

ON THE REGULARITY OF CRITICAL AND MINIMAL SETS OF A FREE INTERFACE PROBLEM

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ABSTRACT. We study a free interface problem of finding the optimal energy configuration for mixtures of two conducting materials with an additional perimeter penalization of the interface. We employ the regularity theory of linear elliptic equations to study the possible opening angles of Taylor cones and to give a different proof of a partial regularity result by Fan Hua Lin [13].

1. INTRODUCTION

In this paper we consider the functional

$$(1.1) \quad \mathcal{F}(E, v) = \gamma P(E, \Omega) + \int_{\Omega} \sigma_E(x) |Dv|^2 dx,$$

where $\gamma > 0$, $\Omega \subset \mathbb{R}^n$ is an open set, $v \in W^{1,2}(\Omega)$ and $P(E, \Omega)$ stands for the perimeter of E in Ω . Moreover, $\sigma_E(x) = \beta \chi_E(x) + \alpha \chi_{\Omega \setminus E}(x)$, where $0 < \alpha < \beta < \infty$ are given constants.

Given a function $u_0 \in W^{1,2}(\Omega)$ and a measurable $E \subset \Omega$, we denote by u_E , or simply by u if no confusion arises, the corresponding elastic equilibrium, i.e., the minimizer in $W^{1,2}(\Omega)$ of the functional

$$\int_{\Omega} \sigma_E(x) |Dv|^2 dx$$

under the boundary condition $v = u_0$ on $\partial\Omega$. It follows that the function u solves the linear equation

$$(1.2) \quad \int_{\Omega} \langle \sigma_E Du, D\varphi \rangle dx = 0 \quad \text{for every } \varphi \in W_0^{1,2}(\Omega).$$

If we denote by u_{β} and u_{α} the restriction of u on E and $\Omega \setminus E$, respectively, they are harmonic in their domains. Moreover, equation (1.2) implies the transmission condition

$$(1.3) \quad \alpha \partial_{\nu} u_{\alpha}(x) = \beta \partial_{\nu} u_{\beta}(x) \quad \text{for all } x \in \partial E \cap \Omega,$$

where ∂_{ν} denotes the derivative of u in the direction of the exterior normal to ∂E . Note that if (E, u) is a smooth critical point of the functional (1.1) the following Euler-Lagrange equation holds

$$(1.4) \quad \gamma H_{\partial E} + \beta |Du_{\beta}|^2 - \alpha |Du_{\alpha}|^2 = \lambda \quad \text{on } \partial E \cap \Omega,$$

where $H_{\partial E}$ stands for the mean curvature of ∂E and λ is either zero or a Lagrange multiplier (in case of a volume constraint).

In the physical literature critical points of the functional (1.1), i.e, solutions of equations (1.2) and (1.4) are studied to model the shape of liquid drops exposed to an electric or a magnetic field. In the model the set E represents a liquid drop with dielectric permittivity β , surrounded by a fluid with smaller permittivity α , and u stands for the electrostatic potential induced by an applied electric field. At the interface, which is assumed to be in static equilibrium, the normal component of the electric displacement field $\sigma_E Du$ has to be continuous. This implies that u has to satisfy (1.3) or equivalently that it has to

Date: May 26, 2014.

2010 Mathematics Subject Classification. 49Q10, 49N60, 74G40.

be a solution of equation (1.2). On the other hand, on the interface the electric stress and the surface tension has to be balanced, which leads to (1.4).

The occurrence of conical tips at an interface exposed to an electric field has been observed by several authors (see e.g. [19]). Moreover, theoretical investigations ([12], [16], [17], [18]) suggest that conical critical points, the so called *Taylor cones*, may only occur if the ratio $\frac{\beta}{\alpha}$ is sufficiently large and if the opening angle is neither too small nor too close to $\pi/2$, i.e., it belongs to a certain range which is independent of the penalization factor γ .

In order to state our first result we denote by E_θ the right circular cone with opening angle $\theta \in (0, \pi/2)$ and vertex at the origin, i.e.,

$$E_\theta = \{(x', x_n) \in \mathbb{R}^{n-1} \times \mathbb{R} : x_n > \frac{1}{\tan \theta} |x'|\}.$$

Theorem 1.1. *Let $n \geq 3$. There exist two positive constants $\delta_0 = \delta_0(n, \frac{\beta}{\alpha}) > 0$ and $\lambda_0 = \lambda_0(n) > 1$ such that, if E_θ is a right circular cone satisfying (1.4) then $\frac{\beta}{\alpha} \geq \lambda_0$ and*

$$\delta_0 \leq \theta \leq \frac{\pi}{2} - \delta_0.$$

As far as we know, this result is the first rigorous proof of the fact that Taylor cones may occur only for certain angles, and provided that the ratio $\frac{\beta}{\alpha}$ is sufficiently large. We remark that we are able to give explicit estimates of the constants δ_0 and λ_0 . In particular, δ_0 and λ_0 are independent of the penalization factor γ , which is in accordance with the observations and theoretical results reported in the physical literature.

The starting point in the proof of Theorem 1.1 is a rather simple decay estimate for the gradient of a minimizer of the Dirichlet energy (see Proposition 2.4). Roughly speaking, we prove that if u_E is a solution of (1.2) and x_0 is a point in Ω , where either the density of E is close to 0 or 1, or the set E is asymptotically close to a hyperplane, then for sufficiently small ρ we have

$$(1.5) \quad \int_{B_\rho(x_0)} |Du_E|^2 dx \leq C\rho^{n-\delta}$$

for any $\delta > 0$. As a consequence of this estimate one has that if the opening angle of the Taylor cone is not in the above range then the Dirichlet energy around the vertex decays faster than the perimeter thus leading to a contradiction to the criticality condition (1.4).

In the mathematical literature people have also considered both the problem (P) of minimizing (1.1) under the boundary condition $u = u_0$ on $\partial\Omega$ and the constrained problem

$$(P_c) \quad \min\{\mathcal{F}(E, v) : v = u_0 \text{ on } \partial\Omega, |E| = d\}$$

for some given $d < |\Omega|$. The partial regularity of minimizers of the unconstrained problem (P) was proved by Fan-Hua Lin in [13] (see also [1], [14]). In the special case $n = 2$ the result of Lin has been improved by Larsen in [10], [11]. However, the full regularity of the free interface ∂E in two dimensions still remains open.

In the second part of the paper we revisit the proof of the partial regularity of minimizers.

Theorem 1.2. *If (E, u) is a minimizer of either problem (P) or problem (P_c) , then*

- (a) *there exists a relatively open set $\Gamma \subset \partial E$ such that Γ is a $C^{1,\sigma}$ hypersurface for all $0 < \sigma < 1/2$,*
- (b) *there exists $\varepsilon > 0$, depending only on $\frac{\beta}{\alpha}$ and n , such that $\mathcal{H}^{n-1-\varepsilon}((\partial E \setminus \Gamma) \cap \Omega) = 0$.*

The above statement slightly generalizes the regularity result proved in [13], where only the unconstrained problem (P) was considered. The value of ε is greater than or equal to $p - 1$ where $2p$ is the

higher integrability exponent of Du_E which is obtained by a standard application of Gehring's lemma (see Lemma 2.2). In particular, it is independent of the penalization factor γ . We note also that, once the $C^{1,\sigma}$ regularity of $\partial^*E \cap \Omega$ is obtained, then using [13, Lemma 5.3] and a standard bootstrap argument one obtains that $\partial^*E \cap \Omega$ is in fact C^∞ .

As in [13] and in the proof of the regularity of minimizers of the Mumford-Shah functional [3] our proof of Theorem 1.2 is based on an interplay between the perimeter and the Dirichlet integral. Differently from [13] we do not use the heavy machinery of currents and do not derive the monotonicity formula. Instead, our starting point is the same decay estimate (1.5) for the Dirichlet energy used for the study of Taylor cones. This estimate implies that if in a ball $B_r(x_0)$ the perimeter of E is sufficiently small then the total energy in a smaller ball $B_{\tau r}(x_0)$

$$P(E, B_{\tau r}(x_0)) + \int_{B_{\tau r}(x_0)} |Du_E|^2$$

is much smaller than the total energy in $B_r(x_0)$ (Lemma 4.2). In turn this fact leads to a density lower bound for the perimeter.

Another consequence of the estimate (1.5) is that whenever the excess

$$\mathcal{E}(x_0, r) = \inf_{\nu \in \mathcal{S}^{n-1}} \frac{1}{r^{n-1}} \int_{\partial E \cap B_r(x_0)} |\nu_E(x) - \nu|^2 d\mathcal{H}^{n-1} \rightarrow 0 \quad \text{as } r \rightarrow 0,$$

the Dirichlet integral in $B_r(x_0)$ decays as in (1.5). As in the Mumford-Shah case this is one of the two key estimates needed for the regularity proof (see Step 1 of the proof of the Theorem 1.2 at the end of Section 3). Finally the excess decay is proven with a more or less standard argument similar to the one used for the Λ -minimizers of the perimeter.

2. REGULARITY OF ELASTIC MINIMA

In this section we study the regularity of the elastic minimum associated to a set E , i.e., solution of (1.2). In the main result of the section, Proposition 2.4, we prove that, if the density of E is close to 0 or 1 or the set E is asymptotically close to a hyperplane, then the elastic energy $\int_{B_\rho(x_0)} |Du|^2 dx$ decays faster than ρ^{n-1} . We prove Proposition 2.4 with a direct argument and therefore we are able to provide explicit bounds for the relevant constants.

We begin by deriving the Caccioppoli's inequality for solutions of (1.2). Even though the argument is standard, we give the proof in order to keep track of the constants. We denote the cube, centred at x_0 and with side length $2r$, by $Q_r(x_0)$. In the case $x_0 = 0$ we simply write Q_r . We recall the Sobolev-Poincaré inequality, i.e., for every function $u \in W^{1,p}(Q_r)$, $1 \leq p < n$, it holds

$$(2.1) \quad \|u - u_r\|_{L^{p^*}(Q_r)} \leq c(n, p) \|Du\|_{L^p(Q_r)}$$

where $u_r = \int_{Q_r} u dx$ and $p^* = \frac{pn}{n-p}$.

Lemma 2.1. *Let $u \in W^{1,2}(\Omega)$ be a solution of (1.2). Then for every cube $Q_{2r}(x_0) \subset\subset \Omega$ it holds*

$$\int_{Q_r(x_0)} |Du|^2 dx \leq C \left(\int_{Q_{2r}(x_0)} |Du|^{2m} dx \right)^{\frac{1}{m}},$$

where $m = \frac{n}{n+2}$, $C = C_{S,n}^2 2^{n+8} \frac{\beta}{\alpha}$, and $C_{S,n}$ is the constant in the Sobolev-Poincaré inequality (2.1) with $p = \frac{2n}{n+2}$.

Proof. Without loss of generality we may assume that $x_0 = 0$. Let $\zeta \in C_0^\infty(Q_{2r})$ be a cut-off function such that $\zeta \equiv 1$ in Q_r and $|D\zeta| \leq \frac{2}{r}$. We choose a test function $\varphi = (u - u_{2r})\zeta^2$ in (1.2), where $u_{2r} = \int_{Q_{2r}} u \, dx$ and apply Young's inequality to obtain

$$\int_{Q_r} |Du|^2 \, dx \leq \int_{Q_{2r}} |Du|^2 \zeta^2 \, dx \leq \frac{4\beta}{\alpha} \int_{Q_{2r}} |u - u_{2r}|^2 |D\zeta|^2 \, dx \leq \frac{16\beta}{r^2\alpha} \int_{Q_{2r}} |u - u_{2r}|^2 \, dx.$$

We use the Sobolev-Poincaré inequality (2.1) to deduce

$$\int_{Q_{2r}} |u - u_{2r}|^2 \, dx \leq C_{S,n}^2 \left(\int_{Q_{2r}} |Du|^{2m} \, dx \right)^{\frac{1}{m}}.$$

The result then follows from the two inequalities above. \square

We apply Gehring's Lemma to obtain higher integrability for the gradient of u .

Lemma 2.2. *Let $u \in W^{1,2}(\Omega)$ solve (1.2). There exists $p > 1$ such that for any ball $B_{2r}(x_0) \subset\subset \Omega$ it holds*

$$\int_{B_r(x_0)} |Du|^{2p} \, dx \leq C \left(\int_{B_{2r}(x_0)} |Du|^2 \, dx \right)^p.$$

The constants can be estimated explicitly as

$$p = \frac{2C_1 - m}{2C_1 - 1} \quad \text{for } C_1 = C_{S,n}^2 2^{10} \cdot 80^n \frac{\beta}{\alpha} \quad \text{and } C = 2^{2n+1} 5^{np} n^{np/2} \omega_n^{p-1},$$

where ω_n is the volume of the unit ball and $m = \frac{n}{n+2}$.

The above result is well known but it is usually stated without estimates of the constants. In the Appendix we will go through the proof of Lemma 2.2 from [8] and evaluate every constant explicitly.

In the next lemma we prove a monotonicity formula for the elastic minimum in the case when E is a half-space.

Lemma 2.3. *Let $E = \{x \in \mathbb{R}^n \mid \langle x - \bar{x}, e \rangle < 0\} \cap \Omega$ for some unit vector e and a point \bar{x} , and suppose u is a solution of (1.2). Let $x_0 \in \partial E \cap \Omega$ and $r > 0$ be such that $B_r(x_0) \subset \Omega$. Then*

$$\rho \mapsto \int_{B_\rho(x_0)} \sigma_E(x) |Du|^2 \, dx$$

is increasing in $(0, r)$.

Proof. Without loss of generality we may assume that $E = \{x \in \mathbb{R}^n \mid x_n < 0\} \cap \Omega$ and $x_0 = 0$. Let us fix a radius r such that $B_r \subset \Omega$. From standard elliptic regularity theory we know that u is smooth in the upper and in the lower part of the ball B_r with respect to the hyperplane $\partial E = \{x_n = 0\}$. To be more precise, let us denote $\bar{B}_r^+ = \bar{B}_r \setminus \bar{E}$ and $\bar{B}_r^- = \bar{B}_r \cap \bar{E}$. Then $u_\alpha \in C^\infty(\bar{B}_r^+)$ and $u_\beta \in C^\infty(\bar{B}_r^-)$ and they are harmonic in the interior of \bar{B}_r^+ and \bar{B}_r^- , where u_α and u_β are the restrictions of u on $\Omega \setminus E$ and E .

The goal is to show that the function $\varphi : (0, r) \rightarrow \mathbb{R}$

$$\varphi(\rho) := \int_{\partial B_\rho} \sigma_E(x) |Du(x)|^2 \, d\mathcal{H}^{n-1}(x) = \int_{\partial B_1} \sigma_E(\rho y) |Du(\rho y)|^2 \, d\mathcal{H}^{n-1}(y)$$

is increasing. Notice that $\sigma_E(\rho y) = \sigma_E(y)$ since E is a half-space. Denote $v = |Du|^2$, $v_\alpha = |Du_\alpha|^2$ and $v_\beta = |Du_\beta|^2$. Since v_α and v_β are subharmonic in the interior of B_ρ^+ and B_ρ^- we deduce by the divergence

theorem that

$$\begin{aligned}
\varphi'(\rho) &= \int_{\partial B_\rho} \sigma_E(x) \langle Dv(x), \frac{x}{\rho} \rangle d\mathcal{H}^{n-1}(x) \\
&= \frac{1}{\mathcal{H}^{n-1}(\partial B_\rho)} \left(\alpha \int_{\partial B_\rho^+} \partial_\nu v_\alpha d\mathcal{H}^{n-1} + \beta \int_{\partial B_\rho^-} \partial_\nu v_\beta d\mathcal{H}^{n-1} + \int_{\partial E \cap B_\rho} \alpha \partial_{x_n} v_\alpha - \beta \partial_{x_n} v_\beta d\mathcal{H}^{n-1} \right) \\
&= \frac{1}{\mathcal{H}^{n-1}(\partial B_\rho)} \left(\alpha \int_{B_\rho^+} \Delta v_\alpha d\mathcal{H}^{n-1} + \beta \int_{B_\rho^-} \Delta v_\beta d\mathcal{H}^{n-1} + \int_{\partial E \cap B_\rho} \alpha \partial_{x_n} v_\alpha - \beta \partial_{x_n} v_\beta d\mathcal{H}^{n-1} \right) \\
&\geq \frac{1}{\mathcal{H}^{n-1}(\partial B_\rho)} \int_{\partial E \cap B_\rho} \alpha \partial_{x_n} v_\alpha - \beta \partial_{x_n} v_\beta d\mathcal{H}^{n-1}.
\end{aligned}$$

We will show that $\beta \partial_{x_n} v_\beta = \alpha \partial_{x_n} v_\alpha$ on ∂E from which the claim follows.

Since $u_\alpha = u_\beta$ on ∂E we have

$$(2.2) \quad \partial_{x_i} u_\alpha = \partial_{x_i} u_\beta \quad \text{and} \quad \partial_{x_i x_i} u_\alpha = \partial_{x_i x_i} u_\beta \quad \text{for } i = 1, 2, \dots, n-1.$$

The transmission condition (1.3) reads as

$$(2.3) \quad \alpha \partial_{x_n} u_\alpha = \beta \partial_{x_n} u_\beta \quad \text{on } \partial E.$$

Differentiating (2.3) with respect to x_i , for $i = 1, 2, \dots, n-1$, yields

$$\alpha \partial_{x_i x_n} u_\alpha = \beta \partial_{x_i x_n} u_\beta \quad \text{on } \partial E.$$

On the other hand, since u_α and u_β are harmonic, we have by (2.2) that

$$\partial_{x_n x_n} u_\alpha = - \sum_{i=1}^{n-1} \partial_{x_i x_i} u_\alpha = - \sum_{i=1}^{n-1} \partial_{x_i x_i} u_\beta = \partial_{x_n x_n} u_\beta \quad \text{on } \partial E.$$

Therefore on ∂E it holds

$$\alpha \partial_{x_n} v_\alpha = 2 \sum_{i=1}^n \alpha \partial_{x_i} u_\alpha \partial_{x_i x_n} u_\alpha = 2 \sum_{i=1}^n \beta \partial_{x_i} u_\beta \partial_{x_i x_n} u_\beta = \beta \partial_{x_n} v_\beta$$

which implies $\varphi'(\rho) \geq 0$. □

The main result of this section is the following decay estimate for elastic minimum.

Proposition 2.4. *Let $u \in W^{1,2}(\Omega)$ be a solution of (1.2). For all $\tau \in (0, 1)$ there exists $\varepsilon_0 = \varepsilon_0(\tau) > 0$ such that if $B_r(x_0) \subset\subset \Omega$ and one of the following conditions hold*

- (i) $\frac{|E \cap B_r(x_0)|}{|B_r|} < \varepsilon_0$,
- (ii) $\frac{|B_r(x_0) \setminus E|}{|B_r|} < \varepsilon_0$,
- (iii) *there exists a half-space H such that $\frac{|(E \Delta H) \cap B_r(x_0)|}{|B_r|} < \varepsilon_0$,*

then

$$\int_{B_{\tau r}(x_0)} |Du|^2 dx \leq C_0 \tau^n \int_{B_r(x_0)} |Du|^2 dx$$

for some constant C_0 depending only on $\frac{\beta}{\alpha}$ and n .

Proof. Without loss of generality we may assume that $\tau < 1/2$. We first treat the cases (i) and (ii). We fix a ball $B_r(x_0) \subset\subset \Omega$ and assume without loss of generality that $x_0 = 0$. Choose v to be the

harmonic function in $B_{r/2}$ with the boundary condition $v = u$ on $\partial B_{r/2}$. We choose the test function $\varphi = v - u \in W_0^{1,2}(B_{r/2})$ in the equations

$$\int_{B_{r/2}} Dv \cdot D\varphi \, dx = 0$$

and

$$(2.4) \quad \alpha \int_{B_{r/2} \setminus E} Du \cdot D\varphi \, dx + \beta \int_{B_{r/2} \cap E} Du \cdot D\varphi \, dx = 0.$$

We write the latter equation as

$$\int_{B_{r/2}} Du \cdot (Dv - Du) \, dx = -\frac{\beta - \alpha}{\alpha} \int_{B_{r/2} \cap E} Du \cdot (Dv - Du) \, dx.$$

We subtract to this $\int_{B_{r/2}} Dv \cdot (Dv - Du) \, dx = 0$ and use Hölder's inequality to deduce

$$\int_{B_{r/2}} |Dv - Du|^2 \, dx \leq \frac{(\beta - \alpha)^2}{\alpha^2} \int_{B_{r/2} \cap E} |Du|^2 \, dx.$$

By the higher integrability stated in Lemma 2.2 we have

$$(2.5) \quad \begin{aligned} \int_{B_{\tau r}} |Dv - Du|^2 \, dx &\leq \frac{(\beta - \alpha)^2}{\alpha^2} |E \cap B_r|^{1-1/p} |B_r|^{1/p} \left(\int_{B_r} |Du|^{2p} \right)^{1/p} \\ &\leq C^{1/p} \frac{(\beta - \alpha)^2}{\alpha^2} \left(\frac{|E \cap B_r|}{|B_r|} \right)^{1-1/p} \int_{B_r} |Du|^2 \, dx \end{aligned}$$

where C and $p > 1$ are from Lemma 2.2. Similarly we deduce

$$(2.6) \quad \int_{B_{\tau r}} |Dv - Du|^2 \, dx \leq C^{1/p} \frac{(\beta - \alpha)^2}{\beta^2} \left(\frac{|E \setminus B_r|}{|B_r|} \right)^{1-1/p} \int_{B_r} |Du|^2 \, dx.$$

On the other hand, since v is harmonic, we have

$$\int_{B_{\tau r}} |Dv|^2 \, dx \leq 2^n \tau^n \int_{B_{r/2}} |Dv|^2 \, dx \leq 2^n \tau^n \int_{B_{r/2}} |Du|^2 \, dx.$$

Hence, we may estimate

$$(2.7) \quad \begin{aligned} \int_{B_{\tau r}} |Du|^2 \, dx &\leq 2 \int_{B_{\tau r}} |Dv - Du|^2 \, dx + 2 \int_{B_{\tau r}} |Dv|^2 \, dx \\ &\leq 2 \int_{B_{\tau r}} |Dv - Du|^2 \, dx + 2^{n+1} \tau^n \int_{B_r} |Du|^2 \, dx. \end{aligned}$$

If ε_0 is such that $\varepsilon_0^{1-\frac{1}{p}} = \tau^n$ then (2.5), (2.6) and (2.7) yield

$$\int_{B_{\tau r}} |Du|^2 \, dx \leq C \tau^n \int_{B_r} |Du|^2 \, dx$$

for a constant C depending only on β/α and n .

We are left with the case (iii). Let H be the half-space from the assumption. We choose v which minimizes the energy $\int_{B_{r/2}} \sigma_H(x) |Dv|^2 \, dx$ with the boundary condition $v = u$ on $\partial B_{r/2}$. Hence

$$(2.8) \quad \beta \int_{B_{r/2} \cap H} Dv \cdot D\varphi \, dx + \alpha \int_{B_{r/2} \setminus H} Dv \cdot D\varphi \, dx = 0$$

for every $\varphi \in W_0^{1,2}(B_{r/2})$. Lemma 2.3 yields

$$\int_{B_{\tau r}} |Dv|^2 \, dx \leq \frac{\beta}{\alpha} 2^n \tau^n \int_{B_{r/2}} |Dv|^2 \, dx.$$

Moreover, from the minimality of v it follows

$$\int_{B_{r/2}} |Dv|^2 dx \leq \frac{\beta}{\alpha} \int_{B_{r/2}} |Du|^2 dx.$$

Hence, we have

$$(2.9) \quad \int_{B_{\tau r}} |Du|^2 dx \leq 2 \int_{B_{\tau r}} |Dv - Du|^2 dx + 2^{n+1} \left(\frac{\beta}{\alpha}\right)^2 \tau^n \int_{B_r} |Du|^2 dx.$$

Let us now rewrite the equation (2.4) satisfied by u as

$$\begin{aligned} \beta \int_{B_{r/2} \cap H} Du \cdot D\varphi dx + \alpha \int_{B_{r/2} \setminus H} Du \cdot D\varphi dx \\ = (\beta - \alpha) \int_{B_{r/2} \cap (H \setminus E)} Du \cdot D\varphi dx - (\beta - \alpha) \int_{B_{r/2} \cap (E \setminus H)} Du \cdot D\varphi dx. \end{aligned}$$

Then, subtracting (2.8) from this equation and choosing $\varphi = u - v$ we get at once

$$\int_{B_{\tau r}} |Du - Dv|^2 dx \leq \left(\frac{\beta}{\alpha} - 1\right)^2 \int_{B_{r/2} \cap (E \Delta H)} |Du|^2 dx$$

and from Lemma 2.2 we deduce

$$(2.10) \quad \int_{B_{\tau r}} |Du - Dv|^2 dx \leq C^{1/p} \left(\frac{\beta}{\alpha} - 1\right)^2 \left(\frac{|(E \Delta H) \cap B_r|}{|B_r|}\right)^{1-1/p} \int_{B_r} |Du|^2 dx.$$

The conclusion then follows as in the previous cases. \square

3. TAYLOR CONES

In this section we study critical configurations (E, u) , i.e., they satisfy (1.2) and (1.4). In particular, we are interested in those which are circular cones satisfying (1.4) outside the vertex. It was shown in [12] and [16] that there exist circular cones in \mathbb{R}^3

$$E_{\theta_0} = \left\{ x \in \mathbb{R}^3 \mid x_3 > \frac{1}{\tan \theta_0} \sqrt{x_1^2 + x_2^2} \right\},$$

for $\theta_0 \in (0, \pi/2)$, which are critical. Indeed, one may find an associated elastic minimum given in spherical coordinates by

$$u(\rho, \theta) = \sqrt{\rho} \cdot f(\theta),$$

where $\rho = \sqrt{x_1^2 + x_2^2 + x_3^2}$ and θ is the angle formed by the vector $x \in \mathbb{R}^3$ with the positive x_3 semi-axis. Denote by $P_{\frac{1}{2}}$ the Legendre function of the first kind of order 1/2 which solves the equation

$$P''(t)(1-t^2) - 2tP'(t) + \frac{3}{4}P(t) = 0, \quad t \in (-1, 1).$$

Then the function f is given by

$$f(\theta) = \begin{cases} P_{\frac{1}{2}}(-\cos \theta_0)P_{\frac{1}{2}}(\cos \theta), & \theta \in [0, \theta_0] \\ P_{\frac{1}{2}}(\cos \theta_0)P_{\frac{1}{2}}(-\cos \theta), & \theta \in [\theta_0, \pi]. \end{cases}$$

The transmission condition (1.3) then reads as

$$\beta P_{\frac{1}{2}}(-\cos \theta_0)P'_{\frac{1}{2}}(\cos \theta_0) + \alpha P_{\frac{1}{2}}(\cos \theta_0)P'_{\frac{1}{2}}(-\cos \theta_0) = 0.$$

It can be proved that there exists a critical threshold $\lambda_1 \approx 17.59$ such that this equation has no solutions if $\frac{\beta}{\alpha} < \lambda_1$ and it has two solutions in $(0, \pi/2)$ if $\frac{\beta}{\alpha} > \lambda_1$.

In [17] estimates of the angles corresponding to critical cones are given by a different approach. Although the known results give sharp estimates for the angles which allow existence of critical spherical

cones, they do not give any rigorous answer whether there exists a range of angles where no critical spherical cones appear.

We apply the regularity from Section 2 to prove that only cones with certain angles are possible. This estimate is independent of γ , which reflects the fact that the perimeter has only regularizing effect. This result rigorously answers to the question connected to Taylor Cones, of why cones of certain angles do not appear. The result also generalizes to convex cones E whose base is uniformly convex and C^2 -regular. Since the result is purely local, with no loss of generality we set $\Omega = \mathbb{R}^n$.

We begin by revisiting the decay estimate proved in Proposition 2.4 with an explicit choice of the constants.

Proposition 3.1. *Let $u \in W_{loc}^{1,2}(\mathbb{R}^n)$ be a solution of (1.2). There exist $\delta_1, \sigma > 0$ and $\vartheta \in (0, 1)$, depending only on the dimension n and the ratio $\frac{\beta}{\alpha}$, such that if one of the following conditions hold*

- (i) $|E \cap B_1| < \delta_1 |B_1|$,
- (ii) $|B_1 \setminus E| < \delta_1 |B_1|$,
- (iii) *there exists a half-space H such that $|(E \Delta H) \cap B_1| < \delta_1 |B_1|$,*

then we have

$$\int_{B_\vartheta} |Du|^2 dx \leq \vartheta^{n-1+\sigma} \int_{B_1} |Du|^2 dx.$$

Proof. We first deal the case (i). We recall that from (2.5) and (2.7) it follows that if $\tau \in (0, \frac{1}{2})$ then we have

$$\int_{B_\tau} |Du|^2 dx \leq 2 \left(C^{1/p} \frac{(\beta - \alpha)^2}{\alpha^2} \delta_1^{1-1/p} + 2^n \tau^n \right) \int_{B_1} |Du|^2 dx$$

where C and p are the constants from Lemma 2.2. Note that this inequality is trivially satisfied if $\tau \in (\frac{1}{2}, 1)$. Let us denote by χ the largest number such that the equation

$$(3.1) \quad 2\chi + 2^{n+1}\vartheta^n = \vartheta^{n-1}$$

has a solution for some $\vartheta > 0$. We may easily solve ϑ

$$\vartheta = \frac{n-1}{2^{n+1}n}.$$

Then if

$$C^{1/p} \frac{(\beta - \alpha)^2}{\alpha^2} \delta_1^{1-1/p} < \chi$$

we get

$$(3.2) \quad \int_{B_\vartheta} |Du|^2 dx \leq \vartheta^{n-1+2\sigma} \int_{B_1} |Du|^2 dx,$$

for some $\sigma > 0$. Similarly in the case (ii) we have (3.2) provided that we choose χ and ϑ as in (3.1) and δ_1 such that

$$C^{1/p} \frac{(\beta - \alpha)^2}{\beta^2} \delta_1^{1-1/p} < \chi.$$

In the case (iii) we choose χ the largest number such that the equation

$$2\chi + 2^{n+1} \left(\frac{\beta}{\alpha} \right)^2 \vartheta^n = \vartheta^{n-1}$$

has a solution $\vartheta > 0$. We may again solve ϑ

$$(3.3) \quad \vartheta = \frac{n-1}{2^{n+1}n} \left(\frac{\alpha}{\beta} \right)^2.$$

Arguing as before, the estimate (3.2) follows from (2.9) and (2.10) if we choose δ_1 such that

$$(3.4) \quad C^{1/p} \left(\frac{\beta}{\alpha} \right)^2 \delta_1^{1-1/p} < \chi.$$

By comparing the above three cases the claim obviously follows by choosing

$$(3.5) \quad \delta_1^{1-1/p} < \left(\frac{\alpha}{\beta} \right)^2 C^{-1/p} \chi$$

where $\chi = \frac{1}{2} \vartheta^{n-1} - 2^n \left(\frac{\beta}{\alpha} \right)^2 \vartheta^n$ and ϑ is from (3.3). \square

We recall the following result (see the proof of [6, Theorem 3.1])

Proposition 3.2. *Let $u \in W^{1,2}(\Omega)$ be a solution of (1.2). Let $x_0 \in \Omega$ and $r > 0$ such that $B_r(x_0) \subset\subset \Omega$. There exist $\lambda_0 > 1, c > 0$ and $\sigma > 0$, depending only on n , such that if*

$$\frac{\beta}{\alpha} < \lambda_0$$

then for every $\rho < r$

$$\int_{B_\rho(x_0)} |Du|^2 dx \leq c \left(\frac{\rho}{r} \right)^{n-1+\sigma} \int_{B_r(x_0)} |Du|^2 dx.$$

We remark that we may estimate the number λ_0 (see [6, (20)-(21)]) by

$$(3.6) \quad \lambda_0 = \frac{n^n + n(n-1)^{n-1} - (n-1)^n}{n^n - n(n-1)^{n-1} + (n-1)^n}.$$

Proof of Theorem 1.1. Let us recall the Euler-Lagrange equation for the critical set E_θ

$$\gamma H_{\partial E_\theta} + \beta |Du_\beta|^2 - \alpha |Du_\alpha|^2 = \lambda \quad \text{on } \partial E_\theta \setminus \{0\}.$$

This can be rewritten as

$$(3.7) \quad \gamma H_{\partial E_\theta} + \beta |\partial_\nu u_\beta|^2 - \alpha |\partial_\nu u_\alpha|^2 + (\beta - \alpha) |D_\tau u|^2 = \lambda \quad \text{on } \partial E_\theta \setminus \{0\},$$

where $D_\tau u$ is the tangential gradient of u on $\partial E_\theta \setminus \{0\}$. From the transmission condition (1.3) we deduce

$$(3.8) \quad \beta |\partial_\nu u_\beta|^2 < \alpha |\partial_\nu u_\alpha|^2 \quad \text{on } \partial E_\theta \setminus \{0\}.$$

Since $H_{\partial E_\theta}(x) = (n-1) \cos \theta \cdot |x|^{-1}$ for $x \in \partial E \setminus \{0\}$, we obtain from the Euler-Lagrange equation (3.7), from (3.8), and from the transmission condition (1.3) that

$$(3.9) \quad |\partial_\nu u_\beta(x)| \geq \frac{c}{\sqrt{|x|}} \quad \text{on } \partial E \setminus \{0\},$$

for some constant $c > 0$. In particular, since the set $\partial E \setminus \{0\}$ is connected, this implies that $\partial_\nu u_\beta$ does not change sign on $\partial E \setminus \{0\}$, and we may thus assume it to be positive.

Let us fix $\rho > 0$ and choose a cut-off function $\zeta \in C_0^\infty(B_\rho)$ such that $\zeta \equiv 1$ in $B_{\rho/2}$ and $|D\zeta| \leq 4/\rho$. Since u_β is harmonic in E , we obtain from (3.9) and by integrating by parts that

$$\int_{E \cap B_\rho} \langle Du_\beta, D\zeta \rangle dx = \int_{\partial E \cap B_\rho} \partial_\nu u_\beta \zeta d\mathcal{H}^{n-1} \geq c \int_{\partial E \cap B_{\rho/2}} |x|^{-1/2} \mathcal{H}^{n-1} \geq \tilde{c} \rho^{n-3/2}.$$

On the other hand Hölder's inequality implies

$$\begin{aligned} \int_{E \cap B_\rho} \langle Du_\beta, D\zeta \rangle dx &\leq \left(\int_{E \cap B_\rho} |D\zeta|^2 dx \right)^{1/2} \left(\int_{E \cap B_\rho} |Du_\beta|^2 dx \right)^{1/2} \\ &\leq C \rho^{n/2-1} \left(\int_{E \cap B_\rho} |Du_\beta|^2 dx \right)^{1/2}. \end{aligned}$$

Therefore

$$\int_{E \cap B_\rho} |Du_\beta|^2 dx \geq c_0 \rho^{n-1}$$

for some constant $c_0 > 0$. The claim now follows by a standard iteration argument from Proposition 3.1 (i) and (iii) and from Proposition 3.2. \square

Remark 3.3. The constant δ_0 can be explicitly estimated in terms of the constant δ_1 from Proposition 3.1, since the spherical sector has the volume

$$|E_\theta \cap B_1| = \omega_{n-1} \left(\int_0^\theta \sin^n t dt + \frac{\sin^{n-1} \theta \cos \theta}{n} \right),$$

where ω_{n-1} is the volume of the $(n-1)$ -dimensional unit ball. The formula for δ_1 is given by (3.5). The constant λ_0 is estimated in (3.6).

Note that in dimension 2 the Euler-Lagrange equation (1.4) reduces to

$$\beta |Du_\beta|^2 - \alpha |Du_\alpha|^2 = \lambda \quad \text{on } \partial E_\theta \setminus \{0\}$$

and it is not clear to us if this weaker information is enough to establish Theorem 1.1 also in this case. However, as we already mentioned earlier, for $n \geq 3$, the proof of that theorem can be easily generalized to more general cones.

Remark 3.4. Let $K \subset \mathbb{R}^{n-1}$ be an open, uniformly convex and C^2 -regular set such that $0 \in K$. Theorem 1.1 can be generalized to conical sets of the form

$$E = \{\lambda(x', 1) \in \mathbb{R}^n \mid x' \in K, \lambda \geq 0\}.$$

4. REGULARITY OF MINIMIZERS

Throughout this section the dimension n and the constants α, β and γ will remain fixed. Thus we denote by C a generic constant depending on these quantities and whose value may change from line to line. On the other hand special constants will be numbered and their dependence on other quantities will be explicitly mentioned. Before proceeding in the regularity proof we recall the following result which was proved in [6, Theorem 1].

Theorem 4.1. *Let (E, u) be a minimizer of problem (P_c) . There exists a constant Λ such that (E, u) is also a minimizer of the penalized functional*

$$\mathcal{F}_\Lambda(F, v) = \gamma P(F, \Omega) + \int_\Omega \sigma_F(x) |Dv|^2 dx + \Lambda ||F| - |E||$$

among all (F, v) such that $v = u$ on $\partial\Omega$.

Motivated by the previous theorem we give the following definition. To this aim we denote for a set E with finite perimeter in Ω and a function $u \in W^{1,2}(\Omega)$ by

$$\mathcal{F}(E, u; U) = \gamma P(E, U) + \int_U \sigma_E |Du|^2 dx$$

the energy of the pair (E, u) in an open set $U \subset\subset \Omega$. We say that a pair (E, u) is a Λ -minimizer of \mathcal{F} in Ω if for every $B_r(x_0) \subset \Omega$ and every pair (F, v) , where F is a set of finite perimeter with $F \Delta E \subset\subset B_r(x_0)$ and $v - u \in W_0^{1,2}(B_r(x_0))$, we have

$$\mathcal{F}(E, u; B_r(x_0)) \leq \mathcal{F}(F, v; B_r(x_0)) + \Lambda |F \Delta E|.$$

Note that any minimizing pair (E, u_E) of the constrained problem (P_c) is a Λ -minimizer of \mathcal{F} for some Λ . Since Λ will be fixed from now on the dependence of the constants on Λ will not be highlighted.

Lemma 4.2. *Let (E, u) be a Λ -minimizer of \mathcal{F} in Ω . For every $\tau \in (0, 1)$ there exists $\varepsilon_1 = \varepsilon_1(\tau) > 0$ such that if $B_r(x_0) \subset \Omega$ and $P(E; B_r(x_0)) < \varepsilon_1 r^{n-1}$ then*

$$\mathcal{F}(E, u; B_{\tau r}(x_0)) \leq C_1 \tau^n (\mathcal{F}(E, u; B_r(x_0)) + r^n)$$

for some constant C_1 independent of τ and r .

Proof. If we replace E by $\frac{E-x_0}{r}$, u by $y \rightarrow r^{-1/2}u(x_0 + ry)$ and Λ by Λr we may assume that $r = 1$, $x_0 = 0$ and (E, u) is a Λr -minimizer in $\frac{\Omega-x_0}{r}$. After this rescaling we are left to prove that for a given $\tau \in (0, 1)$ there exists $\varepsilon_1 = \varepsilon_1(\tau) > 0$ such that if $P(E, B_1) < \varepsilon_1$ then

$$\mathcal{F}(E, u; B_\tau) \leq C_1 \tau^n (\mathcal{F}(E, u; B_1) + \Lambda r).$$

Without loss of generality we may assume that $\tau < 1/2$.

If the perimeter of E in B_1 is small then, by the relative isoperimetric inequality, either $|B_1 \cap E|$ or $|B_1 \setminus E|$ is small. Assume the latter is true. Then we have

$$|B_1 \setminus E| \leq c(n)P(E, B_1)^{\frac{n}{n-1}}.$$

By Fubini's theorem and choosing as a representative of E the set of points of density one, we have

$$|B_1 \setminus E| \geq \int_\tau^{2\tau} \mathcal{H}^{n-1}(\partial B_\rho \setminus E) d\rho.$$

Therefore we may choose $\rho \in (\tau, 2\tau)$ such that

$$(4.1) \quad \mathcal{H}^{n-1}(\partial B_\rho \setminus E) \leq \frac{c(n)}{\tau} P(E, B_1)^{\frac{n}{n-1}} \leq \frac{c(n)\varepsilon_1^{\frac{1}{n-1}}}{\tau} P(E, B_1).$$

Set $F = E \cup B_\rho$ and observe that by the choice of the representative of E it holds

$$P(F; B_1) \leq P(E, B_1 \setminus \bar{B}_\rho) + \mathcal{H}^{n-1}(\partial B_\rho \setminus E).$$

Choosing (F, u) as a competing pair and using the Λr -minimality of (E, u) we get

$$\begin{aligned} \gamma P(E, B_1) + \int_{B_1} \sigma_E |Du|^2 dx &\leq \gamma P(F; B_1) + \int_{B_1} \sigma_F |Du|^2 dx + \Lambda r |F \setminus E| \\ &\leq \gamma (P(E, B_1 \setminus \bar{B}_\rho) + \mathcal{H}^{n-1}(\partial B_\rho \setminus E)) + \int_{B_1} \sigma_F |Du|^2 dx + \Lambda r |B_\rho|. \end{aligned}$$

Hence, choosing ε_1 such that $c(n)\varepsilon_1^{\frac{1}{n-1}} \leq \tau^{n+1}$ and recalling that $\rho \in (\tau, 2\tau)$ we get from (4.1) that

$$\gamma P(E, B_\tau) + \int_{B_\tau} \sigma_E |Du|^2 dx \leq \gamma \tau^n P(E, B_1) + \beta \int_{B_{2\tau}} |Du|^2 dx + \Lambda r |B_{2\tau}|.$$

If we choose ε_1 such that $c(n)\varepsilon_1^{\frac{n}{n-1}} \leq \varepsilon_0(2\tau)|B_1|$, where ε_0 is from Proposition 2.4, we get

$$\int_{B_{2\tau}} |Du|^2 dx \leq 2^n C_0 \tau^n \int_{B_1} |Du|^2 dx$$

and the result follows. □

The next result is contained in [13, Theorem 2], where it is proven for local minimizers of \mathcal{F} . However, in the case of Λ -minimizers the same proof applies without changes once one observes that by comparing (E, u) with (F, u) , where $F = E \setminus B_r(x_0)$ and $B_r(x_0) \subset\subset \Omega$, one gets

$$\gamma P(E, B_r(x_0)) + (\beta - \alpha) \int_{E \cap B_r(x_0)} |Du|^2 dx \leq \gamma \mathcal{H}^{n-1}(\partial B_r \cap E) + \Lambda |B_r| \leq Cr^{n-1}.$$

Theorem 4.3. *Let (E, u) be a Λ -minimizer of \mathcal{F} in Ω . For every open set $U \subset\subset \Omega$ there exists a constant C_2 , depending on U and $\|Du\|_{L^2(\Omega)}$, such that for every $B_r(x_0) \subset U$ it holds*

$$\mathcal{F}(E, u; B_r(x_0)) \leq C_2 r^{n-1}.$$

By combining the energy decay Lemma 4.2 with the energy upper bound given by Theorem 4.3 we obtain the following density lower bound by a standard iteration argument. From now on the topological boundary ∂E must be understood by considering the correct representative of the set (see [15, Proposition 12.19]). In particular, for such representative of E it holds $\overline{\partial^* E} = \partial E$.

Proposition 4.4. *Let (E, u) be a Λ -minimizer of \mathcal{F} and $U \subset\subset \Omega$. There exists a constant $c > 0$, depending on U and $\|Du\|_{L^2(\Omega)}$, such that for every $B_r(x_0) \subset U$ with $x_0 \in \partial E$ we have*

$$P(E, B_r(x_0)) \geq cr^{n-1}.$$

Moreover $\mathcal{H}^{n-1}((\partial E \setminus \partial^* E) \cap \Omega) = 0$.

Proof. The proof is as in [3, Theorem 7.21] and in fact, even simpler. Fix $x_0 \in \partial^* E$. Without loss of generality we may assume that $x_0 = 0$. Fix $\tau \in (0, 1)$ such that $2C_1\tau^{\frac{1}{2}} < 1$ and $\sigma \in (0, 1)$ such that $2C_1C_2\sigma < \varepsilon_1(\tau)\gamma$ and r_0 such that $r_0 < \min\{\varepsilon_1(\tau)\gamma, C_2\}$, where C_1, ε_1 and C_2 are the constants from Lemma 4.2 and Theorem 4.3. Assume by contradiction that for some $B_r \subset U$ with $r < r_0$ we have $P(E, B_r) < \varepsilon_1(\sigma)r^{n-1}$. Then by induction we deduce that

$$(4.2) \quad \mathcal{F}(E, u; B_{\sigma\tau^h r}) \leq \varepsilon_1(\tau)\gamma\tau^{\frac{h}{2}}(\sigma\tau^h r)^{n-1}$$

for all $h = 0, 1, 2, \dots$. Indeed, when $h = 0$, using Lemma 4.2 and Theorem 4.3 and recalling our choice of σ and r_0 we obtain

$$\mathcal{F}(E, u; B_{\sigma r}) \leq C_1\sigma^n(\mathcal{F}(E, u; B_r) + r^n) \leq 2C_1C_2\sigma(\sigma r)^{n-1} < \varepsilon_1(\tau)\gamma(\sigma r)^{n-1}.$$

If (4.2) holds for h to deduce that it also holds for $h + 1$ it is enough to apply Lemma 4.2 and to recall that $2C_1\tau^{\frac{1}{2}} < 1$ and that $r < r_0 < \varepsilon_1(\tau)\gamma$. In particular, this implies

$$\lim_{r \rightarrow 0} \frac{P(E, B_r)}{r^{n-1}} = 0,$$

which is a contradiction since $x_0 \in \partial^* E$. This proves that for every $x_0 \in \partial^* E$ and $r < r_0$ it holds $P(E, B_r(x_0)) \geq \varepsilon_1(\sigma)r^{n-1}$. The claim follows from the fact that $\overline{\partial^* E} = \partial E$. The fact that $\mathcal{H}^{n-1}((\partial E \setminus \partial^* E) \cap \Omega) = 0$ follows from the density lower bound and [3, (2.42)]. \square

Let (E, u) be a Λ -minimizer of \mathcal{F} . We introduce *the excess of E at the point $x \in \partial E$ at the scale $r > 0$ in direction $\nu \in S^{n-1}$*

$$\mathcal{E}(x, r, \nu) := \frac{1}{2r^{n-1}} \int_{\partial E \cap B_r(x)} |\nu_E(y) - \nu|^2 d\mathcal{H}^{n-1}(y)$$

and

$$\mathcal{E}(x, r) := \min_{\nu \in S^{n-1}} \mathcal{E}(x, r, \nu)$$

and the rescaled Dirichlet integral of u

$$\mathcal{D}(x, r) := \frac{1}{r^{n-1}} \int_{B_r(x)} |Du|^2 dy.$$

The density lower bound implies the following important properties of Λ -minimizers of \mathcal{F} . The first one is the so called height bound lemma.

Lemma 4.5 (Height bound). *Let (E, u) be a Λ -minimizer of \mathcal{F} in $B_r(x_0)$. There exist C and $\varepsilon > 0$, depending on $\|Du\|_{L^2(B_r(x_0))}$, such that if $x_0 \in \partial E$ and*

$$\mathcal{E}(x_0, r, \nu) < \varepsilon$$

for some $\nu \in S^{n-1}$ then

$$\sup_{y \in \partial E \cap B_{r/2}(x_0)} \frac{|\nu \cdot (y - x_0)|}{r} \leq C \mathcal{E}(x_0, r, \nu)^{\frac{1}{2(n-1)}}.$$

Proof. The proof of this result can be obtained arguing exactly as in the case of Λ -minimizers of the perimeter [15, Theorem 22.8]. Indeed, it is based only on the relative isoperimetric inequality, the density lower bound, i.e., Proposition 4.4 and the compactness result below. \square

The next result is proved as in the case of the Λ -minimizers of the perimeter with the obvious changes due to the presence of the Dirichlet integral.

Lemma 4.6 (Compactness). *Let (E_h, u_h) be a sequence of Λ_h -minimizers of \mathcal{F} in Ω such that $\Lambda_h \rightarrow \Lambda \in [0, \infty)$ and $\sup_h \mathcal{F}(E_h, u_h; \Omega) < \infty$. Then there exist a subsequence, not relabelled, and a Λ -minimizer (E, u) of \mathcal{F} such that for every open set $U \subset\subset \Omega$, $E_h \rightarrow E$ in $L^1(U)$, $P(E_h, U) \rightarrow P(E, U)$, $u_h \rightarrow u$ in $W^{1,2}(U)$ and moreover*

- (i) if $x_h \in \partial E_h \cap U$ and $x_h \rightarrow x \in U$, then $x \in \partial E \cap U$.
- (ii) if $x \in \partial E \cap U$ there exists $x_h \in \partial E_h \cap U$ such that $x_h \rightarrow x$.

If in addition, $Du_h \rightarrow 0$ weakly in $L^2_{loc}(\Omega, \mathbb{R}^n)$ and $\Lambda_h \rightarrow 0$, then E is a local minimizer of the perimeter, i.e., for every F such that $F \Delta E \subset\subset B_r(x_0) \subset \Omega$ it holds

$$P(E, B_r(x_0)) \leq P(F, B_r(x_0)).$$

Proof. We start by proving the Λ -minimality of (E, u) . Let us fix $B_r(x_0) \subset\subset \Omega$ and assume that $x_0 = 0$. Without loss of generality we may assume that $E_h \rightarrow E$ in $L^1(B_r)$, $u_h \rightarrow u$ weakly in $W^{1,2}(B_r)$ and strongly in $L^2(B_r)$. Let (F, v) be a pair such that $F \Delta E \subset\subset B_r$ and $\text{supp}(u - v) \subset\subset B_r$. By Fubini's theorem and passing to a subsequence if necessary we may choose $\rho < r$ such that $F \Delta E \subset\subset B_\rho$, $\text{supp}(u - v) \subset\subset B_\rho$,

$$\mathcal{H}^{n-1}(\partial^* F \cap \partial B_\rho) = \mathcal{H}^{n-1}(\partial^* E_h \cap \partial B_\rho) = 0 \quad \text{and} \quad \lim_{h \rightarrow 0} \mathcal{H}^{n-1}(\partial B_\rho \cap E \Delta E_h) = 0$$

where it is understood that E and E_h stand for the set of points of density one of E and E_h , respectively. Fix a cut-off function $\zeta \in C_0^1(B_r)$ such that $\zeta \equiv 1$ in B_ρ , and choose $F_h = (F \cap B_\rho) \cup (E_h \setminus B_\rho)$ and $v_h = \zeta v + (1 - \zeta)u_h$. Then by the Λ_h -minimality of (E_h, u_h) , the choice of ρ , the strong convergence of $u_h \rightarrow u$ in L^2 and by convexity, we have

(4.3)

$$\begin{aligned} \gamma P(E_h, B_r) + \int_{B_r} \sigma_{E_h} |Du_h|^2 dx &\leq \gamma P(F_h, B_r) + \int_{B_r} \sigma_{F_h} |Dv_h|^2 dx + \Lambda_h |F_h \Delta E_h| \\ &\leq \gamma [P(F, B_\rho) + P(E_h, B_r \setminus \bar{B}_\rho)] + \int_{B_r} \sigma_{F_h} \zeta |Dv|^2 dx + \int_{B_r} \sigma_{E_h} (1 - \zeta) |Du_h|^2 dx + \varepsilon_h + \Lambda_h |F_h \Delta E_h| \end{aligned}$$

for some $\varepsilon_h \rightarrow 0$. Thus by a simple lower semicontinuity argument we have

$$\gamma P(E, B_\rho) + \int_{B_r} \sigma_E \zeta |Du|^2 dx \leq \gamma P(F, B_\rho) + \int_{B_r} \sigma_F \zeta |Dv|^2 dx + \Lambda |F \Delta E|.$$

Letting $\zeta \downarrow \chi_{B_\rho}$ in the previous inequality we conclude that

$$(4.4) \quad \gamma P(E, B_\rho) + \int_{B_\rho} \sigma_E |Du|^2 dx \leq \gamma P(F, B_\rho) + \int_{B_\rho} \sigma_F |Dv|^2 dx + \Lambda |F \Delta E|$$

thus proving the Λ -minimality of (E, u) . Similarly, choosing $F = E$ and $v = u$ in (4.3) and arguing as before we get

$$\limsup_{h \rightarrow \infty} \left(\gamma P(E_h, B_\rho) + \int_{B_r} \sigma_{E_h} \zeta |Du_h|^2 dx \right) \leq \gamma P(E, B_\rho) + \int_{B_r} \sigma_F \zeta |Du|^2 dx.$$

Letting $\zeta \downarrow \chi_{B_\rho}$ we conclude that

$$\lim_{h \rightarrow \infty} P(E_h, B_\rho) = P(E, B_\rho) \quad \text{and} \quad \lim_{h \rightarrow \infty} \int_{B_\rho} \sigma_{E_h} |Du_h|^2 dx = \int_{B_\rho} \sigma_E |Du|^2 dx.$$

A standard argument then implies that $P(E_h, U) \rightarrow P(E, U)$, $u_h \rightarrow u$ in $W^{1,2}(U)$ for every open set $U \subset\subset \Omega$.

Finally we note that if $Du = 0$ we can choose $v = u$ in (4.4) thus proving that E is a Λ -minimizer of the perimeter. The claim (i) and (ii) follow exactly as in [15, Theorem 21.14]. \square

The next consequence of the density lower bound is the Lipschitz approximation. Its proof is based only on the height bound estimate and can be obtained by following word by word the proof given in [15, Theorem 23.7]. To that aim we use the notation $x = (x', x_n) \in \mathbb{R}^{n-1} \times \mathbb{R}$ for a generic point in \mathbb{R}^n and with a slight abuse of notation we still denote by Df the gradient of a function $f : \mathbb{R}^{n-1} \rightarrow \mathbb{R}$.

Proposition 4.7 (Lipschitz approximation). *Let (E, u) be a Λ -minimizer of \mathcal{F} in $B_r(x_0)$. There exist C_3 and $\varepsilon_3 > 0$, depending on $\|Du\|_{L^2(B_r(x_0))}$, such that if $x_0 \in \partial E$ and*

$$\mathcal{E}(x_0, r, e_n) < \varepsilon_3$$

then there exists a Lipschitz function $f : \mathbb{R}^{n-1} \rightarrow \mathbb{R}$ such that

$$\sup_{x' \in \mathbb{R}^{n-1}} \frac{|f(x')|}{r} \leq C_3 \mathcal{E}(x_0, r, e_n)^{\frac{1}{2(n-1)}}, \quad \|Df\|_{L^\infty} \leq 1$$

and

$$\frac{1}{r^{n-1}} \mathcal{H}^{n-1}(\partial E \Delta \Gamma_f \cap B_{r/2}(x_0)) \leq C_3 \mathcal{E}(x_0, r, e_n)$$

where Γ_f is the graph of f . Moreover

$$\frac{1}{r^{n-1}} \int_{B_{r/2}^{n-1}(x'_0)} |Df|^2 dx' \leq C_3 \mathcal{E}(x_0, r, e_n).$$

Finally we state the following reverse Poincaré inequality which plays the role of the classical Caccioppoli inequality in the elliptic regularity theory. The proof can be obtained exactly as in the case of Λ -minimizers of the perimeter by constructing a suitable squashing-deformation F of the set E and comparing the energies at (E, u) and (F, u) , see [15, Theorem 24.1].

Theorem 4.8. *Let (E, u) be a Λ -minimizer of \mathcal{F} in $B_r(x_0)$. There exist two constants C_4 and $\varepsilon_4 > 0$ such that if $x_0 \in \partial E$ and $\mathcal{E}(x_0, r, \nu) < \varepsilon_4$ then*

$$\mathcal{E}(x_0, r/2, \nu) \leq C_4 \left(\frac{1}{r^{n+1}} \int_{\partial E \cap B_r(x_0)} |\nu \cdot (x - x_0) - c|^2 d\mathcal{H}^{n-1} + \mathcal{D}(x_0, r) + r \right)$$

for every $c \in \mathbb{R}$.

While in our case the proofs for the height bound, the compactness, the reverse Poincaré and the Lipschitz approximation are exactly as in the case of Λ -minimizers of the perimeter, the next step in the regularity proof, i.e., the excess decay, is different. This is due to the interplay between the excess and the rescaled Dirichlet integral. We follow an argument similar to the one used in proving the flatness decay for the minimizers of the Mumford-Shah functional, see [3, Theorem 8.15]. To this aim we first prove the following weaker form of the Euler-Lagrange equation.

Proposition 4.9. *Let (E, u) be a Λ -minimizer of \mathcal{F} in $B_r(x_0)$. For every vector field $X \in C_0^1(B_r(x_0); \mathbb{R}^n)$ it holds*

$$(4.5) \quad \gamma \int_{\partial E} \operatorname{div}_\tau X d\mathcal{H}^{n-1} + \int_\Omega \sigma_E (|Du|^2 \operatorname{div} X - 2\langle DX Du, Du \rangle) dx \leq \Lambda \int_{\partial E} |X| d\mathcal{H}^{n-1},$$

where div_τ denotes the tangential divergence on ∂E .

Proof. The argument is similar to the one in [3, Theorem 7.35] and therefore we only give the sketch of the proof. For a given $X \in C_0^1(B_r(x_0); \mathbb{R}^n)$ we set for every small $t > 0$, $\Phi_t(x) = x - tX(x)$, $E_t = \Phi_t(E)$ and $u_t(y) = u(\Phi_t^{-1}(y))$. From the Λ -minimality it follows

$$\gamma [P(E_t, B_r(x_0)) - P(E, B_r(x_0))] + \int_{B_r(x_0)} (\sigma_{E_t} |Du_t|^2 - \sigma_E |Du|^2) dx + \Lambda |E_t \Delta E| \geq 0.$$

The conclusion then follows with the same standard calculations used to derive the first variation of the perimeter and the Dirichlet integral (see [3, Theorem 7.35]) and observing that

$$\lim_{t \rightarrow 0} \frac{|E_t \Delta E|}{t} \leq \int_{\partial E} |X \cdot \nu_E| d\mathcal{H}^{n-1},$$

see for instance [9, Theorem 3.2]. □

We are ready to prove the excess improvement.

Proposition 4.10. *Let (E, u) be a Λ -minimizer of \mathcal{F} in $B_r(x_0)$. For every $\tau \in (0, 1/2)$ and M there exists $\varepsilon_5 = \varepsilon_5(\tau, M) \in (0, 1)$ such that if $x_0 \in \partial E$ and*

$$\mathcal{E}(x_0, r) \leq \varepsilon_5 \quad \text{and} \quad \mathcal{D}(x_0, r) + r \leq M\mathcal{E}(x_0, r)$$

then

$$\mathcal{E}(x_0, \tau r) \leq C_5(\tau^2 \mathcal{E}(x_0, r) + \mathcal{D}(x_0, 2\tau r) + \tau r)$$

for some constant C_5 depending on $\mathcal{F}(E, u; B_r(x_0))$.

Proof. Without loss of generality we may assume that $\tau < 1/8$. We argue by contradiction. After performing the same translation and rescaling used in the proof of Lemma 4.2, we may assume that there exist an infinitesimal sequence ε_h , a sequence $r_h \in \mathbb{R}$ and a sequence (E_h, u_h) of Λr_h -minimizers of \mathcal{F} in B_1 , with equibounded energies such that, denoting by \mathcal{E}_h the excess of E_h and by \mathcal{D}_h the rescaled Dirichlet integral of u_h , we have

$$(4.6) \quad \mathcal{E}_h(0, 1) := \varepsilon_h, \quad \mathcal{D}_h(0, 1) + r_h \leq M\varepsilon_h,$$

and

$$(4.7) \quad \mathcal{E}_h(0, \tau) > C_5(\tau^2 \mathcal{E}_h(0, 1) + \mathcal{D}_h(0, 2\tau) + \tau r_h)$$

Moreover, up to rotating each E_h if necessary we may also assume, that for all h

$$\mathcal{E}_h(0, 1) = \mathcal{E}_h(0, 1, e_n) = \frac{1}{2} \int_{\partial E_h \cap B_1} |\nu_{E_h}(x) - e_n|^2 d\mathcal{H}^{n-1}.$$

Step 1. Recalling Proposition 4.7, we have that for every h sufficiently large there exists a 1-Lipschitz function $f_h : \mathbb{R}^{n-1} \rightarrow \mathbb{R}$ such that

$$(4.8) \quad \sup_{\mathbb{R}^{n-1}} |f_h| \leq \varepsilon_h^{\frac{1}{2(n-1)}}, \quad \mathcal{H}^{n-1}(\partial E_h \Delta \Gamma_{f_h} \cap B_{1/2}) \leq C_3 \varepsilon_h, \quad \int_{B_{1/2}^{n-1}} |Df_h|^2 dx' \leq C_3 \varepsilon_h.$$

Therefore, setting

$$g_h(x') := \frac{f_h(x') - a_h}{\sqrt{\varepsilon_h}}, \quad \text{where } a_h = \int_{B_{1/2}^{n-1}} f_h dx',$$

we may assume, up to a not relabelled subsequence, that the functions g_h converge weakly in $H^1(B_{1/2}^{n-1})$ and strongly in $L^2(B_{1/2}^{n-1})$ to a function g .

We claim that g is harmonic in $B_{1/2}$. To prove this it is enough to show that for any $\varphi \in C_0^1(B_{1/2}^{n-1})$

$$(4.9) \quad \lim_{h \rightarrow \infty} \frac{1}{\sqrt{\varepsilon_h}} \int_{B_{1/2}^{n-1}} Df_h \cdot D\varphi dx' = 0.$$

In order to prove this equality we fix $\delta > 0$ so that $\text{supp } \varphi \times [-2\delta, 2\delta] \subset B_{1/2}$, choose a cut-off function $\psi : \mathbb{R} \rightarrow [0, 1]$ with support in $(-2\delta, 2\delta)$ and $\psi \equiv 1$ in $[-\delta, \delta]$, and apply to E_h the weak Euler-Lagrange equation (4.5) with $X = (0, \dots, 0, \varphi\psi)$. By the height bound (Lemma 4.5) for h large it holds $\partial E_h \cap B_{1/2} \subset B_{1/2}^{n-1} \times (-\delta, \delta)$. Therefore by denoting ν'_{E_h} the vector made up by the first $n-1$ components of ν_{E_h} , we have

$$\begin{aligned} -\gamma \int_{\partial E_h \cap B_{1/2}} \nu_{E_h} \cdot e_n D\varphi \cdot \nu'_{E_h} d\mathcal{H}^{n-1} + \int_{B_{1/2}} \sigma_{E_h} (|Du_h|^2 \varphi\psi' - 2D_n u_h Du_h \cdot D(\varphi\psi)) dx \\ \leq \Lambda r_h \int_{\partial E_h \cap B_{1/2}} |\varphi\psi| d\mathcal{H}^{n-1}. \end{aligned}$$

Thus, using the energy upper bound and recalling the inequality in (4.6), we have

$$-\gamma \int_{\partial E_h \cap B_{1/2}} \nu_{E_h} \cdot e_n D\varphi \cdot \nu'_{E_h} d\mathcal{H}^{n-1} \leq C \varepsilon_h$$

for some constant C depending on Λ , φ and ψ , but independent of h . Therefore, dividing the above inequality by $\sqrt{\varepsilon_h}$, letting $h \rightarrow \infty$ and replacing φ by $-\varphi$ we conclude that

$$\lim_{h \rightarrow \infty} \frac{1}{\sqrt{\varepsilon_h}} \int_{\partial E_h \cap B_{1/2}} \nu_{E_h} \cdot e_n D\varphi \cdot \nu'_{E_h} d\mathcal{H}^{n-1} = 0.$$

From this equation we get (4.9) by observing that

$$\begin{aligned} - \int_{\partial E_h \cap B_{1/2}} \nu_{E_h} \cdot e_n D\varphi \cdot \nu'_{E_h} d\mathcal{H}^{n-1} &= - \int_{\Gamma_{f_h} \cap B_{1/2}} \nu_{E_h} \cdot e_n D\varphi \cdot \nu'_{E_h} d\mathcal{H}^{n-1} \\ &\quad - \int_{(\partial E_h \setminus \Gamma_{f_h}) \cap B_{1/2}} \nu_{E_h} \cdot e_n D\varphi \cdot \nu'_{E_h} d\mathcal{H}^{n-1} + \int_{(\Gamma_{f_h} \setminus \partial E_h) \cap B_{1/2}} \nu_{E_h} \cdot e_n D\varphi \cdot \nu'_{E_h} d\mathcal{H}^{n-1}. \end{aligned}$$

Indeed, recalling the second inequality in (4.8), we have that

$$0 = \lim_{h \rightarrow \infty} \frac{-1}{\sqrt{\varepsilon_h}} \int_{\Gamma_{f_h} \cap B_{1/2}} \nu_{E_h} \cdot e_n D\varphi \cdot \nu'_{E_h} d\mathcal{H}^{n-1} = \lim_{h \rightarrow \infty} \frac{1}{\sqrt{\varepsilon_h}} \int_{B_{1/2}^{n-1}} \frac{Df_h \cdot D\varphi}{\sqrt{1 + |Df_h|^2}} dx',$$

from which (4.9) immediately follows using the third inequality in (4.8).

Step 2. Since g is harmonic we have for $\tau \in (0, 1/8)$ that

$$\begin{aligned} \int_{B_{2\tau}^{n-1}} |g(x') - g(0) - Dg(0) \cdot x'|^2 dx' &\leq c(n)\tau^{n+3} \sup_{B_{1/4}} |D^2g|^2 \\ &\leq c(n)\tau^{n+3} \int_{B_{1/2}^{n-1}} |Dg|^2 dx'. \end{aligned}$$

Since by (4.8) we have that

$$\int_{B_{1/2}^{n-1}} |Dg|^2 dx \leq \liminf_{h \rightarrow \infty} \int_{B_{1/2}^{n-1}} |Dg_h|^2 dx \leq C_3,$$

and by the mean value property $(g)_r := \int_{B_r^{n-1}} g dx' = g(0)$ and $(Dg)_r = Dg(0)$, we may conclude that

$$\lim_{h \rightarrow \infty} \int_{B_{2\tau}^{n-1}} |g_h(x') - (g_h)_{2\tau} - (Dg_h)_{2\tau} \cdot x'|^2 dx' = \int_{B_{2\tau}^{n-1}} |g(x') - (g)_{2\tau} - (Dg)_{2\tau} \cdot x'|^2 dx' \leq \widehat{C}\tau^{n+3},$$

for some constant \widehat{C} depending only on C_3 and n . In turn, recalling the definition of g_h , this inequality is equivalent to

$$\lim_{h \rightarrow \infty} \frac{1}{\varepsilon_h} \int_{B_{2\tau}^{n-1}} |f_h(x') - (f_h)_{2\tau} - (Df_h)_{2\tau} \cdot x'|^2 dx' \leq \widehat{C}\tau^{n+3}.$$

From this inequality, recalling that $|Df_h| \leq 1$ and setting

$$c_h := \frac{(f_h)_{2\tau}}{\sqrt{1 + |(Df_h)_{2\tau}|^2}}, \quad \nu_h = \frac{(-(Df_h)_{2\tau}, 1)}{\sqrt{1 + |(Df_h)_{2\tau}|^2}},$$

we easily have

$$\begin{aligned} \limsup_{h \rightarrow \infty} \frac{1}{\varepsilon_h} \int_{\partial E_h \cap \Gamma_{f_h} \cap B_{2\tau}} |\nu_h \cdot x - c_h|^2 d\mathcal{H}^{n-1} \\ \leq \lim_{h \rightarrow \infty} \frac{\sqrt{2}}{\varepsilon_h} \int_{B_{2\tau}^{n-1}} |f_h(x') - (f_h)_{2\tau} - (Df_h)_{2\tau} \cdot x'|^2 dx' \leq \sqrt{2}\widehat{C}\tau^{n+3}. \end{aligned}$$

On the other hand, arguing as in Step 1, we immediately get from the height bound and from the first two inequalities in (4.8) that

$$\lim_{h \rightarrow \infty} \frac{1}{\varepsilon_h} \int_{(\partial E_h \setminus \Gamma_{f_h}) \cap B_{2\tau}} |\nu_h \cdot x - c_h|^2 d\mathcal{H}^{n-1} = 0.$$

Hence, we conclude that

$$(4.10) \quad \limsup_{h \rightarrow \infty} \frac{1}{\varepsilon_h} \int_{\partial E_h \cap B_{2\tau}} |\nu_h \cdot x - c_h|^2 d\mathcal{H}^{n-1} \leq \sqrt{2}\widehat{C}\tau^{n+3}.$$

Note that

$$\begin{aligned} \int_{\partial E_h \cap B_{2\tau}} |\nu_{E_h} - \nu_h|^2 d\mathcal{H}^{n-1} &\leq 2 \int_{\partial E \cap B_{2\tau}} |\nu_{E_h} - e_n|^2 d\mathcal{H}^{n-1} + 2|e_n - \nu_h|^2 \mathcal{H}^{n-1}(\partial E_h \cap B_{2\tau}) \\ &\leq 4\varepsilon_h + C \int_{B_{1/2}^{n-1}} |Df_h|^2 dx' \leq C\varepsilon_h \end{aligned}$$

by the third inequality from (4.8). In particular, this shows that $\mathcal{E}_h(0, 2\tau, \nu_h) \rightarrow 0$ as $h \rightarrow \infty$. Therefore applying Theorem 4.8 and (4.10) we have for h large that

$$\mathcal{E}_h(0, \tau) \leq \mathcal{E}_h(0, \tau, \nu_h) \leq C_4(\widehat{C}\tau^2 \mathcal{E}_h(0, 1) + D_h(0, 2\tau) + 2\tau r_h)$$

which is a contradiction to (4.7) if we choose $C_5 > C_4 \max\{\widehat{C}, 2\}$. \square

Finally we give the proof of the regularity theorem (Theorem 1.2).

Proof of Theorem 1.2. Step 1. We begin by proving that for every $\tau \in (0, 1)$ there exists $\varepsilon_6 = \varepsilon_6(\tau) > 0$ such that if $\mathcal{E}(x, r) \leq \varepsilon_6$ then

$$\mathcal{D}(x, \tau r) \leq C_0 \tau \mathcal{D}(x, r)$$

where C_0 is from Proposition 2.4. We argue by contradiction. After performing the same translation and rescaling used in the proof of Lemma 4.2 we may assume that there exist sequences $\varepsilon_h, r_h > 0$ and a sequence E_h of Λr_h -minimizers of \mathcal{F} in B_1 with equibounded energies, such that, denoting by \mathcal{E}_h the excess of E_h and by \mathcal{D}_h the rescaled Dirichlet integral of u_{E_h} , we have that $0 \in \partial E_h$,

$$(4.11) \quad \mathcal{E}_h(0, 1) = \varepsilon_h \rightarrow 0 \quad \text{and} \quad \mathcal{D}_h(0, \tau) > C_0 \tau \mathcal{D}_h(0, 1).$$

By the energy upper bound (Theorem 4.3) and the compactness lemma (Lemma 4.6) we may assume that $E_h \rightarrow E$ in $L^1(B_1)$ and $0 \in \partial E$. Moreover by the lower semicontinuity of the excess we have that $\mathcal{E}(0, 1) = 0$ where $\mathcal{E}(0, 1)$ is the excess of E at 0. Thus it follows that E is a half space, say H , in B_1 , see [15, Proposition 22.2]. In particular, for h large it holds

$$|(E_h \Delta H) \cap B_1| < \varepsilon_0(\tau) |B_1|$$

where ε_0 is from Proposition 2.4 which gives a contradiction with the inequality in (4.11).

Step 2. Let $U \subset\subset \Omega$ be an open set. We show that for every $\tau \in (0, 1)$ there exists $\varepsilon = \varepsilon(\tau, U) > 0$ such that if $x_0 \in \partial E$, $B_r(x_0) \subset U$ and $\mathcal{E}(x_0, r) + \mathcal{D}(x_0, r) + r < \varepsilon$ then

$$(4.12) \quad \mathcal{E}(x_0, \tau r) + \mathcal{D}(x_0, \tau r) + \tau r \leq C_6 \tau (\mathcal{E}(x_0, r) + \mathcal{D}(x_0, r) + r).$$

First of all, if $\varepsilon \leq \varepsilon_6(\tau)$, then Step 1 implies

$$(4.13) \quad \mathcal{D}(x_0, \tau r) \leq C_0 \tau \mathcal{D}(x_0, r).$$

In order to prove (4.12) we may assume that $\tau < 1/2$. Assume first that $D(x_0, r) + r \leq \tau^{-n} \mathcal{E}(x_0, r)$. Then if $\mathcal{E}(x_0, r) < \min\{\varepsilon_5(\tau, M), \varepsilon_6(2\tau)\}$, for $M = \tau^{-n}$, it follows from Proposition 4.10 that

$$\begin{aligned} \mathcal{E}(x_0, \tau r) &\leq C_5 (\tau^2 \mathcal{E}(x_0, r) + \mathcal{D}(x_0, 2\tau r) + \tau r) \\ &\leq C_5 (\tau^2 \mathcal{E}(x_0, r) + 2C_0 \tau \mathcal{D}(x_0, r) + \tau r) \end{aligned}$$

where the last inequality follows from (4.13) applied to 2τ . On the other hand if $\mathcal{E}(x_0, r) \leq \tau^n (D(x_0, r) + r)$ we immediately obtain

$$\mathcal{E}(x_0, \tau r) \leq \tau^{1-n} \mathcal{E}(x_0, r) \leq \tau (D(x_0, r) + r).$$

Therefore (4.13) implies (4.12) by choosing $\varepsilon = \min\{\varepsilon_5(\tau, M), \varepsilon_6(2\tau), \varepsilon_6(\tau)\}$.

Step 3. Let us fix $\sigma \in (0, 1/2)$. We choose $\tau_0 \in (0, 1)$ such that $C_6 \tau_0 \leq \tau_0^{2\sigma}$ where C_6 is the constant in (4.12). Let $U \subset\subset \Omega$ be an open set. We define

$$\Gamma \cap U = \{x \in \partial E \cap U : \mathcal{E}(x, r) + \mathcal{D}(x, r) + r < \varepsilon(\tau_0, U) \text{ for some } r > 0 \text{ such that } B_r(x) \subset U\}$$

where $\varepsilon(\tau_0)$ is from Step 2. Note that $\Gamma \cap U$ is a relatively open in ∂E . We show that $\Gamma \cap U$ is $C^{1, \sigma}$ -hypersurface.

Indeed (4.12) implies via standard iteration argument that if $x_0 \in \Gamma \cap U$ there exist $r_0 > 0$ and a neighborhood V of x_0 such that for every $x \in \partial E \cap V$ it holds

$$\mathcal{E}(x, \tau_0^k r_0) + \mathcal{D}(x, \tau_0^k r_0) + \tau_0^k r_0 \leq \tau_0^{2\sigma k} \quad \text{for } k = 0, 1, 2, \dots$$

In particular

$$\mathcal{E}(x, \tau_0^k r_0) \leq \tau_0^{2\sigma k}.$$

From this estimate and the density lower bound, arguing exactly as in [7, Theorem 8.2], we obtain that for every $x \in \partial E \cap V$ and $0 < s < t < r_0$ it holds

$$(4.14) \quad |(\nu_E)_s(x) - (\nu_E)_t(x)| \leq ct^\sigma$$

for a constant c depending on τ_0, r_0 and n . Here

$$(\nu_E)_t(x) = \int_{\partial E \cap B_t(x)} \nu_E(y) d\mathcal{H}^{n-1}.$$

The estimate (4.14) first implies that $\Gamma \cap U$ is C^1 (see for instance [7, Theorem 8.4]). By a standard argument we then deduce again from (4.14) that $\Gamma \cap U$ is $C^{1,\sigma}$ -hypersurface. Finally we define $\Gamma := \cup_i (\Gamma \cap U_i)$ where (U_i) is an increasing sequence of open sets such that $U_i \subset\subset \Omega$ and $\Omega = \cup_i U_i$

Step 4. Finally we prove that exists $\eta > 0$, depending only on $\frac{\beta}{\alpha}$ and n , such that

$$\mathcal{H}^{n-1-\eta}(\partial E \setminus \Gamma) = 0.$$

Since the argument is fairly standard we only give the sketch of the proof, see e.g. Section 5 in [2], [4] and [5]. We set

$$\Sigma = \{x \in \partial E \setminus \Gamma : \lim_{r \rightarrow 0} D(x, r) = 0\}.$$

Since by Lemma 2.2 $Du \in L_{loc}^{2p}(\Omega)$ for some $p > 1$, depending only on $\frac{\beta}{\alpha}$ and n , we have that

$$\dim_{\mathcal{H}} \left(\{x \in \Omega : \limsup_{r \rightarrow 0} D(x, r) > 0\} \right) \leq n - p,$$

where $\dim_{\mathcal{H}}$ denotes the Hausdorff dimension. The conclusion will follow if we show that $\Sigma = \emptyset$ when $n \leq 7$ and $\dim_{\mathcal{H}}(\Sigma) \leq n - 8$ otherwise.

Let us first treat the case $n \leq 7$. We argue by contradiction and assume, up to a translation, that $0 \in \Sigma$. Let us take a sequence $r_h \rightarrow 0$ and set $E_h = \frac{E}{r_h}$ and $u_h(x) = r_h^{-1/2}u(r_h x)$. Then (E_h, u_h) is Λr_h -minimizer of \mathcal{F} . Since $Du_h \rightarrow 0$ in $L^2(B_1)$ Lemma 4.6 implies that, up to a subsequence, E_h converges to a minimizer of the perimeter E_∞ and moreover $\lim_{h \rightarrow \infty} P(E_h, U) = P(E_\infty, U)$ for every open set $U \subset B_1$ and $0 \in \partial E_\infty$. Since $n \leq 7$, we know that ∂E_∞ is a smooth manifold. In particular, for any $\varepsilon > 0$ there exists $r > 0$ such that $\mathcal{E}(0, E_\infty, r) < \varepsilon$. However the above convergence of the perimeter implies that $\mathcal{E}(0, E_h, r_h) < \varepsilon$ when h is large enough. By the definition of Γ this contradicts the fact that $0 \in \Sigma$.

In the case $n \geq 8$ we claim that if $s > n - 8$ then it holds $\mathcal{H}^s(\Sigma) = 0$. The proof of this can be achieved arguing exactly as in the proof of [2, Theorem 5.6]. \square

5. APPENDIX

We conclude by going through the proof of Lemma 2.2 and estimate all the relevant constants in the statement.

Proof of Lemma 2.2. Without loss of generality we may assume that $Q_R(x_0)$ is the unit cube Q . Denote $d(x) = \text{dist}(x, \partial Q)$ and define

$$\mathcal{C}_k = \{x \in Q \mid \frac{3}{4}2^{-k-1} \leq d(x) \leq \frac{3}{4}2^{-k}\}.$$

Each \mathcal{C}_k can be divided into cubes of side $\frac{3}{4}2^{-k-1}$. We call this collection \mathcal{G}_k . By Lemma 2.1 we have for $F(x) = d(x)^n |Du(x)|^2$ that

$$\int_P F dx \leq C_0 \left(\int_{\bar{P}} F^m dx \right)^{\frac{1}{m}},$$

where \tilde{P} is the concentric cube to $P \subset \mathcal{C}_k$ or $P \subset Q_{1/4}$, for a constant $C_0 = 4^n C$, where C is the constant from Lemma 2.1.

Denote next $\Phi_t = \{x \in Q \mid F(x) > t\}$, where $t > a := \int_Q |Du|^2 dx$. Applying Calderón-Zygmund decomposition we obtain (in the proof [8, Lemma 6.2] choose $\lambda = 2^{1/m} C_0$)

$$\int_{\Phi_t} F dx \leq C_1 t^{1-m} \int_{\Phi_t} F^m dx$$

for

$$C_1 = 5^n 2^n \lambda = 2^{1/m} 10^n C_0 \leq 4 \cdot 40^n C = C_{S,n}^2 2^{10} \cdot 80^n \frac{\beta}{\alpha}.$$

The result of [8, Proposition 6.1] now follows with the constants $A = C_1$ and $r = p > 1$ such that

$$C_1(p-1) = \frac{p-m}{2}$$

that is

$$p = \frac{2C_1 - m}{2C_1 - 1}.$$

This leads to the inequality

$$\int_Q F^p dx \leq 2a^{p-1} \int_Q F dx$$

for $a = \int_Q |Du|^2 dx$. Recalling the definition of F we finally obtain

$$(5.1) \quad \int_{Q_{1/2}} |Du|^{2p} dx \leq 2^{n+pn+1} \left(\int_Q |Du|^2 dx \right)^p.$$

Let $B_1 \subset\subset \Omega$. Observe that for any integer $h > 1$, $Q_{1/2}$ can be covered by h^n cubes of side length $1/h$. Hence, $B_{1/2}$ can be covered by N_h cubes $Q_{1/2h}(x_i)$ having non-empty intersection with $B_{1/2}$ and $N_h \leq h^n$. Using the rescaled analogue of the inequality (5.1) we get

$$\begin{aligned} \int_{B_{1/2}} |Du|^{2p} dx &\leq \frac{2^n}{\omega_n h^n} \sum_{i=1}^{N_h} \int_{Q_{1/2h}(x_i)} |Du|^{2p} dx \\ &\leq 2^{n+pn+1} \frac{2^n}{\omega_n h^n} \sum_{i=1}^{N_h} \left(\int_{Q_{1/h}(x_i)} |Du|^2 dx \right)^p \\ &\leq 2^{n+pn+1} \frac{2^n}{\omega_n h^n} \sum_{i=1}^{N_h} \left(\frac{h^n \omega_n}{2^n} \int_{B_1} |Du|^2 dx \right)^p \end{aligned}$$

provided $h > 4\sqrt{n}$, in which case $Q_{1/h}(x_i) \subset B_1$ for every $i = 1, \dots, N_h$. We may choose $h \leq 5\sqrt{n}$ and thus we get

$$\int_{B_{1/2}} |Du|^{2p} dx \leq 2^{2n+1} 5^{np} n^{np/2} \omega_n^{p-1} \left(\int_{B_1} |Du|^2 dx \right)^p.$$

□

ACKNOWLEDGMENT

The research of NF was partially supported by the 2008 ERC Advanced Grant 226234 ‘‘Analytic Techniques for Geometric and Functional Inequalities’’ and by the 2013 FiDiPro project ‘‘Quantitative Isoperimetric Inequalities’’. The research of VJ was supported by the Academy of Finland grant 268393. The authors want also to thank R. Kohn for pointing out to them the phenomenon of Taylor cones and G. De Philippis and A. Figalli for finding a mistake in the first version of the paper.

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