# DOMAIN VARIATIONS OF THE FIRST EIGENVALUE VIA A STRICT FABER-KRAHN TYPE INEQUALITY 

T. V. ANOOP AND K. ASHOK KUMAR


#### Abstract

For $d \geq 2$ and $\frac{2 d+2}{d+2}<p<\infty$, we prove a strict Faber-Krahn type inequality for the first eigenvalue $\lambda_{1}(\Omega)$ of the $p$-Laplace operator on a bounded Lipschitz domain $\Omega \subset \mathbb{R}^{d}$ (with mixed boundary conditions) under the polarizations. We apply this inequality to the obstacle problems on the domains of the form $\Omega \backslash \mathcal{O}$, where $\mathcal{O} \subset \subset \Omega$ is an obstacle. Under some geometric assumptions on $\Omega$ and $\mathcal{O}$, we prove the strict monotonicity of $\lambda_{1}(\Omega \backslash \mathcal{O})$ with respect to certain translations and rotations of $\mathcal{O}$ in $\Omega$.


## 1. Introduction

In 1877, Lord Rayleigh [26] conjectured that 'the disk is the only planar domain that minimizes the first Dirichlet eigenvalue of the Laplace operator among all planar domains of fixed area.' Nearly after 45 years, this conjecture was proved by Faber [18] and Krahn [23] for the planar domains (in 1923), and it is extended for higher dimensional domains by Krahn [24] (in 1925). This result is known as the Faber-Krahn inequality which is also available for the first Dirichlet eigenvalue of the $p$-Laplace operator $\Delta_{p}$, defined by $\Delta_{p} u=\operatorname{div}\left(|\nabla u|^{p-2} \nabla u\right)$ with $p \in(1, \infty)$, see for example [25, page 191] and [20, II.4]. For a domain $\Omega \subset \mathbb{R}^{d}$, the Faber-Krahn inequality states that

$$
\begin{equation*}
\lambda_{1}\left(\Omega^{*}\right) \leq \lambda_{1}(\Omega) \tag{1.1}
\end{equation*}
$$

where $\lambda_{1}(D)$ denotes the first Dirichlet eigenvalue of the $p$-Laplace operator on a domain $D$ and $\Omega^{*}$ is the open ball centred at the origin in $\mathbb{R}^{d}$ with the same Lebesgue measure as that of $\Omega$. If $\Omega$ is a ball, then the equality holds in (1.1). The question 'for which domains the strict inequality holds in (1.1)?' is addressed in [2, 9, 14, 16, 21].

Noting that $\Omega^{*}$ is the Schwarz symmetrization of $\Omega$, the inequality (1.1) asserts that the first Dirichlet eigenvalue decreases under the Schwarz symmetrization. Next, we see that a similar result easily holds under the polarization as well. The polarization is one of the simplest rearrangements on $\mathbb{R}^{d}$ that was first introduced for sets by Wolontis [32], and for functions by Ahlfors [1] (for $d=2$ ) and Baernstein and Taylor [6] (for $d \geq 2$ ). We refer to [ $5,10,12,28,31]$ for further reading on polarizations and their applications. Now, we define the polarization of measurable sets and functions with respect to an open affine-halfspace in $\mathbb{R}^{d}$. Let $H$ be an open affine-halfspace in $\mathbb{R}^{d}$ (called a polarizer), and let $\sigma_{H}$ be the reflection with respect to the boundary $\partial H$ in $\mathbb{R}^{d}$. We denote the set of all polarizers in $\mathbb{R}^{d}$ by $\mathcal{H}$.

Definition 1.1 (Polarization). Let $H \in \mathcal{H}$ and $\Omega \subseteq \mathbb{R}^{d}$. The polarization $P_{H}(\Omega)$ and the dual-polarization $P^{H}(\Omega)$ of $\Omega$ with respect to $H$ are defined as:

$$
\begin{aligned}
P_{H}(\Omega) & =\left[\left(\Omega \cup \sigma_{H}(\Omega)\right) \cap H\right] \cup\left[\Omega \cap \sigma_{H}(\Omega)\right], \\
P^{H}(\Omega) & =\left[\left(\Omega \cup \sigma_{H}(\Omega)\right) \cap H^{c}\right] \cup\left[\Omega \cap \sigma_{H}(\Omega)\right] .
\end{aligned}
$$

For a measurable function $u: \mathbb{R}^{d} \longrightarrow \mathbb{R}$, the polarization $P_{H}(u)$ with respect to $H$ is defined as

$$
P_{H}(u)(x)= \begin{cases}\max \left\{u(x), u\left(\sigma_{H}(x)\right)\right\}, & \text { for } x \in H, \\ \min \left\{u(x), u\left(\sigma_{H}(x)\right)\right\}, & \text { for } x \in \mathbb{R}^{d} \backslash H .\end{cases}
$$

Now, for $u: \Omega \longrightarrow \mathbb{R}$ let $\widetilde{u}$ be the zero extension of $u$ to $\mathbb{R}^{d}$. The polarization $P_{H}(u)$ is defined as the restriction of $P_{H}(\widetilde{u})$ to $P_{H}(\Omega)$.

[^0]Remark 1.2. The polarization of the sets and the functions satisfy the following relation:

$$
P_{H}\left(\mathbb{1}_{\Omega}\right)=\mathbb{1}_{P_{H}(\Omega)}, \text { for any } \Omega \subseteq \mathbb{R}^{d}
$$

where $\mathbb{1}_{\Omega}$ denotes the characteristic function of $\Omega$.
In Figure 1, the dark shaded regions on the right side represent the polarization $P_{H}(\Omega)$ of $\Omega$ with respect to $H$.


Figure 1. Polarization of an ellipse and a square.

For $H \in \mathcal{H}$, the polarization $P_{H}$ is a rearrangement (preserves the inclusion order and the measure) on $\mathbb{R}^{d}$, see $\left[13\right.$, Section 3]. Further, $P_{H}$ takes an open set to an open set and a closed set to a closed set in $\mathbb{R}^{d}$. Throughout this article, we consider $p \in(1, \infty)$, unless otherwise specified. For a non-negative function $u \in W_{0}^{1, p}(\Omega)$ the polarization $P_{H}(u) \in W_{0}^{1, p}\left(P_{H}(\Omega)\right)$ and the norms are preserved, see [13, Corollary 5.1]:

$$
\|u\|_{p, \Omega}=\left\|P_{H}(u)\right\|_{p, P_{H} \Omega} \text { and }\|\nabla u\|_{p, \Omega}=\left\|\nabla P_{H}(u)\right\|_{p, P_{H} \Omega} .
$$

Therefore, we have the equality in the Pólya-Szëgo type inequality for the polarizations on $\mathbb{R}^{d}$. As an immediate consequence, the variational characterization of $\lambda_{1}(\Omega)$ yields the following Faber-Krahn type inequality:

$$
\begin{equation*}
\lambda_{1}\left(P_{H}(\Omega)\right) \leq \lambda_{1}(\Omega) \tag{1.2}
\end{equation*}
$$

Clearly, if $P_{H}(\Omega)=\Omega$ or $P_{H}(\Omega)=\sigma_{H}(\Omega)$ then the equality holds in (1.2). In this article, we identify the domains for which the strict inequality holds in (1.2) for certain values of $p$. More precisely, we show that, if $p>\frac{2 d+2}{d+2}$ and the equality holds in (1.2) then $P_{H}(\Omega)=\Omega$ or $P_{H}(\Omega)=\sigma_{H}(\Omega)$. We prove this result for the first eigenvalue of the $p$-Laplace operator with mixed boundary conditions on the multiply connected domains of the following form:
$\left(\mathbf{A}_{\mathbf{0}}\right) \Omega_{\mathrm{out}} \backslash \overline{\Omega_{\mathrm{in}}} \subset \mathbb{R}^{d}$ is a bounded Lipschitz domain with $\Omega_{\mathrm{in}} \subset \subset \Omega_{\mathrm{out}}$, and $\Omega_{\mathrm{in}}=\bigcup_{j=1}^{m} \Omega_{j}$, where $\Omega_{j}$ is simply connected and $\overline{\Omega_{i}} \cap \overline{\Omega_{j}}=\emptyset$ for $i, j \in\{1,2, \ldots, m\}$ with $i \neq j$.
For $\Omega_{\text {out }} \backslash \overline{\Omega_{\text {in }}}$ as in $\left(\mathbf{A}_{\mathbf{0}}\right)$, we consider the following family of admissible polarizers

$$
\mathcal{H}_{\mathrm{ad}}:=\left\{H \in \mathcal{H}: \sigma_{H}\left(\Omega_{\mathrm{in}}\right) \subset \subset \Omega_{\mathrm{out}}\right\} .
$$

Since $\Omega_{\mathrm{in}} \subset \subset \Omega_{\mathrm{out}}$, the set $\mathcal{H}_{\mathrm{ad}}$ is always non-empty, and for $H \in \mathcal{H}_{\mathrm{ad}}$ we have (see Proposition 2.12):

$$
P_{H}\left(\Omega_{\mathrm{out}} \backslash \overline{\Omega_{\mathrm{in}}}\right)=P_{H}\left(\Omega_{\mathrm{out}}\right) \backslash P^{H}\left(\overline{\Omega_{\mathrm{in}}}\right) \text { and } \partial P_{H}\left(\Omega_{\mathrm{out}} \backslash \overline{\Omega_{\mathrm{in}}}\right)=\partial P_{H}\left(\Omega_{\mathrm{out}}\right) \sqcup \partial P^{H}\left(\Omega_{\mathrm{in}}\right)
$$

For $H \in \mathcal{H}_{\text {ad }}$, we consider the following eigenvalue problems for the $p$-Laplace operator on both $\Omega$ and $P_{H}(\Omega)$ with mixed boundary conditions:

Neumann condition on $\partial \Omega_{\mathrm{in}}$ :

$$
\left.\left.\left.\begin{array}{rlrl}
-\Delta_{p} u & =\nu|u|^{p-2} u \text { in } \Omega_{\mathrm{out}} \backslash \overline{\Omega_{\mathrm{in}}} \\
u & =0 \text { on } \partial \Omega_{\mathrm{out}}, \\
\frac{\partial u}{\partial n} & =0 \text { on } \partial \Omega_{\mathrm{in}} ;
\end{array}\right\} \begin{array}{rl}
-\Delta_{p} v & =\nu|v|^{p-2} v \text { in } P_{H}\left(\Omega_{\mathrm{out}} \backslash \overline{\Omega_{\mathrm{in}}}\right)  \tag{1.4}\\
v & =0 \text { on } \partial P_{H}\left(\Omega_{\mathrm{out}}\right) \\
\end{array}\right\} \quad(1.3) \quad \begin{array}{l}
\frac{\partial v}{\partial n}
\end{array}\right)=0 \text { on } \partial P^{H}\left(\Omega_{\mathrm{in}}\right) ;
$$

Neumann condition on $\partial \Omega_{\text {out }}$ :

$$
\left.\begin{array}{rl}
-\Delta_{p} u & =\tau|u|^{p-2} u \text { in } \Omega_{\mathrm{out}} \backslash \overline{\Omega_{\mathrm{in}}}  \tag{1.5}\\
u & =0 \text { on } \partial \Omega_{\mathrm{in}} \\
\frac{\partial u}{\partial n} & =0 \text { on } \partial \Omega_{\mathrm{out}} ;
\end{array}\right\}
$$

$$
\left.\begin{array}{rl}
-\Delta_{p} v & =\tau|v|^{p-2} v \text { in } P_{H}\left(\Omega_{\mathrm{out}} \backslash \overline{\Omega_{\mathrm{in}}}\right) \\
v & =0 \text { on } \partial P^{H}\left(\Omega_{\mathrm{in}}\right),  \tag{1.6}\\
\frac{\partial v}{\partial n} & =0 \text { on } \partial P_{H}\left(\Omega_{\mathrm{out}}\right),
\end{array}\right\}
$$

where $\nu, \tau \in \mathbb{R}$.
The above eigenvalue problems can be collectively expressed as the following problem:

$$
\left.\begin{array}{rl}
-\Delta_{p} u & =\gamma|u|^{p-2} u \text { in } \Omega, \\
u & =0 \text { on } \Gamma_{D},  \tag{E}\\
\frac{\partial u}{\partial n} & =0 \text { on } \Gamma_{N},
\end{array}\right\}
$$

where $\Omega \subset \mathbb{R}^{d}$ is a bounded Lipschitz domain with $\partial \Omega=\Gamma_{N} \sqcup \Gamma_{D}$, and $\gamma \in \mathbb{R}$. Let

$$
\mathcal{C}_{\Gamma_{D}}^{0,1}(\Omega):=\left\{v \text { is a Lipschitz continuous function on } \Omega \text { with } \operatorname{supp}(v) \cap \Gamma_{D}=\emptyset\right\}
$$

and we consider the following Sobolev space:

$$
W_{\Gamma_{D}}^{1, p}(\Omega)=\text { the closure of } \mathcal{C}_{\Gamma_{D}}^{0,1}(\Omega) \text { in } W^{1, p}(\Omega)
$$

If $\Gamma_{N}=\emptyset$ (equivalently $\Gamma_{D}=\partial \Omega$ ) then $W_{\Gamma_{D}}^{1, p}(\Omega)=W_{0}^{1, p}(\Omega)$. A real number $\gamma$ is said to be an eigenvalue of $(\mathcal{E})$ if there exists a non-zero function $u \in W_{\Gamma_{D}}^{1, p}(\Omega)$ such that

$$
\int_{\Omega}|\nabla u|^{p-2} \nabla u \cdot \nabla v \mathrm{~d} x-\gamma \int_{\Omega}|u|^{p-2} u v \mathrm{~d} x=0 \text { for every } v \in W_{\Gamma_{D}}^{1, p}(\Omega)
$$

and the function $u$ is called as an eigenfunction corresponding to the eigenvalue $\gamma$. The standard variational arguments establish the existence of an infinite subset of eigenvalues tending to infinity (see [3, Proposition A.1]). The first eigenvalue $\gamma_{1}(\Omega)$ of $(\mathcal{E})$ is simple (the dimension of the eigenspace is one) and the corresponding eignfunctions have constant sign in $\Omega$ (see [3, Proposition A.2]). Moreover, the first eigenvalue $\gamma_{1}(\Omega)$ of $(\mathcal{E})$ has the following variational characterization:

$$
\gamma_{1}(\Omega)=\inf \left\{\int_{\Omega}|\nabla u|^{p} \mathrm{~d} x: u \in W_{\Gamma_{D}}^{1, p}(\Omega) \text { with } \int_{\Omega}|u|^{p} \mathrm{~d} x=1\right\} .
$$

Now, we state a Faber-Krahn type inequality for the first eigenvalues of the eigenvalue problems (1.3) and (1.4), and similarly for the first eigenvalues of the eigenvalue problems (1.5) and (1.6).
Theorem 1.3. Let $p \in(1, \infty)$, $\Omega_{\text {out }} \backslash \overline{\Omega_{\mathrm{in}}} \subset \mathbb{R}^{d}$ be a domain as given in $\left(\mathbf{A}_{\mathbf{0}}\right)$, and $H \in \mathcal{H}_{\mathrm{ad}}$.
(i) If $\sigma_{H}\left(\Omega_{\mathrm{in}}\right)=\Omega_{\mathrm{in}}$, then

$$
\begin{equation*}
\nu_{1}\left(P_{H}\left(\Omega_{\text {out }} \backslash \overline{\Omega_{\mathrm{in}}}\right)\right) \leq \nu_{1}\left(\Omega_{\text {out }} \backslash \overline{\Omega_{\mathrm{in}}}\right) \tag{1.7}
\end{equation*}
$$

(ii) If $\Omega_{\mathrm{in}} \neq \emptyset$ and $\sigma_{H}\left(\Omega_{\mathrm{out}}\right)=\Omega_{\mathrm{out}}$, then

$$
\begin{equation*}
\tau_{1}\left(P_{H}\left(\Omega_{\text {out }} \backslash \overline{\Omega_{\text {in }}}\right)\right) \leq \tau_{1}\left(\Omega_{\text {out }} \backslash \overline{\Omega_{\text {in }}}\right) . \tag{1.8}
\end{equation*}
$$

(iii) If $\frac{2 d+2}{d+2}<p<\infty$, and the equality holds in (1.7) or (1.8) then

$$
P_{H}\left(\Omega_{\text {out }} \backslash \overline{\Omega_{\text {in }}}\right)=\Omega_{\text {out }} \backslash \overline{\Omega_{\text {in }}} \text { or } P_{H}\left(\Omega_{\text {out }} \backslash \overline{\Omega_{\text {in }}}\right)=\sigma_{H}\left(\Omega_{\text {out }} \backslash \overline{\Omega_{\text {in }}}\right) .
$$

## Remark 1.4.

(i) If $\Omega_{\mathrm{in}}=\emptyset$, in (i) then $\nu_{1}$ corresponds to the first Dirichlet eigenvalue $\lambda_{1}$ and thus (1.7) gives: for every $H \in \mathcal{H}_{\text {ad }}$,

$$
\lambda_{1}\left(P_{H}\left(\Omega_{\text {out }} \backslash \overline{\Omega_{\text {in }}}\right)\right) \leq \lambda_{1}\left(\Omega_{\text {out }} \backslash \overline{\Omega_{\text {in }}}\right) .
$$

(ii) If $\Omega_{\mathrm{in}}=\emptyset$, then $\tau_{1}\left(\Omega_{\mathrm{out}}\right)=\tau_{1}\left(P_{H}\left(\Omega_{\mathrm{out}}\right)\right)=0$, for every $H \in \mathcal{H}$. This is the reason why we impose the condition $\Omega_{\mathrm{in}} \neq \emptyset$ in (ii).
(iii) The symmetry assumptions in (i) and (ii) of Theorem 1.3 ensure that $\Gamma_{N} \subseteq \partial P_{H}(\Omega)$ and hence the Neumann boundary is unaltered under such polarizations. This fact is crucially used in our proof. Obtaining the same conclusions of Theorem 1.3 without these additional symmetry assumptions seems to be a challenging problem.

Application to the domain variations: Next, we apply Theorem 1.3 for the domains of the form $\Omega \backslash \mathcal{O} \subset$ $\mathbb{R}^{d}$ to study the monotonicity of the first eigenvalue of $(\mathcal{E})$ on $\Omega \backslash \mathcal{O}$ under certain translations and rotations of $\mathcal{O}$ within $\Omega$. We assume the following:
$\left(\mathbf{A}_{\mathbf{1}}\right) \mathcal{O} \subset \Omega$ is a closed set with nonempty interior such that $\Omega \backslash \mathcal{O}$ is a bounded Lipschitz domain.
The set $\mathcal{O}$ in $\left(\mathbf{A}_{\mathbf{1}}\right)$ is called as an obstacle. The main idea is to express the translations and the rotations of $\mathcal{O}$ in terms of polarizations of punctured domain $\Omega \backslash \bigcirc$. Then we apply Theorem 1.3 and get the monotonicity of the eigenvalue.

The monotonicity along a straight line: In this case, we set $\Omega_{\mathrm{in}}=\emptyset$ and $\Omega_{\mathrm{out}}=\Omega \backslash \mathcal{O}$ is a bounded Lipschitz domain in $\mathbb{R}^{d}$. For a given $h \in \mathbb{S}^{d-1}$, we study the monotonicity of the first Dirichlet eigenvalue of the $p$-Laplace operator with respect to the translations of the obstacle $\mathcal{O}$ in the $h$-direction within $\Omega$. Without loss of generality, we may assume that the origin $0 \in \mathcal{O}$. We consider the following family of polarizers:

$$
\begin{equation*}
H_{s}=\left\{x \in \mathbb{R}^{d}: x \cdot h<s\right\} \text {, for } s \in \mathbb{R} \text {. } \tag{1.9}
\end{equation*}
$$

We make the following geometric assumption on $\Omega$ and $\mathcal{O}$ :
$\left(\mathbf{A}_{2}\right) P_{H_{0}}(\Omega)=\Omega$, and $\mathcal{O}$ is Steiner symmetric with respect to the hyperplane $\partial H_{0}$ (see Definition 2.7).
The translations of $\mathcal{O}$ in the directions of $h$ are given by

$$
\begin{equation*}
\mathcal{O}_{s}=s h+\mathcal{O} \text { for } s \in \mathbb{R} . \tag{1.10}
\end{equation*}
$$

For $\Omega$ and $\mathcal{O}$ as given in $\left(\mathbf{A}_{\mathbf{2}}\right)$, define $\mathrm{L}_{\mathcal{O}}=\left\{s \in \mathbb{R}: P_{H_{s}}(\Omega)=\Omega\right.$ and $\left.\mathcal{O}_{s} \subset \Omega\right\}$. Let $\lambda_{1}(s)$ be the first eigenvalue of (1.3) with $\Omega_{\mathrm{in}}=\emptyset$ and $\Omega_{\text {out }}=\Omega \backslash \mathcal{O}_{s}$ for $s \in \mathrm{~L}_{\mathcal{O}}$. For $s \in \mathbb{R}$, let $\Sigma_{s}:=\{x \in \Omega: x \cdot h>s\}$. A set $A \subseteq \mathbb{R}$ is said to be convex in the $h$-direction, if any line segment parallel to the $\mathbb{R} h$-axis with endpoints in $A$ completely lies in $A$. Now, we have the following strict monotonicity result.
Theorem 1.5. Let $\frac{2 d+2}{d+2}<p<\infty$ and $h \in \mathbb{S}^{d-1}$. Assume that $\mathcal{O}, \Omega \subset \mathbb{R}^{d}$ satisfy $\left(\mathbf{A}_{\mathbf{1}}\right)$ and ( $\mathbf{A}_{\mathbf{2}}$ ). If the set $\Sigma_{s_{0}} \cup \sigma_{H_{s_{0}}}\left(\Sigma_{s_{0}}\right)$ is convex in the $h$-direction for some $s_{0} \in \mathrm{~L}_{\mathcal{O}}$, then the set $\left\{s \in \mathrm{~L}_{\mathcal{O}}: s \geq s_{0}\right\}$ is an interval and $\lambda_{1}(\cdot)$ is strictly decreasing on this interval.

Throughout this article, for given $a \in \mathbb{R}^{d}$ and $r \geq 0$, we denote $B_{r}(a)=\left\{x \in \mathbb{R}^{d}:|x-a|<r\right\}$, the open ball centered at $a$ with the radius $r$, and the closure of $B_{r}(a)$ by $\bar{B}_{r}(a)$.

Remark 1.6. In Theorem 1.5, if $\Omega$ itself is convex in the $h$-direction, then $\mathrm{L}_{\mathcal{O}}$ is an interval containing 0 . In particular, if $\Omega=B_{R}(0), \mathcal{O}=\bar{B}_{r}(0)$ for $0<R<r<\infty$ and $h=e_{1}=(1,0, \ldots, 0) \in \mathbb{S}^{d-1}$. Then, both $\Omega$ and $\mathcal{O}$ are Steiner symmetric with respect to $\partial H_{0}$, and $\mathrm{L}_{\mathcal{O}}=[0, R-r)$. Therefore, by Theorem 1.5, the first Dirichlet eigenvalue $\lambda_{1}\left(B_{R}(0) \backslash \bar{B}_{r}\left(s e_{1}\right)\right)$ is strictly decreasing for $s \in[0, R-r)$. Thus, Theorem 1.5 gives an alternate proof for many existing strict monotonicity results that were proved using the shape derivative (Hadamard perturbation) formula. For example, Kesavan [22] and Harrell-Kröger-Kurata [19], and Anoop-Bobkov-Sasi [4] for $p \in\left(\frac{2 d+2}{d+2}, \infty\right)$.
Remark 1.7. Due to the symmetry restrictions on the Neumann boundary in Theorem 1.3, the monotonicity results (similar to that of Dirichlet eigenvalue in Remark 1.6), when the Neumann boundary condition is specified on $\partial B_{R}(0)$ can not be deduced from Theorem 1.5. However, such a monotonicity result is proved for $p=2$, by Anoop-Ashok-Kesavan [5] using the Hadamard perturbation formula and some geometric properties of the first eigenfunctions. This result is open for general $p \neq 2$.

The monotonicity with respect to the rotations about a point: Next, we study the monotonicity of the first eigenvalue of $(\mathcal{E})$ on $\Omega \backslash \mathcal{O}$ with respect to the rotations of the obstacle $\mathcal{O}$ in $\Omega$ about a point $a \in \mathbb{R}^{d}$. We set $\mathbb{R}^{+}=[0, \infty)$, and make the following geometric assumptions on both $\Omega$ and $\mathcal{O}$ :
$\left(\mathbf{A}_{\mathbf{3}}\right)$ The domain $\Omega$ and the obstacle $\mathcal{O}$ are foliated Schwarz symmetric with respect to the ray $a+\mathbb{R}^{+} \eta$, for some $\eta \in \mathbb{S}^{d-1}$ (see Definition 2.7).
For $s \in[-1,1]$, let $\theta_{s}:=\arccos (s) \in[0, \pi]$. For $\xi \in \mathbb{S}^{d-1} \backslash\{\eta\}$, let $R_{s, \xi}$ be the simple rotation on $\mathbb{R}^{d}$ with the plane of rotation is $X_{\xi}:=\operatorname{span}\{\eta, \xi\}$ and the angle of rotation is $\theta_{s}$ from the ray $\mathbb{R}^{+} \eta$ in the counter-clockwise direction. The rotation of the obstacle $\mathcal{O}$ by $R_{s, \xi}$ about the point $a$ is given by

$$
\begin{equation*}
\mathcal{O}_{s, \xi}=a+R_{s, \xi}(-a+\mathcal{O}) \tag{1.11}
\end{equation*}
$$

Now, we observe the following facts (see Proposition 2.9 and Lemma 4.3):
(a) The rotated obstacle $\mathcal{O}_{s, \xi}$ is foliated Schwarz symmetric with respect to the ray $a+\mathbb{R}^{+} R_{s, \xi}(\eta)$.
(b) For any rotation $R$ that fixes $\eta, \mathcal{O}=a+R(-a+\mathcal{O})$ and $\Omega=a+R(-a+\Omega)$.
(c) For any distinct $\xi_{1}, \xi_{2} \in \mathbb{S}^{d-1} \backslash\{\eta\}$, there exists $R$ that fixes $\eta$ such that

$$
R\left(-a+\Omega \backslash \mathcal{O}_{s, \xi_{1}}\right)=-a+\Omega \backslash \mathcal{O}_{s, \xi_{2}}
$$

From the above observations, it is evident that we only need to consider the rotations of the obstacle by $R_{s, \xi}$ with respect to $a$ in a $X_{\xi}$-plane for a fixed $\xi \in \mathbb{S}^{d-1} \backslash\{\eta\}$. Thus for $s \in[-1,1]$, we set $\mathcal{O}_{s}=\mathcal{O}_{s, \xi}$ and consider

$$
\begin{align*}
\mathrm{C}_{\mathcal{O}} & :=\left\{s \in[-1,1]: \mathcal{O}_{s} \subset \Omega\right\}  \tag{1.12}\\
\gamma_{1}(s) & :=\gamma_{1}\left(\Omega \backslash \mathcal{O}_{s}\right), \text { the first eigenvalue of }(\mathcal{E}) \text { on } \Omega \backslash \mathcal{O}_{s} \text { for } s \in \mathrm{C}_{\mathcal{O}} .
\end{align*}
$$

In this article, we consider the following types of $\Omega$ and $\Gamma_{N} \subseteq \partial \Omega$ :


Figure 2. The translations of $\mathcal{O}$ along the $e_{1}$-axis; and rotations of $\mathcal{O}$ about the point $a \in \mathbb{R}^{d}$, here $\theta_{i}=\arccos \left(s_{i}\right)$ for $i=1,2$ with $s_{1}>s_{2}$.
$\left(\mathbf{A}_{4}\right) \Omega=\Omega_{0} \backslash \bar{B}_{\rho_{0}}(a)$, where $\bar{B}_{\rho_{0}}(a) \subsetneq \Omega_{0} \subset \mathbb{R}^{d}, \rho_{0} \geq 0$; and $\Gamma_{N}=\partial B_{\rho_{0}}(a)$.
$\left(\mathbf{A}_{\mathbf{5}}\right) \Omega=B_{R}(a) \backslash \overline{\Omega_{1}}$, where $\overline{\Omega_{1}} \subset B_{R}(a)$, and $\Gamma_{N}=\partial B_{R}(a)$.
Now, we state our monotonicity result for $\gamma_{1}($.$) on \mathbf{C}_{\mathcal{O}}$.
Theorem 1.8. Let $\frac{2 d+2}{d+2}<p<\infty$ and $\Omega \subset \mathbb{R}^{d}$ be a domain. Assume that the pair $\Omega$ and $\Gamma_{N}$ satisfy either $\left(\mathbf{A}_{\mathbf{4}}\right)$ or $\left(\mathbf{A}_{\mathbf{5}}\right)$. If $\Omega$ and $\mathcal{O}$ satisfy $\left(\mathbf{A}_{\mathbf{1}}\right)$ and $\left(\mathbf{A}_{\mathbf{3}}\right)$ for some $a \in \mathbb{R}^{d}$ and $\eta \in \mathbb{S}^{d-1}$, then $\mathrm{C}_{\mathcal{O}}$ is an interval. In addition, if $\Omega$ is not radial with respect to $a$, then $\gamma_{1}(\cdot)$ is strictly increasing on $\mathbf{C}_{\mathcal{O}}$.

Remark 1.9. If $\Omega$ is radial with respect to the point $a$ (see Corollary 2.11), then the first eigenvalue $\gamma_{1}(\cdot)$ remains as a constant on $\mathrm{C}_{\mathcal{O}}$.

The rest of this article is organized as follows. In Section 2, the polarization of measurable sets and functions are introduced, and some of their important properties are discussed. Further, the characterizations of Steiner and foliated Schwarz symmetries using polarizations are given in Section 2. Also, we include a strong comparison principle and a few interior and boundary regularity results that are essential for the development of this article. A proof of Faber-Krahn inequality (Theorem 1.3) is given in Section 3. The proofs of strict monotonicity results (Theorem 1.5 and 1.8) are given in Section 4. Many important remarks and explicit examples are included in Section 5.

## 2. Preliminaries

In this section, we discuss some of the important properties of the polarization of the sets and functions. Further, we give the definitions of Steiner and foliated Schwarz symmetries, and their characterizations in terms of polarizations. Lastly, we give some regularity results and strong comparison principles for the solutions of the $p$-Laplace operator.
2.1. Polarization of sets. We discuss a few simple properties of the polarization of sets.

Proposition 2.1. Let $H \in \mathcal{H}$ and $A, C \subseteq \mathbb{R}^{d}$. Then,
(i) $P_{H}(A)$ is open, if $A$ is open; and $P_{H}(A)$ is closed if $A$ is closed;
(ii) $P_{H}(A) \subseteq P_{H}(C)$, if $A \subseteq C$;
(iii) $P_{H}(A \cap C) \subseteq P_{H}(A) \cap P_{H}(C)$ and $P_{H}(A) \cup P_{H}(C) \subseteq P_{H}(A \cup C)$;
(iv) $P_{H}\left(\sigma_{H}(A)\right)=P_{H}(A), \sigma_{H}\left(P_{H}(A)\right)=P^{H}(A)$, and $\sigma_{H}\left(P^{H}(A)\right)=P_{H}(A)$;
(v) $P_{H}\left(A^{\mathrm{c}}\right)=\left(P^{H}(A)\right)^{\mathrm{c}}$.

Proof. Recall that, for $H \in \mathcal{H}$, the polarizations of a set $A \subseteq \mathbb{R}^{d}$ are given by

$$
P_{H}(A)=\left[\left(A \cup \sigma_{H}(A)\right) \cap H\right] \cup\left[A \cap \sigma_{H}(A)\right], \text { and } P^{H}(A)=\left[\left(A \cup \sigma_{H}(A)\right) \cap H^{\mathrm{c}}\right] \cup\left[A \cap \sigma_{H}(A)\right] .
$$

Since $A \cap \partial H=\sigma_{H}(A) \cap \partial H$, we can also write

$$
P_{H}(A)=\left[\left(A \cup \sigma_{H}(A)\right) \cap \bar{H}\right] \cup\left[A \cap \sigma_{H}(A)\right] \text { and } P^{H}(A)=\left[\left(A \cup \sigma_{H}(A)\right) \cap \bar{H}^{c}\right] \cup\left[A \cap \sigma_{H}(A)\right] .
$$

Now, (i)-(iii) follow easily from the above observations.
(iv) This follows from the fact $\sigma_{H}(H)=\bar{H}^{\mathrm{c}}$ and the above observations.
(v) By the definition, we have $P_{H}\left(A^{\mathrm{c}}\right)=\left[\left(A^{\mathrm{c}} \cup \sigma_{H}\left(A^{\mathrm{c}}\right)\right) \cap H\right] \bigcup\left[A^{\mathrm{c}} \cap \sigma_{H}\left(A^{\mathrm{c}}\right)\right]$, and hence

$$
\begin{aligned}
\left(P_{H}\left(A^{\mathrm{c}}\right)\right)^{\mathrm{c}} & =\left[\left(A \cap \sigma_{H}(A)\right) \cup H^{\mathrm{c}}\right] \cap\left[A \cup \sigma_{H}(A)\right] \\
& =\left[A \cap \sigma_{H}(A)\right] \cup\left[\left(A \cup \sigma_{H}(A)\right) \cap H^{\mathrm{c}}\right]=P^{H}(A) .
\end{aligned}
$$

The following proposition characterizes the invariance of a set under polarizations.
Proposition 2.2. Let $H \in \mathcal{H}$ and $A \subseteq \mathbb{R}^{d}$. Then
(i) $P_{H}(A)=A$ if and only if $\sigma_{H}(A) \cap H \subseteq A$;
(ii) $P^{H}(A)=A$ if and only if $\sigma_{H}(A) \cap H^{c} \subseteq A$;
(iii) $P_{H}(A)=P^{H}(A)$ if and only if $\sigma_{H}(A)=A$.

Proof. (i) From the definition of $P_{H}(A)$, it is clear that

$$
\begin{equation*}
P_{H}(A)=\left[\left(A \cup \sigma_{H}(A)\right) \cap H\right] \cup\left[A \cap \sigma_{H}(A) \cap H^{\mathrm{c}}\right] \tag{2.1}
\end{equation*}
$$

If $P_{H}(A)=A$, then $P_{H}(A) \cap H=A \cap H$. Thus the above equation yields $\left(A \cup \sigma_{H}(A)\right) \cap H \subseteq A$, and hence we must have $\sigma_{H}(A) \cap H \subseteq A$. Conversely, assume that $\sigma_{H}(A) \cap H \subseteq A$. Then, by applying $\sigma_{H}$ on both sides, and using the fact that $A \cap \partial H=\sigma_{H}(A) \cap \partial H$, we obtain $A \cap H^{\text {c }} \subseteq \sigma_{H}(A)$, Therefore, $A \cap H^{c}=A \cap \sigma_{H}(A) \cap H^{\text {c }}$. From the assumption, we also have $A \cap H=\left[\left(A \cup \sigma_{H}(A)\right) \cap H\right]$. Now, using (2.1), we easily conclude that $P_{H}(A)=A$.
(ii) From Proposition 2.1-(iv), we have $\sigma_{H}\left(P^{H}(A)\right)=P_{H}(A)$ and $P_{H}\left(\sigma_{H}(A)\right)=P_{H}(A)$. Therefore, we get $P^{H}(A)=A$ if and only if $P_{H}\left(\sigma_{H}(A)\right)=\sigma_{H}(A)$. Now, from (i) we obtain

$$
P_{H}\left(\sigma_{H}(A)\right)=\sigma_{H}(A) \text { if and only if } A \cap H \subseteq \sigma_{H}(A)
$$

Now applying $\sigma_{H}$ on both sides of last inclusion and using the fact that $\sigma_{H}(A) \cap \partial H=A \cap \partial H$, we get

$$
P^{H}(A)=A \text { if and only if } \sigma_{H}(A) \cap H^{c} \subseteq A
$$

(iii) From the definitions of $P_{H}(A)$ and $P^{H}(A)$, it is clear that

$$
\begin{aligned}
& P_{H}(A)=\left[\left(A \cup \sigma_{H}(A)\right) \cap H\right] \cup\left[A \cap \sigma_{H}(A) \cap H^{\mathrm{c}}\right] \\
& P^{H}(A)=\left[\left(A \cup \sigma_{H}(A)\right) \cap H^{\mathrm{c}}\right] \cup\left[A \cap \sigma_{H}(A) \cap H\right]
\end{aligned}
$$

If $P_{H}(A)=P^{H}(A)$, then $\left(A \cup \sigma_{H}(A)\right) \cap H=A \cap \sigma_{H}(A) \cap H$ and $A \cap \sigma_{H}(A) \cap H^{c}=\left(A \cup \sigma_{H}(A)\right) \cap H^{c}$. Therefore $A \cup \sigma_{H}(A)=A \cap \sigma_{H}(A)$, and hence $\sigma_{H}(A)=A$. Conversely, assume that $\sigma_{H}(A)=A$. Then, from above equations, we get $P_{H}(A)=P^{H}(A)=A$.


Figure 3. The sets $A_{H}$ and $B_{H}$ of $P_{H}(\Omega) \cap H$.

Proposition 2.3. Let $H \in \mathcal{H}$ and $\Omega \subseteq \mathbb{R}^{d}$ be an open set. Then,
(i) $P_{H}(\Omega) \neq \Omega$ if and only if $A_{H}:=\sigma_{H}(\Omega) \cap \Omega^{\mathrm{c}} \cap H$ has non-empty interior;
(ii) $P_{H}(\Omega) \neq \Omega$ if and only if $B_{H}:=\Omega \cap \sigma_{H}\left(\Omega^{\mathrm{c}}\right) \cap H$ has non-empty interior.

Proof. (i) First, we observe that the interior of $A_{H}$ is $\sigma_{H}(\Omega) \cap \bar{\Omega}^{\mathrm{c}} \cap H$. Since $\sigma_{H}(\Omega) \cap H$ is open, from Proposition 2.2, we get $P_{H}(\Omega) \neq \Omega$ if and only if $\sigma_{H}(\Omega) \cap H \nsubseteq \bar{\Omega}$. Clearly, $\sigma_{H}(\Omega) \cap H \nsubseteq \bar{\Omega}$ if and only if $\bar{\Omega}^{\mathrm{c}} \cap \sigma_{H}(\Omega) \cap H \neq \emptyset$.
(ii) For $H \in \mathcal{H}$, we have $\sigma_{H}(H) \in \mathcal{H}$. Then from Proposition 2.1, $P_{H}(\Omega)=\sigma_{H}(\Omega)$ if and only if $P_{\sigma_{H}(H)}(\Omega)=$ $\Omega$. The proof follows from (i) by replacing $H$ with $\sigma_{H}(H)$.

Remark 2.4. For an open set $\Omega \subset \mathbb{R}^{d}$, if $P_{H}(\Omega) \neq \Omega$ then the interior of $P_{H}(\Omega) \backslash \Omega$ is non-empty. Therefore, if $P_{H}(\Omega) \neq \Omega$ then $P_{H}(\Omega)$ can not be equal to $\Omega$ up to a set of measure zero (or up to a set of $p$-capacity zero).

Now, we prove that the set $P_{H}(\Omega) \cap H$ is a domain when $\Omega$ is a domain. For this, we need the following lemma.

Lemma 2.5. Let $H \in \mathcal{H}$ and $\Omega \subseteq \mathbb{R}^{d}$ be a domain. If $\sigma_{H}(\Omega)=\Omega$, then both $\Omega \cap H$ and $\Omega \cap \bar{H}^{\text {c }}$ are connected.

Proof. Let $f: \Omega \cap \bar{H} \longrightarrow\{0,1\}$ be a continuous function. Using the symmetry of $\Omega$, we define

$$
\tilde{f}(x)= \begin{cases}f(x), & \text { for } x \in \Omega \cap \bar{H} \\ f \circ \sigma_{H}(x), & \text { for } x \in \Omega \cap H^{c}\end{cases}
$$

Then $\tilde{f}$ is a continuous function on $\Omega$, since $\sigma_{H}(x)=x$ for $x \in \partial H$. By the connectedness of $\Omega, \tilde{f}$ is constant on $\Omega$. In particular, $f$ is constant on $\Omega \cap \bar{H}$, and hence $\Omega \cap \bar{H}$ is connected. Therefore int $(\Omega \cap \bar{H})=\Omega \cap H$ is connected, and hence $\sigma_{H}(\Omega \cap H)=\Omega \cap \bar{H}^{\text {c }}$ is also connected.

Proposition 2.6. Let $H \in \mathcal{H}$ and $\Omega \subseteq \mathbb{R}^{d}$ be a domain. Then $P_{H}(\Omega) \cap H$ is a domain.
Proof. First, we observe that $P_{H}(\Omega) \cap H=\left(\Omega \cup \sigma_{H}(\Omega)\right) \cap H$ is open. For proving the connectedness, we consider the following two cases: (a) $\Omega \cap \sigma_{H}(\Omega)=\emptyset$, and (b) $\Omega \cap \sigma_{H}(\Omega) \neq \emptyset$.
(a) $\Omega \cap \sigma_{H}(\Omega)=\emptyset$ : In this case, we have $\Omega \cap \partial H=\emptyset$, since $\Omega \cap \partial H \subset \Omega \cap \sigma_{H}(\Omega)$. Therefore, $\Omega$ is the union of two open sets $\Omega \cap H$ and $\Omega \cap \bar{H}^{\mathrm{c}}$. By the connectedness of $\Omega$, one of them is equal to $\Omega$. If $\Omega \cap H=\Omega$ then $P_{H}(\Omega)=\Omega$, and hence $P_{H}(\Omega) \cap H=\Omega$. If $\Omega \cap \bar{H}^{\text {c }}=\Omega$ then $P_{H}(\Omega)=\sigma_{H}(\Omega)$, and hence $P_{H}(\Omega) \cap H=\sigma_{H}(\Omega)$.
(b) $\Omega \cap \sigma_{H}(\Omega) \neq \emptyset$ : In this case, we have $\Omega \cup \sigma_{H}(\Omega)$ is connected, and it is symmetric with respect to $\partial H$. Thus, by Lemma 2.5, $\left(\Omega \cup \sigma_{H}(\Omega)\right) \cap H=P_{H}(\Omega) \cap H$ is connected.
Therefore, in both of the cases, $P_{H}(\Omega) \cap H$ is domain.
The Steiner, axial and foliated Schwarz symmetries: A set in $\mathbb{R}^{d}$ is said to have certain symmetry, if it is invariant under corresponding symmetrization or rearrangement on $\mathbb{R}^{d}$. Here, we directly give the definitions of the Steiner and the foliated Schwarz symmetries without defining the associated symmetrizations (see [30, Definition 3.1 and Definition 3.2]). The foliated Schwarz symmetrization with respect to a ray $a+\mathbb{R}^{+} \eta$ is the cap symmetrization with respect to $a+\mathbb{R}^{+} \eta$, see [30, Definition 3.2].
Definition 2.7. Let $A \subseteq \mathbb{R}^{d}$ be a measurable set.
(1) Steiner symmetry. Let $S$ be an affine-hyperplane in $\mathbb{R}^{d}$. For each $x \in S$, let $L_{x}$ be the line passing through $x$ and orthogonal to $S$. Then $A$ is said to be Steiner symmetric with respect to $S$, if
for each $x \in S, \quad A \cap L_{x}=B_{\rho}(x) \cap L_{x}$ for some $\rho \geq 0$.
(2) Axial symmetry. Let $L$ be a line in $\mathbb{R}^{d}$. For each $x \in L$, let $S_{x}$ be the affine hyperplane passing through $x$ and orthogonal to $L$. Then $A$ is said to be axially symmetric with respect to $L$, if

$$
\text { for each } x \in L, \quad A \cap \partial B_{\rho}(x) \cap S_{x}=\partial B_{\rho}(x) \cap S_{x} \text { for some } \rho \geq 0 \text {. }
$$

(3) Foliated Schwarz symmetry. Let $a+\mathbb{R}^{+} \eta$ be a ray starting for some $a \in \mathbb{R}^{d}$ and $\eta \in \mathbb{S}^{d-1}$. Then $A$ said to be foliated Schwarz symmetric with respect to $a+\mathbb{R}^{+} \eta$, if

$$
\text { for every } r>0, \quad A \cap \partial B_{r}(a)=B_{\rho}(a+r \eta) \cap \partial B_{r}(a) \text { for some } \rho \geq 0
$$

Remark 2.8. We observe that:
(i) a set $A \subseteq \mathbb{R}^{d}$ is Steiner symmetric with respect to an affine-hyperplane $S$, if and only if $A$ is invariant under the reflection with respect to $S$ and convex in the orthogonal direction to $S$;
(ii) a set $A \subseteq \mathbb{R}^{d}$ is axially symmetric with respect to a line $L$, if and only if $A$ is invariant under the reflection with respect to every affine hyperplane containing $L$. In particular, if $L=\mathbb{R} \eta$ and $R$ is any rotation on $\mathbb{R}^{d}$ such that $R(\eta)=\eta$, then $R(A)=A$. This follows from the definition, since the planes of rotation of such $R$ can not contain $\eta$, and hence those planes must be orthogonal to $\eta$;
(iii) let $A \subseteq \mathbb{R}^{d}$ be foliated Schwarz symmetric with respect to $a+\mathbb{R}^{+} \eta$. Let $I_{A}:=\left\{r>0: A \cap \partial B_{r}(a) \neq\right.$ $\emptyset\}$. For $r \in I_{A}$, let $\rho(r)>0$ be such that $A \cap \partial B_{r}(a)=B_{\rho(r)}(a+r \eta) \cap \partial B_{r}(a)$. Then,

$$
\begin{equation*}
A=\bigcup_{r \in I_{A}} B_{\rho(r)}(a+r \eta) \cap \partial B_{r}(a) \tag{2.2}
\end{equation*}
$$

The following proposition provides some properties of the foliated Schwartz symmetric sets.
Proposition 2.9. If $A \subseteq \mathbb{R}^{d}$ is foliated Schwarz symmetric with respect to a ray $a+\mathbb{R}^{+} \eta$ then
(i) $A$ is axially symmetric with respect to $a+\mathbb{R} \eta$,
(ii) for any linear map $T$ and $b \in \mathbb{R}^{d}$, the set $b+T(A)$ is foliated Schwarz symmetric with respect to $b+T(a)+\mathbb{R}^{+} T(\eta)$,
(iii) $R(-a+A)=-a+A$, for any rotation $R$ on $\mathbb{R}^{d}$ that fixes $\eta$.

Proof. (i) Observe that, for every $r \in I_{A}$, the set $B_{\rho(r)}(a+r \eta) \cap \partial B_{r}(a)$ is axially symmetric with respect to $a+\mathbb{R} \eta$. Now, using (2.2), we conclude that $A$ is axially symmetric with respect to $a+\mathbb{R} \eta$.
(ii) Let $r>0$, then

$$
(b+T(A)) \cap \partial B_{r}(b+T(a))=b+T\left(A \cap \partial B_{r}(a)\right)=b+T\left(B_{\rho}(a+r \eta) \cap \partial B_{r}(a)\right), \text { for some } \rho \geq 0,
$$

where the last equality follows from the definition foliated Schwarz symmetry. Thus

$$
\begin{aligned}
(b+T(A)) \cap \partial B_{r}(b+T(a)) & =b+B_{\rho}(T(a)+r T(\eta)) \cap \partial B_{r}(T(a)) \\
& =B_{\rho}(b+T(a)+r T(\eta)) \cap \partial B_{r}(b+T(a))
\end{aligned}
$$

Now, we obtain the required conclusion by the definition of foliated Schwarz symmetry.
(iii) By taking $T=I$ and $b=-a$ in (ii), we get $-a+A$ is foliated Schwarz symmetric with respect to $\mathbb{R}^{+} \eta$. Thus by (i), $-a+A$ is axially symmetric with respect to $\mathbb{R} \eta$. Since $R$ fixes $\eta$, from (ii) of Remark 2.8 , we conclude $R(-a+A)=-a+A$.

Next, we characterize the foliated Schwarz and Steiner symmetric sets using the polarizations. First, we consider the following polarizers: for given $a \in \mathbb{R}^{d}, \eta \in \mathbb{S}^{d-1}$, let

$$
\mathcal{H}_{a, \eta}:=\left\{H \in \mathcal{H}: a+\mathbb{R}^{+} \eta \subset H \text { and } a \in \partial H\right\} .
$$

Some useful characterizations of the Steiner symmetry (from [11, Lemma 2.2]), foliated Schwarz symmetry (from [30, Section 3]) are given in the following proposition.

Proposition 2.10. Let $A \subseteq \mathbb{R}^{d}$ be any set.
(i) Let $H_{s} \in \mathcal{H}$ be as given in (1.9). Then, for $s_{0} \in \mathbb{R}$, the following statements are equivalent:
(a) the set $A$ is Steiner symmetric with respect to the affine-hyperplane $\partial H_{s_{0}}$,
(b) $P_{H_{s}}(A)=A$, for every $s \geq s_{0}$; and $P^{H_{s}}(A)=A$, for every $s \leq s_{0}$.
(ii) Let $a \in \mathbb{R}^{d}$ and $\eta \in \mathbb{S}^{d-1}$. Then the following are equivalent:
(a) the set $A$ is foliated Schwarz symmetric with respect to the ray $a+\mathbb{R}^{+} \eta$,
(b) $P^{H}(A)=\sigma_{H}(A)$, for every $H \in \mathcal{H}_{a, \eta}$.

We have the following corollary.
Corollary 2.11. Let $A \subseteq \mathbb{R}^{d}$ be any set. If $A$ is foliated Schwarz symmetric with respect to both the rays $a+\mathbb{R}^{+} \eta$ and $a-\mathbb{R}^{+} \eta$ for some $a \in \mathbb{R}^{d}$ and $\eta \in \mathbb{S}^{d-1}$. Then $A$ is radial with respect to the point $a$.

Proof. Notice that, $A$ is radial with respect to $a \in \mathbb{R}^{d}$ provided $A \cap \partial B_{r}(a)=\partial B_{r}(a)$ for every $r \in I_{A}$, where $I_{A}=\left\{r \in \mathbb{R}: A \cap \partial B_{r}(a) \neq \emptyset\right\}$. Since $A$ is foliated Schwarz symmetric with respect to both the rays $a+\mathbb{R}^{+} \eta$ and $a-\mathbb{R}^{+} \eta$, for each $r \in I_{A}$ we get:

$$
\begin{equation*}
A \cap \partial B_{r}(a)=B_{\rho_{1}}(a+r \eta) \cap \partial B_{r}(a)=B_{\rho_{2}}(a-r \eta) \cap \partial B_{r}(a) \text { for some } \rho_{1}, \rho_{2} \geq 0 \tag{2.3}
\end{equation*}
$$

Since $|a-(a-r \eta)|=r$, from (2.3) we obtain $a-r \eta \in B_{\rho_{1}}(a+r \eta)$. Thus $\rho_{1} \geq|a-r \eta-(a+r \eta)|=2 r$, and hence $B_{\rho_{1}}(a+r \eta) \cap \partial B_{r}(a)=\partial B_{r}(a)$. Now, from (2.3) we conclude that $A \cap \partial B_{r}(a)=\partial B_{r}(a)$.
2.2. Polarization of punctured domains. We consider the polarization of the punctured domains of the form $A \backslash C$, where $A \subseteq \mathbb{R}^{d}$ is open, and $C \subset A$ is closed. Clearly $\partial(A \backslash C)=\partial A \sqcup \partial C$.

Proposition 2.12. Let $A \subseteq \mathbb{R}^{d}$ be open and $C \subset A$ be closed. If $H \in \mathcal{H}$ is such that $\sigma_{H}(C) \subset A$, then
(i) $P_{H}(A \backslash C)=P_{H}(A) \backslash P^{H}(C)$,
(ii) $P^{H}(C) \subset P_{H}(A)$, in particular $\partial P_{H}(A \backslash C)=\partial P_{H}(A) \sqcup \partial P^{H}(C)$.

Proof. (i) For $A \subseteq \mathbb{R}^{d}$, denote $P_{H}^{+}(A)=P_{H}(A) \cap H$ and $P_{H}^{-}(A)=P_{H}(A) \cap H^{\text {c }}$. Thus $P_{H}(A)=P_{H}^{+}(A) \sqcup$ $P_{H}^{-}(A)$, and

$$
\begin{equation*}
P_{H}(A) \cap P_{H}\left(C^{\mathrm{c}}\right)=\left[P_{H}^{+}(A) \cap P_{H}^{+}\left(C^{\mathrm{c}}\right)\right] \sqcup\left[P_{H}^{-}(A) \cap P_{H}^{-}\left(C^{\mathrm{c}}\right)\right] . \tag{2.4}
\end{equation*}
$$

On the other hand, we have $P_{H}^{-}\left(A \cap C^{\mathrm{c}}\right)=P_{H}^{-}(A) \cap P_{H}^{-}\left(C^{\mathrm{c}}\right)$. Since $C, \sigma_{H}(C) \subseteq A$, we get $\sigma_{H}(A) \cup C^{\mathrm{c}}=$ $A \cup \sigma_{H}\left(C^{\mathrm{c}}\right)=\mathbb{R}^{d}$. Thus, $\left(A \cap C^{\mathrm{c}}\right) \cup \sigma_{H}\left(A \cap C^{\mathrm{c}}\right)=\left(A \cup \sigma_{H}(A)\right) \cap\left(C^{\mathrm{c}} \cup \sigma_{H}\left(C^{\mathrm{c}}\right)\right)$, and hence $P_{H}^{+}\left(A \cap C^{\mathrm{c}}\right)=$ $P_{H}^{+}(A) \cap P_{H}^{+}\left(C^{\mathrm{c}}\right)$. Therefore, from (2.4) and using $P_{H}\left(C^{\mathrm{c}}\right)=\left(P^{H}(C)\right)^{\mathrm{c}}$ (Proposition 2.1-(v)), we obtain

$$
P_{H}(A \backslash C)=P_{H}(A) \cap P_{H}\left(C^{\mathrm{c}}\right)=P_{H}(A) \backslash P^{H}(C)
$$

(ii) Since $C \cup \sigma_{H}(C)$ is a symmetric set in $A$, by the definitions of $P^{H}$ and $P_{H}$, we get

$$
P^{H}(C) \subseteq C \cup \sigma_{H}(C)=P_{H}\left(C \cup \sigma_{H}(C)\right) \subset P_{H}(A)
$$

Moreover, $P^{H}(C)$ is closed and $P_{H}(A)$ is open in $\mathbb{R}^{d}$. Thus,

$$
\partial P_{H}(A \backslash C)=\partial\left(P_{H}(A) \backslash P^{H}(C)\right)=\partial P_{H}(A) \sqcup \partial P^{H}(C)
$$

Remark 2.13. The assumption $\sigma_{H}(C) \subset A$ is essential for the conclusions of the above proposition. To see this, consider $A=B_{R}(0), C=\bar{B}_{r}(0)$, and the polarizers $H_{s}:=\left\{x \in \mathbb{R}^{d}: x_{1}<s\right\}$ for $s \in \mathbb{R}$. For $s>\frac{R-r}{2}$, we have $\left|\sigma_{H_{t}}(C) \cap A^{\mathrm{c}}\right|=\left|\bar{B}_{r}\left(2 t e_{1}\right) \cap B_{R}(0)^{\mathrm{c}}\right|>0$, where $|A|$ is the Lebesgue measure of $A \subseteq \mathbb{R}^{d}$. Then, $\left|P_{H_{t}}(A) \backslash P^{H_{t}}(C)\right|=\left|B_{R}(0) \backslash \bar{B}_{r}\left(2 t e_{1}\right)\right|>\left|B_{R}(0) \backslash \bar{B}_{r}(0)\right|$. Since $P_{H}$ is measure preserving, we conclude that $P_{H}(A) \backslash P^{H}(C) \neq P_{H}(A \backslash C)$.
2.3. Polarization of functions. Now, we consider the polarization of functions defined on a domain $\Omega \subseteq$ $\mathbb{R}^{d}$ and discuss some important properties of polarization of functions, such as Lipschitz continuity, nonexpansivity, norm preserving property. Recall the definition of polarization of functions (from Definition 1.1).
Proposition 2.14. Let $H \in \mathcal{H}$, and $u \in \mathcal{C}\left(\mathbb{R}^{d}\right)$ be a non-negative function. Then

$$
\operatorname{supp}\left(P_{H}(u)\right)=P_{H}(\operatorname{supp}(u))
$$

Proof. Let $F=\operatorname{supp}(u)$. Clearly $u=u \circ \sigma_{H}=0$ on $F^{\mathrm{c}} \cap \sigma_{H}\left(F^{\mathrm{c}}\right)$, and $u=0$ or $u \circ \sigma_{H}=0$ on $F^{\mathrm{c}} \cup \sigma_{H}\left(F^{\mathrm{c}}\right)$. Since $u \geq 0$, by the definition, we get $P_{H}(u)=0$ on $\left[\left(F^{c} \cup \sigma_{H}\left(F^{\mathrm{c}}\right)\right) \cap H^{\mathrm{c}}\right] \cup\left[F^{\mathrm{c}} \cap \sigma_{H}\left(F^{\mathrm{c}}\right)\right]=P^{H}\left(F^{\mathrm{c}}\right)$. Now, since $P^{H}\left(F^{c}\right)=\left(P_{H}(F)\right)^{c}($ from Proposition 2.1-(v) $)$, we get $\operatorname{supp}\left(P_{H}(u)\right) \subseteq P_{H}(F)$. The other way inclusion is easy to see from the definition. Therefore, $\operatorname{supp}\left(P_{H}(u)\right)=P_{H}(\operatorname{supp}(u))$.

Remark 2.15. Similarly, for non-positive function $u \in \mathcal{C}\left(\mathbb{R}^{d}\right), \operatorname{supp}\left(P_{H}(u)\right) \subseteq P^{H}(\operatorname{supp}(u))$. More generally, for any function $u \in \mathcal{C}\left(\mathbb{R}^{d}\right)$ we have $\operatorname{supp}\left(P_{H}(u)\right)=P_{H}\left(\operatorname{supp}\left(u^{+}\right)\right) \cup P^{H}\left(\operatorname{supp}\left(u^{-}\right)\right)$(see [11, Section-2]), where $u^{+}=\max \{0, u\}$ and $u^{-}=\min \{0, u\}$.

The Hölder continuity of polarizations of Hölder continuous functions defined on $\mathbb{R}^{d}$ is given in [13, Corollary 3.1]. The same result holds for the functions defined on a symmetric domain.

Proposition 2.16. Let $\Omega_{0} \subseteq \mathbb{R}^{d}$ be a domain and $H \in \mathcal{H}$ such that $\sigma_{H}\left(\Omega_{0}\right)=\Omega_{0}$. If $u \in \mathcal{C}^{0, \alpha}\left(\Omega_{0}\right)$ for some $\alpha \in(0,1]$, then $P_{H} u \in \mathcal{C}^{0, \alpha}\left(\Omega_{0}\right)$.
Proof. For $u \in \mathcal{C}^{0, \alpha}\left(\Omega_{0}\right)$, there exists $L>0$ such that $|u(x)-u(y)| \leq L|x-y|^{\alpha}$ for any $x, y \in \Omega_{0}$. For simplicity of notation, we denote the reflection $\sigma_{H}(z)$ of $z \in \mathbb{R}^{d}$ with respect to $\partial H$ by $z^{*}$. Let $x, y \in \Omega_{0}$. Since $\sigma_{H}\left(\Omega_{0}\right)=\Omega_{0}$, both $x^{*}, y^{*} \in \Omega_{0}$, and $\operatorname{supp}\left(P_{H}(u)\right) \subseteq \Omega_{0}($ from Remark 2.15). If both $x, y \in H$, then

$$
\begin{aligned}
\left|P_{H} u(x)-P_{H} u(y)\right| & \leq\left|\max \left\{u(x), u\left(x^{*}\right)\right\}-\max \left\{u(y), u\left(y^{*}\right)\right\}\right| \\
& \leq \max \left\{|u(x)-u(y)|,\left|u\left(x^{*}\right)-u\left(y^{*}\right)\right|\right\} \leq L|x-y|^{\alpha} .
\end{aligned}
$$

Similarly, if $x, y \in H^{\text {c }}$ then $\left|P_{H} u(x)-P_{H} u(y)\right| \leq L|x-y|^{\alpha}$. Now, if $x \in H$ and $y \in H^{c}$ then $\left|x-y^{*}\right|=$ $\left|x^{*}-y\right| \leq|x-y|$. Therefore

$$
\begin{aligned}
\left|P_{H} u(x)-P_{H} u(y)\right| & \leq\left|\max \left\{u(x), u\left(x^{*}\right)\right\}-\min \left\{u(y), u\left(y^{*}\right)\right\}\right| \\
& \leq \max \left\{|u(x)-u(y)|,\left|u(x)-u\left(y^{*}\right)\right|,\left|u(y)-u\left(x^{*}\right)\right|,\left|u\left(x^{*}\right)-u\left(y^{*}\right)\right|\right\} \\
& \leq L|x-y|^{\alpha} .
\end{aligned}
$$

Now, we state the following non-expansive property of polarization, see [13, Theorem 3.1] and [15, Theorem 3, Corollary 1].

Proposition 2.17. Let $\Omega_{0} \subset \mathbb{R}^{d}$, and $j$ be any Young function. Then, for any $H \in \mathcal{H}$ and any non-negative measurable functions $u, v$ on $\Omega_{0}$,

$$
\int_{P_{H} \Omega_{0}} j\left(\left|P_{H} u-P_{H} v\right|\right) \mathrm{d} x \leq \int_{\Omega_{0}} j(|u-v|) \mathrm{d} x .
$$

In particular, for $j(t)=t^{p}, 1 \leq p<\infty$,

$$
\left\|P_{H} u-P_{H} v\right\|_{p, P_{H} \Omega_{0}} \leq\|u-v\|_{p, \Omega_{0}} \text { for any non-negative } u, v \in L^{p}\left(\Omega_{0}\right)
$$

We state the following invariance property of polarizations, see [29, Proposition 2.3.] and [31, Lemma 3.1].
Proposition 2.18. Let $\Omega_{0} \subseteq \mathbb{R}^{d}$ be an open set and $H \in \mathcal{H}$ such that $\sigma_{H}\left(\Omega_{0}\right)=\Omega_{0}$. If $u \in W^{1, p}\left(\Omega_{0}\right)$ then $P_{H}(u) \in W^{1, p}\left(\Omega_{0}\right)$, and

$$
\begin{equation*}
\|u\|_{p}=\left\|P_{H} u\right\|_{p} \text { and }\|\nabla u\|_{p}=\left\|\nabla P_{H} u\right\|_{p} \tag{2.5}
\end{equation*}
$$

Proof. Let $u \in W^{1, p}\left(\Omega_{0}\right)$. Since $\Omega_{0}$ is symmetric with respect to $\partial H$, we have $v:=u \circ \sigma_{H} \in W^{1, p}\left(\Omega_{0}\right)$. Moreover, using the standard arguments we can easily show that, $|u-v|, f:=|u-v| \mathbb{1}_{\Omega_{0} \cap H}$, and $g:=$ $-\mid u-v \mathbb{1}_{\Omega_{0} \cap H^{c}}$ are in $W^{1, p}\left(\Omega_{0}\right)$. Thus $P_{H}(u)=\frac{1}{2}(u+v+f+g)$ is also in $W^{1, p}\left(\Omega_{0}\right)$. To prove that the norms are preserved, first observe that

$$
\begin{gathered}
P_{H} u=\left\{\begin{array}{lll}
u & \text { a.e., } & \text { in }\left[\left(\Omega_{0} \cap H\right) \cap\{u>v\}\right] \cup\left[\left(\Omega_{0} \cap H^{\mathrm{c}}\right) \cap\{u<v\}\right], \\
v & \text { a.e., } & \text { in }\left[\left(\Omega_{0} \cap H^{\mathrm{c}}\right) \cap\{u>v\}\right] \cup\left[\left(\Omega_{0} \cap H\right) \cap\{u<v\}\right] ;
\end{array}\right. \\
\nabla P_{H} u=\left\{\begin{array}{lll}
\nabla u & \text { a.e., } & \text { in }\left[\left(\Omega_{0} \cap H\right) \cap\{u>v\}\right] \cup\left[\left(\Omega_{0} \cap H^{\mathrm{c}}\right) \cap\{u<v\}\right], \\
\nabla v & \text { a.e., } & \text { in }\left[\left(\Omega_{0} \cap H^{\mathrm{c}}\right) \cap\{u>v\}\right] \cup\left[\left(\Omega_{0} \cap H\right) \cap\{u<v\}\right] .
\end{array}\right.
\end{gathered}
$$

Now, by integrating $\left|P_{H}(u)\right|^{p}$ and $\left|\nabla P_{H}(u)\right|^{p}$ over $\Omega_{0}$, and using $\sigma_{H}\left(\left(\Omega_{0} \cap H\right) \cap\{u>v\}\right)=\left(\Omega_{0} \cap \bar{H}^{c}\right) \cap\{u<$ $v\}$ we get (2.5).

Recall that, for a domain $\Omega \subseteq \mathbb{R}^{d}$ and $\Gamma_{D} \subseteq \partial \Omega$, the Sobolev space $W_{\Gamma_{D}}^{1, p}(\Omega)$ is defined by

$$
W_{\Gamma_{D}}^{1, p}(\Omega)=\text { the closure of } \mathcal{C}_{\Gamma_{D}}^{0,1}(\Omega) \text { in } W^{1, p}(\Omega)
$$

where $\mathcal{C}_{\Gamma_{D}}^{0,1}(\Omega)=\left\{\varphi \in \mathcal{C}^{0,1}(\Omega): \operatorname{supp}(\phi) \cap \Gamma_{D}=\emptyset\right\}$. We give the analogous result of Proposition 2.18 for the functions in $W_{\Gamma_{D}}^{1, p}(\Omega)$ in the following proposition.
Proposition 2.19. Let $\Omega=\Omega_{\text {out }} \backslash \overline{\Omega_{\mathrm{in}}} \subset \mathbb{R}^{d}$ be as given in $\left(\mathbf{A}_{\mathbf{0}}\right), \Gamma_{D} \subseteq \partial \Omega$ and $H \in \mathcal{H}_{\text {ad }}$. Let $\varphi \in \mathcal{C}_{\Gamma_{D}}^{0,1}(\Omega)$ be any non-negative function.
(i) If $\Gamma_{D}=\partial \Omega_{\text {out }}$ and $\sigma_{H}\left(\Omega_{\mathrm{in}}\right)=\Omega_{\mathrm{in}}$, then $P_{H}(\varphi) \in \mathcal{C}_{\partial P_{H}\left(\Omega_{\mathrm{out})}\right)}^{0,1}\left(P_{H}(\Omega)\right)$.
(ii) If $\Gamma_{D}=\partial \Omega_{\mathrm{in}}$ and $\sigma_{H}\left(\Omega_{\mathrm{out}}\right)=\Omega_{\mathrm{out}}$, then $P_{H}(\varphi) \in \mathcal{C}_{\partial P^{H}\left(\Omega_{\mathrm{in})}\right.}^{0,1}\left(P_{H}(\Omega)\right)$.

In both of the cases (2.5) holds.
Proof. (i) Let $\Omega_{0}=\mathbb{R}^{d} \backslash \overline{\Omega_{\mathrm{in}}}$. Then $\Omega \subset \Omega_{0}$, and $\sigma_{H}\left(\Omega_{0}\right)=\Omega_{0}$. Let $\varphi \in \mathcal{C}_{\Gamma_{D}}^{0,1}(\Omega)$ be a non-negative function, and let $\widetilde{\varphi}$ be its zero extension to $\Omega_{0}$. Then $\widetilde{\varphi} \in \mathcal{C}^{0,1}\left(\Omega_{0}\right)$, and hence by Proposition $2.16, P_{H}(\widetilde{\varphi}) \in \mathcal{C}^{0,1}\left(\Omega_{0}\right)$. Therefore, $P_{H}(\varphi)=P_{H}(\widetilde{\varphi}) \mathbb{1}_{P_{H}(\Omega)} \in \mathcal{C}^{0,1}\left(P_{H} \Omega\right)$. Next, we show that $P_{H}(\varphi)=0$ on $\partial P_{H}\left(\Omega_{\text {out }}\right)$. Let $M=\operatorname{supp}(\varphi) \subsetneq \Omega_{\mathrm{out}}$. Since $\operatorname{supp}\left(P_{H}(\varphi)\right) \subseteq P_{H}(M)$ is closed, $P_{H}\left(\Omega_{\mathrm{out}}\right)$ is open and $P_{H}(M) \subset P_{H}\left(\Omega_{\mathrm{out}}\right)$, we obtain $\operatorname{supp}\left(P_{H}(\varphi)\right) \cap \partial P_{H}\left(\Omega_{\text {out }}\right)=\emptyset$ as required.
(ii) In this case, let $\Omega_{0}=\Omega_{\text {out }}$. For a non-negative function $\varphi \in W_{\Gamma_{D}}^{1, p}(\Omega)$, as before we get $P_{H}(\varphi)=$ $P_{H}(\widetilde{\varphi}) \mathbb{1}_{P_{H}(\Omega)} \in \mathcal{C}^{0,1}\left(P_{H}(\Omega)\right)$. Let $M=\operatorname{supp}(\varphi)$. Then $M \cap \overline{\Omega_{\mathrm{in}}}=\emptyset$ and $M \subset{\overline{\Omega_{\mathrm{in}}}}^{\mathrm{c}}$. Now, using Proposition 2.1 we obtain

$$
P_{H}(M) \subset P_{H}\left({\overline{\Omega_{\mathrm{in}}}}^{\mathrm{c}}\right) \subseteq P_{H}\left(\Omega_{\mathrm{in}}^{\mathrm{c}}\right)=\left(P^{H}\left(\Omega_{\mathrm{in}}\right)\right)^{\mathrm{c}}
$$

Since $\operatorname{supp}\left(P_{H}(\varphi)\right) \subseteq P_{H}(M)$ is closed, and $P^{H}\left(\Omega_{\mathrm{in}}\right)$ is open, we get $\operatorname{supp}\left(P_{H}(\varphi)\right) \cap \partial P^{H}\left(\Omega_{\mathrm{in}}\right)=\emptyset$. Therefore, $P_{H}(\varphi)=0$ on $\partial P^{H}\left(\Omega_{\mathrm{in}}\right)$.

Using the standard approximation techniques and Proposition 2.17 (the non-expansivity of polarizations), we can prove the following analogous result of Proposition 2.18, for the functions in $W_{\Gamma_{D}}^{1, p}(\Omega)$.

Proposition 2.20. Let $\Omega, \Gamma_{D} \subseteq \partial \Omega$ and $H$ be as given in Proposition 2.19. Let $u \in W_{\Gamma_{D}}^{1, p}(\Omega)$ be any non-negative function.
(i) If $\Gamma_{D}=\partial \Omega_{\text {out }}$ and $\sigma_{H}\left(\Omega_{\mathrm{in}}\right)=\Omega_{\mathrm{in}}$, then $P_{H}(u) \in W_{\partial P_{H}\left(\Omega_{\mathrm{out})}\right)}^{1, p}\left(P_{H}(\Omega)\right)$.
(ii) If $\Gamma_{D}=\partial \Omega_{\mathrm{in}}$ and $\sigma_{H}\left(\Omega_{\mathrm{out}}\right)=\Omega_{\mathrm{out}}$, then $P_{H}(u) \in W_{\partial P^{H}\left(\Omega_{\mathrm{in})}\right.}^{1, p}\left(P_{H}(\Omega)\right)$.

In both of the cases (2.5) holds.
2.4. Regularity results and Strong comparison principles. Next, we recall a few regularity results for the eigenfunctions. Using Moser type iteration arguments [7, Proposition 1.2] and the arguments from [8, Remark 2.8], we get that the eigenfunctions are in $L^{q}$ for any $q \in[1, \infty]$. Now, the local $\mathcal{C}^{1, \alpha}$-regularity results of [17, Theorem 1 and 2] give the following boundary regularity of the eigenfunctions.

Proposition 2.21. Let $\Omega \subseteq \mathbb{R}^{d}$ be a Lipschitz domain and $1<p<\infty$. Let $u \in W_{\text {loc }}^{1, p}(\Omega) \cap L_{\text {loc }}^{\infty}(\Omega)$ be a weak solution of $-\Delta_{p} u=\lambda|u|^{p-2} u$ for some $\lambda \in \mathbb{R}$. Then there exists $\alpha \in(0,1)$ such that $u \in \mathcal{C}_{\operatorname{loc}}^{1, \alpha}(\Omega) \cap \mathcal{C}^{0, \alpha}(\bar{\Omega})$.

The following strong comparison principle for the distributional solutions of the $p$-Laplace operator is given in [27, Theorem 1.4].

Proposition 2.22. Let $\Omega \subset \mathbb{R}^{d}$ be a bounded smooth domain and $\frac{2 d+2}{d+2}<p<\infty$. Let $u, v \in C^{1}(\bar{\Omega})$ be positive distributional solutions of $-\Delta_{p} u-g(u)=0$ in $\Omega$, for a non-negative Lipschitz function $g$ on $[0, \infty)$ with $g(s)>0$ for $s>0$. If $u \leq v$ in $\Omega$, then either $u<v$ in $\Omega$ or $u \equiv v$ in $\Omega$.

## 3. Strict Faber-Krahn type inequality under polarization

In this section, we give A proof for Theorem 1.3. Recall the following two subsets of $P_{H}(\Omega) \cap H$ :

$$
A_{H}=\Omega^{\mathrm{c}} \cap \sigma_{H}(\Omega) \cap H \text { and } B_{H}=\Omega \cap \sigma_{H}\left(\Omega^{\mathrm{c}}\right) \cap H
$$

We need the following lemma.
Lemma 3.1. Let $\Omega=\Omega_{\text {out }} \backslash \overline{\Omega_{\mathrm{in}}} \subset \mathbb{R}^{d}$ be a domain as given in $\left(\mathbf{A}_{\mathbf{0}}\right)$, and $H \in \mathcal{H}_{\mathrm{ad}}$. Then $\overline{\Omega \cap H} \cap \overline{A_{H}} \subseteq \partial \Omega$. Furthermore,
(i) if $\sigma_{H}\left(\Omega_{\text {in }}\right)=\Omega_{\text {in }}$ then $\overline{\Omega \cap H} \cap \overline{A_{H}} \subseteq \partial \Omega_{\text {out }}$;
(ii) if $\sigma_{H}\left(\Omega_{\text {out }}\right)=\Omega_{\text {out }}$ then $\overline{\Omega \cap H} \cap \overline{A_{H}} \subseteq \partial \Omega_{\text {in }}$.

Proof. If $A_{H}=\emptyset$, then trivially $\emptyset=\overline{\Omega \cap H} \cap \overline{A_{H}} \subset \partial \Omega$. Let $A_{H} \neq \emptyset$, then from Proposition 2.3, we obtain $P_{H}(\Omega) \neq \Omega$, and $A_{H}$ has non-empty interior. Since $\Omega \cap H$ and $A_{H}$ are disjoint and $P_{H}(\Omega) \cap H=(\Omega \cap H) \cup A_{H}$, using the connectedness of $P_{H}(\Omega) \cap H$ we conclude that $\overline{\Omega \cap H} \cap \overline{A_{H}}=\partial(\Omega \cap H) \cap \partial A_{H} \neq \emptyset$. Clearly $\Omega \cap \partial(\Omega \cap H) \cap \partial A_{H}=\emptyset$ and hence $\partial(\Omega \cap H) \cap \partial A_{H} \subseteq \partial \Omega$.
(i) If $\sigma_{H}\left(\Omega_{\mathrm{in}}\right)=\Omega_{\mathrm{in}}$, then we can write

$$
\begin{aligned}
A_{H}=\Omega^{\mathrm{c}} \cap \sigma_{H}(\Omega) \cap H & =\left(\Omega_{\mathrm{out}}^{\mathrm{c}} \cup \Omega_{\mathrm{in}}\right) \cap \sigma_{H}\left(\Omega_{\text {out }}\right) \cap \Omega_{\mathrm{in}}^{\mathrm{c}} \cap H \\
& =\Omega_{\text {out }}^{\mathrm{c}} \cap \sigma_{H}\left(\Omega_{\text {out }}\right) \cap H .
\end{aligned}
$$

Since $\Omega_{\text {in }} \subset \subset \Omega_{\text {out }}$, we have $\partial \Omega=\partial \Omega_{\text {out }} \sqcup \partial \Omega_{\text {in }}$ and $\partial \Omega_{\text {in }} \cap \partial A_{H}=\emptyset$. Therefore, $\overline{\Omega \cap H} \cap \overline{A_{H}} \subset \partial \Omega_{\text {out }}$. (ii) Similarly, for $\sigma_{H}\left(\Omega_{\text {out }}\right)=\Omega_{\text {out }}$ we have $A_{H}=\Omega_{\mathrm{in}} \cap \sigma_{H}\left(\Omega_{\mathrm{in}}\right) \cap H$. Therefore, we obtain $\overline{\Omega \cap H} \cap \overline{A_{H}} \subset$ $\partial \Omega_{\mathrm{in}}$.

For any non-negative function $u \in \mathcal{C}(\bar{\Omega})$, let $\widetilde{u}$ be its zero extension to $\mathbb{R}^{d}$ and let

$$
M_{u}=\left\{x \in P_{H}(\Omega) \cap H: P_{H}(\widetilde{u})(x)>\widetilde{u}(x)\right\}
$$

Next, we prove a lemma that plays a significant role in our results.
Lemma 3.2. Let $\Omega=\Omega_{\text {out }} \backslash \overline{\Omega_{\mathrm{in}}} \subset \mathbb{R}^{d}$ be a domain as given in $\left(\mathbf{A}_{\mathbf{0}}\right), \Gamma_{D} \subseteq \partial \Omega$, and $H \in \mathcal{H}_{\mathrm{ad}}$. Let $u \in \mathcal{C}(\bar{\Omega})$ be a non-negative function with $u=0$ on $\Gamma_{D}$. If $\Gamma_{D}$ satisfies one of the following assumptions:

$$
\text { (a) } \Gamma_{D}=\partial \Omega, \quad \text { (b) } \Gamma_{D}=\partial \Omega_{\mathrm{out}} \text { and } \sigma_{H}\left(\Omega_{\mathrm{in}}\right)=\Omega_{\mathrm{in}}, \quad\left(\text { c) } \Gamma_{D}=\partial \Omega_{\mathrm{in}} \text { and } \sigma_{H}\left(\Omega_{\mathrm{out}}\right)=\Omega_{\mathrm{out}}\right.
$$

then $\widetilde{u}$ is continuous on $P_{H}(\Omega) \cap H$. Moreover, if $\Omega \neq P_{H}(\Omega) \neq \sigma_{H}(\Omega)$ then there exists a ball $B_{0} \subset \Omega \cap H$ such that

$$
P_{H}(u)>u \text { in } B_{0} \cap M_{u} \text { and } P_{H}(u) \equiv u \text { in } B_{0} \cap M_{u}^{\mathrm{c}} .
$$

Proof. If $P_{H}(\Omega)=\Omega$ then $\widetilde{u}=u$ in $P_{H}(\Omega) \cap H$, and hence it is continuous. If $P_{H}(\Omega) \neq \Omega$, then from Proposition 2.3 we have $A_{H} \neq \emptyset$ and $\Omega \cap H \subsetneq P_{H}(\Omega) \cap H$. Clearly $\widetilde{u}=u$ on $\Omega \cap H$ and $\widetilde{u}=0$ on $A_{H}$, and hence $\widetilde{u}$ is continuous on both $\Omega \cap H$ and $A_{H}$. If $\Gamma_{D}$ satisfies one of the assumptions $(a)$ - $(c)$, then by Lemma 3.1 we get $\emptyset \neq \overline{\Omega \cap H} \cap \overline{A_{H}} \subseteq \Gamma_{D}$. Therefore, $\widetilde{u}=u=0$ on $\overline{\Omega \cap H} \cap \overline{A_{H}}$ and hence $\widetilde{u}$ is continuous on $(\Omega \cap H) \cup A_{H}=P_{H}(\Omega) \cap H$.
Now assume that $\Omega \neq P_{H}(\Omega) \neq \sigma_{H}(\Omega)$. Then from Proposition 2.3, both $A_{H}$ and $B_{H}$ have non-empty interiors. By the definition of $P_{H}(\widetilde{u})$, we get $P_{H}(\widetilde{u}) \geq \widetilde{u}$ in $P_{H}(\Omega) \cap H$, and

$$
\begin{align*}
& \text { in } A_{H}: \widetilde{u}=0, \widetilde{u} \circ \sigma_{H}=u \circ \sigma_{H}>0 \text { and hence } P_{H}(\widetilde{u})=u \circ \sigma_{H}>\widetilde{u} \text {; } \\
& \text { in } B_{H}: \widetilde{u}=u>0, \widetilde{u} \circ \sigma_{H}=0 \text { and hence } P_{H}(\widetilde{u})=\widetilde{u} . \tag{3.1}
\end{align*}
$$

Let $N=\left\{x \in P_{H}(\Omega) \cap H: P_{H}(\widetilde{u})(x)=\widetilde{u}(x)\right\}$. Since $P_{H}(\widetilde{u})$ is also continuous on $P_{H}(\Omega) \cap H$ (Proposition 2.16), from (3.1) we get $N \subsetneq \Omega \cap H$ is a non-empty closed set and $M_{u}=\left(P_{H}(\Omega) \cap H\right) \backslash N$ is a non-empty open set in $P_{H}(\Omega) \cap H$. Now, by the connectedness of $P_{H}(\Omega) \cap H$ we must have $\partial M_{u} \cap N \neq \emptyset$. For $x_{0} \in \partial M_{u} \cap N$, let $B_{0}=B_{r}\left(x_{0}\right) \subset \Omega \cap H$. Then $B_{0}$ has all the the desired properties.

Now, we prove Theorem 1.3.
Proof of Theorem 1.3. Let $1<p<\infty, \Omega_{\text {out }} \backslash \overline{\Omega_{\mathrm{in}}} \subset \mathbb{R}^{d}$ be as given in $\left(\mathbf{A}_{\mathbf{0}}\right)$ and $H \in \mathcal{H}_{\text {ad }}$. Denote $\Omega=\Omega_{\text {out }} \backslash \overline{\Omega_{\text {in }}}$.
(i) Assume that $\sigma_{H}\left(\Omega_{\mathrm{in}}\right)=\Omega_{\mathrm{in}}$. Let $0<u \in \mathcal{C}_{\partial P_{H}\left(\Omega_{\text {out }}\right)}^{0, \alpha}(\bar{\Omega})$ be an eigenfunction corresponding to $\nu_{1}(\Omega)$. Define $v=P_{H}(u)$ in $P_{H}(\Omega)$ then, from Proposition 2.19, we get $v \in \mathcal{C}_{\partial P_{H}\left(\Omega_{\mathrm{out})}\right.}^{0, \alpha}\left(P_{H}(\Omega)\right)$, and

$$
\|u\|_{p, \Omega}=\|v\|_{p, P_{H}(\Omega)} \text { and }\|\nabla u\|_{p, \Omega}=\|\nabla v\|_{p, P_{H}(\Omega)} .
$$

From the variational characterization of $\nu_{1}\left(P_{H}(\Omega)\right)$, we obtain:

$$
\begin{equation*}
\nu_{1}\left(P_{H}(\Omega)\right) \leq \nu_{1}(\Omega) \tag{3.2}
\end{equation*}
$$

(ii) Assume that $\sigma_{H}\left(\Omega_{\mathrm{out}}\right)=\Omega_{\mathrm{out}}$. Let $0<u \in \mathcal{C}_{\partial P^{H}\left(\Omega_{\mathrm{in})}\right.}^{0, \alpha}(\bar{\Omega})$ be an eigenfunction corresponding to $\tau_{1}(\Omega)$. Define $v=P_{H}(u)$ in $P_{H}(\Omega)$ then, from Proposition 2.19, we obtain $v \in \mathcal{C}_{\partial P^{H}\left(\Omega_{\mathrm{in})}\right)}^{0, \alpha}\left(P_{H}(\Omega)\right)$, and

$$
\|u\|_{p, \Omega}=\|v\|_{p, P_{H}(\Omega)} \text { and }\|\nabla u\|_{p, \Omega}=\|\nabla v\|_{p, P_{H}(\Omega)} .
$$

From the variational characterization of $\tau_{1}\left(P_{H}(\Omega)\right)$, we get

$$
\begin{equation*}
\tau_{1}\left(P_{H}(\Omega)\right) \leq \tau_{1}(\Omega) \tag{3.3}
\end{equation*}
$$

(iii) Let $\frac{2 d+2}{d+2}<p<\infty$. Assume that, the equality holds in (3.2). Let $0 \leq u \in \mathcal{C}_{\partial P^{H}\left(\Omega_{\text {out }}\right)}^{0, \alpha^{\prime}}(\bar{\Omega})$ be an eigenfunction corresponding to the eigenvalue $\nu_{1}(\Omega)$. On the contrary, assume that $\Omega \neq P_{H}(\Omega) \neq \sigma_{H}(\Omega)$. Then by Lemma 3.2, there exists a ball $B_{0} \subset \Omega \cap H$ such that

$$
\begin{equation*}
v>u \text { in } B_{0} \cap M_{u} \text { and } v \equiv u \text { in } B_{0} \cap M_{u}^{\mathrm{c}} \tag{3.4}
\end{equation*}
$$

where $M_{u}=\left\{x \in P_{H}(\Omega) \cap H: v(x)>u(x)\right\}$ is a non-empty open set. Then, both $u, v \in \mathcal{C}^{1}\left(\overline{B_{0}}\right)$ are positive distributional solutions for the following problem in $B_{0}$ :

$$
-\Delta_{p} u-\lambda|u|^{p-2} u=0 \text { in } B_{0}
$$

Now, for $\frac{2 d+2}{d+2}<p<\infty$, the strong comparison principle (Proposition 2.22) implies that
either $u<v$ or $u \equiv v$ in $B_{0}$.
This is a contradiction to (3.4), and hence we must have $P_{H}(\Omega)=\Omega$ or $P_{H}(\Omega)=\sigma_{H}(\Omega)$. If the equality holds in (3.3), the proof will follow using a similar set of arguments as given above.

## 4. Strict monotonicity of first eigenvalues via polarization

In this section, we prove Theorem 1.5 and Theorem 1.8. The main idea is to express the translations and the rotations of the obstacle $\mathcal{O}$ in terms of polarizations and apply the Faber-Krahn type inequality to get the desired monotonicity.
4.1. Monotonicity along a straight line. Now, we give a proof for Theorem 1.5. First, we recall that:

$$
\text { for given } h \in \mathbb{S}^{d-1}, \quad H_{s}:=\left\{x \in \mathbb{R}^{d}: x \cdot h<s\right\}, \quad \Sigma_{s}:=\{x \in \Omega: x \cdot h \geq s\}, \text { for } s \in \mathbb{R},
$$

$P_{H_{0}}(\Omega)=\Omega$, the obstacle $\mathcal{O}$ is Steiner symmetric with respect to $\partial H_{0}$, and the translations of $\mathcal{O}$ in the $h$-direction are given by $\mathcal{O}_{s}=s h+\mathcal{O}$ for $s \in \mathbb{R}$, and

$$
\mathrm{L}_{\mathcal{O}}:=\left\{s \in \mathbb{R}: P_{H_{s}}(\Omega)=\Omega \text { and } \mathcal{O}_{s} \subset \Omega\right\} .
$$

We observe the following facts:

$$
\begin{equation*}
\text { for } x \in \mathbb{R}^{d}, \quad \sigma_{H_{0}}(x)=x-2(x \cdot h) h, \text { and } \sigma_{H_{s}}(x)=2 s h+\sigma_{H_{0}}(x) \text { for } s \in \mathbb{R} \tag{4.1}
\end{equation*}
$$

Lemma 4.1. If $\mathcal{O} \subseteq \mathbb{R}^{d}$ satisfies $\sigma_{H_{0}}(\mathcal{O})=\mathcal{O}$, then for any $s, t \in \mathbb{R}, t h+\mathcal{O}_{s}=\mathcal{O}_{s+t}$ and $\sigma_{H_{t}}\left(\mathcal{O}_{s}\right)=\mathcal{O}_{2 t-s}$.
Proof. It is easy to verify that $t h+\mathcal{O}_{s}=(s+t) h+\mathcal{O}=\mathcal{O}_{s+t}$ and $\sigma_{H_{s}}\left(\mathcal{O}_{s}\right)=\mathcal{O}_{s}$. Since $\sigma_{H_{t}}(x)=2 t h+\sigma_{H_{0}}(x)=$ $2(t-s) h+2 s h+\sigma_{H_{0}}(x)=2(t-s) h+\sigma_{H_{s}}(x)$ any $x \in \mathbb{R}^{d}$, we get $\sigma_{H_{t}}\left(\mathcal{O}_{s}\right)=2(t-s) h+\sigma_{H_{s}}\left(\mathcal{O}_{s}\right)=2(t-s) h+\mathcal{O}_{s}=$ $\mathcal{O}_{2 t-s}$.

Proof of Theorem 1.5. Let the set $\Sigma_{s_{0}} \bigcup \sigma_{H_{s_{0}}}\left(\Sigma_{s_{0}}\right)$ is convex in the $h$-direction for some $s_{0} \in L_{\mathcal{O}}$. Let $\mathrm{R}_{\mathcal{O}}:=\sup \mathrm{L}_{\mathcal{O}}$.
The interval $\left[s_{0}, \mathrm{R}_{\mathcal{O}}\right) \subseteq \mathrm{L}_{\mathcal{O}}$ : Let $s \in\left(s_{0}, \mathrm{R}_{\mathcal{O}}\right)$. Clearly, $\Sigma_{s} \subset \Sigma_{s_{0}}$ and the convexity of the set $\Sigma_{s_{0}} \cup \sigma_{H_{s_{0}}}\left(\Sigma_{s_{0}}\right)$ in the $h$-direction implies that $\sigma_{H_{s}}\left(\Sigma_{s}\right) \subset \Sigma_{s_{0}} \bigcup \sigma_{H_{s_{0}}}\left(\Sigma_{s_{0}}\right) \subset \Omega$ (see Proposition 2.10-(i)). Therefore, by Proposition 2.2, we get $P_{H_{s}}(\Omega)=\Omega$. Next, we show $\mathcal{O}_{s} \subset \Omega$. By the definition of $\mathrm{R}_{\mathcal{O}}$, there exists $s_{1} \in \mathrm{~L}_{\mathcal{O}}$ such that $s<s_{1}<\mathrm{R}_{\mathcal{O}}$. Observe that $\mathcal{O}_{s_{0}}, \mathcal{O}_{s_{1}} \subset \Sigma_{s_{0}} \bigcup \sigma_{H_{s_{0}}}\left(\Sigma_{s_{0}}\right)$, and $s=t s_{0}+(1-t) s_{1}$, for some $t \in(0,1)$. Thus $\mathcal{O}_{s}=t \mathcal{O}_{s_{0}}+(1-t) \mathcal{O}_{s_{1}}$, and hence the convexity of $\Sigma_{s_{0}} \cup \sigma_{H_{s_{0}}}\left(\Sigma_{s_{0}}\right)$ in the $h$-direction implies that $\mathcal{O}_{s} \subset \Omega$. Therefore $s \in \mathrm{~L}_{\mathcal{O}}$.
 $\overline{\mathcal{O}_{s}}$ is Steiner symmetric with respect to $\partial H_{s}$ and $\bar{s}>s$, from Proposition 2.10-(i), we get $P^{H_{\bar{s}}}\left(\mathcal{O}_{s}\right)=\sigma_{H_{\bar{s}}}\left(\mathcal{O}_{s}\right)$. From Lemma 4.1 (since $\left.\sigma_{H_{0}}(\mathcal{O})=\mathcal{O}\right)$ we also have $\sigma_{H_{\bar{s}}}\left(\mathcal{O}_{s}\right)=\mathcal{O}_{2 \bar{s}-s}=\mathcal{O}_{t}$. Therefore, from Proposition 2.12 we obtain

$$
P_{H_{\bar{s}}}\left(\Omega \backslash \mathcal{O}_{s}\right)=P_{H_{\bar{s}}}(\Omega) \backslash P^{H_{\bar{s}}}\left(\mathcal{O}_{s}\right)=\Omega \backslash \mathcal{O}_{t} .
$$

For $\frac{2 d+2}{d+2}<p<\infty$, the Faber-Krahn type inequality (Theorem 1.3) implies that $\lambda_{1}(t)<\lambda_{1}(s)$. Therefore, the first Dirichlet eigenvalue $\lambda_{1}(\cdot)$ is strictly decreasing on $\left[s_{0}, \mathrm{R}_{\mathcal{O}}\right)$.

Remark 4.2. If we drop the convexity assumption from Theorem 1.5, then $L_{0}$ might not be an interval. However, for any $s, t \in L_{0}$ with $\frac{s+t}{2} \in L_{0}$ and $s<t$, the above proof still yields $\lambda_{1}(t)<\lambda_{1}(s)$.
4.2. Monotonicity with respect to the rotations about a point. Now, we prove Theorem 1.8. First recall that, for $\xi \in \mathbb{S}^{d-1} \backslash\{\eta\}$ the rotations of the obstacle $\mathcal{O}$ with the plane of rotation is $X_{\xi}=\operatorname{span}\{\eta, \xi\}$ about the point $a \in \mathbb{R}^{d}$ are given by: for $s \in[-1,1]$,

$$
\mathcal{O}_{s, \xi}:=a+R_{s, \xi}(-a+\mathcal{O})
$$

where $R_{s, \xi}$ is the simple rotation in $\mathbb{R}^{d}$ with $X_{\xi}$ as the plane of rotation and $\theta_{s}=\arccos (s) \in[0, \pi]$ as the angle of rotation from the ray $\mathbb{R}^{+} \eta$ in the counter-clockwise direction. We prove the following lemmas.

Lemma 4.3. For any distinct $\xi_{1}, \xi_{2} \in \mathbb{S}^{d-1} \backslash\{\eta\}$, there exists a simple rotation $R$ such that

$$
R\left(-a+\Omega \backslash \mathcal{O}_{s, \xi_{1}}\right)=-a+\Omega \backslash \mathcal{O}_{s, \xi_{2}}
$$

Proof. Let $\xi_{1}, \xi_{2} \in \mathbb{S}^{d-1} \backslash\{\eta\}$, define

$$
\begin{equation*}
\widetilde{\xi}_{i}=\frac{\xi_{i}-\left(\xi_{i} \cdot \eta\right) \eta}{\left\|\xi_{i}-\left(\xi_{i} \cdot \eta\right) \eta\right\|} \text { for } i=1,2 \tag{4.2}
\end{equation*}
$$

Observe that, the rotation of $\eta$ under $R_{s, \xi_{i}}$ is given by $R_{s, \xi_{i}}(\eta)=s \eta+\sqrt{1-s^{2}} \widetilde{\xi}_{i}$ for $i=1,2$. Consider the plane $X=\operatorname{span}\left\{\widetilde{\xi_{1}}, \widetilde{\xi_{2}}\right\}$, that is orthogonal to $\eta$. Let $R$ be the simple rotation such that $R\left(\widetilde{\xi_{1}}\right)=\widetilde{\xi_{2}}$. Thus $R$ must fix $\eta$, and

$$
R \circ R_{s, \xi_{1}}(\eta)=R\left(s \eta+\sqrt{1-s^{2}} \widetilde{\xi_{1}}\right)=s \eta+\sqrt{1-s^{2}} \widetilde{\xi_{2}}=R_{s, \xi_{2}}(\eta)
$$

Therefore, for $r>0, \rho>0$,

$$
R\left(B_{\rho}\left(r R_{s, \xi_{1}}(\eta)\right) \cap \partial B_{r}(0)\right)=B_{\rho}\left(r R \circ R_{s, \xi_{1}}(\eta)\right) \cap \partial B_{r}(0)=B_{\rho}\left(r R_{s, \xi_{2}}(\eta)\right) \cap \partial B_{r}(0),
$$

and from (2.2) we obtain $R\left(R_{s, \xi_{1}}(-a+\mathcal{O})\right)=R_{s, \xi_{2}}(-a+\mathcal{O})$, and hence $R\left(-a+\mathcal{O}_{s, \xi_{1}}\right)=-a+\mathcal{O}_{s, \xi_{2}}$. Since $R$ fixes $\eta$, and $\Omega$ is foliated Schwarz symmetric with respect to $a+\mathbb{R}^{+} \eta$, we get $R(-a+\Omega)=-a+\Omega$. Thus we obtain

$$
\begin{aligned}
R\left(-a+\Omega \backslash \mathcal{O}_{s, \xi_{1}}\right) & =R(-a+\Omega) \backslash R\left(-a+\mathcal{O}_{s, \xi_{1}}\right) \\
& =(-a+\Omega) \backslash\left(-a+\mathcal{O}_{s, \xi_{2}}\right)=-a+\Omega \backslash \mathcal{O}_{s, \xi_{2}}
\end{aligned}
$$

From Lemma 4.3, we only need to consider the rotations of the obstacle by $R_{s, \xi}$ about the point $a$ in a $X_{\xi}$-plane for a fixed $\xi \in \mathbb{S}^{d-1} \backslash\{\eta\}$. Thus for $s \in[-1,1]$, we set $\mathcal{O}_{s}=\mathcal{O}_{s, \xi}$. Recall that:

$$
\text { for } a \in \mathbb{R}^{d} \text { and } \eta \in \mathbb{S}^{d-1}, \quad \mathcal{H}_{a, \eta}=\left\{H \in \mathscr{H}: a \in \partial H \text { and } a+\mathbb{R}^{+} \eta \subset H\right\} .
$$

Lemma 4.4. Let $a \in \mathbb{R}^{d}, \eta \in \mathbb{S}^{d-1}$, and $\mathcal{O} \subset \mathbb{R}^{d}$ is foliated Schwarz symmetric with respect to the ray $a+\mathbb{R}^{+} \eta$. Let $\xi \in \mathbb{S}^{d-1} \backslash\{\eta\}$ and the rotations of $\mathcal{O}$ be as given in (1.11). Then, for any $s<t$ in $[-1,1]$ there exists $H \in \mathcal{H}_{a, \eta}$ such that

$$
\text { (a) } \sigma_{H}\left(\mathcal{O}_{t}\right)=\mathcal{O}_{s}, \quad \text { (b) } P^{H}\left(\mathcal{O}_{t}\right)=\mathcal{O}_{s} \quad \text { and (c) } P_{H}\left(\mathcal{O}_{s}\right)=\mathcal{O}_{t} .
$$

Proof. Let $h=R_{s}(\eta)-R_{t}(\eta)$ and consider the polarizer $H:=\left\{z \in \mathbb{R}^{d}:(z-a) \cdot h<0\right\}$. Observe that $a \in \partial H$, and for $r>0$,

$$
\begin{aligned}
(a+r \eta-a) \cdot h & =r \eta \cdot\left[R_{s}(\eta)-R_{t}(\eta)\right]=r(s-t)<0 \\
\left(a+r R_{t}(\eta)-a\right) \cdot h & =r R_{t}(\eta) \cdot\left[R_{s}(\eta)-R_{t}(\eta)\right]=r\left[R_{s}(\eta) \cdot R_{t}(\eta)-1\right]<0
\end{aligned}
$$

Therefore $H \in \mathcal{H}_{a, \eta} \cap \mathcal{H}_{a, R_{t}(\eta)}$.
(a) Notice that, $\|h\|=2\left[1-R_{s}(\eta) \cdot R_{t}(\eta)\right]$. Now, for $x=a+r R_{t}(\eta), r>0$ we get

$$
\sigma_{H}(x)=x-\frac{2(x-a) \cdot h}{\|h\|^{2}} h=a+r R_{t}(\eta)-\frac{2 r\left[R_{s}(\eta) \cdot R_{t}(\eta)-1\right]}{2\left[1-R_{s}(\eta) \cdot R_{t}(\eta)\right]}\left(R_{s}(\eta)-R_{t}(\eta)\right)=a+r R_{s}(\eta) .
$$

Therefore, $\sigma_{H}\left(a+\mathbb{R}^{+} R_{t}(\eta)\right)=a+\mathbb{R}^{+} R_{s}(\eta)$, and hence from (2.2) we obtain

$$
\sigma_{H}\left(\mathcal{O}_{t}\right)=\bigcup_{r \in I_{\mathcal{O}}} B_{\rho(r)}\left(\sigma_{H}\left(a+r R_{t}(\eta)\right)\right) \cap \partial B_{r}\left(\sigma_{H}(a)\right)=\bigcup_{r \in I_{\mathcal{O}}} B_{\rho(r)}\left(a+r R_{s}(\eta)\right) \cap \partial B_{r}(a)=\mathcal{O}_{s} .
$$

(b) Since $H \in \mathcal{H}_{a, R_{t}(\eta)}$ and $\mathcal{O}_{t}$ is foliated Schwarz symmetric with respect to $a+\mathbb{R}^{+} R_{t}(\eta)$, Proposition 2.10-
(ii) implies that $P^{H}\left(\mathcal{O}_{t}\right)=\sigma_{H}\left(\mathcal{O}_{t}\right)=\mathcal{O}_{s}$.
(c) Since $P_{H}\left(\sigma_{H}\left(\mathcal{O}_{t}\right)\right)=P_{H}\left(\mathcal{O}_{t}\right)$ (Proposition 2.1-(iv)), we get $P_{H}\left(\mathcal{O}_{s}\right)=P_{H}\left(\sigma_{H}\left(\mathcal{O}_{t}\right)\right)=P_{H}\left(\mathcal{O}_{t}\right)=\mathcal{O}_{t}$.

Proof of Theorem 1.8. Given $\frac{2 d+2}{d+2}<p<\infty, \Omega$ and $\mathcal{O}$ are foliated Schwarz symmetric with respect to $a+\mathbb{R}^{+} \eta$.
The set $\mathrm{C}_{\mathcal{O}}$ is interval: We show that, for any $s \in C_{\mathcal{O}}$, the interval $[s, 1] \subseteq C_{\mathcal{O}}$. Let $t \in(s, 1]$. From Lemma 4.4-(c) there exists $H \in \mathcal{H}_{a, \eta}$ such that $P_{H}\left(\mathcal{O}_{s}\right)=\mathcal{O}_{t}$. Since $H \in \mathcal{H}_{a, \eta}$ and $\Omega$ is foliated Schwarz symmetric with respect to $a+\mathbb{R}^{+} \eta$, from Proposition 2.10-(ii), we get $P_{H}(\Omega)=\Omega$. Now, $\mathcal{O}_{s} \subset \Omega$ implies that $\mathcal{O}_{t}=P_{H}\left(\mathcal{O}_{s}\right) \subset P_{H}(\Omega)=\Omega$. Therefore, $t \in C_{\mathcal{O}}$ and hence $[s, 1] \subseteq C_{\mathcal{O}}$.
Monotonicity of $\gamma_{1}(\cdot)$ : Let $s<t$ in $\mathrm{C}_{\mathcal{O}}$. From Lemma 4.4, there exists $H \in \mathcal{H}_{a, \eta}$ such that $P^{H}\left(\mathcal{O}_{t}\right)=\mathcal{O}_{s}$. Since $\Omega$ is foliated Schwarz symmetric with respect to $a+\mathbb{R}^{+} \eta$, from Proposition 2.10-(ii), we have $P_{H}(\Omega)=\Omega$, and from Proposition 2.12 we get

$$
P_{H}\left(\Omega \backslash \mathcal{O}_{t}\right)=P_{H}(\Omega) \backslash P^{H}\left(\mathcal{O}_{t}\right)=\Omega \backslash \mathcal{O}_{s} .
$$

If $\Omega$ satisfies $\left(\mathbf{A}_{4}\right)$ then $\Omega=\Omega_{0} \backslash \bar{B}_{\rho}(a)$ and $\Gamma_{N}=\partial B_{\rho}(a)$. In this case, we set $\Omega_{\text {out }}=\Omega_{0} \backslash \mathcal{O}_{t}$ and $\Omega_{\mathrm{in}}=B_{\rho_{0}}(a)$ (in Theorem 1.3) so that $\Omega_{\text {out }} \backslash \overline{\Omega_{\mathrm{in}}}=\Omega \backslash \mathcal{O}_{t}, \Gamma_{N}=\partial \Omega_{\mathrm{in}}$ and $\Gamma_{D}=\partial \Omega_{\text {out }}$. Therefore, we have $\nu_{1}\left(\Omega_{\text {out }} \backslash \overline{\Omega_{\text {in }}}\right)=\gamma_{1}\left(\Omega \backslash \mathcal{O}_{t}\right)$ and $\nu_{1}\left(P_{H}\left(\Omega_{\text {out }} \backslash \overline{\Omega_{\text {in }}}\right)\right)=\gamma_{1}\left(\Omega \backslash \mathcal{O}_{s}\right)$. Since $\sigma_{H}\left(\Omega_{\text {in }}\right)=\Omega_{\text {in }}$, from Theorem 1.3-(i) we get

$$
\gamma_{1}\left(\Omega \backslash \mathcal{O}_{s}\right) \leq \gamma_{1}\left(\Omega \backslash \mathcal{O}_{t}\right)
$$

Similarly, if $\Omega$ satisfies $\left(\mathbf{A}_{\mathbf{5}}\right)$ then $\Omega=B_{R}(a) \backslash \overline{\Omega_{1}}$ and $\Gamma_{N}=\partial B_{R}(a)$. In this case, we set $\Omega_{\text {out }}=B_{R}(a)$ and $\Omega_{\mathrm{in}}=\Omega_{1} \cup \mathcal{O}_{t}$ (in Theorem 1.3) so that $\Omega_{\mathrm{out}} \backslash \overline{\Omega_{\mathrm{in}}}=\Omega \backslash \mathcal{O}_{t}, \Gamma_{N}=\partial \Omega_{\mathrm{out}}$ and $\Gamma_{D}=\partial \Omega_{\mathrm{in}}$. Therefore, we have $\tau_{1}\left(\Omega_{\text {out }} \backslash \overline{\Omega_{\text {in }}}\right)=\gamma_{1}\left(\Omega \backslash \mathcal{O}_{t}\right)$ and $\tau_{1}\left(P_{H}\left(\Omega_{\text {out }} \backslash \overline{\Omega_{\text {in }}}\right)\right)=\gamma_{1}\left(\Omega \backslash \mathcal{O}_{s}\right)$. Since $\sigma_{H}\left(\Omega_{\text {out }}\right)=\Omega_{\text {out }}$, from Theorem 1.3-(ii) we get

$$
\gamma_{1}\left(\Omega \backslash \mathcal{O}_{s}\right) \leq \gamma_{1}\left(\Omega \backslash \mathcal{O}_{t}\right)
$$

Since $\Omega$ is not radially symmetric with the center $a$, in both cases, we have $\Omega_{\text {out }} \backslash \overline{\Omega_{\text {in }}} \neq P_{H}\left(\Omega_{\text {out }} \backslash \overline{\Omega_{\text {in }}}\right) \neq$ $\sigma_{H}\left(\Omega_{\text {out }} \backslash \overline{\Omega_{\mathrm{in}}}\right)$. Thus, the strict Faber-Krahn type inequality (Theorem 1.3-(iii)) implies

$$
\gamma_{1}(s)<\gamma_{1}(t)
$$

Therefore, $\gamma_{1}(\cdot)$ is strictly increasing on $\mathrm{C}_{\mathcal{O}}$.

## 5. Some remarks and examples

Example 1. Let $\Omega \subset \mathbb{R}^{2}$ is given by $\Omega=\left\{(x, y): x^{2}+y^{2}<R^{2}, x \leq 0\right\} \cup\{(x, y):|x|+|y|<2 R, x \geq 0\}$ for $R>0$, and $\mathcal{O}$ is the rhombus given by $|x|+|y| \leq 2 \ell$ with $\ell<R$ (see Figure 4). Since $\mathcal{O}$ is Steiner symmetric with respect to the hyperplanes $S_{1}:=\left\{(x, y) \in \mathbb{R}^{2}: x=0\right\}$ and $S_{2}:=\left\{(x, y) \in \mathbb{R}^{2}: x+y=0\right\}$, we can consider the translations $\mathcal{O}$ along the $x$-axis, as well as along the straight line $y=x$.


Figure 4. Example 1

Along the $x$-axis, the translations of $\mathcal{O}$ are $\mathcal{O}_{s}=(s, 0)+\mathcal{O} \subset \Omega$ for $|s|<R-\ell$. For $|s|<R-\ell$, let $\lambda_{1}(s)=\lambda_{1}\left(\Omega \backslash \mathcal{O}_{s}\right)$, the first Dirichlet eigenvalue of the $p$-Laplacian, for $\frac{2 d+2}{d+2}<p<\infty$. Let

$$
s_{*}=\sup \left\{s \in(-R+\ell, 0): \sigma_{H_{s}}(\{(x, y) \in \Omega: x<s\}) \subset \Omega\right\}
$$

Now, applying Theorem 1.5
(i) with $h=(1,0)$, we get $\lambda_{1}(s)$ is strictly decreasing for $s \in[0, R-\ell)$,
(ii) with $h=(-1,0)$, we get $\lambda_{1}(s)$ is strictly increasing for $s \in\left(-R+\ell, s_{*}\right)$.

We get similar monotonicity results for the translations of $\mathcal{O}$ along the straight line $y=x$.
Remark 5.1. In Theorem 1.8, we can consider the obstacle $\mathcal{O}$ of the form $\mathcal{O}=\bigcup_{j=1}^{k} \bar{B}_{\rho_{j}}\left(z_{j}\right) \subset \Omega$, a finite union of closed balls, such that the centers $z_{j}$ 's lie on the ray $a+\mathbb{R}^{+} \eta$. Now, the rotations of the obstacle $\mathcal{O}$ about the point $a$ are given by

$$
\mathcal{O}_{s}:=\bigcup_{j=1}^{k} \bar{B}_{\rho_{j}}\left(a+R_{s}\left(-a+z_{j}\right)\right) \text { for } s \in[-1,1], \text { and } \mathrm{C}_{\mathcal{O}}:=\left\{s \in[-1,1]: \mathcal{O}_{s} \subset \Omega\right\}
$$

In this case, also, we have the same conclusions as Theorem 1.8.


Figure 5. An eccentric annular domain with a hole.
5.1. An eccentric annular domain with a spherical hole. Let $\frac{2 d+2}{d+2}<p<\infty$. For given $0<r<R$ and $0 \leq \alpha<R-r$, we consider an eccentric annular domain $\Omega=B_{R}(0) \backslash \overline{B_{r}}\left(-\alpha e_{1}\right) \subset \mathbb{R}^{d}$ (see Figure 5). Let $\rho>0$ be such that $\bar{B}_{\rho}(y) \subset \Omega$ for some $y \in \mathbb{R}^{d}$. Now define

$$
\Omega_{\rho}:=\left\{y \in \Omega: \bar{B}_{\rho}(y) \subset \Omega\right\}, \text { and } \lambda_{1}(y):=\lambda_{1}\left(\Omega \backslash \bar{B}_{\rho}(y)\right) \text { for } y \in \Omega_{\rho}
$$

We want to study the behaviour of $\lambda_{1}(\cdot)$ on $\Omega_{\rho}$. It is easy to observe that
(a) $\bar{B}_{\rho}(y)$ is Steiner symmetric with respect to any affine-hyperplane through $x$;
(b) $\bar{B}_{\rho}(y)$ is foliated Schwarz symmetric with respect to $a+\mathbb{R}^{+}(y-a)$ for any $a \in \mathbb{R}^{d}$;
(c) $\Omega$ is foliated Schwarz symmetric with respect to $t e_{1}+\mathbb{R}^{+} e_{1}$ for $t \in[-\alpha, 0]$;
(d) the sets $\left\{x \in \Omega: x_{1}<\underline{r}\right\} \bigcup \sigma_{H_{\underline{r}}}\left(\left\{x \in \Omega: x_{1}<\underline{r}\right\}\right)$ and $\left\{x \in \Omega: x_{1}>\bar{r}\right\} \bigcup \sigma_{H_{\bar{r}}}\left(\left\{x \in \Omega: x_{1}>\bar{r}\right\}\right)$ are convex in the $e_{1}$-direction, where $\bar{r}=\frac{R+r-\alpha}{2}, \underline{r}=-\frac{R+r+\alpha}{2}$.
For $y \in \mathbb{R}^{d}$, we write $y=(s, z) \in \mathbb{R} \times \mathbb{R}^{d-1}$. Now for a given $z \in \mathbb{R}^{d-1}, \beta>0$, we consider the sets

$$
\begin{aligned}
& \mathrm{L}_{z}:=\left\{s \in(-R, R):(s, z) \in \Omega_{\rho}\right\} \\
& \mathrm{S}_{\beta}\left(t e_{1}\right):=\Omega_{\rho} \cap \partial B_{\beta}\left(t e_{1}\right), \quad t \in[-\alpha, 0]
\end{aligned}
$$

Remark 5.2. Let $\left(s, z_{1}\right) \in \Omega_{\rho}$. Then we have the following:
(i) Using the axial symmetry of $\Omega$, we obtain $\lambda_{1}(s, z)=\lambda_{1}\left(s, z_{1}\right)$, for $z \in \mathbb{R}^{d-1}$ such that $(s, z) \in \Omega_{\rho}$ and $|z|=\left|z_{1}\right|$.
(ii) From (a), (d) and Theorem 1.5 with $h=-e_{1}$ (and $h=e_{1}$ ), we get

$$
\lambda_{1}\left(\cdot, z_{1}\right) \text { is strictly increasing on } \mathrm{L}_{z_{1}} \bigcap\left(-\sqrt{(R-\rho)^{2}-\left|z_{1}\right|^{2}}, \underline{r}\right]
$$

and

$$
\lambda_{1}\left(\cdot, z_{1}\right) \text { is strictly decreasing on } \mathrm{L}_{z_{1}} \bigcap\left[\bar{r}, \sqrt{(R-\rho)^{2}-\left|z_{1}\right|^{2}}\right)
$$

(iii) If $\left(s_{1}, z_{1}\right)$ and $\left(s_{2}, z_{2}\right) \in \Omega_{\rho}$ such that $s_{1}<s_{2}$ and $\left(s_{1}-t\right)^{2}+\left|z_{1}\right|^{2}=\left(s_{2}-t\right)^{2}+\left|z_{2}\right|^{2}$ for some $t \in[-\alpha, 0]$, then by Theorem 1.8, we get

$$
\lambda_{1}\left(s_{1}, z_{1}\right)<\lambda_{1}\left(s_{2}, z_{2}\right)
$$

In particular, for $s_{1}, s_{2} \in[-\alpha, 0]$ with $s_{1}<s_{2}$, by taking $t=\frac{s_{1}+s_{2}}{2}$ we obtain

$$
\lambda_{1}(\cdot, z) \text { is strictly increasing on } \mathrm{L}_{z} \cap[-\alpha, 0] \text {. }
$$

Remark 5.3. In general, for any $h=\left(h_{1}, h^{\prime}\right) \in \mathbb{S}^{d-1}$, let

$$
r^{h}:=\frac{\sqrt{R^{2}-\alpha^{2}\left|h^{\prime}\right|^{2}}+r-\alpha h_{1}}{2}
$$

Then, the set $\left\{x \in \Omega: x \cdot h>r^{h}\right\} \bigcup \sigma_{H_{r_{h}}}\left(\left\{x \in \Omega: x \cdot h>r^{h}\right\}\right)$ is convex in the $h$-direction. For $y \in \Omega_{\rho}$, define $\mathrm{L}_{y}=\left\{s \in[0, R): y+s h \in \Omega_{\rho}\right\}$. Now, from (a) and Theorem 1.5 we get

$$
\lambda_{1}(y+s h) \text { is strictly decreasing for } s \in \mathrm{~L}_{y} \cap\left[r^{h}, R\right)
$$



Figure 6. Monotonicity along certain paths in $\Omega$.
Remark 5.4. Let $C:(-R, R) \longrightarrow \Omega$ be a continuous path in $\Omega$ such that (see Figure 6):

- on $(-R,-\alpha), C$ is a circular arc centered at $t_{1} e_{1}$ with $t_{1} \in[-\alpha, 0]$;
- on $[-\alpha, 0], C$ is either a circular arc centered at $t_{2} e_{1}$ with $t_{2} \in[-\alpha, 0]$ or a line segment parallel to the $e_{1}$-axis;
- on $(0, R), C$ is a circular arc centered at $t_{3} e_{1}$ with $t_{3} \in[-\alpha, 0]$.

Now, from Remark 5.2, we see that $\lambda_{1}(\cdot)$ is strictly increasing along the path $C$; i.e., for any $\left(s_{1}, z_{1}\right),\left(s_{2}, z_{2}\right) \in$ $C$ with $s_{1}<s_{2}$ we have $\lambda_{1}\left(s_{1}, z_{1}\right)<\lambda_{1}\left(s_{2}, z_{2}\right)$.

Remark 5.5 (Optimal placement of the obstacle). For $y \in \Omega_{\rho}$, we have $y,|y| e_{1} \in \mathrm{~S}_{|y|}(0)$. If $y_{1}<|y|$, then using Theorem 1.8 we obtain $\lambda_{1}(y)<\lambda\left(|y| e_{1}\right)$. Thus

$$
\sup \left\{\lambda_{1}(y): y \in \Omega_{\rho}\right\}=\sup \left\{\lambda_{1}(s, 0): s \in \mathrm{~L}_{0} \cap[0, \bar{r}]\right\}
$$

If $0<\rho<\alpha-r$, then $0, \bar{r} \in \mathrm{~L}_{0}$, and hence by (iii) of Remark 5.2 we get

$$
\sup \left\{\lambda_{1}(y): y \in \Omega_{\rho}\right\}=\max \left\{\lambda_{1}(s, 0): s \in[0, \bar{r}]\right\}
$$

On the other hand, if $\alpha \leq r$ or $\rho>\alpha-r$, then $0 \notin \mathrm{~L}_{0}$. Thus, the above arguments fail to conclude that the supremum is attained in $\Omega_{\rho}$. However, from a Mathematica 12 plot of $\lambda_{1}(\cdot)$ on $[0, R) \cap L_{0}$ (for the various values of $\alpha, r$, and $\rho$ ), we observed that the maximum is attained at a unique point in $(0, \bar{r}) \cap L_{0}$. Giving an analytic explanation of this behaviour of $\lambda_{1}(\cdot, 0)$ in $[0, \bar{r}] \cap \mathrm{L}_{0}$ seems to be an interesting problem to explore.
5.2. The symmetries of the first eigenfunctions: In this subsection, we take $p \in(1, \infty)$ and $\Omega_{\text {out }} \backslash \overline{\Omega_{\text {in }}} \subset$ $\mathbb{R}^{d}$ is a domain as given in $\left(\mathbf{A}_{\mathbf{0}}\right)$. We establish that the first eigenfunctions of (1.3) and (1.6) inherit some of the symmetries of the underlying domains.
Remark 5.6. Let $H \in \mathcal{H}_{\text {ad }}$ be such that $P_{H}\left(\Omega_{\text {out }} \backslash \overline{\Omega_{\text {in }}}\right)=\Omega_{\text {out }} \backslash \overline{\Omega_{\text {in }}}$.
(i) Let $u$ be an eigenfunction corresponding to the first eigenvalue $\nu_{1}\left(\Omega_{\text {out }} \backslash \overline{\Omega_{\mathrm{in}}}\right)$ of (1.3). Assume that $\sigma_{H}\left(\Omega_{\mathrm{in}}\right)=\Omega_{\mathrm{in}}$. If $u$ is positive, then from Proposition 2.19 and (2.5), we see that $P_{H}(u)$ is also an eigenfunction corresponding to $\nu_{1}\left(\Omega_{\text {out }} \backslash \overline{\Omega_{\mathrm{in}}}\right)$. Since the norms of $u$ and $P_{H}(u)$ are same, by the simplicity of $\nu_{1}$, we get $P_{H}(u)=u$. If $u$ is negative then, we get $P^{H}(u)=-P_{H}(-u)=u$.
(ii) Similarly, if $\sigma_{H}\left(\Omega_{\mathrm{out}}\right)=\Omega_{\mathrm{out}}$ and $v$ is a positive eigenfunction corresponding to the first eigenvalue $\tau_{1}\left(\Omega_{\text {out }} \backslash \overline{\Omega_{\text {in }}}\right)$ of (1.5), then $P_{H}(v)=v$ and $P^{H}(-v)=-v$.
Definition 5.7. Let $\Omega$ be foliated Schwarz symmetric with respect to the ray $a+\mathbb{R}^{+} \eta$. Then a function $u: \Omega \longrightarrow \mathbb{R}$ is said to be foliated Schwarz symmetric with respect to the same ray, if $P_{H}(u)=u$ for every $H \in \mathcal{H}_{a, \eta}$ (see [13, Lemma 6.3] and [30, Section 3]).

Remark 5.8. Let $\Omega_{\text {out }} \backslash \overline{\Omega_{\text {in }}}$ be foliated Schwarz symmetric with respect to the ray $a+\mathbb{R}^{+} \eta$ for some $a \in \mathbb{R}^{d}$ and $\eta \in \mathbb{S}^{d-1}$. Then, we have $P_{H}\left(\Omega_{\text {out }} \backslash \overline{\Omega_{\mathrm{in}}}\right)=\Omega_{\text {out }} \backslash \overline{\Omega_{\mathrm{in}}}$ for every $H \in \mathcal{H}_{a, \eta}$.
(i) Assume that $\Omega_{\mathrm{in}}=B_{r}(a)$ for some $r \geq 0$. Then $\sigma_{H}\left(\Omega_{\mathrm{in}}\right)=\Omega_{\mathrm{in}}$ for every $H \in \mathcal{H}_{a, \eta}$. Hence, for any positive eigenfunction $u$ corresponding to $\nu_{1}\left(\Omega_{\mathrm{out}} \backslash \overline{\Omega_{\mathrm{in}}}\right)$, from Remark 5.6 , we obtain $P_{H}(u)=u$ for every $H \in \mathcal{H}_{a, \eta}$. Thus, $u$ is foliated Schwarz symmetric with respect to the ray $a+\mathbb{R}^{+} \eta$.
(ii) Similarly, Assume that $\Omega_{\mathrm{out}}=B_{R}(a)$ for some $R>0$. Then, any positive eigenfunction $u$ corresponding to the first eigenvalue $\tau_{1}\left(\Omega_{\text {out }} \backslash \overline{\Omega_{\mathrm{in}}}\right)$ of (1.5) is foliated Schwarz symmetric with respect to $a+\mathbb{R}^{+} \eta$.
Remark 5.9. Let $\Omega=B_{R}(0) \backslash \bar{B}_{r}\left(-\alpha e_{1}\right) \subset \mathbb{R}^{d}$ be the eccentric annular domain as given in Subsection 5.1. Then $\Omega$ is foliated Schwarz symmetric with respect to $t e_{1}+\mathbb{R}^{+} e_{1}$ for $t \in[-\alpha, 0]$.
(i) If, we take $\Omega_{\text {out }}=\Omega$ and $\Omega_{\mathrm{in}}=\emptyset$, then any positive eigenfunction corresponding to the first Dirichlet eigenvalue $\lambda_{1}(\Omega)$ is foliated Schwarz symmetric with respect to $t e_{1}+\mathbb{R}^{+} e_{1}$ for every $t \in[-\alpha, 0]$.
(ii) If, we take $\Omega_{\text {out }}=B_{R}(0)$ and $\Omega_{\mathrm{in}}=B_{r}\left(-\alpha e_{1}\right)$, then

- any positive eigenfunction corresponding to the first eigenvalue $\nu_{1}(\Omega)$ of (1.3) is foliated Schwarz symmetric with respect to the ray $-\alpha e_{1}+\mathbb{R}^{+} e_{1}$;
- any positive eigenfunction corresponding to the first eigenvalue $\tau_{1}(\Omega)$ of (1.5) is foliated Schwarz symmetric with respect to the ray $\mathbb{R}^{+} e_{1}$.


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T. V. Anoop, Department of Mathematics, Indian Institute of Technology Madras, Chennai 600 036, India Email address: anoop@iitm.ac.in
K. Ashok Kumar, Department of Mathematics, Indian Institute of Technology Madras, Chennai 600 036, India

Email address: s.r.asoku@gmail.com


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