# ON THE HORIZONTAL MEAN CURVATURE FLOW FOR AXISYMMETRIC SURFACES IN THE HEISENBERG GROUP 

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#### Abstract

We study the horizontal mean curvature flow in the Heisenberg group by using the level-set method. We prove the uniqueness, existence and stability of axisymmetric viscosity solutions of the level-set equation. An explicit solution is given for the motion starting from a subelliptic sphere. We also give several properties of the level-set method and the mean curvature flow in the Heisenberg group


## 1. Introduction

We are interested in a family of compact hypersurfaces $\left\{\Gamma_{t}\right\}_{t \geq 0}$ in the Heisenberg group parametrized by time $t \geq 0$. The motion of the hypersurfaces is governed by the following law:

$$
\begin{equation*}
V_{H}=\kappa_{H}, \tag{1.1}
\end{equation*}
$$

where $V_{H}$ denotes its horizontal normal velocity and $\kappa_{H}$ stands for the horizontal mean curvature in the Heisenberg group. The geometric motion (1.1) is thus called horizontal mean curvature flow. The objective of this work is to investigate the evolution of the surface $\Gamma_{t}$ for $t>0$ for a general class of initial surface $\Gamma_{0}$.

We implement a version of the level-set method adapted to the Heisenberg group. Let us assume, for the moment, that $\Gamma_{t}$ is smooth for any $t \geq 0$. If there exists $u \in C^{2}(\mathcal{H} \times[0, \infty))$ such that

$$
\Gamma_{t}=\{p \in \mathcal{H}: u(p, t)=0\}
$$

for $t \geq 0$, then one may represent the horizontal normal velocity $V_{H}$ as

$$
V_{H}=\frac{u_{t}}{\left|\nabla_{H} u\right|}
$$

and the horizontal mean curvature $\kappa_{H}$ as

$$
\kappa_{H}=\operatorname{div}_{H}\left(\frac{\nabla_{H} u}{\left|\nabla_{H} u\right|}\right)=\frac{1}{\left|\nabla_{H} u\right|} \operatorname{tr}\left[\left(I-\frac{\nabla_{H} u \otimes \nabla_{H} u}{\left|\nabla_{H} u\right|^{2}}\right)\left(\nabla_{H}^{2} u\right)^{*}\right] .
$$

Here $u_{t}, \nabla_{H} u$ and $\left(\nabla_{H}^{2} u\right)^{*}$ respectively denote the derivative in $t$, the horizontal gradient and the (symmetrized) horizontal Hessian of $u$, and $\operatorname{div}_{H}$ is the horizontal divergence

[^0]operator. The horizontal gradient of $u$ is given by $\nabla_{H} u=\left(X_{1} u, X_{2} u\right)$, where
\[

$$
\begin{aligned}
& X_{1}=\frac{\partial}{\partial p_{1}}-\frac{p_{2}}{2} \frac{\partial}{\partial p_{3}} \\
& X_{2}=\frac{\partial}{\partial p_{2}}+\frac{p_{1}}{2} \frac{\partial}{\partial p_{3}} .
\end{aligned}
$$
\]

In order to understand the law of motion by curvature (1.1), it therefore suffices to solve (MCF) $u_{t}-\operatorname{tr}\left[\left(I-\frac{\nabla_{H} u \otimes \nabla_{H} u}{\left|\nabla_{H} u\right|^{2}}\right)\left(\nabla_{H}^{2} u\right)^{*}\right]=0$ in $\mathcal{H} \times(0, \infty)$, $u(p, 0)=u_{0}(p)$ in $\mathcal{H}$. with a given function $u_{0} \in C(\mathcal{H})$ satisfying

$$
\Gamma_{0}=\left\{p \in \mathcal{H}: u_{0}(p)=0\right\} .
$$

We refer the reader to [CGG, ES, G] for a detailed derivation of (MCF) in the Euclidean spaces and to [CDPT, CC] for the analogue in the Heisenberg group.

In this work, we aim to establish the uniqueness, existence and stability of the solutions of (MCF) that are spatially axisymmetric about the third coordinate axis. Namely, we are interested in the solutions $u$ satisfying

$$
\begin{equation*}
u\left(p_{1}, p_{2}, p_{3}, t\right)=u\left(p_{1}^{\prime}, p_{2}^{\prime}, p_{3}, t\right) \text { when }\left(p_{1}^{\prime}\right)^{2}+\left(p_{2}^{\prime}\right)^{2}=p_{1}^{2}+p_{2}^{2} . \tag{1.2}
\end{equation*}
$$

The symmetric structure of the functions is useful to obtain positive results. We thus consider our contribution as a first step in order to prove more general results. Consult [AAG, SS] for the results on motion by mean curvature for axisymmetric surfaces in the Euclidean spaces.

The symmetry with respect to the third axis in the Heisenberg group is not accidental. Indeed it is well known that this coordinate plays a key role in the Heisenberg group in several cases. In particular, we recall, for example, that $\left\{\left(0,0, p_{3}\right) \in \mathcal{H}: p_{3} \in \mathbb{R}\right\}$ is the center of the Heisenberg group and moreover the points along the $p_{3}$-axis correspond to conjugate points for the exponential map [Mo]. We warn the reader that, in general, our results do not apply to functions with different axes of symmetry.

Hereafter the property (1.4) is sometimes referred to as "spatial symmetry about the vertical axis" or simply as "axisymmetric".

Since the general regularity of $u$ is not known a priori, we discuss the problem in the framework of viscosity solutions [CIL]. As it is easily observed from the equation, a key difficulty lies at the characteristic set of the level set $\Gamma_{t}$, i.e., at the points where $\nabla_{H} u=0$.
1.1. Uniqueness. Even in the Euclidean case [CGG, ES, S, G], the proof of the comparison principle and the uniqueness of solutions for this type of degenerate equations need special techniques to deal with the characteristic set. The comparison principle we expect is as follows: for any upper semicontinuous subsolution $u$ and lower semicontinuous supersolution $v$ defined on $\mathcal{H} \times[0, \infty)$ satisfying $u(p, 0) \leq v(p, 0)$ for all $p \in \mathcal{H}$, we have $u(p, t) \leq v(p, t)$ for any $t \geq 0$. L. Capogna and G. Citti [CC] extended the results of [ES] and proved a comparison principle by excluding the characteristic points. Their comparison principle further required that (i) either $u$ or $v$ be uniformly continuous and (ii) the initial surfaces are completely separated in the horizontal directions, i.e., $u(p, 0) \leq v(q, 0)$ for all $p=\left(p_{1}, p_{2}, p_{3}\right), q=\left(q_{1}, q_{2}, q_{3}\right) \in \mathcal{H}$ such that $p_{i}=q_{i}$ for $i=1,2$. The general comparison principle, as stated above, remains an open question.

In this paper, we follow [CGG, G] and give a comparison principle without assuming those two conditions above but requiring that either $u$ or $v$ be axisymmetric. We also restrict ourselves to the case of compact surfaces for simplicity. The comparison theorem we present is as follows.

Theorem 1.1 (Comparison theorem). Let $u$ and $v$ be respectively an upper semicontinuous subsolution and a lower semicontinuous supersolution of

$$
u_{t}-\operatorname{tr}\left[\left(I-\frac{\nabla_{H} u \otimes \nabla_{H} u}{\left|\nabla_{H} u\right|^{2}}\right)\left(\nabla_{H}^{2} u\right)^{*}\right]=0
$$

in $\mathcal{H} \times(0, T)$ for any $T>0$. Assume that there is a compact set $K \subset \mathcal{H}$ and $a, b \in \mathbb{R}$ with $a \leq b$ such that $u(p, t)=a$ and $v(p, t)=b$ for all $p \in \mathcal{H} \backslash K$ and $t \in[0, T]$. Assume that either $u$ or $v$ is spatially axisymmetric about the vertical axis. If $u(p, 0) \leq v(p, 0)$ for all $p \in \mathcal{H}$, then $u \leq v$ for all $(x, t) \in \mathcal{H} \times[0, T]$.

The uniqueness of the axisymmetric solutions follows immediately from the theorem above. It is worth remarking that when showing comparison principles involving viscosity solutions, one usually needs to double the variables and maximize

$$
u(p, t)-v(q, s)-\frac{\phi(p, q)+|t-s|^{2}}{\varepsilon}
$$

where $\varepsilon>0, p, q \in \mathcal{H}, t, s \in[0, \infty)$ and $\phi$ is a smooth penalty function on $\mathcal{H} \times \mathcal{H}$, and argues by contradiction.

The typical choice of $\phi$ in the Euclidean spaces, as discussed in [CIL] and [G], is a quadratic function $\phi(x, y)=|x-y|^{2}$ usually or a quartic function $\phi(x, y)=|x-y|^{4}$ for mean curvature flow equation (for $x, y \in \mathbb{R}^{n}$ ). The advantages of these choices are:
(a) The derivatives of $\phi$ with respect to $x$ and $y$ are opposite, i.e., $\nabla_{x} \phi=-\nabla_{y} \phi$. We would plug these derivatives in the viscosity inequalities, since they serve as semidifferentials for the compared functions. This construction enables us to derive a contradiction.
(b) When discussing (mild) singular equations such as curvature flow equations, it will be convenient to have the second derivatives be 0 whenever the first derivatives are 0 , as in the case of $|x-y|^{4}$.
The analogue of the choice $|x-y|^{4}$ is not immediate in the Heisenberg group. Since the group multiplication is not commutative, the two natural options $f(p, q)=\left|q^{-1} \cdot p\right|^{4}$ and $g(p, q)=\left|p \cdot q^{-1}\right|^{4}$ are different. It seems that we have more options but it turns out that neither of them satisfies both conditions above. By direct calculation, we may find that $g$ fulfills the requirement (a) above but its derivatives do not satisfy (b). The function $f$ is good for our requirement (b) but unfortunately fails to have the property (a). Hence, the main difficulty of the uniqueness argument in the Heisenberg group consists in a wise choice of the penalty function $\phi$.

Our approach combines both choices $f$ and $g$. On one hand, we use $f$ to derive a relaxed definition (Definition 3.2) of solutions of (3.1), facilitating us to overcome the singularity. On the other hand, under the help of axial symmetry, we obtain the property (b) when employing $g$ type of penalty functions in the proof of the comparison principle. The symmetry plays an important role since it largely simplifies the structure of characteristic points; see [FLM2] for some geometric details.

Roughly speaking, when a smooth function $u(p, t)$ is spatially symmetric about the vertical axis, i.e., $u=u\left(r, p_{3}, t\right)$, where $r=\left(p_{1}^{2}+p_{2}^{2}\right)^{1 / 2}$, we get

$$
\begin{aligned}
& X_{1} u=\frac{p_{1}}{r} \frac{\partial}{\partial r} u-\frac{p_{2}}{2} \frac{\partial}{\partial p_{3}} u \\
& X_{2} u=\frac{p_{2}}{r} \frac{\partial}{\partial r} u+\frac{p_{1}}{2} \frac{\partial}{\partial p_{3}} u
\end{aligned}
$$

Then $\nabla_{H} u(p, t)=0$ implies that either $\partial u / \partial r=\partial u / \partial p_{3}=0$ or $p_{1}^{2}+p_{2}^{2}=0$. This observation enables us to obtain property (b) for a large power of the function $g$.

Our definition of viscosity solutions is actually an extension of that introduced in [CGG, G] to the Heisenberg group. In Section 3, we discuss the equivalence of this definition and the others.
1.2. Existence. Generally speaking, there are at least three possible approaches to get the existence of solutions of (MCF). One may follow [ES] to use the uniformly parabolic theory by considering a regularized equation $\mathrm{u}_{t}-\operatorname{tr}\left[\left(I-\frac{\nabla_{H} u \otimes \nabla_{H} u}{\left|\nabla_{H} u\right|^{2}+\varepsilon^{2}}\right)\left(\nabla_{H}^{2} u\right)^{*}\right]=0$ in $\mathcal{H} \times(0, \infty)$, $u(p, 0)=u_{0}(p)$ in $\mathcal{H}$. and take the limit of its solution as $\varepsilon \rightarrow 0$; see [CC] for results in the Carnot groups with this method. Another possible option is to employ Perron's method by considering the supremum of all subsolutions or the infimum of all supersolutions, as is shown in [CGG, G] for the Euclidean case. We refer to [I, CIL] for a general introduction of this method in the framework of viscosity solutions.

A third method for existence is based on the representation theorem involving optimal control or game theory, which recently generated a spur of activity. Consult the works [CSTV, KS1, KS2, MPR1, MPR2, PSSW, PS, ST] for the development of this approach to various equations in Euclidean spaces. For the mean curvature flow in the sub-Riemannian geometry, a stochastic control-based formulation analogous to [ST] is addressed in [DDR], where the authors found a solution via a suitable optimal stochastic control problem.

In this work, we adapt the deterministic game-theoretic approach of R. V. Kohn and S. Serfaty [KS1] to the Heisenberg group. For any given axisymmetric continuous function $u_{0}$, we set up a family of games, whose value functions $u^{\varepsilon}$ converge to the solution $u$ to the mean curvature flow equation. We not only get the existence of solutions but also obtain a game interpretation of the equation in the Heisenberg group. The proof is based on the dynamic programming principle, which can be regarded as a (nonlinear) semigroup. Our convergence theorem relies on the comparison principle given in Theorem 1.1. More precisely, taking the half relaxed limits, defined on $\mathcal{H} \times[0, \infty)$,

$$
\begin{align*}
\bar{u}(p, t): & =\limsup _{\varepsilon \rightarrow 0}^{*} u^{\varepsilon}(p, t) \\
& =\lim _{\delta \rightarrow 0} \sup \left\{u^{\varepsilon}(q, s): s \geq 0,0<\varepsilon<\delta,|p-q|+|t-s|<\delta\right\} \tag{1.3}
\end{align*}
$$

and

$$
\begin{align*}
\underline{u}(p, t): & =\liminf _{\varepsilon \rightarrow 0} u^{\varepsilon}(p, t) \\
& =\liminf _{\delta \rightarrow 0}\left\{u^{\varepsilon}(q, s): s \geq 0,0<\varepsilon<\delta,|p-q|+|t-s|<\delta\right\}, \tag{1.4}
\end{align*}
$$

we show that $\bar{u}$ and $\underline{u}$ are respectively a subsolution and a supersolution of (1.2) using the dynamic programming principle. We also show that $\bar{u}(p, 0) \leq u_{0}(p) \leq \underline{u}(p, 0)$ and that $u^{\varepsilon}, \bar{u}$ and $\underline{u}$ are spatially axisymmetric about the vertical axis. Our game approximation then follows immediately from the comparison theorem. See Section 5 for more details on the game setting and the existence theorem.

We discuss asymptotic mean value properties related to random tug-of-war games for $p$-harmonic functions on the Heisenberg group in [FLM1].
1.3. Stability and uniqueness of the evolution. We give a stability theorem, which is used to show that the equation (1.2) is invariant under the change of dependent variable. We prove that for any continuous function $\theta: \mathbb{R} \rightarrow \mathbb{R}$, the composition $\theta \circ u$ is a solution provided that $u$ is a solution. Note that this is clear if $\theta$ is smooth and strictly monotone, since the mean curvature flow equation is geometric and orientation-free; see [G] for more explanation. Our stability result is applied so as to weaken the regularity of $\theta$.

It follows from the invariance property that any axisymmetric evolution $\Gamma_{t}$ does not depend on the particular choice of $u_{0}$ but depends on $\Gamma_{0}$ only, which is important for the level-set method.
1.4. Evolution of spheres. Our uniqueness and existence results enable us to discuss motion by mean curvature with a variety of initial hypersurfaces including spheres, tori and other compact surfaces. We are particularly interested in the motion of a subelliptic sphere. It turns out that if $u_{0}$ is a defining function of the sphere centered at 0 with radius $r$, say

$$
u_{0}(p)=\min \left\{\left(p_{1}^{2}+p_{2}^{2}\right)^{2}+16 p_{3}^{2}-r^{4}, M\right\}
$$

with $p=\left(p_{1}, p_{2}, p_{3}\right) \in \mathcal{H}$ and $M>0$ large, then the unique solution of (MCF) is

$$
u(p, t)=\min \left\{\left(p_{1}^{2}+p_{2}^{2}\right)^{2}+12 t\left(p_{1}^{2}+p_{2}^{2}\right)+16 p_{3}^{2}+12 t^{2}-r^{4}, M\right\}
$$

for any $t \geq 0$. We need to truncate the initial function and the solution by a constant $M$ because all of our wellposedness results are for solutions that are constant outside a compact set. It is obvious that the zero level set $\Gamma_{t}$ of $u$ vanishes after time $t=r^{2} / \sqrt{12}$, which, by Theorem 1.1, indicates that all compact surfaces under the motion by horizontal mean curvature disappear in finite time.

To understand the asymptotic profile at the extinction time, we normalize the evolution $\Gamma_{t}$ initialized from the sphere and find that the normalized surface $\Gamma_{t} / \sqrt{r^{4}-12 t^{2}}$ converges to an ellipsoid given by the following equation:

$$
\sqrt{12} r^{2}\left(p_{1}^{2}+p_{2}^{2}\right)+16 p_{3}^{2}=1
$$

The asymptotic profile above depends on $r$, the size of the initial surface, which is quite different from the Euclidean case.

The paper is organized in the following way. We present an introduction in Section 2.1 about the Heisenberg group including calculations of some particular functions we will use later. In Section 3, we discuss various kinds of definitions of solution to (1.2). We propose a new definition and show its equivalence with the others. An explicit solution related to the evolution of a subelliptic sphere is given at the end of this section. The comparison principle, Theorem 1.1, is proved in Section 4. We establish the games and show the existence theorem in Section 5 . Section 6 is devoted to the stability results and Section 7 is intended to show further properties of the evolution including the uniqueness and finite extinction with the interesting asymptotic profile.

## 2. Tools from Calculus in $\mathcal{H}$

Good references for this section are the course notes [M] and the monograph [CDPT].
2.1. Preliminaries. Recall the that Heisenberg group $\mathcal{H}$ is $\mathbb{R}^{3}$ endowed with the noncommutative group multiplication

$$
\left(p_{1}, p_{2}, p_{3}\right) \cdot\left(q_{1}, q_{2}, q_{3}\right)=\left(p_{1}+q_{1}, p_{2}+q_{2}, p_{3}+q_{3}+\frac{1}{2}\left(p_{1} q_{2}-q_{1} p_{2}\right)\right)
$$

for all $p=\left(p_{1}, p_{2}, p_{3}\right)$ and $q=\left(q_{1}, q_{2} \cdot q_{3}\right)$ in $\mathcal{H}$. The Haar measure if $\mathcal{H}$ is the usual Lebesgue measure in $\mathbb{R}^{3}$. The Korányi gauge is given by

$$
|p|=\left(\left(p_{1}^{2}+p_{2}^{2}\right)^{2}+16 p_{3}^{2}\right)^{1 / 4}
$$

and the left-invariant Korányi or gauge metric is

$$
d(p, q)=\left|q^{-1} \cdot p\right|
$$

The Korányi ball of radius $r>0$ centered at $p$ is

$$
B_{r}(p):=\{q \in \mathcal{H}: d(p, q)<r\} .
$$

The Lie Algebra of $\mathcal{H}$ is generated by the left-invariant vector fields

$$
\begin{aligned}
X_{1} & =\frac{\partial}{\partial p_{1}}-\frac{p_{2}}{2} \frac{\partial}{\partial p_{3}} \\
X_{2} & =\frac{\partial}{\partial p_{2}}+\frac{p_{1}}{2} \frac{\partial}{\partial p_{3}} \\
X_{3} & =\frac{\partial}{\partial p_{3}}
\end{aligned}
$$

One may easily verify the commuting relation $X_{3}=\left[X_{1}, X_{2}\right]=X_{1} X_{2}-X_{2} X_{1}$.
For any smooth real valued function $u$ defined in an open subset of $\mathcal{H}$, the horizontal gradient of $u$ is

$$
\nabla_{H} u=\left(X_{1} u, X_{2} u\right)
$$

while the complete gradient of $u$ is

$$
\nabla u=\left(X_{1} u, X_{2} u, X_{3} u\right)
$$

For further details about the relation between sub-Riemannian metrics in Carnot group and Riemaniann metrics see [AFM].

The symmetrized second horizontal Hessian $\left(\nabla_{H}^{2} u\right)^{*}$ is the $2 \times 2$ symmetry matrix given by

$$
\left(\nabla_{H}^{2} u\right)^{*}:=\left(\begin{array}{cc}
X_{1}^{2} u & \left(X_{1} X_{2} u+X_{2} X_{1} u\right) / 2 \\
\left(X_{1} X_{2} u+X_{2} X_{1} u\right) / 2 & X_{2}^{2} u
\end{array}\right)
$$

We will also consider the symmetrized complete Hessian $\left(\nabla^{2} u\right)^{*}$ defined as the $3 \times 3$ symmetric matrix

$$
\left(\nabla^{2} u\right)^{*}:=\left(\begin{array}{ccc}
X_{1}^{2} u & \left(X_{1} X_{2} u+X_{2} X_{1} u\right) / 2 & \left(X_{1} X_{3} u+X_{3} X_{1} u\right) / 2 \\
\left(X_{1} X_{2} u+X_{2} X_{1} u\right) / 2 & X_{2}^{2} u & \left(X_{2} X_{3} u+X_{3} X_{2} u\right) / 2 \\
\left(X_{1} X_{3} u+X_{3} X_{1} u\right) / 2 & \left(X_{2} X_{3} u+X_{3} X_{2} u\right) / 2 & X_{3}^{2} u
\end{array}\right)
$$

2.2. Derivatives of auxiliary functions. Here we include several basic calculations for some test functions related to the Korányi distance, which will be used in the proof of comparison theorem for generalized horizontal mean curvature flow.

We are interested in the first and second horizontal derivatives of

$$
\begin{aligned}
f(p, q): & =d(p, q)^{4} \\
& =\left(\left(p_{1}-q_{1}\right)^{2}+\left(p_{2}-q_{2}\right)^{2}\right)^{2}+16\left(p_{3}-q_{3}-\frac{1}{2} q_{1} p_{2}+\frac{1}{2} q_{2} p_{1}\right)^{2}
\end{aligned}
$$

We use the super index $p$ to denote derivatives with respect to the $p$ variable and follow the same convention for derivatives with respect to $q$.

Let us record the results of our calculation:

$$
\begin{align*}
& X_{1}^{p} f(p, q)=4\left(\left(p_{1}-q_{1}\right)^{2}+\left(p_{2}-q_{2}\right)^{2}\right)\left(p_{1}-q_{1}\right) \\
& \quad-16\left(p_{2}-q_{2}\right)\left(p_{3}-q_{3}+\frac{1}{2}\left(q_{2} p_{1}-q_{1} p_{2}\right)\right)  \tag{2.1}\\
& \begin{aligned}
X_{2}^{p} f(p, q)=4\left(\left(p_{1}-q_{1}\right)^{2}+\left(p_{2}-q_{2}\right)^{2}\right)\left(p_{2}-q_{2}\right)
\end{aligned} \\
& \quad+16\left(p_{1}-q_{1}\right)\left(p_{3}-q_{3}+\frac{1}{2}\left(q_{2} p_{1}-q_{1} p_{2}\right)\right)  \tag{2.2}\\
& X_{1}^{q} f(p, q)=-4\left(\left(p_{1}-q_{1}\right)^{2}+\left(p_{2}-q_{2}\right)^{2}\right)\left(p_{1}-q_{1}\right) \\
&  \tag{2.3}\\
& -16\left(p_{2}-q_{2}\right)\left(p_{3}-q_{3}+\frac{1}{2}\left(q_{2} p_{1}-q_{1} p_{2}\right)\right)
\end{align*}
$$

$$
\begin{align*}
& X_{2}^{q} f(p, q)=-4\left(\left(p_{1}-q_{1}\right)^{2}+\left(p_{2}-q_{2}\right)^{2}\right)\left(p_{2}-q_{2}\right) \\
&+16\left(p_{1}-q_{1}\right)\left(p_{3}-q_{3}+\frac{1}{2}\left(q_{2} p_{1}-q_{1} p_{2}\right)\right) \tag{2.4}
\end{align*}
$$

It is clear that in general $\nabla_{H}^{p} f(p, q) \neq-\nabla_{H}^{q} f(p, q)$, which is not the case in the Euclidean case. But the following Euclidean property still holds here.
Proposition 2.1. If either $\nabla_{H}^{p}\left(\left|q^{-1} \cdot p\right|^{4}\right)=0$ or $\nabla_{H}^{q}\left(\left|q^{-1} \cdot p\right|^{4}\right)=0$, then the horizontal components of $p$ and $q$ are equal, i.e., $p_{1}=q_{1}$ and $p_{2}=q_{2}$.

Proof. Set

$$
\begin{aligned}
A & :=4\left(\left(p_{1}-q_{1}\right)^{2}+\left(p_{2}-q_{2}\right)^{2}\right) \\
B & :=16\left(p_{3}-q_{3}+\frac{1}{2}\left(q_{2} p_{1}-q_{1} p_{2}\right)\right) .
\end{aligned}
$$

When $\nabla_{H}^{p}\left(\left|q^{-1} \cdot p\right|^{4}\right)=0$, the calculations (2.1) and (2.2) read

$$
\left\{\begin{array}{r}
A\left(p_{1}-q_{1}\right)-B\left(p_{2}-q_{2}\right)=0  \tag{2.5}\\
B\left(p_{1}-q_{1}\right)+A\left(p_{2}-q_{2}\right)=0
\end{array}\right.
$$

with $\operatorname{det}\left(\begin{array}{cc}A & -B \\ B & A\end{array}\right)=A^{2}+B^{2} \geq 0$. Since $A^{2}+B^{2}=0$ implies that $p_{i}=q_{i}$ for $i=1,2$, the desired result is trivial if $A^{2}+B^{2}=0$. If the determinant is not zero, then we also obtain $q_{1}=p_{1}$ and $q_{2}=p_{2}$ by solving the linear system (2.5). The same argument applies to the case when $\nabla_{H}^{q}\left(\left|q^{-1} \cdot p\right|^{4}\right)=0$.

We next calculate the second horizontal derivatives.

$$
\begin{gather*}
X_{1}^{2, p} f(p, q)=X_{1}^{2, q} f(p, q)=12\left(p_{1}-q_{1}\right)^{2}+12\left(p_{2}-q_{2}\right)^{2}  \tag{2.6}\\
X_{2}^{2, p} f(p, q)=X_{2}^{2, q} f(p, q)=12\left(p_{2}-q_{2}\right)^{2}+12\left(p_{1}-q_{1}\right)^{2}  \tag{2.7}\\
X_{2}^{p} X_{1}^{p} f(p, q)=X_{1}^{q} X_{2}^{q} f(p, q)=-16\left(p_{3}-q_{3}+\frac{1}{2}\left(q_{2} p_{1}-q_{1} p_{2}\right)\right)=-B  \tag{2.8}\\
X_{1}^{p} X_{2}^{p} f(p, q)=X_{2}^{q} X_{1}^{q} f(p, q)=16\left(p_{3}-q_{3}+\frac{1}{2}\left(q_{2} p_{1}-q_{1} p_{2}\right)\right)=B \tag{2.9}
\end{gather*}
$$

It is clear that

$$
\frac{1}{2}\left(X_{1}^{p} X_{2}^{p} f+X_{2}^{p} X_{1}^{p} f\right)=\frac{1}{2}\left(X_{1}^{q} X_{2}^{q} f+X_{2}^{q} X_{1}^{q} f\right)=0
$$

For later use, let us investigate the derivatives of another function. Take

$$
\begin{align*}
& g(p, q):=\left|p \cdot q^{-1}\right|^{4} \\
& =\left(\left(p_{1}-q_{1}\right)^{2}+\left(p_{2}-q_{2}\right)^{2}\right)^{2}+16\left(p_{3}-q_{3}-\frac{1}{2} p_{1} q_{2}+\frac{1}{2} p_{2} q_{1}\right)^{2} \tag{2.10}
\end{align*}
$$

Then

$$
\begin{align*}
X_{1}^{p} g(p, q)=4\left(\left(p_{1}-q_{1}\right)^{2}\right. & \left.+\left(p_{2}-q_{2}\right)^{2}\right)\left(p_{1}-q_{1}\right) \\
& -16\left(p_{2}+q_{2}\right)\left(p_{3}-q_{3}-\frac{1}{2} p_{1} q_{2}+\frac{1}{2} p_{2} q_{1}\right) \tag{2.11}
\end{align*}
$$

$$
\begin{align*}
X_{2}^{p} g(p, q)=4\left(\left(p_{1}-q_{1}\right)^{2}\right. & \left.+\left(p_{2}-q_{2}\right)^{2}\right)\left(p_{2}-q_{2}\right) \\
& +16\left(p_{1}+q_{1}\right)\left(p_{3}-q_{3}-\frac{1}{2} p_{1} q_{2}+\frac{1}{2} p_{2} q_{1}\right) ;  \tag{2.12}\\
X_{1}^{q} g(p, q)=-4\left(\left(p_{1}-q_{1}\right)^{2}\right. & \left.+\left(p_{2}-q_{2}\right)^{2}\right)\left(p_{1}-q_{1}\right) \\
& +16\left(p_{2}+q_{2}\right)\left(p_{3}-q_{3}-\frac{1}{2} p_{1} q_{2}+\frac{1}{2} p_{2} q_{1}\right) ;  \tag{2.13}\\
X_{2}^{q} g(p, q)=-4\left(\left(p_{1}-q_{1}\right)^{2}\right. & \left.+\left(p_{2}-q_{2}\right)^{2}\right)\left(p_{2}-q_{2}\right) \\
& -16\left(p_{1}+q_{1}\right)\left(p_{3}-q_{3}-\frac{1}{2} p_{1} q_{2}+\frac{1}{2} p_{2} q_{1}\right) .
\end{align*}
$$

Remark 2.1. In this case, we do have $\nabla_{H}^{p} g(p, q)=-\nabla_{H}^{q} g(p, q)$. But the property in Proposition 2.1 does not hold in general.

The second derivatives are given below.

$$
\begin{align*}
& X_{1}^{2, p} g(p, q)=X_{1}^{2, q} g(p, q)= 12\left(p_{1}-q_{1}\right)^{2}+4\left(p_{2}-q_{2}\right)^{2}+8\left(p_{2}+q_{2}\right)^{2} ;  \tag{2.15}\\
& X_{2}^{2, p} g(p, q)=X_{2}^{2, q} g(p, q)= 4\left(p_{1}-q_{1}\right)^{2}+12\left(p_{2}-q_{2}\right)^{2}+8\left(p_{1}+q_{1}\right)^{2} ;  \tag{2.16}\\
& X_{1}^{p} X_{2}^{p} g(p, q)=X_{2}^{q} X_{1}^{q} g(p, q)=8\left(p_{1}-q_{1}\right)\left(p_{2}-q_{2}\right)-8\left(p_{1}+q_{1}\right)\left(p_{2}+q_{2}\right) \\
&+16\left(p_{3}-q_{3}-\frac{1}{2} p_{1} q_{2}+\frac{1}{2} p_{2} q_{1}\right) ;  \tag{2.17}\\
& \begin{aligned}
X_{2}^{p} X_{1}^{p} g(p, q)=X_{1}^{q} X_{2}^{q} g(p, q)= & 8\left(p_{1}-q_{1}\right)\left(p_{2}-q_{2}\right)-8\left(p_{1}+q_{1}\right)\left(p_{2}+q_{2}\right) \\
& -16\left(p_{3}-q_{3}-\frac{1}{2} p_{1} q_{2}+\frac{1}{2} p_{2} q_{1}\right) ; \\
\frac{1}{2}\left(X_{1}^{p} X_{2}^{p} g+X_{2}^{p} X_{1}^{p} g\right)= & \frac{1}{2}\left(X_{1}^{q} X_{2}^{q} g+X_{2}^{q} X_{1}^{q} g\right) \\
= & 8\left(p_{1}-q_{1}\right)\left(p_{2}-q_{2}\right)-8\left(p_{1}+q_{1}\right)\left(p_{2}+q_{2}\right) .
\end{aligned}
\end{align*}
$$

2.3. Extrema in the Heisenberg group. As $|p|^{2} \approx p_{1}^{2}+p_{2}^{2}+\left|p_{3}\right|$ in Heisenberg group, the Taylor formula reads

$$
\begin{equation*}
u(p)=u(\hat{p})+\left\langle\hat{p}^{-1} \cdot p, \nabla u(\hat{p})\right\rangle+\frac{1}{2}\left\langle\left(\nabla_{H}^{2} u\right)^{*}(\hat{p}) h, h\right\rangle+o\left(\left|\hat{p}^{-1} \cdot p\right|^{2}\right), \tag{2.20}
\end{equation*}
$$

where $h=\left(p_{1}-\hat{p}_{1}, p_{2}-\hat{p}_{2}\right)$ is the horizontal projection of $\hat{p}^{-1} \cdot p$.
The following proposition follows easily from the Euclidean analog.
Proposition 2.2 (Maxima on Heisenberg group). Suppose $\mathcal{O}$ is an open subset of $\mathcal{H}$. Let $u \in C^{2}(\mathcal{O})$ and $\hat{p} \in \mathcal{O}$. If $u(p) \leq u(\hat{p})$ for all $p \in \mathcal{O}$, then $\nabla u(\hat{p})=0$ and $\left(\nabla_{H}^{2} u\right)^{*}(\hat{p}) \leq 0$.

Analogously, for minima we have that if $u(p) \geq u(\hat{p})$ for all $p \in \mathcal{O}$, then $\nabla u(\hat{p})=0$ and $\left(\nabla_{H}^{2} u\right)^{*}(\hat{p}) \geq 0$.

## 3. Definitions of solutions

3.1. General definitions. For a vector $\eta \in \mathbb{R}^{2}$ and a $2 \times 2$ symmetric matrix $Y \in \mathbf{S}^{2}$ we define

$$
F(\eta, Y)=-\operatorname{tr}\left(\left(I-\frac{\eta \otimes \eta}{|\eta|^{2}}\right) Y\right)
$$

In any open subset $\mathcal{O} \subset \mathcal{H} \times(0, \infty)$ the mean curvature flow equation

$$
\begin{equation*}
u_{t}-\operatorname{tr}\left[\left(I-\frac{\nabla_{H} u \otimes \nabla_{H} u}{\left|\nabla_{H} u\right|^{2}}\right)\left(\nabla_{H}^{2} u\right)^{*}\right]=0 \text { in } \mathcal{O} \tag{3.1}
\end{equation*}
$$

can be written as

$$
u_{t}+F\left(\nabla_{H} u,\left(\nabla_{H}^{2} u\right)^{*}\right)=0 \text { in } \mathcal{O}
$$

We next define the semicontinuous envelopes in the following way: for any function $h$ defined on a set $\mathcal{O}$ of a metric space $\mathcal{M}$ with values in $\mathbb{R} \cup\{ \pm \infty\}$, we take

$$
\begin{equation*}
h^{\star}(x)=\lim _{r \rightarrow 0} \sup \left\{h(y): y \in \mathcal{O} \cap B_{r}(x)\right\} \tag{3.2}
\end{equation*}
$$

and

$$
\begin{equation*}
h_{\star}(x)=\lim _{r \rightarrow 0} \inf \left\{h(y): y \in \mathcal{O} \cap B_{r}(x)\right\} \tag{3.3}
\end{equation*}
$$

for any $x \in \overline{\mathcal{O}}$, where $B_{r}(x)$ denotes the ball with radius $r>0$ centered at $x$. It is easily seen that

$$
\begin{gathered}
F^{\star}(0,0)=F_{\star}(0,0)=0 \\
F^{\star}(\eta, X)=F_{\star}(\eta, X)=F(\eta, X) \text { for all }(\eta, X) \in \mathbb{R}^{2} \backslash\{0\} \times \mathbf{S}^{2}
\end{gathered}
$$

One type of definition of viscosity solutions of (3.1) is as follows.
Definition 3.1. An upper (resp., lower) semicontinuous function $u$ defined on $\mathcal{O} \subset \mathcal{H} \times$ $(0, \infty)$ is a subsolution (resp., supersolution) of (3.1) if
(i) $u<\infty$ (resp., $u>-\infty$ ) in $\mathcal{O}$;
(ii) for any smooth function $\phi$ such that

$$
\begin{gathered}
\max _{\mathcal{O}} u-\phi=(u-\phi)(\hat{p}, \hat{t}) \\
\left(\text { resp., } \min _{\mathcal{O}} u-\phi=(u-\phi)(\hat{p}, \hat{t}),\right)
\end{gathered}
$$

it satisfies

$$
\begin{gathered}
\phi_{t}+F_{\star}\left(\nabla_{H} \phi,\left(\nabla_{H}^{2} \phi\right)^{*}\right) \leq 0 \text { at }(\hat{p}, \hat{t}) \\
\text { (resp., } \left.\phi_{t}+F^{\star}\left(\nabla_{H} \phi,\left(\nabla_{H}^{2} \phi\right)^{*}\right) \geq 0 \text { at }(\hat{p}, \hat{t})\right) .
\end{gathered}
$$

A function $u$ is called a solution of (3.1) if it is both a subsolution and a supersolution.
We now propose another definition for the horizontal mean curvature flow equation following Giga [G].

Definition 3.2. An upper (resp., lower) semicontinuous function $u$ defined on $\mathcal{O} \subset \mathcal{H} \times$ $(0, \infty)$ is a subsolution (resp., supersolution) of (3.1) if
(i) $u<\infty$ (resp., $u>-\infty$ ) in $\mathcal{O}$;
(ii) for any smooth function $\phi$ such that

$$
\begin{gathered}
\max _{\mathcal{O}} u-\phi=(u-\phi)(\hat{p}, \hat{t}) \\
\text { (resp., } \left.\min _{\mathcal{O}} u-\phi=(u-\phi)(\hat{p}, \hat{t}),\right)
\end{gathered}
$$

it satisfies

$$
\begin{aligned}
& \quad \phi_{t}+F\left(\nabla_{H} \phi,\left(\nabla_{H}^{2} \phi\right)^{*}\right) \leq 0 \text { at }(\hat{p}, \hat{t}) \\
& \left(\text { resp. }, \phi_{t}+F\left(\nabla_{H} \phi,\left(\nabla_{H}^{2} \phi\right)^{*}\right) \geq 0 \text { at }(\hat{p}, \hat{t}),\right)
\end{aligned}
$$

when $\nabla_{H} \phi(\hat{p}, \hat{t}) \neq 0$ and

$$
\begin{gathered}
\phi_{t}(\hat{p}, \hat{t}) \leq 0 \\
\left(\operatorname{resp} ., \phi_{t}(\hat{p}, \hat{t}) \geq 0,\right)
\end{gathered}
$$

when $\nabla_{H} \phi(\hat{p}, \hat{t})=0$ and $\left(\nabla_{H}^{2} \phi\right)^{*}(\hat{p}, \hat{t})=0$.
A function $u$ is called a solution of (3.1) if it is both a subsolution and a supersolution.
Remark 3.1. One may replace the maximum (resp., minimum) in condition (ii) of the above definitions with a strict maximum by adding a positive (resp., negative) smooth gauge to $\phi$.

The definition using subelliptic semijets is as follows.
Definition 3.3. An upper (resp., lower) semicontinuous function $u$ defined on $\mathcal{O} \subset \mathcal{H} \times$ $(0, \infty)$ is a subsolution (resp., supersolution) of (3.1) if
(1) $u<\infty$ (resp., $u>-\infty$ ) in $\mathcal{O}$;
(2) for any $(\tau, \eta, \mathcal{X}) \in \bar{J}_{H}^{2,+} u(\hat{p}, \hat{t})$ (resp., $\left.(\tau, \eta, \mathcal{X}) \in \bar{J}_{H}^{2,-} u(\hat{p}, \hat{t})\right)$ with $(\hat{p}, \hat{t}) \in \mathcal{O}$, we have

$$
\begin{gathered}
\phi_{t}+F_{\star}\left(\nabla_{H} \phi,\left(\nabla_{H}^{2} \phi\right)^{*}\right) \leq 0 \text { at }(\hat{p}, \hat{t}), \\
\left(\text { resp. }, \phi_{t}+F^{\star}\left(\nabla_{H} \phi,\left(\nabla_{H}^{2} \phi\right)^{*}\right) \geq 0 \text { at }(\hat{p}, \hat{t}),\right. \text { ) }
\end{gathered}
$$

A function $u$ is called a solution of (3.1) if it is both a subsolution and a supersolution.
It is not hard to see that Definition 3.3 is equivalent to Definition 3.1. Roughly speaking, in Definition 3.2 and Definition 3.3 we restrict the test function space to the following

$$
\mathcal{A}_{0}=\left\{\phi \in C^{\infty}(\mathcal{H}): \nabla_{H} \phi(p)=0 \text { implies }\left(\nabla_{H}^{2}\right)^{*} \phi(p)=0\right\}
$$

The next result, which is actually a variant of [G, Proposition 2.2.8] for the Heisenberg group, indicates the equivalence between this new definition and the known one in spite of the restriction on the test functions.
Proposition 3.1 (Equivalence of definitions). An upper (resp., lower) semicontinuous function $u: \mathcal{O} \rightarrow \mathbb{R}$ is a subsolution (resp., supersolution) of (3.1) defined as in Definition 3.2 (in $\mathcal{O}$ ) if and only if it is a subsolution (resp., superolution) in $\mathcal{O}$ in the sense of Definition 3.1.
Proof. It is obvious that Definition 3.2 is a relaxation of Definition 3.1. We prove the reverse implication only for subsolutions. The statement for supersolutions can be proved similarly. Suppose there are a smooth function $\phi$ and $(\hat{p}, \hat{t}) \in \mathcal{O}$ such that

$$
\max _{\mathcal{O}}(u-\phi)=(u-\phi)(\hat{p}, \hat{t})
$$

By usual modification in the definition of viscosity solutions, we may assume it is a strict maximum. We construct

$$
\Psi_{\varepsilon}(p, q, t):=u(p, t)-\frac{1}{\varepsilon}\left|q^{-1} \cdot p\right|^{4}-\phi(q, t)
$$

It is clear that

$$
\Psi^{*}(p, q, t):=\limsup _{\varepsilon \rightarrow 0}^{*} \Psi_{\varepsilon}(p, q, t)= \begin{cases}u(p, t)-\phi(p, t) & \text { if } p=q \\ -\infty & \text { if } p \neq q\end{cases}
$$

attains a strict maximum at $(\hat{p}, \hat{p}, \hat{t})$. By the convergence of maximizers ([G, Lemma 2.2.5]), we may take $p^{\varepsilon}, q^{\varepsilon}, t^{\varepsilon}$ converging to $\hat{p}, \hat{p}, \hat{t}$ respectively as $\varepsilon \rightarrow 0$ such that $\Psi_{\varepsilon}$ attains a maximum at $\left(p^{\varepsilon}, q^{\varepsilon}, t^{\varepsilon}\right)$. It follows that $q \mapsto-\frac{1}{\varepsilon}\left|q^{-1} \cdot p^{\varepsilon}\right|^{4}-\phi(q, t)$ has a maximum at $q^{\varepsilon}$, which, by Proposition 2.2 , implies that

$$
\begin{equation*}
-\frac{1}{\varepsilon} \nabla_{H}^{q} f\left(p^{\varepsilon}, q^{\varepsilon}\right)=\nabla_{H} \phi\left(q^{\varepsilon}, t^{\varepsilon}\right) \tag{3.4}
\end{equation*}
$$

$$
\begin{equation*}
-\frac{1}{\varepsilon}\left(\nabla_{H}^{2, q} f\right)^{*}\left(p^{\varepsilon}, q^{\varepsilon}\right) \leq\left(\nabla_{H}^{2} \phi\right)^{*}\left(q^{\varepsilon}, t^{\varepsilon}\right), \tag{3.5}
\end{equation*}
$$

where $f(p, q)=\left|q^{-1} \cdot p\right|^{4}$.
We next discuss the following two cases.
Case A. $\nabla_{H} \phi\left(q^{\varepsilon}, t^{\varepsilon}\right) \neq 0$ for a subsequence of $\varepsilon \rightarrow 0$. (We still use $\varepsilon$ to denote the subsequence.)
Since the maximality of $\Psi$ at $\left(p^{\varepsilon}, q^{\varepsilon}, t^{\varepsilon}\right)$ implies that

$$
(p, t) \mapsto u(p, t)-\frac{1}{\varepsilon} f\left(p^{\varepsilon}, q^{\varepsilon}\right)-\phi\left(p \cdot\left(p^{\varepsilon}\right)^{-1} \cdot q^{\varepsilon}, t\right)
$$

attains a maximum at $\left(p^{\varepsilon}, t^{\varepsilon}\right) \in \mathcal{O}$. Denote $\phi^{\varepsilon}(p, t)=\phi\left(p \cdot\left(p^{\varepsilon}\right)^{-1} \cdot q^{\varepsilon}, t\right)$. We apply Definition 3.2 to get

$$
\begin{equation*}
\phi_{t}+F\left(\nabla_{H} \phi^{\varepsilon},\left(\nabla_{H}^{2} \phi^{\varepsilon}\right)^{*}\right) \leq 0 \text { at }\left(p^{\varepsilon}, t^{\varepsilon}\right) \tag{3.6}
\end{equation*}
$$

Since the derivative of the right multiplication tends to 0 as $\varepsilon \rightarrow 0$ and its second derivatives are 0 , we have

$$
\nabla_{H} \phi^{\varepsilon}\left(p^{\varepsilon}, t^{\varepsilon}\right) \rightarrow \nabla_{H} \phi(\hat{p}, \hat{t}) \text { and }\left(\nabla_{H}^{2} \phi^{\varepsilon}\right)^{*}\left(p^{\varepsilon}, t^{\varepsilon}\right) \rightarrow\left(\nabla_{H}^{2} \phi\right)^{*}(\hat{p}, \hat{t}) \text { as } \varepsilon \rightarrow 0 \text {. }
$$

It follows immediately that

$$
\phi_{t}+F_{\star}\left(\nabla_{H} \phi,\left(\nabla_{H}^{2} \phi\right)^{*}\right) \leq 0 \text { at }(\hat{p}, \hat{t}) .
$$

Case B. $\nabla_{H} \phi\left(q^{\varepsilon}, t^{\varepsilon}\right)=0$ for all sufficiently small $\varepsilon>0$.
It follows from (3.4) that $\nabla_{H}^{q} f\left(p^{\varepsilon}, q^{\varepsilon}\right)=0$, which by Proposition 2.1 yields that

$$
\begin{equation*}
p_{i}^{\varepsilon}=q_{i}^{\varepsilon} \text { for } i=1,2 . \tag{3.7}
\end{equation*}
$$

In terms of (2.1)-(2.2) and (2.6)-(2.9), we have

$$
\begin{equation*}
\nabla_{H}^{p} f\left(p^{\varepsilon}, q^{\varepsilon}\right)=0 \text { and } \nabla_{H}^{2, p} f\left(p^{\varepsilon}, q^{\varepsilon}\right)=0 . \tag{3.8}
\end{equation*}
$$

Since $\left(p^{\varepsilon}, t^{\varepsilon}\right)$ is a maximizer of

$$
(p, t) \mapsto u(p, t)-\frac{1}{\varepsilon} f\left(p, q^{\varepsilon}\right)-\phi\left(q^{\varepsilon}, t\right),
$$

applying Definition 3.2 and sending the limit, we obtain

$$
\begin{equation*}
\phi_{t}(\hat{p}, \hat{t}) \leq 0, \tag{3.9}
\end{equation*}
$$

On the other hand, by passing to the limit in (3.4) and (3.5), we have

$$
\begin{equation*}
\nabla_{H} \phi(\hat{p}, \hat{t})=0 \tag{3.10}
\end{equation*}
$$

and

$$
\begin{equation*}
\nabla_{H}^{2} \phi(\hat{p}, \hat{t}) \geq 0 . \tag{3.11}
\end{equation*}
$$

By (3.10), (3.9) is equivalent to

$$
\phi_{t}(\hat{p}, \hat{t})+F_{\star}\left(\nabla_{H} \phi(\hat{p}, \hat{t}), 0\right) \leq 0,
$$

which, thanks to (3.11) and the ellipticity of $F$, implies that

$$
\phi_{t}(\hat{p}, \hat{t})+F_{\star}\left(\nabla_{H} \phi(\hat{p}, \hat{t}), \nabla_{H}^{2} \phi(\hat{p}, \hat{t})\right) \leq 0 .
$$

3.2. An explicit solution. We provide an example of solutions of (3.1) when the initial value is the fourth power of a smooth gauge of the Heisenberg group. We can actually express a solution explicitly.

Proposition 3.2. For any $p=\left(p_{1}, p_{2}, p_{3}\right) \in \mathcal{H}$, let

$$
\begin{equation*}
G(p)=|p|^{4}=\left(p_{1}^{2}+p_{2}^{2}\right)^{2}+16 p_{3}^{2} . \tag{3.12}
\end{equation*}
$$

Then

$$
\begin{equation*}
w(p, t)=\left(p_{1}^{2}+p_{2}^{2}\right)^{2}+12 t\left(p_{1}^{2}+p_{2}^{2}\right)+16 p_{3}^{2}+12 t^{2} \tag{3.13}
\end{equation*}
$$

is a continuous solution of (1.2) and $w(p, 0)=G(p)$.
Proof. Since $w$ is smooth, the proof is based on a straightforward calculation of the first derivatives of $w$

$$
\begin{align*}
& w_{t}=12\left(p_{1}^{2}+p_{2}^{2}\right)+24 t, \\
& X_{1} w=K p_{1}-16 p_{2} p_{3}, \quad X_{2} w=K p_{2}+16 p_{1} p_{3}, \tag{3.14}
\end{align*}
$$

where $K:=4\left(p_{1}^{2}+p_{2}^{2}\right)+24 t$ and the second derivatives

$$
\begin{align*}
& X_{1}^{2} w=X_{2}^{2} w=12 p_{1}^{2}+12 p_{2}^{2}+24 t, \\
& X_{1} X_{2} w=16 p_{3}, X_{2} X_{1} w=-16 p_{3}, \\
& \left(\nabla_{H}^{2} w\right)^{*}=\left(\begin{array}{cc}
12 p_{1}^{2}+12 p_{2}^{2}+24 t & 0 \\
0 & 12 p_{1}^{2}+12 p_{2}^{2}+24 t
\end{array}\right) . \tag{3.15}
\end{align*}
$$

Noting that $\left(\nabla_{H}^{2} w\right)^{*}$ is constant multiple of the identity, we easily conclude from our calculation that

$$
\begin{aligned}
& F^{\star}\left(\nabla_{H} w, \nabla_{H}^{2} w\right)=F_{\star}\left(\nabla_{H} w, \nabla_{H}^{2} w\right) \\
= & \operatorname{tr}\left[\left(I-\frac{\nabla_{H} w \otimes \nabla_{H} w}{\left|\nabla_{H} w\right|^{2}}\right)\left(\nabla_{H}^{2} w\right)^{*}\right]=12\left(p_{1}^{2}+p_{2}^{2}\right)+24 t=w_{t},
\end{aligned}
$$

which means that $w$ satisfies (3.1) by Definition 3.1.
Remark 3.2. There is another way to understand that $w$ is a solution of (3.1) by adopting Definition 3.2 when $\nabla_{H} w=0$ at $(p, t) \in \mathcal{H} \times(0, \infty)$. If $\nabla_{H} w(p, t)=0$, we have $p_{1}=p_{2}=0$ by solving a linear system

$$
\left(\begin{array}{cc}
K & -16 p_{3} \\
16 p_{3} & K
\end{array}\right)\binom{p_{1}}{p_{2}}=\binom{0}{0}
$$

with

$$
\operatorname{det}\left(\begin{array}{cc}
K & -16 p_{3} \\
16 p_{3} & K
\end{array}\right)=K^{2}+16 p_{3}^{2}>0 .
$$

In addition,

$$
\left(\nabla_{H}^{2} w\right)^{*}=\left(\begin{array}{cc}
24 t & 0 \\
0 & 24 t
\end{array}\right) .
$$

Note that, by Proposition 2.2, it is not possible to take a smooth function $\phi$ touching $w$ from above at ( $p, t$ ) with

$$
\begin{equation*}
\nabla_{H} \phi(p, t)=0 \text { and }\left(\nabla_{H}^{2} \phi\right)^{*}(p, t)=0 . \tag{3.16}
\end{equation*}
$$

Therefore $w$ is a subsolution of (1.2) at $(p, t)$ by Definition 3.2. On the other hand, whenever a test function $\phi$ touches $w$ from below at ( $p, t$ ) with (3.16), we get $\phi_{t}(p, t)=$ $w_{t}(p, t)=24 t>0$, which implies that $w$ is also a supersolution due to Definition 3.2.

Remark 3.3. A basic transformation keeps the solution (3.13) being a solution. To be more precise, for any fixed $c \in \mathbb{R}, L>0$ and $\hat{p} \in \mathcal{H}$, we define $\hat{w}(p, t)=\operatorname{Lw}\left(\hat{p}^{-1} \cdot p, t\right)+c$ for all $(p, t) \in \mathcal{H} \times[0, \infty)$. Then we claim that $\hat{w}$ is a solution of (1.2). Indeed, our calculation above extends to

$$
\begin{aligned}
& X_{1}^{2} \hat{w}=X_{2}^{2} \hat{w}=12 L\left(p_{1}-\hat{p}_{1}\right)^{2}+12 L\left(p_{2}-\hat{p}_{2}\right)^{2}+24 L t ; \\
& X_{1} X_{2} \hat{w}=-X_{2} X_{1} \hat{w}=16 L\left(p_{3}-\hat{p}_{3}-\frac{1}{2} \hat{p}_{1} p_{2}+\frac{1}{2} p_{1} \hat{p}_{2}\right) .
\end{aligned}
$$

The conclusion follows immediately as in the proof of Proposition 3.2.
A primary and geometric observation for the explicit solution $u$ in (3.13) is as follows. For any fixed $\mu>0$, the $\mu$-level set,

$$
\Gamma_{t}^{\mu}=\{p \in \mathcal{H}: w(p, t)=\mu\}
$$

describes the position of surface at time $t \geq 0$. It is obvious that even if $\Gamma_{0}^{\mu} \neq \emptyset, \Gamma_{t}$ will vanish when $t$ is sufficiently large, which agrees with the usual extinction of mean curvature flows. We will revisit this property in Section 7.2.

A natural question now is whether the explicit solution we found is the only solution of (MCF) with the initial data (3.12). This is related to the open question on the uniqueness of solutions of (MCF). In the following sections we will give an affirmative answer for the case when the initial data are cylindrically symmetric about the vertical axis.

## 4. Comparison principle

4.1. Cylindrically symmetric solutions. Before presenting the proof of Theorem 1.1, let us investigate the properties for the solutions of (MCF) that are axisymmetric with respect to the vertical axis; in other words, we consider solutions of the form $u=u(r, z, t)$ where $r=\left(x^{2}+y^{2}\right)^{1 / 2}$.
Lemma 4.1 (Tests for axisymmetric solutions). Let u be a subsolution (resp., supersolution) of (3.1). Suppose that there exists $(\hat{p}, \hat{t}) \subset \mathcal{H} \times(0, \infty)$ and $\phi \in C^{2}(\mathcal{O})$ such that

$$
\max _{\mathcal{O}}(u-\phi)=(u-\phi)(\hat{p}, \hat{t}) \quad\left(\text { resp. }, \min _{\mathcal{O}}(u-\phi)=(u-\phi)(\hat{p}, \hat{t})\right) .
$$

If $\hat{p}=\left(\hat{p}_{1}, \hat{p_{2}}, \hat{p_{3}}\right)$ satisfies $\hat{p}_{1}^{2}+\hat{p}_{2}^{2} \neq 0$ and $u$ is axisymmetric about the vertical axis, then there exists $k \in \mathbb{R}$ such that

$$
\begin{equation*}
\frac{\partial}{\partial p_{1}} \phi(\hat{p}, \hat{t})=\hat{p_{1}} k \quad \text { and } \frac{\partial}{\partial p_{2}} \phi(\hat{p}, \hat{t})=\hat{p_{2}} k . \tag{4.1}
\end{equation*}
$$

Remark 4.1. It is clear that $k=\frac{\partial}{\partial r} \phi\left(\sqrt{\hat{p}_{1}^{2}+\hat{p}_{2}^{2}}, \hat{p}_{3}, \hat{t}\right)$ provided that $\phi=\phi\left(r, p_{3}, t\right)$, i.e., $\phi$ is also axisymmetric about the vertical axis.
Proof. Denote $\hat{r}=\sqrt{\hat{p}_{1}^{2}+\hat{p}_{2}^{2}}$. We only prove the situation when $u$ is a subsolution. By the symmetry of $u, u\left(p_{1}, p_{2}, \hat{p}_{3}, \hat{t}\right)=u\left(\hat{p}_{1}, \hat{p}_{2}, \hat{p}_{3}, \hat{t}\right)$ for all $p_{1}^{2}+p_{2}^{2}=\hat{r}^{2}$. By assumption, we have

$$
(u-\phi)\left(p_{1}, p_{2}, \hat{p}_{3}, \hat{t}\right) \leq(u-\phi)\left(\hat{p}_{1}, \hat{p}_{2}, \hat{p}_{3}, \hat{t}\right) \text { for all }\left(p_{1}, p_{2}, p_{3}, t\right) \in \mathcal{O},
$$

which implies that

$$
\phi\left(p_{1}, p_{2}, \hat{p}_{3}, \hat{t}\right) \geq \phi\left(\hat{p}_{1}, \hat{p}_{2}, \hat{p}_{3}, \hat{t}\right)
$$

for all ( $p_{1}, p_{2}$ ) close to ( $\hat{p}_{1}, \hat{p}_{2}$ ) with $p_{1}^{2}+p_{2}^{2}=r^{2}$. Applying the method of Lagrange's multiplier, we get $k \in \mathbb{R}$ such that

$$
\begin{gathered}
\frac{\partial}{\partial p_{1}}\left(\phi\left(p_{1}, p_{2}, \hat{p}_{3}, \hat{t}\right)-\frac{k}{2}\left(p_{1}^{2}+p_{2}^{2}-\hat{r}^{2}\right)\right)=0, \\
\frac{\partial}{\partial p_{2}}\left(\phi\left(p_{1}, p_{2}, \hat{p}_{3}, \hat{t}\right)-\frac{k}{2}\left(p_{1}^{2}+p_{2}^{2}-\hat{r}^{2}\right)\right)=0
\end{gathered}
$$

at ( $\hat{p}_{1}, \hat{p}_{2}$ ). We conclude (4.1) by straightforward calculations.
4.2. Proof of the comparison theorem. We are now in a position to prove Theorem 1.1.

Proof of Theorem 1.1. Let us assume $u$ is axisymmetric about the vertical axis. The same argument applies to the case when $v$ is axisymmetric. Suppose by contradiction that there exists $(\hat{p}, \hat{t}) \in \mathcal{H} \times(0, T)$ such that

$$
(u-v)(\hat{p}, \hat{t})>0
$$

We may assume that $(\hat{p}, \hat{t})$ satisfies

$$
\begin{equation*}
u(\hat{p}, \hat{t})-v(\hat{p}, \hat{t})-\frac{\sigma}{T-\hat{t}}=\max _{\mathcal{H} \times[0, T)}\left(u(p, t)-v(p, t)-\frac{\sigma}{T-t}\right)=\mu>0, \tag{4.2}
\end{equation*}
$$

when $\sigma>0$ is small. We fix such $\sigma$, double the variables and set up an auxiliary function

$$
\Phi^{\varepsilon}(p, t, q, s)=u(p, t)-v(q, s)-\frac{1}{\varepsilon} g^{2}(p, q)-\frac{1}{2 \varepsilon}(t-s)^{2}-\frac{\sigma}{T-t},
$$

where $g(p, q)=\left|p \cdot q^{-1}\right|^{4}$. Let $\left(p^{\varepsilon}, t^{\varepsilon}, q^{\varepsilon}, s^{\varepsilon}\right) \in(\mathcal{H} \times[0, T))^{2}$ be a maximizer of $\Phi^{\varepsilon}$, then it is clear that

$$
\Phi^{\varepsilon}\left(p^{\varepsilon}, t^{\varepsilon}, q^{\varepsilon}, s^{\varepsilon}\right)=\sup _{(\mathcal{H} \times[0, T))^{2}} \Phi^{\varepsilon}>\Phi^{\varepsilon}(\hat{p}, \hat{t}, \hat{p}, \hat{t}),
$$

which implies that

$$
\begin{equation*}
\frac{1}{\varepsilon} g^{2}\left(p^{\varepsilon}, q^{\varepsilon}\right)+\frac{1}{2 \varepsilon}\left(t^{\varepsilon}-s^{\varepsilon}\right)^{2} \leq u\left(p^{\varepsilon}, t^{\varepsilon}\right)-v\left(q^{\varepsilon}, s^{\varepsilon}\right)-u(\hat{p}, \hat{t})+v(\hat{p}, \hat{t})+\frac{\sigma}{T-\hat{t}}-\frac{\sigma}{T-t^{\varepsilon}} . \tag{4.3}
\end{equation*}
$$

By the boundedness of $u$ and $v$, we have

$$
\left|p^{\varepsilon} \cdot\left(q^{\varepsilon}\right)^{-1}\right| \rightarrow 0 \text { and }\left|t^{\varepsilon}-s^{\varepsilon}\right| \rightarrow 0 \text { as } \varepsilon \rightarrow 0
$$

Since $u=a$ and $v=b$ with $a \leq b$ outside $K \times[0, \infty)$, we may take a subsequence of $\varepsilon$, still indexed by $\varepsilon$, such that $p^{\varepsilon}, q^{\varepsilon} \rightarrow \bar{p} \in \mathcal{H}$ and $t^{\varepsilon}, s^{\varepsilon} \rightarrow \bar{t} \in[0, T)$ as $\varepsilon \rightarrow 0$. Sending the limit in (4.3) and applying (4.2), we get

$$
\limsup _{\varepsilon \rightarrow 0}\left(\frac{1}{\varepsilon} g^{2}\left(p^{\varepsilon}, q^{\varepsilon}\right)+\frac{1}{2 \varepsilon}\left(t^{\varepsilon}-s^{\varepsilon}\right)^{2}\right) \leq 0 .
$$

In other words, we have

$$
\begin{equation*}
\frac{1}{\varepsilon} g^{2}\left(p^{\varepsilon}, q^{\varepsilon}\right) \rightarrow 0 \text { and } \frac{1}{2 \varepsilon}\left(t^{\varepsilon}-s^{\varepsilon}\right)^{2} \rightarrow 0 \text { as } \varepsilon \rightarrow 0 . \tag{4.4}
\end{equation*}
$$

We next claim that $\bar{t} \neq 0$. Indeed, if $\bar{t}=0$, then, since $u(p, 0) \leq v(p, 0)$ for all $p \in \mathcal{H}$, we are led to

$$
\Phi^{\varepsilon}\left(p^{\varepsilon}, t^{\varepsilon}, q^{\varepsilon}, s^{\varepsilon}\right) \rightarrow u(p, 0)-v(p, 0)-\frac{\sigma}{T}<0
$$

which contradicts the fact that $\Phi^{\varepsilon}\left(p^{\varepsilon}, t^{\varepsilon}, q^{\varepsilon}, s^{\varepsilon}\right) \geq \mu$. We next apply the Crandall-Ishii lemma and get

$$
\begin{aligned}
& \left(\frac{\sigma}{\left(T-t^{\varepsilon}\right)^{2}}+\frac{1}{\varepsilon}\left(t^{\varepsilon}-s^{\varepsilon}\right), \frac{1}{\varepsilon} \nabla^{p} g^{2}\left(p^{\varepsilon}, q^{\varepsilon}\right), \mathcal{X}^{\varepsilon}\right) \in \bar{J}_{H}^{2,+} u\left(p^{\varepsilon}, t^{\varepsilon}\right) ; \\
& \left(\frac{1}{\varepsilon}\left(t^{\varepsilon}-s^{\varepsilon}\right),-\frac{1}{\varepsilon} \nabla^{q} g^{2}\left(p^{\varepsilon}, q^{\varepsilon}\right), \mathcal{Y}^{\varepsilon}\right) \in \bar{J}_{H}^{2,-} v\left(q^{\varepsilon}, s^{\varepsilon}\right),
\end{aligned}
$$

where $\bar{J}_{H}^{2,+}$ and $\bar{J}_{H}^{2,-}$ denote the closure of the semijets in Heisenberg group and $\mathcal{X}^{\varepsilon}, \mathcal{Y}^{\varepsilon} \in \mathbf{S}^{2}$ satisfy

$$
\left\langle\mathcal{X}^{\varepsilon} \xi, \xi\right\rangle-\left\langle\mathcal{Y}^{\varepsilon} \xi, \xi\right\rangle \leq \frac{C}{\varepsilon} g\left(p^{\varepsilon}, q^{\varepsilon}\right)\left|p^{\varepsilon} \cdot\left(q^{\varepsilon}\right)^{-1}\right|^{4}|\xi|^{2}=\frac{C}{\varepsilon} g^{2}\left(p^{\varepsilon}, q^{\varepsilon}\right)|\xi|^{2}
$$

for some $C>0$ and all $\xi \in \mathbb{R}^{2}$. See $[\mathrm{B}, \mathrm{M}]$ for more details on the semijets and the Crandall-Ishii lemma on the Heisenberg group. It follows from (4.4) that

$$
\begin{equation*}
\limsup _{\varepsilon \rightarrow 0}\left(\left\langle\mathcal{X}^{\varepsilon} \xi, \xi\right\rangle-\left\langle\mathcal{Y}^{\varepsilon} \xi, \xi\right\rangle\right) \leq 0 \tag{4.5}
\end{equation*}
$$

uniformly for all bounded $\xi \in \mathbb{R}^{2}$. Moreover, as is derived from Remark 2.1, the following gradient relation holds:

$$
\frac{1}{\varepsilon} \nabla_{H}^{p} g^{2}\left(p^{\varepsilon}, q^{\varepsilon}\right)=-\frac{1}{\varepsilon} \nabla_{H}^{q} g^{2}\left(p^{\varepsilon}, q^{\varepsilon}\right) .
$$

Let $\eta^{\varepsilon}$ denote $\frac{1}{\varepsilon} \nabla_{H}^{p} g\left(p^{\varepsilon}, q^{\varepsilon}\right)$.
Finally, we adopt Definition 3.3 to derive a contradiction.
Case A. If $\eta^{\varepsilon} \neq 0$ for all $\varepsilon>0$ small, then

$$
\begin{equation*}
\frac{\sigma}{\left(T-t^{\varepsilon}\right)^{2}}+\frac{1}{\varepsilon}\left(t^{\varepsilon}-s^{\varepsilon}\right)+F\left(\eta^{\varepsilon}, \mathcal{X}^{\varepsilon}\right) \leq 0 \tag{4.6}
\end{equation*}
$$

and

$$
\begin{equation*}
\frac{1}{\varepsilon}\left(t^{\varepsilon}-s^{\varepsilon}\right)+F\left(\eta^{\varepsilon}, \mathcal{Y}^{\varepsilon}\right) \geq 0 \tag{4.7}
\end{equation*}
$$

Taking the difference of (4.6) and (4.7) yields

$$
\frac{\sigma}{\left(T-t^{\varepsilon}\right)^{2}} \leq \operatorname{tr}\left(I-\frac{\eta^{\varepsilon} \otimes \eta^{\varepsilon}}{\left|\eta^{\varepsilon}\right|^{2}}\right)\left(\mathcal{X}^{\varepsilon}-\mathcal{Y}^{\varepsilon}\right) .
$$

Passing to the limit as $\varepsilon \rightarrow 0$ with an application of (4.5), we end up with

$$
\frac{\sigma}{(T-\bar{t})^{2}} \leq 0,
$$

which is clearly a contradiction.
Case B. If $\eta^{\varepsilon_{j}}=\frac{1}{\varepsilon} \nabla_{H}^{p} g^{2}\left(p^{\varepsilon}, q^{\varepsilon}\right)=\frac{1}{\varepsilon} \nabla_{H}^{q} g^{2}\left(p^{\varepsilon}, q^{\varepsilon}\right)=0$ for a subsequence $\varepsilon_{j} \rightarrow 0$, we obtain, by computation, that

$$
\begin{align*}
& 2 g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right) X_{1}^{p} g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)=2 g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)\left(\frac{\partial}{\partial p_{1}} g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)-\frac{p_{2}^{\varepsilon_{j}}}{2} \frac{\partial}{\partial p_{3}} g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)\right)=0 ; \\
& 2 g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right) X_{2}^{p} g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)=2 g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)\left(\frac{\partial}{\partial p_{2}} g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)+\frac{p_{1}^{\varepsilon_{j}}}{2} \frac{\partial}{\partial p_{3}} g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)\right)=0 \tag{4.8}
\end{align*}
$$

and

$$
\begin{align*}
& 2 g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right) X_{1}^{q} g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)=2 g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)\left(\frac{\partial}{\partial q_{1}} g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)-\frac{q_{2}^{\varepsilon_{j}}}{2} \frac{\partial}{\partial q_{3}} g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)\right)=0 ; \\
& 2 g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right) X_{2}^{q} g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)=2 g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)\left(\frac{\partial}{\partial q_{2}} g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)+\frac{q_{1}^{\varepsilon_{j}}}{2} \frac{\partial}{\partial q_{3}} g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)\right)=0 . \tag{4.9}
\end{align*}
$$

We further discuss two sub-cases.
Case 1. When $g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)=0$, we get $p^{\varepsilon_{j}}=q^{\varepsilon_{j}}$, which implies that

$$
\nabla_{H}^{p} g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)=\nabla_{H}^{q} g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)=0 .
$$

Since

$$
\begin{array}{ll}
X_{1}^{2, p} g^{2}=2\left(X_{1}^{p} g\right)^{2}+2 g X_{1}^{2, p} g, & X_{2}^{2, p} g^{2}=2\left(X_{2}^{p} g\right)^{2}+2 g X_{2}^{2, p} g, \\
X_{1}^{p} X_{2}^{p} g^{2}=2 X_{1}^{p} g X_{2}^{p} g+2 g X_{1}^{p} X_{2}^{p} g, & X_{2}^{p} X_{1}^{p} g^{2}=2 X_{2}^{p} g X_{1}^{p} g+2 g X_{2}^{p} X_{1}^{p} g, \tag{4.10}
\end{array}
$$

We have $\left(\nabla_{H}^{2, p} g^{2}\right)^{*}\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)=0$. Similarly, we can deduce $\left(\nabla_{H}^{2, q} g^{2}\right)^{*}\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)=0$. By Definition 3.2, the viscosity inequalities read

$$
\begin{equation*}
\frac{\sigma}{\left(T-t^{\varepsilon_{j}}\right)^{2}}+\frac{1}{\varepsilon_{j}}\left(t^{\varepsilon_{j}}-s^{\varepsilon_{j}}\right) \leq 0 \tag{4.11}
\end{equation*}
$$

and

$$
\begin{equation*}
\frac{1}{\varepsilon_{j}}\left(t^{\varepsilon_{j}}-s^{\varepsilon_{j}}\right) \geq 0 \tag{4.12}
\end{equation*}
$$

whose difference implies that $\sigma /\left(T-t^{\varepsilon_{j}}\right)^{2} \leq 0$. This is certainly a contradiction.
Case 2. When $g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right) \neq 0$, we get $X_{1}^{p} g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)=X_{2}^{p} g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)=0$. We first claim that $p_{1}^{\varepsilon_{j}}=p_{2}^{\varepsilon_{j}}=0$. Suppose by contradiction that $\left(p_{1}^{\varepsilon_{j}}\right)^{2}+\left(p_{2}^{\varepsilon_{j}}\right)^{2} \neq 0$. In terms of Lemma 4.1, there is $k \in \mathbb{R}$ such that (4.8) reduces to

$$
p_{1}^{\varepsilon_{j}} k-\frac{p_{2}^{\varepsilon_{j}}}{2} \frac{\partial}{\partial p_{3}} g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)=0 \text { and } p_{2}^{\varepsilon_{j}} k+\frac{p_{1}^{\varepsilon_{j}}}{2} \frac{\partial}{\partial p_{3}} g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)=0
$$

which yields that $k=0$ and

$$
\frac{\partial}{\partial p_{3}} g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)=p_{3}^{\varepsilon_{j}}-q_{3}^{\varepsilon_{j}}-\frac{1}{2} p_{1}^{\varepsilon_{j}} q_{2}^{\varepsilon_{j}}+\frac{1}{2} p_{2}^{\varepsilon_{j}} q_{1}^{\varepsilon_{j}}=0 .
$$

It follows from (4.8), (2.11) and (2.12) that $p^{\varepsilon_{j}}=q^{\varepsilon_{j}}$, which contradicts the assumption that $g\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right) \neq 0$. This completes the proof of our claim.

As $p_{1}^{\varepsilon_{j}}=p_{2}^{\varepsilon_{j}}=0$, we apply (4.8), (2.11) and (2.12) again and get

$$
\begin{aligned}
& \left.4\left(\left(q_{1}^{\varepsilon_{j}}\right)^{2}+\left(q_{2}^{\varepsilon_{j}}\right)\right)^{2}\right)\left(-q_{1}^{\varepsilon_{j}}\right)-16 q_{2}^{\varepsilon_{j}}\left(p_{3}^{\varepsilon_{j}}-q_{3}^{\varepsilon_{j}}\right)=0 ; \\
& 4\left(\left(q_{1}^{\varepsilon_{j}}\right)^{2}+\left(q_{2}^{\varepsilon_{j}}\right)^{2}\right)\left(-q_{2}^{\varepsilon_{j}}\right)+16 q_{1}^{\varepsilon_{j}}\left(p_{3}^{\varepsilon_{j}}-q_{3}^{\varepsilon_{j}}\right)=0 .
\end{aligned}
$$

We are then led to $q_{1}^{\varepsilon_{j}}=q_{2}^{\varepsilon_{j}}=0$. Now simplifying the second derivatives of $g^{2}$ in (4.10) by using (2.15)-(2.19), we obtain $\left(\nabla_{H}^{2, p} g^{2}\right)^{*}\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)=0$. An analog of calculation yields that $\left(\nabla_{H}^{2, q} g^{2}\right)^{*}\left(p^{\varepsilon_{j}}, q^{\varepsilon_{j}}\right)=0$. The proof is complete since Definition 3.2 can be adopted once again to get (4.11)-(4.12) and deduce a contradiction.

## 5. Existence theorem by games

The game setting is as follows. A marker, representing the game state, is initialized at a state $p \in \mathcal{H}$ from time 0 . The maturity time given is denoted by $t$. Let the step size for space be $\varepsilon>0$. Time $\varepsilon^{2}$ is consumed for every step. Then the total number of game steps $N$ can be regarded as $\left[t / \varepsilon^{2}\right]$. The game states for all steps are denoted in order by $\zeta^{0}, \zeta^{1}, \ldots, \zeta^{N}$ with $\zeta^{0}=p$. Two players, Player I and Player II participate the game. Player I intends to minimize at the final state an objective function, which in our case is $u_{0}: \mathcal{H} \rightarrow \mathbb{R}$, while Player II is to maximize it. At the $(k+1)$-th round $(k<N)$,
(1) Player I chooses in $\mathcal{H}$ a unit horizontal vector $v^{k}$, i.e., $v^{k}=\left(v_{1}^{k}, v_{2}^{k}, 0\right)$ satisfying $\left|v^{k}\right|^{2}=\left(v_{1}^{k}\right)^{2}+\left(v_{2}^{k}\right)^{2}=1$. We denote by $S_{h}^{1}$ the set of all unit horizontal vectors.
(2) Carol has the right to reverse Paul's choice, which determines $b^{k}= \pm 1$;
(3) The marker is moved from the present state $\zeta^{k}$ to $\zeta^{k} \cdot\left(\sqrt{2} \varepsilon b^{k} v^{k}\right)$.

Then the state equation is written inductively as

$$
\left\{\begin{array}{l}
\zeta^{k+1}=\zeta^{k} \cdot\left(\sqrt{2} \varepsilon b^{k} v^{k}\right), \quad k=0,1, \ldots, N-1 ;  \tag{5.1}\\
\zeta^{0}=p
\end{array}\right.
$$

The value function is defined to be

$$
\begin{equation*}
u^{\varepsilon}(p, t):=\min _{v^{1}} \max _{b^{1}} \ldots \min _{v^{N}} \max _{b^{N}} u_{0}\left(\zeta^{N}\right), \tag{5.2}
\end{equation*}
$$

By the dynamic programming:

$$
\begin{equation*}
u^{\varepsilon}(p, t)=\min _{v \in S_{h}^{1}} \max _{b= \pm 1} u^{\varepsilon}\left(p \cdot(\sqrt{2} \varepsilon b v), t-\varepsilon^{2}\right) \tag{5.3}
\end{equation*}
$$

with $u^{\varepsilon}(p, 0)=u_{0}(p)$.
Our main result of this section is given below.

Theorem 5.1 (Existence theorem by games). Assume that $u_{0}$ is uniformly continuous function in $\mathcal{H}$ and is constant $C \in \mathbb{R}$ outside a compact set. Assume also that $u_{0}$ is spatially axisymmetric about the vertical axis. Let $u^{\varepsilon}$ be the value function defined as in (5.2). Then $u^{\varepsilon}$ converges, as $\varepsilon \rightarrow 0$, to the unique axisymmetric viscosity solution of (MCF) uniformly on compact subsets of $\mathcal{H} \times[0, \infty)$. Moreover, $u=C$ in $(\mathcal{H} \backslash K) \times(0, \infty)$ for some compact set $K \subset \mathcal{H}$.

Before presenting the proof of Theorem 5.1, we first give bounds for the game trajectories under some particular strategies.

Lemma 5.2 (Lower bound of the game trajectories). For any $p \in \mathcal{H}$ and $t \geq 0$ with $N=\left[t / \varepsilon^{2}\right]$, let $\zeta_{k}$ be defined as in (5.1) for all $k=0,1, \ldots, N$. Then the following statements hold.
(i) There exists a strategy of Player I such that

$$
\begin{equation*}
\left(\left|\zeta_{1}^{N}\right|^{2}+\left|\zeta_{2}^{N}\right|^{2}\right)^{2}+16\left|\zeta_{3}^{N}\right|^{2} \geq\left(\left|p_{1}\right|^{2}+\left|p_{2}\right|^{2}\right)^{2}+16\left|p_{3}\right|^{2} \tag{5.4}
\end{equation*}
$$

under this strategy regardless of Player II's choices.
(ii) There exists a strategy of Player II such that (5.4) holds under this strategy regardless of Player I's choices.

Proof. (i) By direct calculation, we have

$$
\begin{align*}
& \left(\left(p_{1}+\sqrt{2} \varepsilon b v_{1}\right)^{2}+\left(p_{2}+\sqrt{2} \varepsilon b v_{2}\right)^{2}\right)^{2}+16\left(p_{3}+\frac{1}{2} \sqrt{2} \varepsilon b\left(p_{1} v_{2}-p_{2} v_{1}\right)\right)^{2} \\
= & \left(p_{1}^{2}+p_{2}^{2}+2 \varepsilon^{2}\right)^{2}+8 \varepsilon^{2}\left(p_{1}^{2}+p_{2}^{2}\right)+16 p_{3}^{2}  \tag{5.5}\\
& +4 \sqrt{2} \varepsilon b\left(\left(p_{1}^{2}+p_{2}^{2}+2 \varepsilon^{2}\right)\left(p_{1} v_{1}+p_{2} v_{2}\right)+4\left(p_{1} p_{3} v_{2}-p_{2} p_{3} v_{1}\right)\right)
\end{align*}
$$

It is clear that Player I may take $v=\left(v_{1}, v_{2}, 0\right) \in S_{h}^{1}$ satisfying

$$
\begin{aligned}
& v_{1}=\frac{1}{\rho}\left(\left(p_{1}^{2}+p_{2}^{2}+2 \varepsilon^{2}\right) p_{2}+4 p_{1} p_{3}\right) \\
& v_{2}=-\frac{1}{\rho}\left(\left(p_{1}^{2}+p_{2}^{2}+2 \varepsilon^{2}\right) p_{1}+4 p_{2} p_{3}\right),
\end{aligned}
$$

with

$$
\rho=\left(p_{1}^{2}+p_{2}^{2}\right)^{1 / 2}\left(\left(p_{1}^{2}+p_{2}^{2}+2 \varepsilon^{2}\right)^{2}+16 p_{3}^{2}\right)^{1 / 2}
$$

so that, no matter which $b$ is picked, we have

$$
b\left(\left(p_{1}^{2}+p_{2}^{2}+2 \varepsilon^{2}\right)\left(p_{1} v_{1}+p_{2} v_{2}\right)+4\left(p_{1} p_{3} v_{2}-p_{2} p_{3} v_{1}\right)\right)=0
$$

and, furthermore by (5.5),

$$
\begin{align*}
& \left(\left(p_{1}+\sqrt{2} \varepsilon b v_{1}\right)^{2}+\left(p_{2}+\sqrt{2} \varepsilon b v_{2}\right)^{2}\right)^{2}+16\left(p_{3}+\frac{1}{2} \sqrt{2} \varepsilon b\left(p_{1} v_{2}-p_{2} v_{1}\right)\right)^{2}  \tag{5.6}\\
= & \left(p_{1}^{2}+p_{2}^{2}+2 \varepsilon^{2}\right)^{2}+8 \varepsilon^{2}\left(p_{1}^{2}+p_{2}^{2}\right)+16 p_{3}^{2} \geq\left(p_{1}^{2}+p_{2}^{2}\right)^{2}+16 p_{3}^{2} .
\end{align*}
$$

We can iterate (5.6) to get

$$
\left(\left|\zeta_{1}^{k}\right|^{2}+\left|\zeta_{2}^{k}\right|^{2}\right)^{2}+16\left|\zeta_{3}^{k}\right|^{2} \geq\left(\left|\zeta_{1}^{k-1}\right|^{2}+\left|\zeta_{2}^{k-1}\right|^{2}\right)^{2}+16\left|\zeta_{3}^{k-1}\right|^{2}
$$

for all $k=1,2, \ldots, N$ and (5.4) follows easily.
(ii) The proof of (ii) is similar and even easier. Note that Player II may take a proper $b= \pm 1$ so that

$$
b\left(\left(p_{1}^{2}+p_{2}^{2}+2 \varepsilon^{2}\right)\left(p_{1} v_{1}+p_{2} v_{2}\right)+4\left(p_{1} p_{3} v_{2}-p_{2} p_{3} v_{1}\right)\right) \geq 0
$$

and therefore (5.6) holds immediately. We then complete the proof by iteration again.

Lemma 5.3 (Upper bound of the game trajectories). For any $p \in \mathcal{H}$ and $t \geq 0$ with $N=\left[t / \varepsilon^{2}\right]$, let $\zeta^{k}$ be defined as in (5.1) for all $k=0,1, \ldots, N$. Then the following statements hold.
(i) There exists a strategy of Player I such that

$$
\begin{equation*}
\left(\left|\zeta_{1}^{N}\right|^{2}+\left|\zeta_{2}^{N}\right|^{2}\right)^{2}+16\left|\zeta_{3}^{N}\right|^{2} \leq\left(\left|p_{1}\right|^{2}+\left|p_{2}\right|^{2}+6 N \varepsilon^{2}\right)^{2}+16\left|p_{3}\right|^{2} \tag{5.7}
\end{equation*}
$$

under this strategy regardless of Player II's choices.
(ii) There exists a strategy of Player II such that (5.7) holds under this strategy regardless of Player I's choices.

Remark 5.1. With the notation of the gauge $G$ in (3.12), the inequality (5.7) can be simplified into

$$
G\left(\zeta^{N}\right) \leq\left(\left|p_{1}\right|^{2}+\left|p_{2}\right|^{2}+6 t\right)^{2}+16\left|p_{3}\right|^{2}
$$

which is intuitively natural, since the explicit solution given in (3.13) satisfies

$$
w(p, t) \leq\left(\left|p_{1}\right|^{2}+\left|p_{2}\right|^{2}+6 t\right)^{2}+16\left|p_{3}\right|^{2}
$$

Proof. By iteration, it suffices to show there exist strategies of Player I or Player II such that

$$
\begin{align*}
& \left(\left(p_{1}+\sqrt{2} \varepsilon b v_{1}\right)^{2}+\left(p_{2}+\sqrt{2} \varepsilon b v_{2}\right)^{2}+j \varepsilon^{2}\right)^{2} \\
& +16\left(p_{3}+\frac{1}{2} \sqrt{2} \varepsilon b\left(p_{1} v_{2}-p_{2} v_{1}\right)\right)^{2} \leq\left(p_{1}^{2}+p_{2}^{2}+(j+6) \varepsilon^{2}\right)^{2}+16 p_{3}^{2} \tag{5.8}
\end{align*}
$$

Indeed, the left hand side is calculated to be

$$
\begin{aligned}
& \left(p_{1}^{2}+p_{2}^{2}+(j+2) \varepsilon^{2}\right)^{2}+8 \varepsilon^{2}\left(p_{1}^{2}+p_{2}^{2}\right)+16 p_{3}^{2} \\
& +4 \sqrt{2} \varepsilon b\left(\left(p_{1}^{2}+p_{2}^{2}+2 \varepsilon^{2}\right)\left(p_{1} v_{1}+p_{2} v_{2}\right)+4\left(p_{1} p_{3} v_{2}-p_{2} p_{3} v_{1}\right)\right)
\end{aligned}
$$

As in the proof of Lemma 5.3, either Player I or Player II may let

$$
b\left(\left(p_{1}^{2}+p_{2}^{2}+2 \varepsilon^{2}\right)\left(p_{1} v_{1}+p_{2} v_{2}\right)+4\left(p_{1} p_{3} v_{2}-p_{2} p_{3} v_{1}\right)\right) \leq 0
$$

with no regard for their opponents strategies. Hence, by a strategy of either Player I or Player II, we have

$$
\begin{aligned}
\left(\left(p_{1}+\sqrt{2} \varepsilon b v_{1}\right)^{2}\right. & \left.+\left(p_{2}+\sqrt{2} \varepsilon b v_{2}\right)^{2}+j \varepsilon^{2}\right)^{2}+16\left(p_{3}+\frac{1}{2} \sqrt{2} \varepsilon b\left(p_{1} v_{2}-p_{2} v_{1}\right)\right)^{2} \\
& \leq\left(p_{1}^{2}+p_{2}^{2}+(j+2) \varepsilon^{2}\right)^{2}+8 \varepsilon^{2}\left(p_{1}^{2}+p_{2}^{2}\right)+16 p_{3}^{2} \\
& \leq\left(p_{1}^{2}+p_{2}^{2}+(j+6) \varepsilon^{2}\right)^{2}+16 p_{3}^{2}
\end{aligned}
$$

which proves (5.8).
Remark 5.2. For any $\hat{p} \in \mathcal{H}, c \in \mathbb{R}$ and $L>0$, let

$$
\begin{equation*}
\hat{G}(p)=c+L G\left(\hat{p}^{-1} \cdot p\right) \tag{5.9}
\end{equation*}
$$

Our proof above can be directly generalized to show that

$$
\hat{G}\left(\zeta^{N}\right) \leq c+L\left(\left|p_{1}-\hat{p}_{1}\right|^{2}+\left|p_{2}-\hat{p}_{2}\right|^{2}+6 N \varepsilon^{2}\right)^{2}+16 L\left|p_{3}-\hat{p}_{3}+\frac{1}{2}\left(p_{1} \hat{p}_{2}-p_{2} \hat{p}_{1}\right)\right|^{2}
$$

with either a strategy of Player I or a strategy of Player II.
We now return to the proof of Theorem 5.1, which actually rests on showing that $\bar{u}$ and $\underline{u}$, as defined in (1.5) and (1.6), are respectively a subsolution and a supersolution of (MCF). (Note that our definitions are valid since the game value $u_{\varepsilon}$ are bounded uniformly for all $\varepsilon>0$ by its definition.) Moreover, we show that $\bar{u}(p, 0) \leq \underline{u}(p, 0)$ and $\bar{u}$ and $\underline{u}$ are constant outside a compact set. Then it follows immediately from the comparison principle (Theorem 1.1) that $\bar{u} \leq \underline{u}$ and therefore $u^{\varepsilon} \rightarrow u$ locally uniformly as $\varepsilon \rightarrow 0$.

Proposition 5.4 (Constant value outside a compact set). Assume that $u_{0}$ is uniformly continuous function in $\mathcal{H}$ and is a constant $C \in \mathbb{R}$ outside a compact set. Let $u^{\varepsilon}$ be the value function defined by (5.2). Then for any $T>0, \bar{u}(p, t)=\underline{u}(p, t)=C$ for all $p \in \mathcal{H}$ outside a compact set and for all $t \in[0, T]$.

Proof. Suppose there exists $B_{r}$ such that $u_{0}(p)=C$ for any $p \in \mathcal{H} \backslash B_{r}$. Then for any $\hat{p} \in \mathcal{H} \backslash B_{r}$ and $t \geq 0$, we use the strategy of Player I introduced in Lemma 5.2, we get $\zeta^{N} \in \mathcal{H} \backslash B_{r}$ regardless of Player II's choices, which implies that

$$
u^{\varepsilon}(p, t) \leq u_{0}\left(\zeta^{N}\right)=C .
$$

Similarly, we may use the strategy of Player II to deduce that

$$
u^{\varepsilon}(p, t) \geq C .
$$

Hence, $u^{\varepsilon}=C$ and $\bar{u}=\underline{u}=C$ in $\mathcal{H} \backslash B_{r}$.
Proposition 5.5. Assume that $u_{0}$ is uniformly continuous function in $\mathcal{H}$ and is constant outside a compact set. Let $u^{\varepsilon}$ be the value function defined by (5.2). Then $\bar{u}(p, 0) \leq u_{0}(p)$ and $\underline{u}(p, 0) \geq u_{0}(p)$ for all $p \in \mathcal{H}$.

In order to prove this result, we first need to regularize the initial data with the smooth gauge $G$ in (3.12). We define

$$
\begin{equation*}
\psi^{L}(p)=\sup _{q \in \mathcal{H}}\left\{u_{0}(q)-L G\left(p^{-1} \cdot q\right)\right\} \tag{5.10}
\end{equation*}
$$

and

$$
\begin{equation*}
\psi_{L}(p)=\inf _{q \in \mathcal{H}}\left\{u_{0}(q)+L G\left(p^{-1} \cdot q\right)\right\} \tag{5.11}
\end{equation*}
$$

for any $p \in \mathcal{H}$ and fixed $L>0$. These two functions are called the sup-convolution and inf-convolution of $u_{0}$ respectively. Our definitions here are slightly different from those in [W] in that we plug $p^{-1} \cdot q$ instead of $q \cdot p^{-1}$ in $G$. However, the properties remain the same. We present one of the important properties for our use.

Lemma 5.6 (Approximation by semi-convolutions). Assume that $u_{0}$ is uniformly continuous on $\mathcal{H}$ and is constant outside a compact set. Let $\psi^{L}$ and $\psi_{L}$ be respectively defined as in (5.10) and (5.11). Then $\psi^{L}$ and $\psi_{L}$ converge to $u_{0}$ uniformly in $\mathcal{H}$ as $L \rightarrow \infty$.
Proof. We only show the statement for $\psi^{L}$. The proof for the statement on $\psi_{L}$ is symmetric.

It is easily seen that

$$
\begin{equation*}
\psi^{L} \geq u_{0} \text { in } \mathcal{H} \tag{5.12}
\end{equation*}
$$

On the other hand, since $u_{0}$ is uniformly continuous, for any $p \in \mathcal{H}$, we may find $q_{L} \in \mathcal{H}$ such that

$$
\psi^{L}(p)=\sup _{q \in \mathcal{H}}\left\{u_{0}(p)-L G\left(p^{-1} \cdot q\right)\right\}=u_{0}\left(q_{L}\right)-L G\left(p^{-1} \cdot q_{L}\right) .
$$

By (5.12), we have

$$
\begin{equation*}
L G\left(p^{-1} \cdot q_{L}\right) \leq u_{0}\left(q_{L}\right)-u_{0}(p), \tag{5.13}
\end{equation*}
$$

which, by the boundedness of $u_{0}$, implies that

$$
\left|p^{-1} \cdot q_{L}\right| \leq\left(2 K_{0} / L\right)^{1 / 4}
$$

where $K_{0}=\sup _{H}\left|u_{0}\right|$. By the uniform continuity of $u_{0}$, for any $\delta>0$, there exists $\varepsilon>0$ such that $\left|u_{0}(p)-u_{0}(q)\right| \leq \delta$ for any $p, q \in \mathcal{H}$ satisfying $\left|p^{-1} \cdot q\right| \leq \varepsilon$. Then we may let $L>0$ be sufficiently large such that $\left(2 K_{0} / L\right)^{1 / 4} \leq \varepsilon$ and therefore

$$
u_{0}\left(q_{L}\right)-u_{0}(p) \leq \delta,
$$

which, combined with (5.12), yields

$$
\left|\psi^{L}(p)-u_{0}(p)\right| \leq \delta \text { for all } p \in \mathcal{H}
$$

Proof of Proposition 5.5. We arbitrarily fix $\hat{p} \in \mathcal{H}$. By Lemma 5.6, for any $\delta>0$, there exists $L>0$ such that

$$
\psi^{L}(\hat{p}) \leq u_{0}(\hat{p})+\delta
$$

which implies that

$$
u_{0}(p) \leq u_{0}(\hat{p})+\delta+L G\left(\hat{p}^{-1} \cdot p\right)
$$

Let us use the right hand side, which is exactly $\hat{G}$ in (5.9) with $c=u_{0}(\hat{p})+\delta$, as the objective function of the games. Suppose the game value is $w^{\varepsilon}$. Then by using the special strategy of Player I given in Lemma 5.3 and Remark 5.2, we obtain a game estimate

$$
\begin{aligned}
w^{\varepsilon}(p, t) \leq & u_{0}(\hat{p})+\delta+L G\left(\hat{p}^{-1} \cdot \zeta^{N}\right) \\
\leq & u_{0}(\hat{p})+\delta+L\left(\left|p_{1}-\hat{p}_{1}\right|^{2}+\left|p_{2}-\hat{p}_{2}\right|^{2}+6 N \varepsilon^{2}\right)^{2} \\
& +16 L\left|p_{3}-\hat{p}_{3}+\frac{1}{2}\left(p_{1} \hat{p}_{2}-p_{2} \hat{p}_{1}\right)\right|^{2}
\end{aligned}
$$

no matter what choices are made by Player II during the game. On the other hand, since it is clear that $u^{\varepsilon} \leq w^{\varepsilon}$ and $N \varepsilon^{2} \leq t$, we get

$$
\begin{aligned}
u^{\varepsilon}(p, t) \leq u_{0}(\hat{p})+\delta+L\left(\left|p_{1}-\hat{p}_{1}\right|^{2}+\right. & \left.\left|p_{2}-\hat{p}_{2}\right|^{2}+6 t\right)^{2} \\
& +16 L\left|p_{3}-\hat{p}_{3}+\frac{1}{2}\left(p_{1} \hat{p}_{2}-p_{2} \hat{p}_{1}\right)\right|^{2}
\end{aligned}
$$

Taking the relaxed limit of $u^{\varepsilon}$ at $(\hat{p}, 0)$ as $\varepsilon \rightarrow 0$, we have

$$
\bar{u}(\hat{p}, 0) \leq u_{0}(\hat{p})+\delta
$$

We finally send $\delta \rightarrow 0$ and get $\bar{u}(\hat{p}, 0) \leq u_{0}(\hat{p})$ for any $\hat{p} \in \mathcal{H}$.
The proof for the statement that $\underline{u}(p, 0) \geq u_{0}(p)$ for all $p \in \mathcal{H}$ is symmetric. In fact, the key is to use the strategy of Player II introduced in Lemma 5.3 and Remark 5.2 to deduce

$$
\begin{aligned}
u^{\varepsilon}(p, t) \geq u_{0}(\hat{p})-\delta-L\left(\left|p_{1}-\hat{p}_{1}\right|^{2}+\right. & \left.\left|p_{2}-\hat{p}_{2}\right|^{2}+6 t\right)^{2} \\
& -16 L\left|p_{3}-\hat{p}_{3}+\frac{1}{2}\left(p_{1} \hat{p}_{2}-p_{2} \hat{p}_{1}\right)\right|^{2}
\end{aligned}
$$

Proposition 5.7 (Axial symmetry of the game values). Suppose that $u_{0}$ is uniformly continuous on $\mathcal{H}$ and is spatially axisymmetric with respect to the vertical axis. Let $u^{\varepsilon}$ be the value function defined as in (5.2) Then $u^{\varepsilon}, \bar{u}$ and $\underline{u}$ are also spatially axisymmetric about the vertical axis.

Proof. We argue by induction. Assume that $u^{\varepsilon}(p, t)=u^{\varepsilon}\left(p^{\prime}, t\right)$ for some $t \geq 0$ and for any $p, p^{\prime} \in \mathcal{H}$ such that

$$
\begin{equation*}
p_{1}^{2}+p_{2}^{2}=\left(p_{1}^{\prime}\right)^{2}+\left(p_{2}^{\prime}\right)^{2} \text { and } p_{3}=p_{3}^{\prime} \tag{5.14}
\end{equation*}
$$

We aim to show $u^{\varepsilon}\left(p, t+\varepsilon^{2}\right)=u^{\varepsilon}\left(p^{\prime}, t+\varepsilon^{2}\right)$ for all $p, p^{\prime} \in \mathcal{H}$ satisfying the condition (5.14).
Since the dynamic programming principle (5.3) gives

$$
u^{\varepsilon}\left(p, t+\varepsilon^{2}\right)=\min _{v \in S_{h}^{1}} \max _{b= \pm 1} u^{\varepsilon}(p \cdot(\sqrt{2} \varepsilon b v), t)
$$

there exists $v \in S_{h}^{1}$ such that

$$
\begin{equation*}
u^{\varepsilon}\left(p, t+\varepsilon^{2}\right)=\max _{b= \pm 1} u^{\varepsilon}(p \cdot(\sqrt{2} \varepsilon b v), t) \tag{5.15}
\end{equation*}
$$

We claim that there is $v^{\prime} \in S_{h}^{1}$ such that the coordinates of $p \cdot(\sqrt{2} \varepsilon b v)$ and $p^{\prime} \cdot\left(\sqrt{2} \varepsilon b v^{\prime}\right)$ satisfy (5.14) as well. Indeed, as

$$
p \cdot(\sqrt{2} \varepsilon b v)=\left(p_{1}+\sqrt{2} \varepsilon b v_{1}, p_{2}+\sqrt{2} \varepsilon b v_{2}, p_{3}+\frac{1}{2} \sqrt{2} \varepsilon b\left(p_{1} v_{2}-p_{2} v_{1}\right)\right)
$$

and

$$
p^{\prime} \cdot\left(\sqrt{2} \varepsilon b v^{\prime}\right)=\left(p_{1}^{\prime}+\sqrt{2} \varepsilon b v_{1}^{\prime}, p_{2}^{\prime}+\sqrt{2} \varepsilon b v_{2}^{\prime}, p_{3}^{\prime}+\frac{1}{2} \sqrt{2} \varepsilon b\left(p_{1}^{\prime} v_{2}^{\prime}-p_{2}^{\prime} v_{1}^{\prime}\right)\right),
$$

we are looking for $v_{1}^{\prime}, v_{2}^{\prime} \in S_{h}^{1}$ such that

$$
\left\{\begin{array}{l}
\left(p_{1}^{\prime}+\sqrt{2} \varepsilon b v_{1}^{\prime}\right)^{2}+\left(p_{2}^{\prime}+\sqrt{2} \varepsilon b v_{2}^{\prime}\right)^{2}=\left(p_{1}+\sqrt{2} \varepsilon b v_{1}\right)^{2}+\left(p_{2}+\sqrt{2} \varepsilon b v_{2}\right)^{2} \\
p_{3}^{\prime}+\frac{1}{2} \sqrt{2} \varepsilon b\left(p_{1}^{\prime} v_{2}^{\prime}-p_{2}^{\prime} v_{1}^{\prime}\right)=p_{3}+\frac{1}{2} \sqrt{2} \varepsilon b\left(p_{1} v_{2}-p_{2} v_{1}\right) .
\end{array}\right.
$$

Since $p$ and $p^{\prime}$ satisfy (5.14), it suffices to solve the linear system

$$
\left\{\begin{array}{l}
p_{1}^{\prime} v_{1}^{\prime}+p_{2}^{\prime} v_{2}^{\prime}=p_{1} v_{1}+p_{2} v_{2}, \\
\\
-p_{2}^{\prime} v_{1}^{\prime}+p_{1}^{\prime} v_{2}^{\prime}=-p_{2} v_{1}+p_{1} v_{2} .
\end{array}\right.
$$

The problem is trivial if $p_{1}^{2}+p_{2}^{2}=\left(p_{1}^{\prime}\right)^{2}+\left(p_{2}^{\prime}\right)^{2}=0$. When $p_{1}^{2}+p_{2}^{2}=\left(p_{1}^{\prime}\right)^{2}+\left(p_{2}^{\prime}\right)^{2} \neq 0$, we get a unique pair of solutions

$$
\begin{aligned}
& v_{1}^{\prime}=\frac{1}{\left(p_{1}^{\prime}\right)^{2}+\left(p_{2}^{\prime}\right)^{2}}\left(\left(p_{1} p_{1}^{\prime}+p_{2} p_{2}^{\prime}\right) v_{1}+\left(p_{1}^{\prime} p_{2}-p_{1} p_{2}^{\prime}\right) v_{2}\right), \\
& v_{2}^{\prime}=\frac{1}{\left(p_{1}^{\prime}\right)^{2}+\left(p_{2}^{\prime}\right)^{2}}\left(\left(p_{1} p_{2}^{\prime}-p_{1}^{\prime} p_{2}\right) v_{1}+\left(p_{1} p_{1}^{\prime}+p_{2} p_{2}^{\prime}\right) v_{2}\right) .
\end{aligned}
$$

Thanks to the relation (5.14), it is easy to verify that $v^{\prime}=\left(v_{1}^{\prime}, v_{2}^{\prime}, 0\right) \in S_{h}^{1}$, i.e., $\left(v_{1}^{\prime}\right)^{2}+$ $\left(v_{2}^{\prime}\right)^{2}=1$. We complete the proof of the claim.

In view of the induction hypothesis, we obtain

$$
u^{\varepsilon}\left(p^{\prime} \cdot\left(\sqrt{2} \varepsilon b v^{\prime}\right), t\right)=u^{\varepsilon}(p \cdot(\sqrt{2} \varepsilon b v), t) \text { for both } b= \pm 1
$$

which, together with the dynamic programming (5.3) and (5.15), yields

$$
u^{\varepsilon}\left(p^{\prime}, t+\varepsilon^{2}\right) \leq \max _{b= \pm 1} u^{\varepsilon}\left(p^{\prime} \cdot\left(\sqrt{2} \varepsilon b v^{\prime}\right), t\right) \leq u^{\varepsilon}\left(p, t+\varepsilon^{2}\right) .
$$

We may similarly prove that $u^{\varepsilon}\left(p^{\prime}, t+\varepsilon^{2}\right) \geq u^{\varepsilon}\left(p, t+\varepsilon^{2}\right)$ and therefore $u^{\varepsilon}\left(p^{\prime}, t+\varepsilon^{2}\right)=$ $u^{\varepsilon}\left(p, t+\varepsilon^{2}\right)$ for all $p, p^{\prime} \in \mathcal{H}$ satisfying (5.14).

It follows from the definitions (1.5)-(1.6) of half relaxed limits that the same results for $\bar{u}$ and $\underline{u}$ hold.

Proposition 5.8. Assume that $u^{\varepsilon}$ satisfies the dynamic programming principle (5.3). Let $\bar{u}$ be the upper relaxed limit defined as in (1.5). Then $\bar{u}$ is a subsolution of (1.2).
Proof. Assume that there exists $(\hat{p}, \hat{t}) \in \mathcal{H} \times(0, \infty)$ and $\phi \in C^{2}(\mathcal{H} \times(0, \infty))$ such that $\bar{u}-\phi$ attains a strict maximum at $(\hat{p}, \hat{t})$. Then by definitions of $\bar{u}$, we may take a sequence, still indexed by $\varepsilon,\left(p^{\varepsilon}, t^{\varepsilon}\right) \in \mathcal{H} \times(0, \infty)$ such that $\left(p^{\varepsilon}, t^{\varepsilon}\right) \rightarrow(\hat{p}, \hat{t})$ and $u^{\varepsilon}\left(p^{\varepsilon}, t^{\varepsilon}\right) \rightarrow \bar{u}(\hat{p}, \hat{t})$ as $\varepsilon \rightarrow 0$ and

$$
\begin{equation*}
u^{\varepsilon}\left(p^{\varepsilon}, t^{\varepsilon}\right)-\phi\left(p^{\varepsilon}, t^{\varepsilon}\right)=\max _{B_{r}(\hat{p}, t)}\left(u^{\varepsilon}-\phi\right) \tag{5.16}
\end{equation*}
$$

Applying the dynamic programming principle (5.3) with $(p, t)=\left(p^{\varepsilon}, t^{\varepsilon}\right)$, we have

$$
u^{\varepsilon}\left(p^{\varepsilon}, t^{\varepsilon}\right)=\min _{v} \max _{b} u^{\varepsilon}\left(p^{\varepsilon} \cdot(\sqrt{2} \varepsilon b v), t^{\varepsilon}-\varepsilon^{2}\right),
$$

which, combined with (5.16), implies that

$$
\phi\left(p^{\varepsilon}, t^{\varepsilon}\right) \leq \min _{v} \max _{b} \phi\left(p^{\varepsilon} \cdot(\sqrt{2} \varepsilon b v), t^{\varepsilon}-\varepsilon^{2}\right)
$$

We next use the Taylor expansion for the right hand side at $\left(p^{\varepsilon}, t^{\varepsilon}\right)$ and obtain

$$
\begin{equation*}
\varepsilon^{2} \phi_{t}\left(p^{\varepsilon}, t^{\varepsilon}\right)-\min _{v} \max _{b}\left(\left\langle\sqrt{2} \varepsilon b v, \nabla \phi\left(p^{\varepsilon}, t^{\varepsilon}\right)\right\rangle+\varepsilon^{2}\left\langle\left(\nabla_{H}^{2} \phi\right)^{*}\left(p^{\varepsilon}, t^{\varepsilon}\right) v_{h}, v_{h}\right\rangle\right) \leq o\left(\varepsilon^{2}\right) \tag{5.17}
\end{equation*}
$$

where $v_{h}$ is the horizontal projection of $v$, i.e., $v_{h}=\left(v^{1}, v^{2}\right)$ for any $v=\left(v^{1}, v^{2}, v^{3}\right)$. Since $v=\left(v^{1}, v^{2}, 0\right)$, we may rewrite (5.17) as

$$
\begin{equation*}
\varepsilon^{2} \phi_{t}\left(p^{\varepsilon}, t^{\varepsilon}\right)-\min _{v} \max _{b}\left(\left\langle\sqrt{2} \varepsilon b v_{h}, \nabla_{H} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right)\right\rangle+\varepsilon^{2}\left\langle\left(\nabla_{H}^{2} \phi\right)^{*}\left(p^{\varepsilon}, t^{\varepsilon}\right) v_{h}, v_{h}\right\rangle\right) \leq o\left(\varepsilon^{2}\right) \tag{5.18}
\end{equation*}
$$

We discuss two cases:
Case A: $\nabla_{H} \phi(\hat{p}, \hat{t}) \neq 0$. Then $\nabla_{H} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right) \neq 0$ for all sufficiently small $\varepsilon>0$. Letting

$$
\tilde{v}=\frac{1}{\left|\nabla_{H} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right)\right|}\left(X_{2} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right),-X_{1} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right), 0\right)
$$

with

$$
\tilde{v}_{h}=\frac{1}{\left|\nabla_{H} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right)\right|}\left(X_{2} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right),-X_{1} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right)\right)
$$

we have from (5.17)

$$
\begin{equation*}
\left.\phi_{t}\left(p^{\varepsilon}, t^{\varepsilon}\right)-\left\langle\left(\nabla_{H}^{2} \phi\right)^{*}\left(p^{\varepsilon}, t^{\varepsilon}\right) \tilde{v}_{h}, \tilde{v}_{h}\right\rangle\right) \leq o(1) \tag{5.19}
\end{equation*}
$$

Noticing that

$$
\tilde{v}_{h} \otimes \tilde{v}_{h}=I-\frac{\nabla_{H} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right) \otimes \nabla_{H} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right)}{\left|\nabla_{H} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right)\right|^{2}}
$$

we are thus led from (5.19) to

$$
\begin{equation*}
\phi_{t}\left(p^{\varepsilon}, t^{\varepsilon}\right)-\operatorname{tr}\left[\left(I-\frac{\nabla_{H} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right) \otimes \nabla_{H} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right)}{\left|\nabla_{H} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right)\right|^{2}}\right)\left(\nabla_{H}^{2} \phi\right)^{*}\left(p^{\varepsilon}, t^{\varepsilon}\right)\right] \leq o(1) \tag{5.20}
\end{equation*}
$$

Sending $\varepsilon \rightarrow 0$, we get

$$
\phi_{t}(\hat{p}, \hat{t})-\operatorname{tr}\left[\left(I-\frac{\nabla_{H} \phi(\hat{p}, \hat{t}) \otimes \nabla_{H} \phi(\hat{p}, \hat{t})}{\left|\nabla_{H} \phi(\hat{p}, \hat{t})\right|^{2}}\right)\left(\nabla_{H}^{2} \phi\right)^{*}(\hat{p}, \hat{t})\right] \leq 0
$$

Case B: $\nabla_{H} \phi(\hat{p}, \hat{t})=0$. In this case, we have by Definition 3.2 that

$$
\begin{equation*}
\nabla_{H}^{2} \phi^{*}(\hat{p}, \hat{t})=0 \tag{5.21}
\end{equation*}
$$

If $\nabla_{H} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right) \neq 0$ for all $\varepsilon>0$, we may follow the same argument as in Case A, passing to the limit for (5.20) as $\varepsilon \rightarrow 0$ with an application of (5.21), and get

$$
\begin{equation*}
\phi_{t}(\hat{p}, \hat{t}) \leq 0 \tag{5.22}
\end{equation*}
$$

as desired.
If there exists a subsequence $\varepsilon_{j}$ such that $\nabla_{H} \phi\left(p^{\varepsilon_{j}}, t^{\varepsilon_{j}}\right)=0$ for all $j$, then it follows from (5.17) that

$$
\left.\phi_{t}\left(p^{\varepsilon_{j}}, t^{\varepsilon_{j}}\right)-\left\langle\left(\nabla_{H}^{2} \phi\right)^{*}\left(p^{\varepsilon_{j}}, t^{\varepsilon_{j}}\right) v_{h}, v_{h}\right\rangle\right) \leq o(1) \text { for some } v
$$

which again implies (5.22) as the limit when $\varepsilon_{j} \rightarrow 0$.
Proposition 5.9. Assume that $u^{\varepsilon}$ satisfies the dynamic programming principle (5.3). Let $\underline{u}$ be the lower relaxed limit defined as in (1.6). Then $\underline{u}$ is a supersolution of (1.2).

In order to facilitate the proof, let us present an elementary result.

Lemma 5.10 (Lemma 4.1 in [GL]). Suppose $\xi$ is a unit vector in $\mathbb{R}^{2}$ and $X$ is a real symmetric $2 \times 2$ matrix, then there exists a constant $M>0$ that depends only on the norm of $X$, such that for any unit vector $v \in \mathbb{R}^{2}$,

$$
\begin{equation*}
\left|\left\langle X \xi^{\perp}, \xi^{\perp}\right\rangle-\langle X v, v\rangle\right| \leq M|\langle\xi, v\rangle|, \tag{5.23}
\end{equation*}
$$

where $\xi^{\perp}$ denotes a unit orthonormal vector of $\xi$.
Proof. Let $\cos \theta=\left\langle\xi^{\perp}, v\right\rangle$ and $\sin \theta=\langle\xi, v\rangle$. Then we have

$$
\begin{aligned}
& \left|\left\langle X \xi^{\perp}, \xi^{\perp}\right\rangle-\langle X v, v\rangle\right| \\
= & \left|\operatorname{tr}\left(X\left(\xi^{\perp} \otimes \xi^{\perp}-v \otimes v\right)\right)\right| \\
\leq & \|X\|\left\|\xi^{\perp} \otimes \xi^{\perp}-\left(\xi \sin \theta+\xi^{\perp} \cos \theta\right) \otimes\left(\xi \sin \theta+\xi^{\perp} \cos \theta\right)\right\| \\
= & \|X\|\left\|\sin ^{2} \theta \xi^{\perp} \otimes \xi^{\perp}-\sin \theta \cos \theta\left(\xi \perp \xi^{\perp}+\xi^{\perp} \otimes \xi\right)\right\| \\
\leq & M|\sin \theta|
\end{aligned}
$$

where $M>0$ depends on $\|X\|$.
We refer the reader to [L, Lemma 2.3] for a higher dimensional extension of this lemma.
Proof of Proposition 5.9. Assume that there exists $(\hat{p}, \hat{t}) \in \mathcal{H} \times(0, \infty)$ and $\phi \in C^{2}(\mathcal{H} \times$ $(0, \infty))$ such that $\underline{u}-\phi$ attains a strict minimum at $(\hat{p}, \hat{t})$. We may again take a sequence $\left(p^{\varepsilon}, t^{\varepsilon}\right) \in \mathcal{H} \times(0, \infty)$ such that $\left(p^{\varepsilon}, t^{\varepsilon}\right) \rightarrow(\hat{p}, \hat{t})$ and $u^{\varepsilon}\left(p^{\varepsilon}, t^{\varepsilon}\right) \rightarrow \underline{u}(\hat{p}, \hat{t})$ as $\varepsilon \rightarrow 0$ and

$$
\begin{equation*}
u^{\varepsilon}\left(p^{\varepsilon}, t^{\varepsilon}\right)-\phi\left(p^{\varepsilon}, t^{\varepsilon}\right)=\min _{B_{r}(\hat{p}, \hat{t})}\left(u^{\varepsilon}-\phi\right) \tag{5.24}
\end{equation*}
$$

Applying the dynamic programming principle (5.3) with $(p, t)=\left(p^{\varepsilon}, t^{\varepsilon}\right)$, we have

$$
u^{\varepsilon}\left(p^{\varepsilon}, t^{\varepsilon}\right)=\min _{v} \max _{b} u^{\varepsilon}\left(p^{\varepsilon} \cdot \sqrt{2} \varepsilon b v, t^{\varepsilon}-\varepsilon^{2}\right) .
$$

It then follows from (5.24) that

$$
\phi\left(p^{\varepsilon}, t^{\varepsilon}\right) \geq \min _{v} \max _{b} \phi\left(p^{\varepsilon} \cdot \sqrt{2} \varepsilon b v, t^{\varepsilon}-\varepsilon^{2}\right) .
$$

As an analogue of (5.18), the Taylor expansion at $\left(p^{\varepsilon}, t^{\varepsilon}\right)$ yields

$$
\begin{equation*}
\phi_{t}\left(p^{\varepsilon}, t^{\varepsilon}\right)-\min _{v} \max _{b}\left(\frac{1}{\varepsilon}\left\langle\sqrt{2} b v_{h}, \nabla_{H} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right)\right\rangle+\left\langle\left(\nabla_{H}^{2} \phi\right)^{*}\left(p^{\varepsilon}, t^{\varepsilon}\right) v_{h}, v_{h}\right\rangle\right) \geq o(1), \tag{5.25}
\end{equation*}
$$

as $\varepsilon \rightarrow 0$.
We again divide our discussion into two situations.
Case A: $\nabla_{H} \phi(\hat{p}, \hat{t}) \neq 0$. Then $\nabla \phi\left(p^{\varepsilon}, t^{\varepsilon}\right) \neq 0$ for all sufficiently small $\varepsilon>0$. We adopt Lemma 5.10 and get

$$
\begin{aligned}
& \max _{b}\left(\frac{1}{\varepsilon}\left\langle\sqrt{2} b v_{h}, \nabla_{H} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right)\right\rangle+\left\langle\left(\nabla_{H}^{2} \phi\right)^{*}\left(p^{\varepsilon}, t^{\varepsilon}\right) v_{h}, v_{h}\right\rangle\right) \\
\leq & \left\langle\left(\nabla_{H}^{2} \phi\right)^{*}\left(p^{\varepsilon}, t^{\varepsilon}\right) \tilde{v}_{h}, \tilde{v}_{h}\right\rangle+\left(M+\frac{\sqrt{2}}{\varepsilon}\right)\left|\left\langle v_{h}, \phi_{H}\left(p^{\varepsilon}, t^{\varepsilon}\right)\right\rangle\right|,
\end{aligned}
$$

where

$$
\tilde{v}_{h}=\frac{1}{\left|\nabla_{H} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right)\right|}\left(X_{2} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right),-X_{1} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right)\right),
$$

as given in the proof of Proposition 5.8. It is now clear, by taking $v_{h}=\tilde{v}_{h}$, that

$$
\begin{aligned}
& \min _{v} \max _{b}\left(\frac{1}{\varepsilon}\left\langle\sqrt{2} b v_{h}, \nabla_{H} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right)\right\rangle+\left\langle\left(\nabla_{H}^{2} \phi\right)^{*}\left(p^{\varepsilon}, t^{\varepsilon}\right) v_{h}, v_{h}\right\rangle\right) \\
\leq & \left\langle\left(\nabla_{H}^{2} \phi\right)^{*}\left(p^{\varepsilon}, t^{\varepsilon}\right) \tilde{v}_{h}, \tilde{v}_{h}\right\rangle,
\end{aligned}
$$

which implies through (5.25) that

$$
\phi_{t}\left(p^{\varepsilon}, t^{\varepsilon}\right)-\left\langle\left(\nabla_{H}^{2} \phi\right)^{*}\left(p^{\varepsilon}, t^{\varepsilon}\right) \tilde{v}_{h}, \tilde{v}_{h}\right\rangle \geq o(1)
$$

Letting $\varepsilon \rightarrow 0$, we obtain

$$
\phi_{t}(\hat{p}, \hat{t})-\operatorname{tr}\left[\left(I-\frac{\nabla_{H} \phi(\hat{p}, \hat{t}) \otimes \nabla_{H} \phi(\hat{p}, \hat{t})}{\left|\nabla_{H} \phi(\hat{p}, \hat{t})\right|^{2}}\right)\left(\nabla_{H}^{2} \phi\right)^{*}(\hat{p}, \hat{t})\right] \geq 0
$$

Case B: $\nabla_{H} \phi(\hat{p}, \hat{t})=0$. We may further assume (5.21) again in this case. We may apply the same argument above and get

$$
\begin{equation*}
\phi_{t}(\hat{p}, \hat{t}) \geq 0 \tag{5.26}
\end{equation*}
$$

provided that $\nabla_{H} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right) \neq 0$ for all $\varepsilon>0$. It remains to show (5.26) when there is a subsequence $\varepsilon_{j}$ such that $\nabla_{H} \phi\left(p^{\varepsilon_{j}}, t^{\varepsilon_{j}}\right)=0$. By (5.25), we have on this occasion

$$
\left.\phi_{t}\left(p^{\varepsilon_{j}}, t^{\varepsilon_{j}}\right)-\left\langle\left(\nabla_{H}^{2} \phi\right)^{*}\left(p^{\varepsilon_{j}}, t^{\varepsilon_{j}}\right) v_{h}, v_{h}\right\rangle\right) \geq o(1) \text { for some } v
$$

Sending $\varepsilon \rightarrow 0$, we get (5.26).
We are now in a position to prove Theorem 5.1.
Proof of Theorem 5.1. In terms of Proposition 5.7, Proposition 5.8 and Proposition 5.9, $\bar{u}$ and $\underline{u}$ are respectively a subsolution and a supersolution of (1.2) that are axisymmetric with respect to the vertical axis. For any $T>0, \bar{u}(p, t)$ and $\underline{u}(p, t)$ are constant outside a compact set of $\mathcal{H}$ for all $t \in[0, T]$, owing to Proposition 5.4. Also, since $\bar{u}(p, 0) \leq u_{0}(p)$ and $\underline{u}(p, 0) \geq u_{0}(p)$ for all $p \in \mathcal{H}$, we may apply Theorem 1.1 to get $\bar{u} \leq \underline{u}$ in $\mathcal{H} \times[0, T]$. As it is obvious that $\bar{u} \geq \underline{u}$, we get $\bar{u}=\underline{u}$ in $\mathcal{H} \times[0, T]$ with $u(\cdot, 0)=u_{0}(\cdot)$. In conclusion, $u=\bar{u}=\underline{u}$ is the unique continuous solution of (MCF) and the locally uniform convergence $u^{\varepsilon} \rightarrow u$ follows immediately.

## 6. Stability

The following stability result is standard in the theory of viscosity solutions.
Theorem 6.1 (Stability under the uniform convergence). Let $u^{\varepsilon}$ be solutions of (1.2) and $u^{\varepsilon} \rightarrow u$ locally uniformly in $\mathcal{H} \times[0, \infty)$. Then $u$ is also a solution of (1.2).
Lemma 6.2. If $u^{\varepsilon}$ is a subsolution (resp., supersoution) of (1.2) for all small $\varepsilon>0$, then

$$
\bar{u}=\limsup _{\varepsilon \rightarrow \infty}^{*} u^{\varepsilon} \quad\left(\text { resp., } \underline{u}=\liminf _{\varepsilon \rightarrow \infty} u^{\varepsilon}\right)
$$

is also a subsolution (resp., supersoution) of (1.2).
Proof. Suppose there exists $\phi \in C^{2}(\mathcal{H} \times[0, \infty))$ and $(\hat{p}, \hat{t}) \in \mathcal{H} \times(0, \infty)$ such that $\bar{u}-\phi$ attains a strict maximum at $(\hat{p}, \hat{t})$. Then by the convergence of maximizers as shown in [G, Lemma 2.2.5], we can take subsequences of $p^{\varepsilon}, t^{\varepsilon}$ and $u^{\varepsilon}$, still indexed by $\varepsilon$, satisfying $\left(p^{\varepsilon}, t^{\varepsilon}\right) \rightarrow(\hat{p}, \hat{t})$ as $\varepsilon \rightarrow 0$ and

$$
\left(u^{\varepsilon}-\phi\right)\left(p^{\epsilon}, t^{\epsilon}\right)=\max _{\mathcal{H} \times[0, \infty)}\left(u^{\varepsilon}-\phi\right) .
$$

We discuss two cases.
Case 1: $\nabla_{H} \phi(\hat{p}, \hat{t}) \neq 0$. Then $\nabla_{H} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right) \neq 0$ for all $\varepsilon>0$ small. We apply Definition 3.2 and get

$$
\phi_{t}-\operatorname{tr}\left[\left(I-\frac{\nabla_{H} \phi \otimes \nabla_{H} \phi}{\left|\nabla_{H} \phi\right|^{2}}\right)\left(\nabla_{H}^{2} \phi\right)^{*}\right] \leq 0 \text { at }\left(p^{\varepsilon}, t^{\varepsilon}\right) .
$$

Sending $\varepsilon \rightarrow 0$, we get the desired inequality

$$
\phi_{t}-\operatorname{tr}\left[\left(I-\frac{\nabla_{H} \phi \otimes \nabla_{H} \phi}{\left|\nabla_{H} \phi\right|^{2}}\right)\left(\nabla_{H}^{2} \phi\right)^{*}\right] \leq 0 \text { at }(\hat{p}, \hat{t}) .
$$

Case 2: $\nabla_{H} \phi(\hat{p}, \hat{t})=0$. Then by Definition 3.2 we only need to discuss the situation when $\left(\nabla_{H}^{2} \phi\right)^{*}(\hat{p}, \hat{t})=0$ also holds. Hence, $\nabla_{H} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right) \rightarrow 0$ and $\left(\nabla_{H}^{2} \phi\right)^{*}\left(p^{\varepsilon}, t^{\varepsilon}\right) \rightarrow 0$ as $\varepsilon \rightarrow 0$. If there exists a subsequence $\varepsilon_{j}$ such that $\nabla_{H} \phi\left(p^{\varepsilon_{j}}, t^{\varepsilon_{j}}\right) \neq 0$, then we have

$$
\phi_{t}-\operatorname{tr}\left[\left(I-\frac{\nabla_{H} \phi \otimes \nabla_{H} \phi}{\left|\nabla_{H} \phi\right|^{2}}\right)\left(\nabla_{H}^{2} \phi\right)^{*}\right] \leq 0 \text { at }\left(p^{\varepsilon_{j}}, t^{\varepsilon_{j}}\right) .
$$

Passing to the limit $j \rightarrow \infty$, we obtain $\phi_{t}(\hat{p}, \hat{t}) \leq 0$.
If, on the other hand, $\nabla_{H} \phi\left(p^{\varepsilon}, t^{\varepsilon}\right)=0$ for all $\varepsilon>0$ small, then we get $\phi_{t}\left(p^{\varepsilon}, t^{\varepsilon}\right) \leq 0$ and the limit immediately yields $\phi_{t}(\hat{p}, \hat{t}) \leq 0$, which completes our proof.

One may similarly prove that $\underline{u}=\liminf _{* \varepsilon \rightarrow 0} u^{\varepsilon}$ is a supersolution provided that $u^{\varepsilon}$ is a supersolution for all $\varepsilon>0$ small.

Proof of Theorem 6.1. Let

$$
\bar{u}=\underset{\varepsilon \rightarrow 0}{\limsup }{ }^{*} u^{\varepsilon} \text { and } \underline{u}=\liminf _{\varepsilon \rightarrow 0} u^{\varepsilon} .
$$

Then in virtue of Lemma 6.2, $\bar{u}$ is a subsolution of (1.2) and $\underline{u}$ is a supersolution of (1.2). Noting that $u^{\varepsilon} \rightarrow u$ locally uniformly, we must have $u=\bar{u}=\underline{u}$ and therefore $u$ is a solution of (1.2).

## 7. Properties of the evolution

We have shown that there is a unique solution $u$ of (MCF) for any given continuous function $u_{0}$ which is axisymmetric with respect to the vertical axis and attains constant value outside a compact set. Let us turn to discuss the surface evolution described by the level-set equation (MCF). More precisely, given an axisymmetric compact surface $\Gamma_{0} \subset \mathcal{H}$, we choose $u_{0} \in C(\mathcal{H})$ such that it is axisymmetric constant outside a compact set and satisfies

$$
\begin{equation*}
\Gamma_{0}=\left\{p: \mathcal{H}: u_{0}(p)=0\right\} . \tag{7.1}
\end{equation*}
$$

We then solve (MCF) for the unique solution $u$ and get the surface

$$
\begin{equation*}
\Gamma_{t}=\{p \in \mathcal{H}: u(p, t)=0\} \text { for any } t \geq 0 . \tag{7.2}
\end{equation*}
$$

In what follows, we first show that the surface represented by the level-set $\Gamma_{t}$ of $u$ does not depend on the particular choice of $u_{0}$.

### 7.1. Uniqueness of the surface evolution.

Theorem 7.1 (Invariance). Assume that $\theta: \mathbb{R} \rightarrow \mathbb{R}$ is continuous. If $u$ is a solution of (1.2). Then $w=\theta \circ u$ is also a solution of (1.2).

Proof. We prove the theorem in several steps.
Step 1. We first give the proof in the case that $\theta \in C^{2}(\mathbb{R})$ and $\theta^{\prime}>0$. Suppose that there is $\phi \in C^{2}(\mathcal{H} \times[0, \infty))$ and $(\hat{p}, \hat{t}) \in \mathcal{H} \times(0, \infty)$ such that $\theta \circ u-\phi$ attains a maximum at $(\hat{p}, \hat{t})$. Then it is clear that $u-h(\phi)$ attains a maximum at $(\hat{p}, \hat{t})$, where $h=\theta^{-1} \in C^{2}(\mathbb{R})$ with $h^{\prime}>0$. Denote $\psi=h(\phi)$. Since $u$ is a subsolution of (1.2), we have

$$
\psi_{t}-\operatorname{tr}\left[\left(I-\frac{\nabla_{H} \psi \otimes \nabla_{H} \psi}{\left|\nabla_{H} \psi\right|^{2}}\right)\left(\nabla_{H}^{2} \psi\right)^{*}\right] \leq 0 \text { at }(\hat{p}, \hat{t}) .
$$

Note that $\psi_{t}=h^{\prime} \phi_{t}, \nabla_{H} \psi=h^{\prime} \nabla_{H} \phi$ and

$$
\left(\nabla_{H}^{2} \psi\right)^{*}=h^{\prime \prime} \nabla_{H} \phi \otimes \nabla_{H} \phi+h^{\prime}\left(\nabla_{H}^{2} \phi\right)^{*} .
$$

It follows that

$$
\phi_{t}-\operatorname{tr}\left[\left(I-\frac{\nabla_{H} \phi \otimes \nabla_{H} \phi}{\left|\nabla_{H} \phi\right|^{2}}\right)\left(\nabla_{H}^{2} \phi\right)^{*}\right] \leq 0 \text { at }(\hat{p}, \hat{t}) \text {, }
$$

which shows that $\theta \circ u$ is a subsolution of (1.2). An analogue of this argument yields that $\theta \circ u$ is also a supersolution.

We also claim that $\theta \circ u$ remains being a solution when $\theta \in C^{2}(\mathbb{R})$ and $\theta^{\prime}<0$. Indeed, when $\theta$ is a decreasing function, $-\theta$ is increasing. We obtain that $-\theta \circ u$ is a solution of (1.2). Thanks to the fact that the mean curvature flow is orientation-free or (1.2) is homogeneous in all of the derivatives, we easily see that $\theta \circ u$ is a solution as well. In particular, we note that $-u$ is a solution when $u$ is a solution.
Step 2. We generalize the consequence obtained in Step 1 for a continuous nondecreasing or nonincreasing function. Indeed, for any continuous nondecreasing function $\theta$, we may take $\theta_{n} \in C^{2}(\mathbb{R})$ with $\theta_{n}^{\prime}>0$ for all $n=1,2, \ldots$ such that

$$
\underset{n \rightarrow \infty}{\limsup ^{*}} \theta_{n} \circ u=\theta \circ u
$$

We refer the reader to [G, Lemma 4.2.3] for details about the construction of $\theta_{n}$. Since $\theta_{n} \circ u$ is a solution of (1.2) for all $n$, as shown in Step $1, \theta \circ u$ is a subsolution, due to Lemma 6.2.

To show that $\theta \circ u$ is a supersolution, we define $\tilde{\theta}(x)=\theta(-x)$ for any $x \in \mathbb{R}$ and observe that $\theta(u)=\tilde{\theta}(-u)$. Since $\tilde{\theta}$ is nonincreasing and $-u$ is a solution, we may apply a symmetric version of [G, Lemma 4.2.3] to get $\theta(u)=\tilde{\theta}(-u)$ is a supersoluiton.

When $\theta$ is a continuous nonincreasing function, $-\theta$ is nondecreasing. We apply again the homogeneity of (1.2) to obtain that $\theta \circ u$ is a solution. Since the verification of definition of (1.2) is pointwise, one can further relax the monotonicity condition on $\theta$ to a local monotonicity condition.

To conclude this step, we notice that $\max \{\min \{u, C\}-C\}$ is a solution for any $C>0$ provided that $u$ is a solution.
Step 3. We finally discuss the situation when $\theta$ is assumed to be continuous only. By Theorem 6.1, it suffices to discuss the bounded function $\max \{\min \{u, C\}-C\}$ instead of $u$ for arbitrarily large $C>0$. We approximate $\theta$ uniformly by polynomials $\theta_{m}$ in $[-C-1, C+1]$. Since polynomials only have finitely many maximizers and minimizers, we may also assume each $\theta_{m}$ is constant near all of its local maximizers and minimizers.

In fact, if, for instance, $\theta_{m}$ attains a local maximum at $x_{0} \in \mathbb{R}$, we take $\min \left\{\theta_{m}(x), \theta\left(x_{0}\right)-\right.$ $\left.\varepsilon_{m}\right\}$, where $\varepsilon_{m}>0$ is sufficiently small $\left(\varepsilon_{m} \rightarrow 0\right.$ as $\left.m \rightarrow \infty\right)$ such that $\theta_{m}$ is continuous.

Now $\theta_{m}$ is locally nonincreasing or nondecreasing. We apply the result in Step 2 and find that $\theta_{m} \circ u$ is a solution of (1.2). Since $\theta_{m} \rightarrow \theta$ uniformly, by the stability result given in Theorem 6.1, we see that $\theta \circ u$ is a solution by sending $m \rightarrow \infty$.

An immediate consequence of the theorem above is that our generalized surface evolution does not depend on the choice of the initial level-set function $u_{0}$.

Corollary 7.2 (Independence of the choice of the initial function). Suppose that $u_{0}$ and $\tilde{u}_{0}$ are continuous functions in $\mathcal{H}$ axisymmetric about the vertical axis and are constant outside a compact set $K \subset \mathcal{H}$. Let $\Gamma_{0}=\left\{p \in \mathcal{H}: u_{0}(p)=0\right\}=\left\{p \in \mathcal{H}: \tilde{u}_{0}(p)=0\right\}$ be bounded. Let $u$ and $\tilde{u}$ be the unique continuous solutions of (1.2) with the initial conditions $u_{0}$ and $\tilde{u}_{0}$ respectively. For any $t \geq 0$, set

$$
\Gamma_{t}=\{p \in \mathcal{H}: u(p, t)=0\} \text { and } \tilde{\Gamma}_{t}=\{p \in \mathcal{H}: \tilde{u}(p, t)=0\} .
$$

Then $\Gamma_{t}=\tilde{\Gamma}_{t}$ for all $t \geq 0$.
Proof. We follow the proof of [ES, Theorem 5.1]. It is obvious, from Theorem 1.1 and Theorem 5.1, that $u$ and $\tilde{u}$ are axisymmetric about the vertical axis.

We may assume $u_{0} \geq 0$ without changing the zero level set of $u_{0}$, since $|u|$ is a solution of (1.2) with the initial condition $u(p, 0)=\left|u_{0}\right|$ by Theorem 7.1. Similarly, let us also assume that $\tilde{u}_{0} \geq 0$ and $\tilde{u} \geq 0$.

For any $k=1,2, \ldots$ let $E_{0}=\emptyset$ and $E_{k}=\left\{p \in \mathcal{H}: u_{0}(p)>l / k\right\}$ such that $E_{k}$ is nondecreasing and $\mathcal{H} \subset \Gamma_{0}=\cup_{k} E_{k}$. Define

$$
a_{k}=\max _{\mathcal{H} \backslash E_{k-1}} \tilde{u}_{0}(k=1,2, \ldots)
$$

Then we have $\lim _{k \rightarrow \infty} a_{k}=0$. We then construct a continuous function $\theta$ satisfying $\theta(0)=0, \theta(1 / k)=a_{k}$ for all $k$ and $\theta=a_{1}$ in $[1, \infty)$.

Now it is clear that $\theta \circ u$ is an axisymmetric solution of (1.2) with initial data $\theta \circ u_{0}$, again due to Theorem 7.1. By our construction of $\theta$, we easily see that $\theta \circ u_{0} \geq \tilde{u}_{0}$. Applying Theorem 1.1 for all $T>0$, we get $\theta \circ u \geq \tilde{u}$. This means that $\Gamma_{t} \subset \tilde{\Gamma}_{t}$ for any $t \geq 0$. Indeed, for any $p \in \Gamma_{t}$, we have $u(p, t)=0$, which implies that $\theta \circ u(p, t)=0$ and therefore $\tilde{u}(p, t)=0$.

We conclude the proof by similarly showing the inclusion $\tilde{\Gamma}_{t} \subset \Gamma_{t}$ for any $t \geq 0$.
7.2. Finite time extinction. We give a simple geometric property of the mean curvature flow. The following result shows that an axisymmetric compact surface evolving by its mean curvature shrinks and disappears in finite time.

Theorem 7.3 (Finite time extinction for bounded evolution). Suppose that $\left\{\Gamma_{t}\right\}_{t \geq 0}$ denotes an axisymmetric surface evolution of the mean curvature flow. If $\Gamma_{0} \subset B_{r}$, for $r>0$, then $\Gamma_{t}=\emptyset$ when $t>r^{2} / \sqrt{12}$.

Proof. We may take an axisymmetric $u_{0} \in C(\mathcal{H})$ with a constant value $C>0$ outside $B_{r}$ satisfying (7.1) and

$$
u_{0} \geq \min \left\{|p|^{4}-r^{4}, C\right\}
$$

Taking

$$
w(p, t)=\left(p_{1}^{2}+p_{2}^{2}\right)^{2}+12 t\left(p_{1}^{2}+p_{2}^{2}\right)+16 p_{3}^{2}+12 t^{2}
$$

as in (3.13), we easily see that $w^{C}(p, t):=\min \left\{w(p, t)-r^{4}, C\right\}$ is a solution of (1.2) with initial data $w^{C}(p, 0)=\min \left\{|p|^{4}-r^{4}, C\right\}$, by Theorem 7.1 with $\theta(x)=\min \{x, C\}$. We are therefore led to $u \geq w-r^{4}$ by Theorem 1.1.

It is clear that $w^{C}(p, t)>0$ when $t>r^{2} / \sqrt{12}$ for all $p \in \mathcal{H}$, which implies that $u>0$ when $t>r^{2} / \sqrt{12}$. Hence $\Gamma_{t}$ defined in (7.2) is empty when $t>r^{2} / \sqrt{12}$. Note that the conclusion does not depend on the particular choice of $u_{0}$, as explained in Corollary 7.2.

Remark 7.1. Theorem 7.3 indicates that a bounded axisymmetric mean curvature flow encounters singularities at a certain time $T>0$.

Remark 7.2. The following result stronger than Theorem 7.3 holds: For any continuous solution $u$ of (MCF) with zero level set $\Gamma_{t}$ for any $t \geq 0$, if $\Gamma_{0} \subset B_{r}$ with some $r>0$, then $\Gamma_{t}=\emptyset$ when $t>r^{2} / \sqrt{12}$. Here we do not need to assume the axial symmetry of $\Gamma_{0}$ but we must specify the solution $u$ since it is not known in general whether or not $\Gamma_{t}$ depends on the choice of $u_{0}$.

Definition 7.1. We say $T \geq 0$ is the extinction time of the mean curvature flow $\Gamma_{t}$ in the Heisenberg group, if $\Gamma_{t} \neq \emptyset$ when $t \leq T$ and $\Gamma_{t}=\emptyset$ when $t>T$.

We next proceed to investigate the asymptotic profile after normalization for a sphere in the Heisenberg group. It is well-known that in the Euclidean space any normalized compact convex surface converges to a sphere as $t$ tends to the extinction time [H]. However, the normalized curvature flow from a sphere of radius $r$ in the Heisenberg group looks like an ellipsoid

$$
\begin{equation*}
E_{T}:=\left\{P \in \mathcal{H}: 12 T\left(P_{1}^{2}+P_{2}^{2}\right)+16 P_{3}^{2}=1\right\} \tag{7.3}
\end{equation*}
$$

at the extinction time $T=r^{2} / \sqrt{12}$.
Proposition 7.4. Suppose that $\Gamma_{t} \subset \mathcal{H}(t \geq 0)$ is the horizontal mean curvature flow as defined in (7.2) with $\Gamma_{0}=\{p \in \mathcal{H}:|p|=r\}$, where $r>0$ is a given radius. Then the extinction time $T=r^{2} / \sqrt{12}$ and the normalized flow $\Gamma_{t} / \sqrt{r^{4}-12 t^{2}} \rightarrow E_{T}$ as $t \rightarrow T$, where $E_{T}$ is given in (7.3).

Proof. We take

$$
w^{C}(p, t)=\min \left\{\left(p_{1}^{2}+p_{2}^{2}\right)^{2}+12 t\left(p_{1}^{2}+p_{2}^{2}\right)+16 p_{3}^{2}+12 t^{2}-r^{4}, C\right\}
$$

with $C>0$. It is easily seen that $w^{C}(p, 0)=0$ if and only if $p \in \Gamma_{0}$. We have also shown that $w^{C}$ is a solution of (1.2). We track the evolution by setting $\Gamma_{t}=\left\{p \in \mathcal{H}: w^{C}(p, t)=\right.$ $0\}$ for all $t \geq 0$. It is clear that $\Gamma_{t}=\emptyset$ when $t>r^{2} / \sqrt{12}$ and $\Gamma_{t} \neq \emptyset$ when $t \leq r^{2} / \sqrt{12}$. For any $p(t)=\left(p_{1}(t), p_{2}(t), p_{3}(t)\right) \in \Gamma_{t}$, we have

$$
\begin{equation*}
12 t\left(p_{1}^{2}(t)+p_{2}^{2}(t)\right)+16 p_{3}^{2}(t) \leq r^{4}-12 t^{2} \tag{7.4}
\end{equation*}
$$

We normalize the flow by letting $P(t)=p(t) / \sqrt{r^{4}-12 t^{2}}$ for any $p(t) \in \Gamma_{t}$. Then (7.4) is written as

$$
\begin{equation*}
12 t\left(P_{1}^{2}(t)+P_{2}^{2}(t)\right)+16 P_{3}^{2}(t) \leq 1 \tag{7.5}
\end{equation*}
$$

By setting

$$
U(P, t)=\left(P_{1}^{2}+P_{2}^{2}\right)^{2}\left(r^{4}-12 t^{2}\right)+12 t\left(P_{1}^{2}+P_{2}^{2}\right)+16 P_{3}^{2}-1
$$

we get

$$
\left.0=\frac{1}{r^{4}-12 t^{2}} w^{C}\left(\sqrt{r^{4}-12 t^{2}} P(t), t\right)=U(P(t), t)\right) .
$$

Sending the limit as $t \rightarrow T$ with (7.5) taken into account, we obtain

$$
12 T\left(P_{1}^{2}(T)+P_{2}^{2}(T)\right)+16 P_{3}^{2}(T)=1
$$

for the limit $P(T)$ of any subsequence of $P(t)$ as $t \rightarrow \infty$. The consequence above amounts to saying that the limit of the set $\Gamma_{t} / \sqrt{r^{4}-12 t^{2}}$ is contained in $E_{T}$.

On the other hand, for any $P=\left(P_{1}, P_{2}, P_{3}\right) \in E_{T}$, we have

$$
w^{C}\left(\sqrt{r^{4}-12 t^{2}} \lambda P\right)=\left(r^{4}-12 t^{2}\right) W(\lambda, P, t)
$$

where $\lambda>0$ and

$$
W(\lambda, P, t)=\left(\lambda^{4}\left(P_{1}^{2}+P_{2}^{2}\right)^{2}\left(r^{4}-12 t^{2}\right)+\lambda^{2}-1+12 \lambda^{2}(t-T)\left(P_{1}^{2}+P_{2}^{2}\right) .\right.
$$

One may take $\lambda(t)>0$ such that $w^{C}\left(\sqrt{r^{4}-12 t^{2}} \lambda(t) P, t\right)=0$; in other words, $\lambda(t) P \in$ $\Gamma_{t} / \sqrt{r^{4}-12 t^{2}}$. Moreover, $\lambda(t) \rightarrow 1$ as $t \rightarrow T$, which implies that $P$ belongs to the limit of a sequence of elements in $\Gamma_{t} / \sqrt{r^{4}-12 t^{2}}$.

In conclusion, we obtain $\Gamma_{t} / \sqrt{r^{4}-12 t^{2}} \rightarrow E_{T}$ as $t \rightarrow T$.
We stress that this result is very different from that in the Euclidean space. The normalized asymptotic shape of horizontal mean curvature flow in the Heisenberg group starting from a ball is an ellipsoid. Moreover, the shape of the ellipsoid depends on the extinction time $T$ and therefore the size of the initial surface. It would be interesting to show this result for a general compact and convex initial surface.

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