and transported following the law $\pi(\mathrm{d}x\mathrm{d}y)$ so that "the destination is in the causal (resp. timelike) future of the source". Clearly, for $\mu = \delta_{x_0}$ and $\nu = \delta_{x_1}$, it holds that $\pi = \mu \otimes \nu \in \Pi_\leq(\mu, \nu)$ (resp. $\Pi_\ll(\mu, \nu)$) if and only if $x_0 \leq x_1$ (resp. $x_0 \ll x_1$). $\qquad \square$

Definition 1.4. Let $(X, \mathsf{d}, \ll, \leq, \tau)$ be a Lorentzian pre-length space and let $p \in (0, 1]$. Given $\mu, \nu \in \mathcal{P}(X)$, the *p-Lorentz-Wasserstein distance* is defined by

$$\ell_p(\mu, \nu) := \sup_{\pi \in \Pi_\leq(\mu, \nu)} \left(\int_{X \times X} \tau(x, y)^p \, \pi(\mathrm{d}x\mathrm{d}y) \right)^{1/p}. \tag{1.6}$$

When $\Pi_\leq(\mu, \nu) = \emptyset$, we set $\ell_p(\mu, \nu) := -\infty$.

Definition 1.4 extends to Lorentzian pre-length spaces the corresponding notion given in the smooth Lorentzian setting in [30] (see also [71, 75], and [89] for $p = 1$). When $\Pi_\leq(\mu, \nu) = \emptyset$, we adopt the convention of McCann [71] (note that [30] set $\ell_p(\mu, \nu) = 0$ in this case).

The function ℓ_p inherits the properties of the time-separation function τ. Indeed, by using the classical technique of gluing (used for instance to prove the triangular inequality for the Wasserstein distance), one can prove that ℓ_p verifies the reverse triangular inequality:

$$\ell_p(\mu_0, \mu_1) + \ell_p(\mu_1, \mu_2) \leq \ell_p(\mu_0, \mu_2), \quad \forall \, \mu_0, \mu_1, \mu_2 \in \mathcal{P}(X) \tag{1.7}$$

with the convention on the left hand side that $\infty - \infty = -\infty$.

To invoke the classical theory of Optimal Transport it is more convenient to move the causal constraint on the transport plans to the cost function. In other words, it is useful to reformulate the variational problem in the right hand side of (1.6) with the following cost function:

$$\ell(x, y) := \begin{cases} \tau(x, y), & \text{if } x \leq y, \\ -\infty, & \text{otherwise.} \end{cases} \tag{1.8}$$

Notice that

$$\int_{X \times X} \tau(x, y)^p \, \pi(\mathrm{d}x\mathrm{d}y) = \int_{X \times X} \ell(x, y)^p \, \pi(\mathrm{d}x\mathrm{d}y) \in \mathbb{R}_{\geq 0}, \quad \forall \pi \in \Pi_\leq(\mu, \nu).$$

Moreover, if $\pi \in \Pi(\mu, \nu)$ satisfies $\int_{X \times X} \ell(x, y)^p \, \pi(\mathrm{d}x\mathrm{d}y) > -\infty$ then $\pi \in \Pi_{\leq}(\mu, \nu)$. Thus the maximization problem (1.6) is equivalent (i.e. the sup and the set of maximisers coincide) to the maximisation problem

$$\sup_{\pi \in \Pi(\mu, \nu)} \left(\int_{X \times X} \ell(x, y)^p \, \pi(\mathrm{d}x\mathrm{d}y) \right)^{1/p}. \tag{1.9}$$

A $\pi \in \Pi_{\leq}(\mu, \nu)$ maximising in (1.6) is said ℓ_p-*optimal*. The set of ℓ_p-*optimal* plans from μ to ν is denoted by $\Pi_{\leq}^{p\text{-opt}}(\mu, \nu)$.

The advantage of the formulation (1.9) is that, when $(X, \mathsf{d}, \ll, \leq, \tau)$ is causally closed (so that $\{(x, y) \mid x \leq y\} \subset X \times X$ is a closed subset) and globally hyperbolic geodesic (so that τ is continuous) then ℓ is upper semi-continuous on $X \times X$. Similarly, when X is *locally* causally closed globally hyperbolic geodesic, if μ and ν have compact support then ℓ is upper semi-continuous on $\operatorname{supp} \mu \times \operatorname{supp} \nu$.

In both cases, one can apply standard optimal transport techniques (e.g. [93]) to the Monge-Kantorovich problem (1.9).

Proposition 1.5 (Existence of optimal plans). *Let $(X, \mathsf{d}, \ll, \leq, \tau)$ be a causally closed (resp. locally causally closed) globally hyperbolic Lorentzian geodesic space and let $\mu, \nu \in \mathcal{P}(X)$ (resp. $\mathcal{P}_c(X)$). If $\Pi_{\leq}(\mu, \nu) \neq \emptyset$ and if there exist measurable functions $a, b : X \to \mathbb{R}$, with $a \oplus b \in L^1(\mu \otimes \nu)$ such that $\ell^p \leq a \oplus b$ on $\operatorname{supp} \mu \times \operatorname{supp} \nu$ (e.g. when μ and ν are compactly supported) then the $\sup$ in (1.6) (and henceforth in (1.9)) is attained and finite.*

We fix the following notation

$$\Pi_{\leq}^{p\text{-opt}}(\mu, \nu) := \{\pi \in \Pi_{\leq}(\mu, \nu) \text{ is } p\text{-optimal}\},$$
$$\Pi_{\ll}^{p\text{-opt}}(\mu, \nu) := \{\pi \in \Pi_{\ll}(\mu, \nu) \text{ is } p\text{-optimal}\}.$$

1.3 The Kantorovich duality

Existence of optimal plans $\pi \in \Pi_{\leq}(\mu, \nu)$ easily follows by the direct method of Calculus of Variations. The rough picture behind a transport plan is that π is moving the mass $\mu(\mathrm{d}x)$ "at x" in other "suitable points" of $\operatorname{supp} \nu$. A natural question is then: which are such "suitable points" chosen by x? Or, better said, how can we detect $\Gamma \subset X_{\leq}^2$ such that $\pi(\Gamma) = 1$?

These questions can answered via two key notions in optimal transport theory: cyclical monotonicity and Kantorovich duality. Optimal transport problems possess indeed a rich duality theory, firstly discovered by Kantorovich and valid for a large family of cost functions. The dual variational problem permits, among other things, to describe the geometry of the optimizers of the original variational problem (1.9).

A set $\Gamma \subset X^2$ is cyclically monotone with respect to a cost function c if it is c-optimal with respect to perturbation by finitely many points (for the precise definition we refer to [93, Sec. 5]). Kantorovich duality gives a recipe to construct such cyclically monotone sets. The relevance of this condition can be understood by mentioning that, in the smooth Riemannian setting, cyclical monotonicity of a transport plan is equivalent to its optimality.

However, when the cost function is not real valued, for instance in our setting with $\ell(x, y)$ as in (1.8), more attention is required. In particular, Kantorovich duality is more subtle ([14], [54], [71]) and the equivalence between the optimality of a transport plan and the cyclical monotonicity of its support breaks down.

To have a full duality theory we have to consider a more regular family of measures. The following definition relaxes the notion of q-separated introduced by McCann [71, Definition 4.1] in the smooth Lorentzian setting.

In the following, given two functions $a, b : X \to \mathbb{R}$, we denote by $a \oplus b : X \times X \to \mathbb{R}$ the function on the product defined by

$$(a \oplus b)(x, y) := a(x) + b(y), \quad \forall x, y \in X.$$

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Definition 1.6 (Timelike p-dualisable)**.** Let $(X, \mathsf{d}, \ll, \leq, \tau)$ be a Lorentzian pre-length space and let $p \in (0, 1]$. We say that $(\mu, \nu) \in \mathcal{P}(X)^2$ is *timelike p-dualisable (by $\pi \in \Pi_{\ll}(\mu, \nu)$)* if

1. $\ell_p(\mu, \nu) \in (0, \infty)$;

2. $\pi \in \Pi_{\ll}^{p\text{-opt}}(\mu, \nu)$;

3. there exist measurable functions $a, b : X \to \mathbb{R}$, with $a \oplus b \in L^1(\mu \otimes \nu)$ such that $\ell^p \leq a \oplus b$ on $\operatorname{supp} \mu \times \operatorname{supp} \nu$.

If X is globally hyperbolic, (μ, ν) have compact support and admit a timelike p-optimal plan $\pi \in \Pi_{\ll}^{p\text{-opt}}(\mu, \nu)$, then (μ, ν) is timelike p-dualisable by π. In particular, the notion of timelike p-dualisability relaxes the condition of q-separation introduced by McCann [71]. Moreover timelike p-dualisabiliy ensures a weak form of duality.

Proposition 1.7 (Weak Kantorovich duality)**.** *Fix $p \in (0, 1]$. Let $(X, \mathsf{d}, \ll, \leq, \tau)$ be a (resp. locally) causally closed globally hyperbolic Lorentz geodesic space. If $(\mu, \nu) \in \mathcal{P}(X)^2$ (resp. $\mathcal{P}_c(X)^2$) is timelike p-dualisable, then Kantorovich duality holds:*

$$\ell_p(\mu, \nu)^p = \inf \left\{ \int_X u \, \mu + \int_X v \, \nu \right\}, \tag{1.10}$$

where the $\inf$ is taken over all measurable functions $u : \operatorname{supp} \mu \to \mathbb{R} \cup \{+\infty\}$ and $v : \operatorname{supp} \nu \to \mathbb{R} \cup \{+\infty\}$ with $u \oplus v \geq \ell^p$ on $\operatorname{supp} \mu \times \operatorname{supp} \nu$ and $u \oplus v \in L^1(\mu \otimes \nu)$. Furthermore, the value of the right hand side does not change if one restricts the $\inf$ to bounded and continuous functions.

We next discuss the validity of the *strong* Kantorovich duality, i.e. the existence of optimal functions (called Kantorovich potentials) achieving the infimum in the right hand side of (1.10).

Since the dual minimization problem runs over the couple of functions verifying $u \oplus v \geq \ell^p$ on $\operatorname{supp} \mu \times \operatorname{supp} \nu$, one can consider special couple of functions. The following definition is indeed tailored on this duality principle.

Definition 1.8 (ℓ^p-concave functions, ℓ^p-transform and ℓ^p-subdifferential)**.** Fix $p \in (0, 1]$ and let $U, V \subset X$. A measurable function $\varphi : U \to \mathbb{R}$ is *ℓ^p-concave relatively to (U, V)* if there exists a function $\psi : V \to \mathbb{R}$ such that

$$\varphi(x) = \inf_{y \in V} \psi(y) - \ell^p(x, y), \quad \forall x \in U.$$

The function

$$\varphi^{(\ell^p)} : V \to \mathbb{R} \cup \{-\infty\}, \quad \varphi^{(\ell^p)}(y) := \sup_{x \in U} \varphi(x) + \ell^p(x, y) \tag{1.11}$$

is called *ℓ^p-transform of φ*. The *ℓ^p-subdifferential $\partial_{\ell^p} \varphi \subset (\operatorname{supp} \mu \times \operatorname{supp} \nu) \cap X_{\leq}^2$* is defined by

$$\partial_{\ell^p} \varphi := \{(x, y) \in (\operatorname{supp} \mu \times \operatorname{supp} \nu) \cap X_{\leq}^2 \ : \ \varphi^{(\ell^p)}(y) - \varphi(x) = \ell^p(x, y)\}.$$

Then the strong form of duality can be defined as follows.

Definition 1.9 (Strong Kantorovich duality)**.** Fix $p \in (0, 1]$. We say that $(\mu, \nu) \in \mathcal{P}(X)^2$ satisfies *strong ℓ^p-Kantorovich duality* if

1. $\ell_p(\mu, \nu) \in (0, \infty)$;

2. there exists $\varphi : \operatorname{supp} \mu \to \mathbb{R}$ which is ℓ^p-concave relatively to $(\operatorname{supp} \mu, \operatorname{supp} \nu)$ and satisfying

$$\ell_p(\mu, \nu)^p = \int_X \varphi^{(\ell^p)}(y) \, \nu(\mathrm{d}y) - \int_X \varphi(x) \, \mu(\mathrm{d}x).$$

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It is immediate to check that if $(\mu, \nu) \in \mathcal{P}(X)^2$ satisfies strong ℓ^p-Kantorovich duality, then the following holds: a plan $\pi \in \Pi_{\leq}(\mu, \nu)$ is ℓ_p-optimal if and only if

$$\varphi^{(\ell^p)}(y) - \varphi(x) = \ell^p(x, y) = \tau(x, y)^p, \quad \text{for } \pi\text{-a.e. } (x, y),$$

i.e. if and only if $\pi(\partial_{\ell^p}\varphi) = 1$.

Motivated by this remark, we have devised a set of couples of probability measures larger than those satisfying the strong Kantorovich duality (but still sufficiently regular).

Definition 1.10 (Strongly timelike p-dualisable). A pair $(\mu, \nu) \in (\mathcal{P}(X))^2$ is said to be *strongly timelike p-dualisable* if

1. (μ, ν) is timelike p-dualisable;

2. there exists a measurable ℓ^p-cyclically monotone set $\Gamma \subset X_{\ll}^2 \cap (\operatorname{supp}\mu \times \operatorname{supp}\nu)$ such that a plan $\pi \in \Pi_{\leq}(\mu, \nu)$ is ℓ_p-optimal if and only if π is concentrated on Γ, i.e. $\pi(\Gamma) = 1$.

The next two propositions (for their proof see [24, Cor. 2.29, Cor. 2.30]) show that the notion of strongly timelike p-dualisable measures is non-empty: the first one addresses the "local" behaviour.

Proposition 1.11. *Fix $p \in (0, 1]$. Let $(X, \mathsf{d}, \ll, \leq, \tau)$ be a causally closed (resp. locally causally closed) globally hyperbolic Lorentzian geodesic space and assume that $\mu, \nu \in \mathcal{P}(X)$ (resp. $\mathcal{P}_c(X)$) satisfy:*

1. *there exist measurable functions $a, b : X \to \mathbb{R}$ with $a \oplus b \in L^1(\mu \otimes \nu)$ such that $\tau^p \leq a \oplus b$ on $\operatorname{supp}\mu \times \operatorname{supp}\nu$;*

2. *$\operatorname{supp}\mu \times \operatorname{supp}\nu \subset X_{\ll}^2$.*

Then (μ, ν) is strongly timelike p-dualisable.

The second one shows that in the case ν is a Dirac measure, the strongly timelike p-dualisability is equivalent to the timelike p-dualisability.

Proposition 1.12. *Let $(X, \mathsf{d}, \ll, \leq, \tau)$ be a Lorentzian pre-length space and let $p \in (0, 1]$. Fix $\bar{x} \in X$ and let $\nu := \delta_{\bar{x}}$. Assume that $\mu \in \mathcal{P}(X)$ satisfies:*

$$\tau(\cdot, \bar{x})^p \in L^1(X, \mu) \quad and \quad \tau(\cdot, \bar{x}) > 0 \ \mu\text{-a.e. } . \tag{1.12}$$

Then (μ, ν) is strongly timelike p-dualisable. In other terms, in case ν is a Dirac measure, the strongly timelike p-dualisability is equivalent to the timelike p-dualisability.

1.4 Geodesic structure of the Lorentz-Wasserstein space

A relevant object in the study of the geometry of Lorentz-Wasserstein space are geodesics.

Definition 1.13. Let $(X, \mathsf{d}, \ll, \leq, \tau)$ be a Lorentzian pre-length space.
We say that $(\mu_s)_{s \in [0,1]} \subset \mathcal{P}(X)$ is an ℓ_p-geodesic if and only if

$$\ell_p(\mu_s, \mu_t) = (t - s)\ell_p(\mu_0, \mu_1), \quad \forall \, 0 \leq s \leq t \leq 1.$$

In particular ℓ_p-geodesics are implicitly causal future-directed.

In the next proposition (for the proof see [24, Prop. 2.32]) we collect some useful properties of ℓ_p-geodesics. Before stating it, we introduce the evaluation map

$$\mathrm{e}_t : C([0, 1], X) \to X, \quad \gamma \mapsto \mathrm{e}_t(\gamma) := \gamma_t, \quad \forall t \in [0, 1], \tag{1.13}$$

and the stretching/restriction operator $\mathrm{restr}_{s_1}^{s_2} : C([0, 1], X) \to C([0, 1], X)$

$$(\mathrm{restr}_{s_1}^{s_2}\gamma)_t := \gamma_{(1-t)s_1 + ts_2}, \quad \forall s_1, s_2 \in [0, 1], s_1 < s_2, \forall t \in [0, 1]. \tag{1.14}$$

Proposition 1.14. *Let $(X, \mathsf{d}, \ll, \leq, \tau)$ be a $\mathcal{K}$-globally hyperbolic, Lorentzian geodesic space. Let $\mu_0, \mu_1 \in \mathcal{P}_c(X)$ such that there exists $\pi \in \Pi_{\leq}^{p\text{-}opt}(\mu_0, \mu_1)$ with $\operatorname{supp} \pi \Subset \{\tau > 0\}$ (in particular, if $\operatorname{supp}(\mu_0 \otimes \mu_1) \Subset \{\tau > 0\}$). Then*

1. *There always exists an ℓ_p-geodesic from μ_0 to μ_1.*

2. *For every ℓ_p-geodesic $(\mu_t)_{t \in [0,1]}$ from μ_0 to μ_1 there exists a probability measure $\eta \in \mathcal{P}(C([0,1], X)$ such that $(\mathrm{e}_t)_\sharp \eta = \mu_t$ for every $t \in [0,1]$ and η-a.e. γ is a maximal causal curve from $\gamma_0 \in \operatorname{supp} \mu_0$ to $\gamma_1 \in \operatorname{supp} \mu_1$. Such an η is called ℓ_p-dynamical optimal plan and the set of dynamical optimal plans from μ_0 to μ_1 is denoted by $\operatorname{OptGeo}_{\ell_p}(\mu_0, \mu_1)$.*

3. *If $\eta \in \operatorname{OptGeo}_{\ell_p}(\mu_0, \mu_1)$ then*

$$\eta^{s_1, s_2} := (\operatorname{restr}_{s_1}^{s_2})_\sharp \eta \in \operatorname{OptGeo}_{\ell_p}((\mathrm{e}_{s_1})_\sharp \eta, (\mathrm{e}_{s_2})_\sharp \eta), \quad \forall s_1 < s_2, \; s_1, s_2 \in [0, 1].$$

Proposition 1.14 proves that the ℓ_p-optimal transport is performed along geodesics of the underlying space X. This link with the geometry of X has been crucially used to study Ricci curvature.

1.5 Timelike Ricci curvature lower bounds: smooth setting

In the sequel, we fix a non-negative Borel measure $\mathfrak{m}$ on (X, d) which is finite on bounded sets. Such an $\mathfrak{m}$ will play the role of reference volume measure. In case of a smooth Lorentzian manifold (M, g), a natural choice is given by $\mathfrak{m} = \operatorname{vol}_g$ (i.e. the volume measure associated to g) or $\mathfrak{m} = \exp(f) \operatorname{vol}_g$, where $f \in C^\infty(M)$ plays the role of a weight. Once a reference volume measure $\mathfrak{m}$ is fixed, one can define an entropy functional.

Definition 1.15. Given a probability measure $\mu \in \mathcal{P}(X)$ we define its relative Boltzmann-Shannon entropy by

$$\operatorname{Ent}(\mu|\mathfrak{m}) = \int_M \rho \log(\rho) \, \mathfrak{m}, \tag{1.15}$$

if $\mu = \rho \, \mathfrak{m}$ is absolutely continuous with respect to $\mathfrak{m}$ and $(\rho \log(\rho))_+$ is $\mathfrak{m}$-integrable. Otherwise we set $\operatorname{Ent}(\mu|\mathfrak{m}) = +\infty$.

A simple application of Jensen inequality using the convexity of $(0, \infty) \ni t \mapsto t \log t$ gives

$$\operatorname{Ent}(\mu|\mathfrak{m}) \geq -\log \mathfrak{m}(\operatorname{supp}\mu) > -\infty, \quad \forall \mu \in \mathcal{P}_c(X). \tag{1.16}$$

We set $\operatorname{Dom}(\operatorname{Ent}(\cdot|\mathfrak{m})) := \{\mu \in \mathcal{P}(X) : \operatorname{Ent}(\mu|\mathfrak{m}) \in \mathbb{R}\}$ to be the finiteness domain of the entropy.

An important property of the relative entropy is the lower-semicontinuity under narrow convergence, (for the proof see for instance [86, Lemma 4.1])

$$\mu_n \to \mu_\infty \text{ narrowly and } \mathfrak{m}\left(\bigcup_{n \in \mathbb{N}} \operatorname{supp}\mu_n\right) < \infty \implies \liminf_{n \to \infty} \operatorname{Ent}(\mu_n|\mathfrak{m}) \geq \operatorname{Ent}(\mu_\infty|\mathfrak{m}). \tag{1.17}$$

The connection between Ricci curvature lower bounds and convexity properties of the entropy functional along Wasserstein geodesics (classically called "displacement convexity") is now well understood in the Riemannian setting. It was proved by McCann (in his Ph.D. thesis [70]) that displacement convexity holds in $\mathbb{R}^n$. Then Otto-Villani [77] formally observed that non-negative Ricci implies displacement convexity for smooth Riemannian manifolds, observation then rigorously proved by Cordero Erausquin-McCann-Schmuckenschläger [25]. The circle was then closed by Sturm-von Renesse [82], who proved a complete characterization of Ricci lower bounds for smooth Riemannian manifolds in terms of convexity properties of the entropy.

In the same spirit as in the Riemannian setting, the convexity properties of the entropy functional along ℓ_p-geodesics are equivalent to lower bounds on the Ricci curvature. This was proved first by McCann [71] for non-negative timelike Ricci lower bounds; in an independent slightly subsequent work, M.-Suhr [75] characterised general lower and upper timelike Ricci bounds (for the proof of the exact statement below, see [24, Theorem 3.1]).

Theorem 1.16 (McCann and M.-Suhr)**.** *Let (M^n, g) be a smooth globally hyperbolic spacetime, $0 < p < 1$ and denote by vol_g the associated volume measure.*
Then the following assertions are equivalent:

- $\mathrm{Ric}_g(v, v) \geq -Kg(v, v)$, *for every timelike $v \in TM$.*

- $\forall (\mu_0, \mu_1) \in (\mathrm{Dom}(\mathrm{Ent}(\cdot | \mathrm{vol}_g)))^2$ *timelike p-dualisable, there exists a (unique) ℓ_p-geodesic $(\mu_t)_{t \in [0,1]}$ s.t. $[0,1] \ni t \mapsto e(t) := \mathrm{Ent}(\mu_t | \mathrm{vol}_g)$ is semi-convex (hence locally Lipschitz on $(0, 1)$) and it satisfies:*

$$e''(t) - \frac{1}{n} e'(t)^2 \geq K \int_{M \times M} \tau(x, y)^2 \, \pi(\mathrm{d}x\mathrm{d}y), \qquad (1.18)$$

  in the distributional sense on $[0, 1]$.

- *For any couple $(\mu_0, \mu_1) \in (\mathrm{Dom}(\mathrm{Ent}(\cdot | \mathrm{vol}_g)) \cap \mathcal{P}_c(X))^2$ which is strongly timelike p-dualisable there exists a (unique) ℓ_p-geodesic $(\mu_t)_{t \in [0,1]}$ joining them and satisfying (1.18).*

Remark 1.17 (Timelike Ricci lower bounds and energy conditions in General Relativity)**.**

- **Weak energy condition**. A lower bound on the timelike Ricci curvature of a spacetime (M^n, g), i.e. the first item in Theorem 1.16:

  There exists $K \in \mathbb{R}$ such that $\mathrm{Ric}_g \geq -Kg(v, v)$ for all timelike vectors $v \in TM$, (1.19)

  is quite a natural assumption in General Relativity. Of course, for a C^2-metric g, the lower bound (1.19) is satisfied on compact subsets of the space-time.

  Recalling that the Einstein's equations postulate proportionality of Ric_g and $T - \frac{1}{n-2}\mathrm{tr}_g(T)g$ (where T is the so-called energy-momentum tensor), for a general cosmological constant $\Lambda \in \mathbb{R}$, the lower bound (1.19) is equivalent to require that

  $$T(v, v) \geq -\frac{1}{n-2}\mathrm{tr}_g(T) + \frac{1}{8\pi}\left(K - \frac{2\Lambda}{n-2}\right), \text{ for all } v \in TM \text{ with } g(v, v) = -1.$$

  In particular, if $\inf_M \mathrm{tr}_g(T) > -\infty$ (or, equivalently, $\inf_M \mathrm{R}_g > -\infty$ where R_g is the scalar curvature of g), then the *weak* energy condition $T(v, v) \geq 0$ for all timelike v (which is believed to hold for most physically reasonable T, according to [94, pag. 218]) implies (1.19).

- **Hawking-Penrose's Strong energy condition** (SEC for short). The SEC asserts that, called T the energy-momentum tensor in the Einstein's equations, it holds $T(v, v) \geq \frac{1}{2}\mathrm{tr}_g(T)$ for every time-like vector $v \in TM$ satisfying $g(v, v) = -1$. Assuming that the space-time (M, g) satisfies the Einstein equations with zero cosmological constant, the SEC is equivalent to $\mathrm{Ric}_g(v, v) \geq 0$ for every timelike vector $v \in TM$. This corresponds to the case $K = 0$ in Theorem 1.16.

  The SEC, proposed by Hawking and Penrose [79, 48, 50], plays a key role in General Relativity. For instance, in the presence of trapped surfaces, it implies that the space-time is geodesically incomplete. Fact interpreted as singular behaviour possibly connected to the presence of a black holes (for a general discussion about singularity theorems see for instance the monographs [49, 94, 60]). Moreover, due to the averaged-focusing effect on geodesics, the SEC is sometimes interpreted as the geometric counterpart of the fact that gravity is an attractive force.

Remark 1.18 (A heuristic thermodynamic interpretation of Theorem 1.16)**.** An ℓ_p-geodesic $(\mu_t)_{t \in [0,1]}$ can be interpreted as the evolution [a] of a distribution (in space-time) of gas. Theorem 1.16 says

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[a] strictly speaking t is not the proper time, but only a variable parametrising the evolution

that timelike Ricci lower bounds (which correspond to energy conditions in General Relativity) can be equivalently formulated in terms of the convexity properties of the Boltzmann-Shannon entropy along such evolutions $(\mu_t)_{t \in [0,1]} \subset \mathcal{P}(M)$. Extrapolating a bit more, we might say that the second law of thermodynamics (i.e. in a natural thermodynamic process, the sum of the entropies of the interacting thermodynamic systems decreases, due to our sign convention) concerns the *first* derivative of the Boltzmann-Shannon entropy; gravitation (under the form of Ricci curvature) is instead related to the *second* order derivative of the Boltzmann-Shannon entropy along a natural thermodynamic process.

Remark 1.19 (Some related physics literature). The existence of strong connections between thermodynamics and General Relativity is not new in the physics literature; it has its origins at least in the work Bekenstein [11] and Hawking with collaborators [9] in the mid-1970s about the black hole thermodynamics. These works inspired a new research field in theoretical physics, called entropic gravity (also known as emergent gravity), asserting that gravity is an entropic force rather than a fundamental interaction. Let us give a brief account. In 1995 Jacobson [47] derived the Einstein equations from the proportionality of entropy and horizon area of a black hole, exploiting the fundamental relation $\delta Q = T \delta S$ linking heat Q, temperature T and entropy S. Subsequently, other physicists (let us mention Padmanabhan [78]) have been exploring links between gravity and entropy. More recently, in 2011 Verlinde [90] proposed a heuristic argument suggesting that (Newtonian) gravity can be identified with an entropic force caused by changes in the information associated with the positions of material bodies. A relativistic generalization of those arguments leads to the Einstein equations.

Theorem 1.16 can be seen as an additional strong connection between general relativity and thermodynamics/information theory. It would be interesting to explore this relationship further.

Remark 1.20. Let us stress that the framework (1.19) includes any solution of the Einstein's equations *in vacuum* (i.e. with null stress-energy tensor T) with *possibly non-zero cosmological constant* Λ. Already such a framework is highly interesting as the standard black hole metrics (e.g. Schwartzshild, Kerr) are solutions of the Einstein's vacuum equations, and also the more recent literature on black holes typically focuses on vacuum solutions (see e.g. [26, 28, 29, 56]). A key role in such breakthroughs on black holes is given by a deep analysis of the system of non-linear hyperbolic partial differential equations corresponding to the Einstein's vacuum equations (in a suitable gauge). At least in the smooth setting, it was recently proved by the second author and Suhr [75] that the optimal transport point of view is compatible with the hyperbolic PDEs one (in the sense that it is possible to characterise solutions of the Einstein's equations in terms of optimal transport). We believe that all of this suggests that the proposed optimal transport approach has potential for the future.

2 Synthetic Theory of timelike Ricci lower bounds

2.1 Definitions of $\mathsf{TCD}_p^e(K, N)$, $\mathsf{wTCD}_p^e(K, N)$ and $\mathsf{TMCP}_p^e(K, N)$

Theorem 1.16 proves the equivalence between a condition requiring the smoothness of the Lorentzian metric g and a condition (the ordinary differential inequality for e) that can be formulated in any Lorentzian pre-length space. Inspired by the Riemannian counterpart which led to the flourishing theory of Lott-Sturm-Villani $\mathsf{CD}(K, N)$ metric measure spaces [86, 87, 65], it is then natural to turn it into a definition.

Definition 2.1 (TCD condition). Fix $p \in (0, 1)$, $K \in \mathbb{R}$, $N \in (0, \infty)$. A measured pre-length space $(X, \mathsf{d}, \mathfrak{m}, \ll, \leq, \tau)$ satisfies the $\mathsf{TCD}_p^e(K, N)$ (resp. $\mathsf{wTCD}_p^e(K, N)$) if for any couple $(\mu_0, \mu_1) \in \mathrm{Dom}(\mathrm{Ent}(\cdot | \mathfrak{m}))^2$ which is timelike p-dualisable (resp. $(\mu_0, \mu_1) \in (\mathrm{Dom}(\mathrm{Ent}(\cdot | \mathfrak{m})) \cap \mathcal{P}_c(X))^2$ which is strongly timelike p-dualisable) by some $\pi \in \Pi_{\ll}^{p\text{-opt}}(\mu_0, \mu_1)$, there exists an ℓ_p-geodesic $(\mu_t)_{t \in [0,1]}$

such that $[0,1] \ni t \mapsto e(t) := \mathrm{Ent}(\mu_t|\mathfrak{m})$ is semi-convex (hence locally Lipschitz on $(0,1)$) and it satisfies

$$e''(t) - \frac{1}{N} e'(t)^2 \geq K \int_{X \times X} \tau(x,y)^2 \, \pi(\mathrm{d}x\mathrm{d}y), \tag{2.1}$$

in the distributional sense on $[0,1]$.

The condition (2.1) is a differential, and therefore infinitesimal, condition. One can however formulate an equivalent global condition taking inspiration from the Entropic formulation of the curvature-dimension condition CD obtained by Erbar-Kuwada-Sturm [31]. This is done by noticing that $t \mapsto e(t)$ is semi-convex and satisfies the inequality (2.1) if and only if $t \mapsto u_N(t) := \exp(-e(t)/N)$ is semi-convex and satisfies

$$u_N'' \leq -\frac{K}{N} \|\tau\|_{L^2(\pi)}^2 \, u_N. \tag{2.2}$$

Set

$$\mathfrak{s}_\kappa(\vartheta) := \begin{cases} \frac{1}{\sqrt{\kappa}} \sin(\sqrt{\kappa}\vartheta), & \kappa > 0 \\ \vartheta, & \kappa = 0 \\ \frac{1}{\sqrt{-\kappa}} \sinh(\sqrt{-\kappa}\vartheta), & \kappa < 0 \end{cases}, \qquad \mathfrak{c}_\kappa(\vartheta) := \begin{cases} \cos(\sqrt{\kappa}\vartheta), & \kappa \geq 0 \\ \cosh(\sqrt{-\kappa}\vartheta) & \kappa < 0 \end{cases}, \tag{2.3}$$

and

$$\sigma_\kappa^{(t)}(\vartheta) := \begin{cases} \frac{\mathfrak{s}_\kappa(t\vartheta)}{\mathfrak{s}_\kappa(\vartheta)}, & \kappa\vartheta^2 \neq 0 \text{ and } \kappa\vartheta^2 < \pi^2 \\ t, & \kappa\vartheta^2 = 0 \\ +\infty & \kappa\vartheta^2 \geq \pi^2 \end{cases}.$$

Note that the function $\kappa \mapsto \sigma_\kappa^{(t)}(\vartheta)$ is non-decreasing for every fixed ϑ, t. With the above notation, the differential inequality (2.2) is equivalent to the integrated version (cf. [31, Lemma 2.2]):

$$u_N(t) \geq \sigma_{K/N}^{(1-t)}\left(\|\tau\|_{L^2(\pi)}\right) u_N(0) + \sigma_{K/N}^{(t)}\left(\|\tau\|_{L^2(\pi)}\right) u_N(1). \tag{2.4}$$

We thus proved the following proposition.

Proposition 2.2. *Fix $p \in (0,1)$, $K \in \mathbb{R}$ and $N \in (0,\infty)$. The measured Lorentzian pre-length space $(X, \mathsf{d}, \mathfrak{m}, \ll, \leq, \tau)$ satisfies $\mathsf{TCD}_p^e(K,N)$ (resp. $\mathsf{wTCD}_p^e(K,N)$) if and only if for any couple $(\mu_0, \mu_1) \in \left(\mathrm{Dom}(\mathrm{Ent}(\cdot|\mathfrak{m}))\right)^2$ which is timelike p-dualisable (resp. $(\mu_0, \mu_1) \in [\mathrm{Dom}(\mathrm{Ent}(\cdot|\mathfrak{m})) \cap \mathcal{P}_c(X)]^2$ which is strongly timelike p-dualisable) by some $\pi \in \Pi_\ll^{p\text{-}opt}(\mu_0, \mu_1)$, there exists an ℓ_p-geodesic $(\mu_t)_{t \in [0,1]}$ such that the function $[0,1] \ni t \mapsto u_N(t) := U_N(\mu_t|\mathfrak{m})$ satisfies (2.4).*

By considering (K,N)-convexity properties only of those ℓ_p-geodesics $(\mu_t)_{t \in [0,1]}$ where μ_1 is a Dirac delta one obtains the following weaker condition [24] (inspired by the Riemannian counterparts by Sturm [87] and Ohta [76]).

Definition 2.3. *Fix $p \in (0,1)$, $K \in \mathbb{R}$, $N \in (0,\infty)$. The measured globally hyperbolic Lorentz geodesic space $(X, \mathsf{d}, \mathfrak{m}, \ll, \leq, \tau)$ satisfies $\mathsf{TMCP}_p^e(K,N)$ if and only if for any $\mu_0 \in \mathcal{P}_c(X) \cap \mathrm{Dom}(\mathrm{Ent}(\cdot|\mathfrak{m}))$ and for any $x_1 \in X$ such that $x \ll x_1$ for μ_0-a.e. $x \in X$, there exists an ℓ_p-geodesic $(\mu_t)_{t \in [0,1]}$ from μ_0 to $\mu_1 = \delta_{x_1}$ such that*

$$U_N(\mu_t|\mathfrak{m}) \geq \sigma_{K/N}^{(1-t)}\left(\sqrt{\int_X \tau(x,x_1)^2 \, \mu_0(\mathrm{d}x)}\right) U_N(\mu_0|\mathfrak{m}), \quad \forall t \in [0,1). \tag{2.5}$$

Also the TMCP condition is able to characterize the lower bounds on the time-like Ricci curvature. In particular (see [24, Cor. A.2] if (M^n, g) is a globally hyperbolic smooth spacetime of dimension $n \geq 2$ without boundary then:

1. If $\mathrm{Ric}_g(v,v) \geq -Kg(v,v)$ for every timelike vector $v \in TM$, then the associated Lorentzian geodesic space satisfies $\mathsf{TMCP}_p^e(K', N')$ for every $K' \leq K$ and $N' \geq N$.

2. If the Lorentzian geodesic space associated to (M^n, g) satisfies $\mathsf{TMCP}_p^e(K, N)$, then $n \leq N$.

Hence the notion of TMCP is compatible with the smooth setting. Finally, as expected, the $\mathsf{wTCD}_p^e(K, N)$ condition implies the $\mathsf{TMCP}_p^e(K, N)$ (see [24, Prop. 3.11] for the proof).

Proposition 2.4. *Fix $p \in (0, 1)$, $K \in \mathbb{R}$, $N \in (0, \infty)$. The $\mathsf{wTCD}_p^e(K, N)$ condition implies the $\mathsf{TMCP}_p^e(K, N)$ condition for locally causally closed, d-compatible, $\mathcal{K}$-globally hyperbolic Lorentzian geodesic spaces.*

A number of geometric consequences can be deduced for measured pre-length spaces verifying the wTCD condition (some of them valid also under the weaker TMCP). Referring to [24] for the complete list, here we only mention the timelike Brunn-Minkowski inequality, the timelike Bishop-Gromov inequality, and the timelike Bonnet-Myers.

2.2 Stability of synthetic Ricci curvature lower bounds

One of the main motivation for the huge impact of the Lott-Sturm-Villani theory of synthetic curvature bounds relies in the the stability property of the CD condition: if a sequence of metric measure spaces $(X_n, \mathsf{d}_n, \mathfrak{m}_n)$ verifies $\mathsf{CD}(K, N)$ and it is converging in the measured-Gromov-Hausdorff sense to a limit metric measure space $(X, \mathsf{d}, \mathfrak{m})$, then $(X, \mathsf{d}, \mathfrak{m})$ satisfies the $\mathsf{CD}(K, N)$ condition as well.

We have therefore faced the challenging problem of developing a stability property for the TCD condition. While there is a well established theory of convergence for Riemannian manifolds (and more generally for metric measure spaces), the situation for Lorentzian manifolds (and even more for Lorentzian length spaces) is more open to investigation. Without a well established notion of convergence for Lorentzian length spaces already present in the literature, we have devised the following one. Before stating it, recall that a topological embedding is a map $f : X \to Y$ between two topological spaces X and Y such that f is continuous, injective and with continuous inverse between X and $f(X)$.

Definition 2.5. A sequence of measured Lorentzian geodesic spaces $\{(X_j, \mathsf{d}_j, \mathfrak{m}_j, \ll_j, \leq_j, \tau_j)\}_{j \in \mathbb{N}}$ converges to $(X_\infty, \mathsf{d}_\infty, \mathfrak{m}_\infty, \ll_\infty, \leq_\infty, \tau_\infty)$ if and only if:

1. There exists a locally causally closed, $\mathcal{K}$-globally hyperbolic Lorentzian geodesic space $(\bar{X}, \bar{\mathsf{d}}, \bar{\ll}, \bar{\leq}, \bar{\tau})$ such that each $j \in \mathbb{N} \cup \{\infty\}$, $(X_j, \mathsf{d}_j, \mathfrak{m}_j, \ll_j, \leq_j, \tau_j)$ is isomorphically embedded in it, i.e. there exist topological embedding maps $\iota_j : X_j \to \bar{X}$ such that

   - $x_j^1 \leq_j x_j^2$ if and only if $\iota_j(x_j^1) \bar{\leq} \iota_j(x_j^2)$, for every $j \in \mathbb{N} \cup \{\infty\}$, for every $x_j^1, x_j^2 \in X_j$;
   - $\tau_j(x_j^1, x_j^2) = \bar{\tau}(\iota_j(x_j^1), \iota_j(x_j^2))$ for every $x_j^1, x_j^2 \in X_j$, for every $j \in \mathbb{N} \cup \{\infty\}$;

2. The measures $(\iota_j)_\sharp \mathfrak{m}_j$ converge to $(\iota_\infty)_\sharp \mathfrak{m}_\infty$ weakly in duality with $C_c(\bar{X})$ in $\bar{X}$, i.e.

$$\int \varphi \, (\iota_j)_\sharp \mathfrak{m}_j \to \int \varphi \, (\iota_\infty)_\sharp \mathfrak{m}_\infty \quad \forall \varphi \in C_c(\bar{X}).$$

In Definition 2.5 we use topological embeddings to identify spaces with their image inside a larger space. This procedure to compare different spaces is rather standard in the framework of metric measure spaces and it provides an effective way to define and study the measured Gromov-Haudorff convergence (in the latter, the assumption is actually that each embedding is *isometric*, which of course is much stronger than merely a topological embedding; for an overview of equivalent formulations of convergence for metric measure spaces see for instance [38, Sec. 3]). An important assumption is that the causal structures are preserved under the image of the embeddings.

Even though we haven't specifically listed any topological assumption on the sequence of spaces X_j, they actually inherit them from X_∞ via the topological embeddings ι_j. The map ι_j preserves both the causal relations and τ_j, hence $(X_j, \mathsf{d}_j, \mathfrak{m}_j, \ll_j, \leq_j, \tau_j)$ are locally causally closed and $\mathcal{K}$-globally hyperbolic Lorentzian geodesic (by assumption) spaces.

Theorem 2.6 (Weak stability of TCD). *Let $\{(X_j, \mathsf{d}_j, \mathfrak{m}_j, \ll_j, \leq_j, \tau_j)\}_{j \in \mathbb{N} \cup \{\infty\}}$ be a sequence of measured Lorentzian geodesic spaces converging to $(X_\infty, \mathsf{d}_\infty, \mathfrak{m}_\infty, \ll_\infty, \leq_\infty, \tau_\infty)$ in the sense of Definition 2.5. Assume moreover that there exist $p \in (0,1), K \in \mathbb{R}, N \in (0,\infty)$ such that $(X_j, \mathsf{d}_j, \mathfrak{m}_j, \ll_j, \leq_j, \tau_j)$ satisfies $\mathsf{TCD}_p^e(K, N)$, for each $j \in \mathbb{N}$.*

Then the limit space $(X_\infty, \mathsf{d}_\infty, \mathfrak{m}_\infty, \ll_\infty, \leq_\infty, \tau_\infty)$ satisfies $\mathsf{wTCD}_p^e(K, N)$.

Let us comment on the general strategy to obtain a stability result for the TCD condition; this will highlight the reasons for which full stability is not yet at disposal, and motivates the presence of the two distinct definitions (weak and strong) of the timelike curvature-dimension condition (see Definition 2.1). For the detailed proof see [24, Thm. 3.14]

Some comments on the proof. Start by fixing two probability measures $\mu_{0,\infty}, \mu_{1,\infty} \in \mathcal{P}_c(\bar{X})$ both a.c. with respect to $\mathfrak{m}_\infty$. The goal is to construct an ℓ_p-geodesic $(\mu_{t,\infty})_{t \in [0,1]}$ such that the map $[0,1] \ni t \mapsto u_N(t) := \exp(-\mathrm{Ent}(\mu_{t,\infty}|\mathfrak{m}_\infty)/N)$ satisfies the inequality (2.4).

The classic approach in the metric measure space setting (see for instance [86, 87] would be to use that $\mathfrak{m}_j \to \mathfrak{m}_\infty$ in $W_2^{(\bar{X}, \bar{\mathsf{d}})}$ (this can be assumed without any loss of generality) and to take $\gamma_j \in \mathcal{P}(\bar{X} \times \bar{X})$ optimal plan between $\mathfrak{m}_j$ and $\mathfrak{m}_\infty$.

Then, denoting by $\mu_{i,\infty} = \rho_i \mathfrak{m}_\infty$, for $i = 0, 1$, one defines two sequences of probability measures:

$$\mu_{i,j} := (P_1)_\sharp (\rho_i \gamma_j), \quad i = 0, 1. \quad \implies \quad \mu_{i,j} \ll \mathfrak{m}_j,$$

also verifying $\mu_{i,j} \rightharpoonup \mu_{i,\infty}$ as $j \to \infty$ for $i = 0, 1$.

The general scheme (that in this setting breaks down) would be to consider an ℓ_p-geod $(\mu_{t,j})_{t \in [0,1]}$ from $\mu_{0,j}$ to $\mu_{1,j}$ given by the $\mathsf{TCD}(K, N)$ condition of X_j and then to prove (non-trivial but classical) the upper semi-continuity to the limit on $u_{N,j}(t)$ and continuity for $u_{N,j}(0)$ and $u_{N,j}(1)$.

In the present setting, the major difficulty is that $\mu_{0,j}$ and $\mu_{1,j}$ may or may not be causally related; in particular, one cannot invoke the $\mathsf{TCD}(K, N)$ condition between $\mu_{0,j}$ and $\mu_{1,j}$. This is the main reason why we assume that the spaces in the converging sequence satisfy the $\mathsf{TCD}(K, N)$ (i.e. the slightly stronger version of the condition), and we obtain that the limit space satisfies the $\mathsf{wTCD}(K, N)$ (i.e. the slightly weaker version of the condition).

Alternatively, this issue could be circumvented also by imposing some restrictions on the family of geodesics along which the concavity of the entropy has to be checked. $\qquad\square$

If either μ_0 or μ_1 is a Dirac delta, then full stability is granted. This is the content of the next result (for a detailed proof see [24, Thm 3.12])

Theorem 2.7 (Stability of TMCP). *Let $\{(X_j, \mathsf{d}_j, \mathfrak{m}_j, \ll_j, \leq_j, \tau_j)\}_{j \in \mathbb{N} \cup \{\infty\}}$ be a sequence of measured Lorentzian geodesic spaces converging to $(X_\infty, \mathsf{d}_\infty, \mathfrak{m}_\infty, \ll_\infty, \leq_\infty, \tau_\infty)$ in the sense of Definition 2.5. Assume moreover that there exist $p \in (0,1), K \in \mathbb{R}, N \in (0,\infty)$ such that $(X_j, \mathsf{d}_j, \mathfrak{m}_j, \ll_j, \leq_j, \tau_j)$ satisfies $\mathsf{TMCP}_p^e(K, N)$, for each $j \in \mathbb{N}$.*

Then the limit space $(X_\infty, \mathsf{d}_\infty, \mathfrak{m}_\infty, \ll_\infty, \leq_\infty, \tau_\infty)$ satisfies $\mathsf{TMCP}_p^e(K, N)$ as well.

3 Applications: Hawking singularity Theorem and geometric comparison results in a Lorentzian synthetic setting

3.1 Some useful preliminary concepts

In this section we will address a generalization to the synthetic setting, namely for those Lorentzian geodesic spaces verifying the $\mathsf{TMCP}_p^e(K, N)$ condition, of the Hawking singularity Theorem.

To this aim, we will need an extra regularity assumption on the geodesics, namely we will require the space to be timelike non-branching (recall the definition (1.4) of TGeo(X)).

Definition 3.1 (Timelike non-branching). A Lorentzian pre-length space $(X, \mathsf{d}, \ll, \leq, \tau)$ is said to be *forward timelike non-branching* if and only if for any $\gamma^1, \gamma^2 \in \text{TGeo}(X)$, it holds:

$$\exists\, \bar{t} \in (0, 1) \text{ such that } \forall t \in [0, \bar{t}] \quad \gamma_t^1 = \gamma_t^2 \implies \gamma_s^1 = \gamma_s^2, \quad \forall s \in [0, 1].$$

It is said to be *backward timelike non-branching* if the reversed causal structure is forward timelike non-branching. In case it is both forward and backward timelike non-branching it is said *timelike non-branching*.

Remark 3.2. By Cauchy-Lipschitz Theorem, it is clear that if (M, g) is a space-time whose Christoffel symbols are locally-Lipschitz (e.g. in case $g \in C^{1,1}$) then the associated synthetic structure is timelike non-branching. It is expected that for spacetimes with a metric of lower regularity (e.g. $g \in C^1$ or $g \in C^0$) timelike branching may occur. Similar behaviour could occur in closed cone structures when the Lorentz-Finsler norm is not strictly convex (see [73, Remark 2.8]).

The starting point for the Hawking singularity Theorem is to consider an *achronal* subset of X.

Definition 3.3 (Achronal set). Let $(X, \mathsf{d}, \ll, \leq, \tau)$ be a Lorentzian pre-length space. A set $V \subset X$ is *achronal* if $x \not\ll y$, $\forall x, y \in V$.

Remark 3.4. If V is achronal, then $I^+(V) \cap I^-(V) = \emptyset$.

We will need to give a meaning to lower bounds on the mean curvature of V and we therefore need some additional regularity.
The following compactness property, originally introduced by Galloway [36] in the smooth setting, will play an important role.

Definition 3.5 (Future timelike complete (FTC))). A $V \subset X$ is *future timelike complete* (FTC for short) if for each point $x \in I^+(V)$, the intersection $J^-(x) \cap V \subset V$ has compact closure (w.r.t. d) in V. Analogously for *past timelike complete* (PTC).

A key role in the proof of the Hawking singularity Theorem is played by the signed time-separation function from a FTC subset V. This can be seen as the Lorentzian counterpart of the signed distance function from a compact set in a Riemannian manifold. Let us recall its definition.

Definition 3.6 (Signed time-separation τ_V). Let $V \subset X$ be an achronal subset. The *signed time-separation* from V is the function $\tau_V : X \to [-\infty, +\infty]$ defined by

$$\tau_V(x) := \begin{cases} \sup_{y \in V} \tau(y, x), & \text{for } x \in I^+(V) \\ -\sup_{y \in V} \tau(x, y), & \text{for } x \in I^-(V) \\ 0 & \text{otherwise.} \end{cases} \tag{3.1}$$

Note that τ_V is lower semi-continuous, as supremum of (lower semi-)continuous functions.
In order for these suprema to be attained, global hyperbolicity and geodesic property of X alone are not sufficient; here the additional FTC and PTC assumptions on V are very useful.

Proposition 3.7. *Let $(X, \mathsf{d}, \ll, \leq, \tau)$ be a globally hyperbolic Lorentzian geodesic space and let $V \subset X$ be an achronal FTC (resp. PTC) subset. Then, for all $x \in I^+(V)$ (resp. for all $x \in I^-(V)$) there exists $y_x \in V$ with $\tau_V(x) = \tau(y_x, x) > 0$ (resp. $\tau_V(x) = -\tau(x, y_x) < 0$).*

We refer the reader to [24, Lemma 1.8] for the proof.

3.2 Disintegration and synthetic mean curvature bounds

In order to introduce a weak notion of mean curvature, it is useful to define normal variations to the achronal set V. Since τ_V is also 1-Lipschitz with respect to τ, relying on Proposition 3.7 one can consider, even in the non-smooth setting, integral lines of τ_V saturating the 1-Lipschitz inequality. This will provide, up to a set of measure zero, a (non-smooth) "foliation" of $I^+(V)$ by timelike geodesics that will suffice to our scope.

Indeed, removing sets of measure zero (here the synthetic curvature bound $\mathsf{TMCP}_p^e(K,N)$ is crucially used), the following disintegration formula for the reference measure $\mathfrak{m}$ holds true (see [24, Sec. 4] for the details):

$$\mathfrak{m}{\llcorner}_{I^+(V)} = \int_Q h(\alpha,\cdot)\,\mathcal{L}^1{\llcorner}_{X_\alpha}\,\mathfrak{q}(d\alpha), \qquad X_\alpha \text{ timelike geodesic for every } \alpha \in Q$$

- Q is a (typically uncountable) set of indices, that can be obtained as a Borel quotient set from X under the equivalence relation that identifies points lying on the same integral line of τ_V saturating the 1-Lipschitz (w.r.t. τ) inequality;

- $\mathfrak{q}$ is a probability measure over the set of indices Q (in case the ambient measure $\mathfrak{m}$ is finite, this is obtained as push forward of $\mathfrak{m}$ via the aforementioned quotient map after normalization; the case when $\mathfrak{m}$ is not finite can be treated by a cut-off argument);

- $h(\alpha,\cdot) \in L^1_{loc}(X_\alpha, \mathcal{L}^1{\llcorner}_{X_\alpha})$ for $\mathfrak{q}$-a.e. $\alpha \in Q$;

- $(X_\alpha, |\cdot|, h(\alpha,\cdot)\mathcal{L}^1)$ verifies $\mathsf{MCP}(K,N)$, i.e.

$$\left(\frac{\mathfrak{s}_{K/(N-1)}(b - \tau_V(x_1))}{\mathfrak{s}_{K/(N-1)}(b - \tau_V(x_0))}\right)^{N-1} \leq \frac{h(\alpha, x_1)}{h(\alpha, x_0)} \leq \left(\frac{\mathfrak{s}_{K/(N-1)}(\tau_V(x_1) - a)}{\mathfrak{s}_{K/(N-1)}(\tau_V(x_0) - a)}\right)^{N-1}, \tag{3.2}$$

for all $x_0, x_1 \in X_\alpha$, with $0 \leq a < \tau_V(x_0) < \tau_V(x_1) < b < \pi\sqrt{(N-1)/(K \vee 0)}$. Here $(a,b) \subset \mathbb{R}$ denotes the segment corresponding to the geodesic X_α via the map τ_V, i.e. $(a,b) = \tau_V(X_\alpha)$ (note that the interval can also be closed, or half closed, depending on the geometry of X), and $\mathfrak{s}_{K/(N-1)}(\cdot)$ was defined in (2.3).

By a Fubini-Tonelli argument, one can obtain the following coarea-type formula:

$$\mathfrak{m}{\llcorner}_{I_V^+} = \int_0^\infty \mathcal{H}_t\,dt, \tag{3.3}$$

meaning that for each measurable set $A \subset I_V^+ \cup V$ with $\mathfrak{m}(A) < \infty$, the map $[0,\infty) \ni t \mapsto \mathcal{H}_t(A)$ is measurable and

$$\mathfrak{m}(A) = \int_0^\infty \mathcal{H}_t(A)\,dt = \int_0^\infty \mathcal{H}_t(A \cap \{\tau_V = t\})\,dt.$$

Moreover, by construction, $\mathcal{H}_t$ is concentrated on the level set $\{\tau_V = t\}$ and $\mathcal{H}_0$ is concentrated on V. We refer to [24, Sec. 5.1] for the precise definition of the family of measures $\mathcal{H}_t$ and all the other missing details.

We now review how (3.3) is sufficient to give a meaning to *synthetic mean curvature bounds for V* at least in the smooth setting. In particular we will analyse the disintegration formula and the measures $\mathcal{H}_t$ and their relation to mean curvature bounds.

Let (M^n, g) be a $2 \leq n$-dimensional smooth globally hyperbolic space-time and $V \subset M$ be a smooth compact achronal spacelike hypersurface without boundary. Then, the signed time-separation function τ_V from V is smooth on a neighbourhood U of V and $\nabla \tau_V$ is the smooth timelike past-pointing unit normal vector field along V. More precisely,

$$\nabla \tau_V(x) \perp T_x V, \quad g(\nabla \tau_V(x), \nabla \tau_V(x)) = -1, \quad \forall x \in V.$$

Denote with vol_g the volume measure of (M^n, g) and with vol_V the induced $(n-1)$-dimensional volume measure on V. By compactness of V, there exists $\delta > 0$ such that the g-geodesic $[0, \delta] \ni t \mapsto \exp_x(-t\nabla\tau_V(x))$ is a future pointing maximal geodesic, for every $x \in V$. Define

$$\mathcal{U} := V \times [0, \delta] \subset V \times \mathbb{R}, \quad \Phi : \mathcal{U} \to M, \ \Phi(x, t) := \exp_x(-t\nabla\tau_V(x)).$$

For $\delta > 0$ small enough it is a standard fact (tubular neighbourhood theorem) that Φ is a diffeomorphism onto its image and that the following integration formula holds true:

$$\int_M \varphi \, \mathrm{vol}_g = \int_V \int_0^\delta \varphi \circ \Phi(x, t) \det D\Phi_{(x,t)}|_{T_x V} \, \mathrm{d}t \, \mathrm{vol}_V(\mathrm{d}x), \quad \forall \varphi \in C_c(\Phi(\mathcal{U})). \tag{3.4}$$

In the smooth setting, we can identify the set of indices (which in general is a Borel quotient set) Q with V, and the quotient measure $\mathfrak{q}$ with $\psi \, \mathrm{vol}_V$ defined as follows:

$$\mathfrak{q} := \psi \, \mathrm{vol}_V \ll \mathrm{vol}_V, \quad \text{where } \psi(x) := \left(\int_0^\delta \det D\Phi_{(x,t)}|_{T_x V} \, \mathrm{d}t \right), \forall x \in V.$$

The integration formula (3.4) can be thus rewritten as

$$\int_M \varphi \, \mathrm{vol}_g = \int_V \frac{1}{\psi(x)} \int_0^\delta \varphi \circ \Phi(x, t) \det D\Phi_{(x,t)}|_{T_x V} \, \mathrm{d}t \, \mathfrak{q}(\mathrm{d}x), \quad \forall \varphi \in C_c(\Phi(\mathcal{U})). \tag{3.5}$$

By the uniqueness of the disintegration formula, (3.5) gives:

$$h_\alpha(t) = \frac{1}{\psi(\alpha)} \det D\Phi_{(\alpha,t)}|_{T_\alpha V}, \quad h_\alpha(0) = \frac{1}{\psi(\alpha)}, \quad \forall \alpha \in V, \ \forall t \in [0, \delta].$$

Moreover, it follows that the measure $\mathcal{H}_t$ can be written as

$$\mathcal{H}_t := \Phi(\cdot, t)_\sharp \left(\det D\Phi_{(\alpha,t)}|_{T_\alpha V} \, \mathrm{vol}_V(\mathrm{d}\alpha) \right), \quad \text{for all } t \geq 0,$$

in particular, $\mathcal{H}_0 = \mathrm{vol}_V$ and $\mathcal{H}_t$ is the $(n-1)$-volume measure on the hypersurface $\{\Phi(x, t) : x \in V\}$.

Given any $\phi \in C^\infty(V; \mathbb{R}_{\geq 0})$, one can consider the region $V_{t,\phi}$ as the domain trapped between V and the normal graph of ϕ. The first variation of the volume is thus

$$\frac{d}{dt} \mathrm{vol}_g(V_{t,\phi}) = \mathcal{H}^{n-1}(\{\Phi(x, t\phi(x)) : x \in V\}),$$

where $\mathcal{H}^{n-1}$ is the standard $(n-1)$-volume of the hypersurface $\{\Phi(x, t\phi(x)) : x \in V\}$; in particular, $\frac{d}{dt}\big|_{t=0} \mathrm{vol}_g(V_{t,\phi}) = \mathrm{vol}_V(V) = \mathcal{H}_0(V)$. Taking a further variation, we get

$$\int_V \phi^2 \, g(\vec{H}_V, \nabla\tau_V) \, \mathrm{vol}_V = \frac{d}{dt^2}\bigg|_{t=0} \mathrm{vol}_g(V_{t,\phi}) = \lim_{t \downarrow 0} \frac{\mathfrak{m}(V_{t,\phi}) - t \int_V \phi \mathcal{H}_0}{t^2/2}.$$

The right hand side, corresponding to the second variation of volume, is thus the first variation of the area which gives the mean curvature $\vec{H}_V$ of V. The right hand side can therefore be chosen as a starting point for defining bounds on the mean curvature. Still we need to make precise in the non-smooth setting the definition of variation set $V_{t,\phi}$. We will use the "initial-point projection map" $\mathfrak{a} : I_V^+ \to V$ assigning to each $x \in I_V^+$ the point given by Proposition 3.7. It is not hard to check it is $\mathfrak{m}$-measurable (see [24] for details).

Definition 3.8. [Synthetic mean curvature bounds] The Borel achronal FTC subset $V \subset X$ has *forward mean curvature bounded below by $H_0 \in \mathbb{R}$* if $\mathcal{H}_0$ is a non-negative Radon measure with $(\mathfrak{a})_\sharp \mathfrak{q} \ll \mathcal{H}_0$ and such that for any normal variation

$$V_{t,\phi} := \{x \in \mathcal{T}_V : 0 \leq \tau_V(x) \leq t\phi(\mathfrak{a}(x))\},$$

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the following inequality holds true:

$$\limsup_{t\downarrow 0} \frac{\mathfrak{m}(V_{t,\phi}) - t\int_V \phi\mathcal{H}_0}{t^2/2} \geq H_0 \int_V \phi^2 \mathcal{H}_0,$$

for any bounded Borel function $\phi : V \to [0,\infty)$ with compact support. Analogously, V has *forward mean curvature bounded above by* $H_0 \in \mathbb{R}$ if $\mathcal{H}_0$ is a non-negative Radon measure and for any normal variation $V_{t,\phi}$ as above the following inequality holds true:

$$\liminf_{t\downarrow 0} \frac{\mathfrak{m}(V_{t,\phi}) - t\int_V \phi\mathcal{H}_0}{t^2/2} \leq H_0 \int_V \phi^2 \mathcal{H}_0, \tag{3.6}$$

for any bounded Borel function $\phi : V \to [0,\infty)$ with compact support.

Remark 3.9 (Example of a surface with a conical singularity)**.** The notion of forward mean curvature bound of Definition 3.8 should be compared with the recent related definition proposed by Ketterer [55]. In the notation of [55], in order to have finite bound H_0 one needs an interior & exterior ball condition on V (equivalent, in the smooth setting, to a local L^∞ bound on the full second fundamental form), see [55, Remark 5.9]. The notion proposed above in Definition 3.8 instead works well even if the set V has corners or conical singularities.
For example, one can check that the set

$$V = \{(x,t) \subset \mathbb{R}^{n,1} : t = \alpha|x|\}, \quad \alpha \in (0,1),$$

in the $(n+1)$-dimensional Minkowski space-time $\mathbb{R}^{n,1}$ is an achronal topological hypersurface, smooth outside of the origin (where it is Lipschitz) and having forward mean curvature bounded above by $H_0 = 0$ in the sense of Definition 3.8. Notice that for any compact subset, one could choose the upper bound on the mean curvature to be strictly negative, but such an upper bound approaches zero as $|x| \to \infty$.

3.3 The results

Let us define $D_{H_0,K,N} > 0$ as follows:

$$D_{H_0,K,N} := \begin{cases} \frac{\pi}{2}\sqrt{\frac{N-1}{K}} & \text{if } K > 0, N > 1, H_0 = 0 \\ \sqrt{\frac{N-1}{K}}\cot^{-1}\left(\frac{-H_0}{\sqrt{K(N-1)}}\right) & \text{if } K > 0, N > 1, H_0 \in \mathbb{R}\setminus\{0\} \\ -\frac{N-1}{H_0} & \text{if } K = 0, N > 1, H_0 < 0 \\ \sqrt{-\frac{N-1}{K}}\coth^{-1}\left(\frac{-H_0}{\sqrt{-K(N-1)}}\right) & \text{if } K < 0, N > 1, H_0 < -\sqrt{-K(N-1)}. \end{cases} \tag{3.7}$$

Then the following synthetic version of the Hawking singularity theorem holds true.

Theorem 3.10 (Hawking singularity Theorem for $\mathsf{TMCP}_p^e(K,N)$ spaces)**.** *Let $(X,\mathsf{d},\mathfrak{m},\ll,\leq,\tau)$ be a timelike non-branching, locally causally closed, $\mathcal{K}$-globally hyperbolic, Lorentzian geodesic space satisfying $\mathsf{TMCP}_p^e(K,N)$ for some $p \in (0,1)$, $K \in \mathbb{R}$, $N \in [1,\infty)$ and assume that the causally-reversed structure satisfies the same conditions.*
Let $V \subset X$ be a Borel achronal FTC subset having forward mean curvature bounded above by H_0 in the sense of Definition 3.8. If

1. *$K > 0$, $N > 1$ and $H_0 \in \mathbb{R}$, or*

2. *$K = 0$, $N > 1$ and $H_0 < 0$, or*

3. *$K < 0$, $N > 1$ and $H_0 < \sqrt{-K(N-1)} < 0$,*

then for every $x \in I^+(V)$ it holds $\tau_V(x) \leq D_{H_0,K,N}$. In particular, for every timelike geodesic $\gamma \in$ TGeo(X) with $\gamma_0 \in V$, the maximal (on the right) domain of definition is contained in $\left[0, D_{H_0,K,N}\right]$. In case $N = 1$, $H_0 < 0$, it holds that $I^+(V) = \emptyset$.

Sketch of the proof. The proof of Theorem 3.10 follows the following argument: from the assumption on the mean curvature of V one can deduce with some effort some estimates on the logarithmic derivative of the densities of the disintegration formula $\log h'$. Then, from the disintegration formula and the localization of the curvature bounds, we had that $(X_\alpha, |\cdot|, h(\alpha, \cdot)\mathcal{L}^1)$ verifies MCP(K, N). Putting together this two facts yields bounds on the length of X_α and giving the claim. □

For completeness, we also report the direct application of Theorem 3.10 to the case of continuous, causally plain metrics. Recall that causally plain metrics are a subclass (detected by Chrúsciel-Grant [27]) of the space of continuous metrics, introduced to avoid pathological causal behaviours (such as null curves emanating from a point covering a set with non-empty interior, a phenomenon called "bubbling" and that may happen for Hölder continuous metrics).

Corollary 3.11 (Hawking singularity Theorem for a C^0-Lorentzian metric)**.** *Let (M, g) be a $2 \leq n$-dimensional timelike non-branching, globally hyperbolic, causally plain spacetime with a C^0-Lorentzian metric satisfying TMCP$_p^e(K, N)$ for some $p \in (0,1), K \in \mathbb{R}, N \in (1, \infty)$. Assume the causally-reversed structure to satisfy the same conditions.*
Let $V \subset M$ be a Borel achronal FTC subset (or, more strongly, let V be a Cauchy hypersurface) having forward mean curvature bounded above by $H_0 < 0$ (in the synthetic sense). If

1. *$K > 0$, $N > 1$ and $H_0 \in \mathbb{R}$, or*

2. *$K = 0$, $N > 1$ and $H_0 < 0$, or*

3. *$K < 0$, $N > 1$ and $H_0 < \sqrt{-K(N-1)} < 0$,*

Then for every $x \in I^+(V)$ it holds $\tau_V(x) \leq D_{H_0,K,N}$. In particular, for every timelike geodesic $\gamma \in$ Geo(M) with $\gamma_0 \in V$, the maximal (on the right) domain of definition is contained in $[0, D_{H_0,K,N}]$. Hence (M, g) is not timelike geodesically complete.

Remark 3.12. [Literature about Hawking singularity Theorem] Hawking singularity Theorem was established in [49, Thm. 4, p. 272] for smooth space-times (the proof works for C^2 metrics) assuming that V is a *compact* spacelike slice. The result was extended to $C^{1,1}$ metrics in [59] and to C^1 metrics in [40], by approximating the metric of low regularity with smoother metrics. The extension to non-compact future causally complete V was achieved in [36, Thm. 3.1] (see also [42]) in the smooth setting, and extended to $C^{1,1}$ metrics in [39]. Theorem 3.10 and Corollary 3.11, already in the smooth setting, relax the future causal completeness with the weaker future timelike completeness (in addition to extend the results to a synthetic framework, including C^0 metrics). Hawking (as well as Penrose and Hawking-Penrose) singularity Theorem was also extended to (smooth) closed cone structures [73] and smooth weighted Lorentz-Finsler manifolds [66].

A first synthetic singularity theorem was recently shown in [2] under the stronger assumptions that the space is a synthetic warped product with lower bounds on the sectional curvature in the sense of comparison triangles (á la Alexandrov).

With similar techniques and ideas one can also obtain some important geometric inequalities. We report, only in the case of continuous metrics, the Bishop-Gromov inequality and the Bonnet-Myers upper bound, the very same statements (as well as other comparison results) hold more generally for timelike non-branching, locally causally closed, $\mathcal{K}$-globally hyperbolic, Lorentzian geodesic spaces satisfying TMCP$_p^e(K, N)$ for some $p \in (0,1)$, $K \in \mathbb{R}$, $N \in [1, \infty)$ and whose causally-reversed structure satisfies the same conditions (see [24, Sec. 5.3]).

In order to state the results precisely, let introduce some notation. Given $x \leq y \in X$ we set

$$\mathfrak{I}(x, y, t) := \{\gamma_t \,:\, \gamma \in \mathrm{Geo}(X),\, \gamma_0 = x,\, \gamma_1 = y\}.$$

Fix $x_0 \in X$ and let

$$B^\tau(x_0, r) := \{x \in I^+(x_0) \cup \{x_0\} : \tau(x_0, x) < r\}$$

be the τ-ball of radius r and center x_0. Since typically the volume of a τ-ball is infinite (e.g. in Minkowski space it is the region below an hyperboloid), it is useful to localise volume estimates using star-shaped sets. To this aim, we say that $E \subset I^+(x_0) \cup \{x_0\}$ is τ-star-shaped with respect to x_0 if $\mathfrak{I}(x_0, x, t) \subset E$ for every $x \in E$ and $t \in (0, 1]$. Denote by

$$v(E, r) := \mathfrak{m}(B^\tau(x_0, r) \cap E)$$

the volume of the τ-ball of radius r intersected with the compact subset $E \subset I^+(x_0) \cup \{x_0\}$, τ-star-shaped with respect to x_0.

Corollary 3.13 (Timelike Bishop-Gromov and Bonner-Myers Theorems). *Let (M, g) be a $2 \leq n$-dim. timelike non-branching, globally hyperbolic, causally plain spacetime with C^0-metric satisfying $\mathsf{TMCP}_p^e(K, N)$. Assume the causally-reversed structure satisfes the same conditions. Then the following hold:*

- Timelike Bishop-Gromov inequality*: for every $x_0 \in M$, every compact subset $E \subset I^+(x_0) \cup \{x_0\}$ τ-star-shaped with respect to x_0 and every $r < R \leq \pi\sqrt{\frac{N-1}{K \vee 0}}$, it holds*

$$\frac{v(E, r)}{v(E, R)} \geq \frac{\int_0^r \mathfrak{s}_{K/(N-1)}(t)^{N-1}\mathrm{d}t}{\int_0^R \mathfrak{s}_{K/(N-1)}(t)^{N-1}\mathrm{d}t}.$$

- Timelike Bonnet-Myers upper bound*: if $K > 0$, then $\sup_{x,y \in X} \tau(x, y) \leq \pi\sqrt{(N-1)/K}$. In particular, for any causal curve γ it holds $\mathrm{L}_\tau(\gamma) \leq \pi\sqrt{(N-1)/K}$ and (M, g) is not timelike geodesically complete.*

4 A synthetic formulation of Einstein's vacuum equations

The goal of this section is to review the synthetic formulation of Einstein's vacuum equations (i.e. zero stress-energy tensor $T \equiv 0$ but possibly non-zero cosmological constant Λ) proposed in [75, App. B], as a natural outcome of the combination of [75] and [24].

The basic idea is that the Einstein vacuum equations (say with zero cosmological constant for convenience of presentation) can be characterised by the condition

$$\mathrm{Ric}_g(v, v) = 0 \qquad \text{for every timelike vector } v \in TM. \tag{4.1}$$

Of course, (4.1) is equivalent to require both the timelike Ricci lower bound

$$\mathrm{Ric}_g(v, v) \geq 0 \qquad \text{for every timelike vector } v \in TM \tag{4.2}$$

and the timelike Ricci upper bound

$$\mathrm{Ric}_g(v, v) \leq 0 \qquad \text{for every timelike vector } v \in TM. \tag{4.3}$$

The idea is to characterise both bounds separately in a synthetic way, and to ask the validity of both the corresponding synthetic characterizations in the non-smooth setting.

The synthetic timelike Ricci lower bound (4.2) was already characterised in a synthetic way via optimal transport in Section 2. Below we treat the upper bound (4.3).

4.1 Synthetic time-like Ricci upper bounds

The following definition is inspired by Sturm's approach to Riemannian/metric Ricci curvature upper bounds [88] and by the characterization of smooth Lorentzian of timelike Ricci upper bounds obtained in [75, Thm 4.7, Rem. 4.8]. We denote the metric ball in (X, d) with center x and radius r by $B_r^{\mathsf{d}}(x)$.

Definition 4.1 (Synthetic time-like Ricci upper bounds)**.** Fix $p \in (0, 1)$, $K \in \mathbb{R}$. We say that a measured Lorentzian pre-length space $(X, \mathsf{d}, \mathfrak{m}, \ll, \leq, \tau)$ has time-like Ricci curvature bounded above by K in a synthetic sense if there exists $r_0 > 0$ and a function $\omega : [0, r_0) \to [0, \infty)$ with $\lim_{r \downarrow 0} \omega(r) = 0$ such that for every $r \in [0, r_0)$ the following holds.

- For every $x, y \in X$ with $\mathsf{d}(x, y) = r > 0$ and such that $B_{r^4}^{\mathsf{d}}(x) \times B_{r^2}^{\mathsf{d}}(y) \subset X_{\ll}^2$,

- for every $\mu_0 \in \mathrm{Dom}(\mathrm{Ent}(\cdot|\mathfrak{m}))$ with $\mathrm{supp}\,\mu_0 \subset B_{r^4}^{\mathsf{d}}(x)$,

there exists an ℓ_p-geodesic $(\mu_t)_{t \in [-1,1]}$ satisfying

- $\mathrm{supp}\,\mu_1 \subset B_{r^2}^{\mathsf{d}}(y)$,

- $\mathrm{supp}\,\mu_{-1} \times \mathrm{supp}\,\mu_1 \subset X_{\leq}^2$,

- $\bigcup_{t \in [-1,1]} \mathrm{supp}\,\mu_t \subset B_{10r_0}^{\mathsf{d}}(x)$,

- $\mathrm{Ent}(\mu_{-1}|\mathfrak{m}) - 2\mathrm{Ent}(\mu_0|\mathfrak{m}) + \mathrm{Ent}(\mu_1|\mathfrak{m}) \leq (K + \omega(r))\, r^2$.

A key property of the above notion of time-like Ricci bounded above is the stability under weak convergence of measured Lorentzian pre-length spaces. The stability of Riemannian/metric Ricci upper bounds via optimal transport was proved by Sturm [88]. The notion of convergence we use below is a slight reinforcement (we ask that the immersion maps are isometries with respect to the metric structures instead of merely topological embedding maps) of the weak convergence used in Section 2.2 to show stability of synthetic timelike Ricci *lower* bounds; in any case it is a natural adaptation to the Lorentzian setting of the mGH convergence used for metric measure spaces (see for instance [38, Sec. 3] for an overview of equivalent formulations of convergence for metric measure spaces).

Theorem 4.2 (Stability of time-like Ricci curvature upper bounds)**.** *Let $\{(X_j, \mathsf{d}_j, \mathfrak{m}_j, \ll_j, \leq_j, \tau_j)\}_{j \in \mathbb{N} \cup \{\infty\}}$ be a sequence of measured Lorentzian geodesic spaces satisfying the following properties :*

1. *There exists a locally causally closed, globally hyperbolic Lorentzian geodesic space $(X, \mathsf{d}, \ll, \leq, \tau)$ such that each $(X_j, \mathsf{d}_j, \mathfrak{m}_j, \ll_j, \leq_j, \tau_j)$, $j \in \mathbb{N} \cup \{\infty\}$, is isomorphically embedded in it, i.e. there exist inclusion maps $\iota_j : X_j \hookrightarrow X$ such that for every $x_j^1, x_j^2 \in X_j$, for every $j \in \mathbb{N} \cup \{\infty\}$, the following holds:*

   - $\mathsf{d}(\iota_j(x_j^1), \iota_j(x_j^2)) = \mathsf{d}_j(x_j^1, x_j^2)$;
   - $x_j^1 \leq_j x_j^2$ *if and only if* $\iota_j(x_j^1) \leq \iota_j(x_j^2)$;
   - $\tau(\iota_j(x_j^1), \iota_j(x_j^2)) = \tau_j(x_j^1, x_j^2)$.

2. *The measures $(\iota_j)_\sharp \mathfrak{m}_j$ converge to $(\iota_\infty)_\sharp \mathfrak{m}_\infty$ weakly in duality with $C_c(X)$ in X, i.e.*

$$\int \varphi\, (\iota_j)_\sharp \mathfrak{m}_j \to \int \varphi\, (\iota_\infty)_\sharp \mathfrak{m}_\infty, \quad \forall \varphi \in C_c(X), \tag{4.4}$$

   where $C_c(X)$ denotes the set of continuous functions with compact support.

3. *Volume non-collapsing: there exists a function $v : (0, \infty) \to (0, \infty)$ such that for every $x_j \in X_j$ it holds $\mathfrak{m}_j(B_r^{\mathsf{d}_j}(x_j)) \geq v(r) > 0$.*

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4. *There exists a function* $\omega : [0, \infty) \to [0, \infty)$ *with* $\lim_{r \downarrow 0} \omega(r) = 0$ *and there exist* $p \in (0, 1)$, $K \in \mathbb{R}$ *such that* $(X_j, \mathsf{d}_j, \mathfrak{m}_j, \ll_j, \leq_j, \tau_j)$ *has time-like Ricci curvature bounded above by* K *with respect to* $p \in (0, 1)$, $r_0 > 0$ *and with remainder function* ω *in the synthetic sense of Definition 4.1.*

Then also the limit space $(X_\infty, \mathsf{d}_\infty, \mathfrak{m}_\infty, \ll_\infty, \leq_\infty, \tau_\infty)$ *has time-like Ricci curvature bounded above by* K *with respect to* $p \in (0, 1), r_0 + 1$ *and with remainder function* ω *in the synthetic sense of Definition 4.1.*

The proof can be performed in the same spirit as the proof of Theorem 2.6, the interested reader is referred to [75, Thm B.6].

4.2 Synthetic Einstein's vacuum equations

Combining the synthetic upper and lower bounds on the time-like Ricci curvature, i.e. Definitions 2.1 and 4.1, the following synthetic version for the vacuum Einstein's equations (with possibly non-zero cosmological constant) is rather natural.

Definition 4.3 (Synthetic vacuum Einstein's equations). Fix $p \in (0, 1)$, $\Lambda \in \mathbb{R}$, $N \in (0, \infty]$. We say that the measured Lorentzian pre-length space $(X, \mathsf{d}, \mathfrak{m}, \ll, \leq, \tau)$ satisfies the (resp. weak) synthetic formulation of the vacuum Einstein equations Ric $\equiv \Lambda$ with cosmological constant $\Lambda \in \mathbb{R}$ and has synthetic dimension $\leq N$ if

- $(X, \mathsf{d}, \mathfrak{m}, \ll, \leq, \tau)$ satisfies the $\mathsf{TCD}_p^e(\Lambda, N)$ (resp. $\mathsf{wTCD}_p^e(\Lambda, N)$) condition;

- There exists $r_0 > 0$ and a function $\omega : [0, r_0) \to [0, \infty)$ with $\lim_{r \downarrow 0} \omega(r) = 0$ such that $(X, \mathsf{d}, \mathfrak{m}, \ll, \leq, \tau)$ has time-like Ricci curvature bounded above by Λ, with respect to $p \in (0, 1), r_0$ and ω.

Combining the stability of time-like Ricci lower and upper bounds (i.e. Theorem 2.6 and Theorem 4.2) gives the stability of the synthetic vacuum Einstein's equations under the aforementioned natural Lorentzian variant of measured Gromov-Hausdorff convergence.

Theorem 4.4 (Weak stability of synthetic vacuum Einstein's equations)**.** *Let* $\{(X_j, \mathsf{d}_j, \mathfrak{m}_j, \ll_j, \leq_j, \tau_j)\}_{j \in \mathbb{N} \cup \{\infty\}}$ *be a sequence of measured Lorentzian geodesic spaces satisfying the following properties:*

1. *There exists a locally causally closed, $\mathcal{K}$-globally hyperbolic Lorentzian geodesic space* $(X, \mathsf{d}, \ll, \leq, \tau)$ *such that each* $(X_j, \mathsf{d}_j, \mathfrak{m}_j, \ll_j, \leq_j, \tau_j)$, $j \in \mathbb{N} \cup \{\infty\}$, *is isomorphically embedded in it (as in (1) of Theorem 4.2).*

2. *The measures* $(\iota_j)_\sharp \mathfrak{m}_j$ *converge to* $(\iota_\infty)_\sharp \mathfrak{m}_\infty$ *weakly in duality with* $C_c(X)$ *in* X, *i.e.* (4.4) *holds.*

3. *Volume non-collapsing: there exists a function* $v : (0, \infty) \to (0, \infty)$ *such that for every* $x_j \in X_j$ *it holds* $\mathfrak{m}_j(B_r^{\mathsf{d}_j}(x_j)) \geq v(r) > 0$.

4. *There exist* $p \in (0, 1), \Lambda \in \mathbb{R}, N \in (0, \infty], r_0 > 0$ *and* $\omega : [0, r_0) \to [0, \infty)$ *with* $\lim_{r \downarrow 0} \omega(r) = 0$ *such that* $(X_j, \mathsf{d}_j, \mathfrak{m}_j, \ll_j, \leq_j, \tau_j)$ *satisfies the synthetic formulation of the vacuum Einstein equations* Ric $\equiv \Lambda$ *with cosmological constant* $\Lambda \in \mathbb{R}$, *with synthetic dimension* $\leq N$ *with respect to* $p \in (0, 1), r_0$ *and* ω *as in Definition 4.3.*

Then the limit space $(X_\infty, \mathsf{d}_\infty, \mathfrak{m}_\infty, \ll_\infty, \leq_\infty, \tau_\infty)$ *satisfies the weak synthetic formulation of the vacuum Einstein equations* Ric $\equiv \Lambda$ *with cosmological constant* $\Lambda \in \mathbb{R}$, *with synthetic dimension* $\leq N$ *with respect to* $p \in (0, 1), r_0 + 1$ *and* ω *as in Definition 4.3.*

Remark 4.5. By [75, Thm 4.7, Rem. 4.8] and Theorem 1.16, if $(X, \mathsf{d}, \mathfrak{m}, \ll, \leq, \tau)$ is a (for simplicity say a compact subset in a) smooth Lorentzian manifold, then $(X, \mathsf{d}, \mathfrak{m}, \ll, \leq, \tau)$ satisfies the Einstein's equations Ric $\equiv \Lambda$ in the smooth classical sense if and only if $(X, \mathsf{d}, \mathfrak{m}, \ll, \leq, \tau)$ satisfies the Einstein's equations in the synthetic sense of Definition 4.3. Therefore, Theorem 4.4 gives that the corresponding limits of smooth solutions to Einstein's equation Ric $\equiv \Lambda$ satisfy the weak synthetic Einstein's equations Ric $\equiv \Lambda$ in the sense of Definition 4.3. In other terms, the vacuum Einstein's equations are stable under the conditions (and with respect to the notion of convergence) of Theorem 4.4.

Let us mention that the stability of the Einstein's equations under various notions of (weak) convergence is a subject of high interest in general relativity.
Classically, the problem is stated in terms of convergence of a sequence of Lorentzian metrics g_j converging to a limit Lorentzian metric g_∞, *on a fixed underlying manifold.*
If g_j solve the vacuum Einstein equations, $g_j \to g_\infty$ in C^0_{loc} and the derivatives of g_j converge in L^2_{loc}, then it is well known that the limit g_∞ satisfies the vacuum Einstein equations as well.
However, if the $g_j \to g_\infty$ in C^0_{loc} and the derivatives of g_j converge *weakly* in L^2_{loc}, explicit examples have been constructed (see [20, 43] for examples in symmetry classes) where the limit g_∞ may satisfy the Einstein equations with a *non-vanishing* stress energy momentum tensor. Burnett [20] conjectured that, if there exist $C > 0$ and $\lambda_j \to 0$ such that

$$ |g_j - g_\infty| \leq \lambda_j, \quad |\partial g_j| \leq C, \quad |\partial^2 g_j| \leq C \lambda_j^{-1}, $$

then g_∞ is isometric to a solution to the Einstein-massless Vlasov system for some suitable Vlasov field. Such a conjecture remains open, although there has been recent progress [51, 52, 46] under symmetry conditions. We also mention the recent work [67] where concentrations (at the level of ∂g_j) are allowed in addition to oscillations.

Theorem 4.4 gives a new point of view on the stability of the vacuum Einstein's equations. Indeed, while in the aforementioned results the metrics g_j are converging *on a fixed underlying manifold*, in Theorem 4.4 also the underlying space X may vary (along the sequence and at the limit), allowing change in topology in the limit, as one may expect in case of formation of singularities. Moreover, the notion of convergence is quite different in spirit: while in the aforementioned results $g_j \to g_\infty$ in a suitable *functional analytic sense*, in Theorem 4.4 the spaces are converging in a *more geometric sense* (inspired by the measured Gromov-Haudorff convergence).

5　Some possible research directions and open problems

Due to the richness of the theory of metric measure spaces satisfying the Lott-Sturm-Villani Curvature-Dimension condition $\mathsf{CD}(K, N)$, it is natural to expect that the Lorentzian counterpart surveyed here will enjoy a rich theory as well. In this section we propose some possible research directions and open problems, without any attempt to be exhaustive; on the contrary, we believe that Poincaré's famous sentence "surprising results shall be obtained" is appropriate as ever.

Convergence of spaces and pre-compactess

A fundamental property of the class of $\mathsf{CD}(K, N)$ spaces (say with a bound on the diameter for simplicity) is the compactness under measured Gromov-Hausdorff (mGH for short) convergence. This is a consequence of two deep results: Gromov pre-compactness theorem [45] and the stability of the $\mathsf{CD}(K, N)$ condition under mGH convergence [86, 87, 65].
In Sections 2.2, 4.1, 4.2, we proposed possible notions of convergence for Lorentzian synthetic spaces and proved stability results for synthetic Ricci bounds (and for a synthetic characterization of the Einstein's vacuum equations), thus ticking the latter key property.

It is an open problem whether the class of $\mathsf{TCD}_p^e(K, N)$ spaces (say with bounded diameter) satisfy some (pre)-compactness result in the same spirit of Gromov pre-compactness theorem (with respect to the convergence proposed in Definition 2.5, or some other suitable notion of convergence).

A key difficulty is that while the CD/MCP conditions for metric measure spaces imply a control on the volume growth of metric balls (more precisely a local volume doubling property), in our setting the τ-balls typically have infinite volume (e.g., in Minkowski space, τ-spheres are hyperboloids). Thus we cannot expect to have compactness in classical pointed-measured-Gromov-Hausdorff topology (which is thus not anymore the clearly natural notion for weak convergence of spaces).

Eulerian Vs Lagrangian approach

The point of view on synthetic Ricci bounds employed in this presentation has been *Lagrangian*, i.e. we analysed the convexity/concavity properties of the entropy along suitable Lorentz-Wasserstein geodesics. In the smooth setting (as well in the metric measure setting, though the non-smoooth framework is more delicate) such a Lagrangian point of view is equivalent to an *Eulerian* approach based on the Bochner inequality (which is also known as Bakry-Émery Γ-calculus), see [5, 6, 7, 31, 8]. It would be very interesting to develop an Eulerian approach to timelike Ricci curvature bounds, and possibly prove the equivalence with the Lagrangian point of view (under suitable assumptions, for instance a natural Lorentzian counterpart of the assumption that the Cheeger energy is a quadratic form).

Building such a bridge between Eulerian and Lagrangian approaches will very likely open the door to a Lorentzian theory for gradient flows, which has been very rich in the smooth Riemannian (and non-smooth metric-measure) setting (see for instance the monograph [4], the aforementioned [5, 6, 7, 31, 8] and the more recent survey [3]).

Compatibility with synthetic sectional curvature lower bounds

In the smooth setting, it is clear that non-negative (or more generally a lower bound on the) sectional curvature implies non-negative (or more generally a lower bound on the) Ricci curvature, as the latter is obtained as a trace of the former.

Synthetic lower bounds on the timelike sectional curvature in terms of comparison triangles (á la Alexandrov) for Lorentzian synthetic spaces have been proposed ans studied in [58] (see also the more recent [10]). It is a natural open problem whether such synthetic lower bounds on the timelike sectional curvature imply the synthetic lower bounds on the timelike Ricci curvature surveyed here. Note that while the former is purely metric, the latter needs to fix a reference volume measure in order to be formulated. Natural candidates are the canonical volume measures on Lorentzian (pre-)length spaces constructed in [72].

Null synthetic Ricci curvature lower bounds

In the present survey, we only treated *timelike* Ricci curvature lower bounds. This was motivated by the applications we had in mind, namely Hawking singularity Theorem, time-like Bishop-Gromov, time-like Bonnet Myers (see Section 3, or [24] for a more complete list) as well as a synthetic formulation of Einstein's equations (see Section 4, or [75] for an optimal transport characterization of the full Einstein's equations possibly with a non-zero stress-energy tensor).

It is however also an interesting direction to find a synthetic characterization of *null* Ricci curvature lower bounds, i.e. $\mathrm{Ric}_g(v, v) \geq K$ for all *null* vectors $v \in TM$. Indeed this is at the heart of Penrose's work in general relativity: it is one of the assumptions of his celebrated singularity theorem [79] which was awarded the Nobel prize, as well as in much of his work (see for instance [60, Sec. 1.9, 6.3, 6.4] or [61]).

Independence of the theory on $p \in (0,1]$

Note that the definition of the $\mathsf{TCD}_p^e(K,N)$ condition (Definition 2.1) depends on the exponent $p \in (0,1]$ chosen to fix the optimal transport cost and hence the geodesic structure on the space of probability measures. In the metric measure setting (for essentially non-branching metric measure spaces) it was recently proved in [1] that the $\mathsf{CD}(K,N)$ theory is independent of the exponent chosen to metrize the space of probability measures. Arguing along the lines of [1], which in turns builts on top of [15, 16, 22, 23], it should be possible to prove the analogous independence of p in the Lorentzian synthethic setting (under suitable timelike non-branching and causal assumptions).

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