

HOMOGENIZATION OF HAMILTON-JACOBI EQUATIONS IN CARNOT GROUPS

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ABSTRACT. We study an homogenization problem for Hamilton-Jacobi equations in the geometry of Carnot Groups. The tiling and the corresponding notion of periodicity are compatible with the dilatations of the Group and use the Lie bracket generating property.

1. BACKGROUND AND MOTIVATIONS

Consider a Hamilton-Jacobi equation:

$$u + H(\xi, \nabla u) = 0 \quad \text{in } \mathbb{R}^N$$

where the Hamiltonian $H(\xi, p) : \mathbb{R}^N \times \mathbb{R}^{N \times N} \rightarrow \mathbb{R}$ is not coercive in p . The lack of coerciveness of the Hamiltonian can be overcome by changing the underlying geometry with a suitable family of vector fields. More precisely, we are able to consider the case when $H(\xi, p) = H(\xi, \sigma(\xi)q)$, where $\sigma(\xi)$ is a $m \times N$ matrix, $m < N$, H is coercive in q . Here the rows of the matrix $\sigma(\xi)$ will be considered as coefficients of vector fields satisfying Hörmander condition and that generate a Carnot Group, therefore $\sigma(\xi)\nabla u$ will be the horizontal gradient in a Carnot Group, denoted by $D_h u$, see section 2.

We shall consider homogenizations problems for Hamilton-Jacobi equations of the form :

$$(1.1) \quad u^\varepsilon(\xi) + H\left(\xi, \frac{\xi}{\varepsilon}, D_h u^\varepsilon(\xi)\right) = 0$$

in a Carnot Group G , where the Hamiltonian H is G -periodic in the second variable and $\frac{\xi}{\varepsilon}$ will be interpreted in the geometry of the group.

Let us revise first some key results in the Euclidean setting, for \mathbb{Z}^N -periodic Hamiltonians. The pioneering paper on homogenization of Hamilton-Jacobi equations is due to P.L. Lions, Papanicolau and Varadhan. They proved that the asymptotic behaviour, as ε tends to zero, of the solutions u^ε is governed by the equation:

$$(1.2) \quad u(\xi) + \bar{H}(\xi, Du(\xi)) = 0$$

The effective Hamiltonian \bar{H} is obtained by solving a cell problem:

$$(1.3) \quad H(x, \xi, p + Dv(\xi)) = \lambda$$

for every (x, p) fixed and putting $\bar{H}(x, p) = \lambda$.

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Evans ([9], [10]) developed the perturbed test functions method for homogenization problems in the framework of the theory of viscosity solutions. The combined use of comparison results and perturbed test functions allow to establish, in a simple way, uniform estimates in the sup norm for the solutions u^ε to the solution of the homogenized equation.

The question of estimating the rate of the uniform convergence of the solutions u^ε to the solution of the homogenized equation has been tackled up by Capuzzo Dolcetta and Ishii([8]).

For recent developments in the framework of stochastic control and second-order Hamilton-Jacobi-Bellmann equations, see the papers [1],[2]. Birindelli and Wigniolle have studied the problem in the Heisenberg Group, [6], giving also an estimate of the rate of convergence.

The aim of this paper is the homogenization in general Carnot Group of step r , with a strong interplay with their geometry. We reinterpret the results obtained “by hand” by Birindelli and Wigniolle, in the geometry of the group, using the Lie bracket generating property, and a geometric construction of the tiling, valid in any Carnot group.

In the Heisenberg group case, the authors were able to give estimates of the rate of convergence, using the homogeneity of test functions of degree two, and, therefore, uniform estimates of the second derivatives; for step r groups, $r > 2$, we will obtain uniform convergence of u^ε towards the solution of the homogenized equation.

Notice that viscosity solutions are Lipschitz continuous with respect to the group distance, but only Hölder continuous with respect to the Euclidean one. Viscosity solutions for Hamilton-Jacobi equations in the Heisenberg group were considered in [17].

2. CARNOT GROUPS

We will consider (\mathbb{R}^N, \cdot) as a Carnot group with a Group operation \cdot and a family of dilations, compatible with the Lie structure.

Let us briefly recall some basic facts about Carnot groups G . A Carnot group G of step $r \geq 1$, is a simply connected nilpotent Lie group whose Lie algebra \mathfrak{g} is stratified. This means that \mathfrak{g} admits a decomposition as a vector space sum

$$\mathfrak{g} = \mathfrak{g}_1 \oplus \mathfrak{g}_2 \oplus \cdots \oplus \mathfrak{g}_r$$

such that

$$[\mathfrak{g}_1, \mathfrak{g}_j] = \mathfrak{g}_{j+1}$$

for $j = 1 \cdots r$ with $\mathfrak{g}_k = \{0\}$ for $k > r$. Note that \mathfrak{g} is generated as a Lie algebra by \mathfrak{g}_1 .

Let $m_j = \dim \mathfrak{g}_j$ and chose a basis of \mathfrak{g}_j formed by left-invariant vector fields $X_{i,j}, i = 1, \cdots, m_j$. The dimension of G as a manifold is $N = m_1 + \cdots + m_r$. The horizontal tangent space at a point $\xi \in G$ is the m_1 -dimensional subspace that is linearly spanned by $X_{1,1}(\xi), \cdots, X_{m_1,1}(\xi)$. From now on, we will indicate simply by m the dimension of G_1 and by X_1, \cdots, X_m a frame of vector fields that span the first layer G_1 . The exponential coordinates are given by the diffeomorphism

$F : \mathbb{R}^N \rightarrow G$ defined by:

$$F(x) = \exp \left(\sum_{j=1}^r \sum_{i=1}^{m_j} x_{i,j} X_{i,j} \right),$$

Denoting by \cdot the group operation in G , the mapping $(\xi, \eta) \rightarrow \xi \cdot \eta$ has polynomial entries, when written in exponential coordinates. Viceversa, given a point $\xi \in G$, we denote by x_s the vector of coordinates $(x_{1,s}, \dots, x_{m_s,s})$ for every $s = 1, \dots, r$. The following formula holds:

$$(2.1) \quad X_\alpha = \frac{\partial}{\partial x_{1,\alpha}} + \sum_{s=2}^r \sum_{\beta=1}^{m_s} b_{s,\alpha}^\beta(x_1(\xi), \dots, x_{s-1}(\xi)) \frac{\partial}{\partial x_{s,\beta}}, \dots, \alpha = 1, \dots, m$$

where each $b_{s,\alpha}^\beta$ is a polynomial of weighted degree $s - 1$. By weighted degree we mean that the layer \mathfrak{g}_s in the stratification of \mathfrak{g} has the degree s . Correspondingly, each homogeneous monomial $x_1^{\alpha_1}, \dots, x_r^{\alpha_r}$, with multi-indices $\alpha_s = (\alpha_{1,s}, \dots, \alpha_{m_s,s})$, $s = 1, \dots, r$ is said to have weighted degree p if

$$\sum_{s=1}^r s \left(\sum_{\beta=1}^{m_s} \alpha_{\beta,s} \right) = p$$

There is a family of dilations that is compatible with the group operation:

$$\delta_\lambda(x_1, \dots, x_N) = (\lambda x_{1,1}, \dots, \lambda x_{1,m}, \lambda^2 x_{2,1}, \dots, \lambda^2 x_{2,m_2}, \dots, \lambda^r x_{r,m_r})$$

With the above notations the horizontal subspace can be identified with the left translation by ξ of G_1 , that is we have

$$\xi \cdot G_1 = \text{Linear-span} \{X_1(\xi), \dots, X_m(\xi)\}.$$

A horizontal curve $\gamma(t)$ is a piece-wise smooth curve whose tangent vector $\gamma'(t)$ is in the horizontal tangent space $(\gamma(t)) \cdot G_1$ whenever it exists. Given two points ξ and η we consider the set of all possible horizontal curves joining these points:

$$\Gamma(\xi, \eta) = \{\gamma \text{ horizontal curve} : \gamma(0) = \xi, \gamma(1) = \eta\}$$

This set is never empty by Chow's accessibility theorem (see for example [5]). For convenience, fix an ambient Riemannian metric in \mathfrak{g} so that $\mathfrak{X} = \{X_{i,j}\}$ is an orthonormal frame and the

Riemannian volume element = Haar measure of G = Lebesgue measure in \mathbb{R}^N .

The Carnot-Carathéodory distance is then defined as the infimum of the length of horizontal curves of the set Γ :

$$d_{CC}(\xi, \eta) = \inf_{\Gamma(\xi, \eta)} \int_0^1 |\gamma'(t)| dt.$$

It depends only of the restriction of the ambient Riemannian metric to the horizontal distribution generated by the horizontal tangent spaces.

The Carnot-Carathéodory ball of radius R centered at a point ξ is given by

$$B(\xi, R) = \{\eta \in G : d_{CC}(\xi, \eta) < R\}.$$

The Carnot-Carathéodory gauge is given by

$$|\xi|_{CC} = d(0, \xi).$$

An important cone-like property of this gauge is that it solves the eikonal equation in the almost everywhere sense. This was proven by Monti and Serra-Cassano [MSc] in general Carnot groups.

Theorem 1. *Consider the horizontal gradient of $|\xi|_{CC}$ given by*

$$X(|\xi|_{CC}) = (X_1(|\xi|_{CC}), \dots, X_m(|\xi|_{CC})).$$

Then, for a.e. $\xi \in \mathbb{R}^N$, we have

$$|X(|\xi|_{CC})| = 1.$$

A smooth gauge in G is defined:

$$|\xi|_G = \left(\sum_{j=1}^r \left(\sum_{i=1}^{m_j} |x_{i,j}|^2 \right)^{\frac{r_i}{j}} \right)^{\frac{1}{2r_1}}$$

Theorem 2. [5]

$$|\xi|_{CC} \sim |\xi|_G \sim \sum_{j=1}^r \sum_{i=1}^{m_j} |x_{i,j}|^{\frac{1}{j}}$$

$$\text{vol}(B(0, R)) \sim r^Q$$

where $Q = \sum_{j=1}^r j m_j$ is the homogeneous dimension of G .

From the point of view of the viscosity theory, we do know that the smooth gauge is a viscosity subsolution of the eikonal equation, see section 3.

For simplicity, we will make the following assumptions on G : denote by X_1, \dots, X_m a basis for the first layer and suppose that the following identities hold:

$$[X_1, X_m] = X_{m+1}, \dots, [X_1, X_{m+r-1}] = X_{m+r}, \quad m+r = N.$$

Following [5],[11], the group operation in G has polynomial entries:

$$\xi \cdot \eta = (x_{1,1} + y_{1,1}, \dots, x_{1,m} + y_{1,m}, x_{m+1} + y_{m+1} + Q_2(x_{1,j}, y_{1,j}),$$

$$x_{m+2} + y_{m+2} + Q_3(x_{1,j}, y_{1,j}, x_{m+1}, y_{m+1}), \dots, x_{m+r} + y_{m+r} + Q_r(x_{1,j}, y_{1,j}, \dots, x_{m+r-1}, y_{m+r-1}))$$

where each $Q_s, s = 1, \dots, r$ is an homogeneous polynomial of degree s in the variables indicated.

2.1. Examples.

2.1.1. *Heisenberg Group.* Such group is the simplest prototype of a Carnot Group of step two. It can be identified with the Euclidean space $\mathbb{C}^n \times \mathbb{R} \equiv \mathbb{R}^{2n+1}$ endowed with the non-abelian group law :

$$(x, y, t) \cdot (x', y', t') = (x + x', y + y', t + t' + \frac{1}{2} \langle x, y' \rangle - \langle x', y \rangle)$$

The Heisenberg algebra is splitted in $V_1 \oplus V_2$, where $V_1 = \mathbb{R}^{2n} \times \{0\}$ and $V_2 = \{0\} \times \mathbb{R}$. It is generated by the vector fields:

$$X_j = \frac{\partial}{\partial x_j} + \frac{1}{2} y_j \frac{\partial}{\partial t}$$

$$Y_j = \frac{\partial}{\partial y_j} - \frac{1}{2} x_j \frac{\partial}{\partial t}$$

The only nontrivial commutator is

$$[X_j, Y_j] = \frac{\partial}{\partial t}$$

The homogeneous dimension is 4.

2.1.2. *Engel Group.* The Engel Group \mathcal{E} can be identified with \mathbb{R}^4 and the Lie algebra is decomposed in:

$$\mathfrak{E} = G_1 \oplus G_2 \oplus G_3$$

The horizontal algebra is spanned by X_1, X_2 and the nontrivial commutators are:

$$[X_1, X_2] = X_3 \quad [X_1, X_3] = [X_1, [X_1, X_2]] = X_4.$$

This is a group of step 3.

The group law is:

$$(x, y, t, s) \cdot (x', y', t', s') = (x + x', y + y', t + t' + Q_2, s + s' + Q_3)$$

where

$$Q_2 = \frac{1}{2}(xy' - yx')$$

$$Q_3 = \frac{1}{2}(xt' - tx') + \frac{1}{12}(x^2y' - xx'(y + y') + yx'^2)$$

The homogeneous dimension is 7. For more examples of Carnot Groups, see the recent monograph [18].

We recall the following property, that in the sequel will be referred as the Lie bracket generating property. Following the notations of [20] and [19], we define the following quantities. Let S_1, \dots, S_ℓ be fields belonging to the family X_1, \dots, X_m and let $a \in R$:

$$C_1(a; S_1) = \exp(aS_1)$$

$$C_2(a; S_1, S_2) = \exp(-aS_2)\exp(-aS_1)\exp(aS_2)\exp(aS_1)$$

and, by induction:

$$C_\ell(a; S_1, \dots, S_\ell) = C_{\ell-1}(a; S_2, \dots, S_\ell)^{-1}\exp(-aS_1)C_{\ell-1}(a; S_2, \dots, S_\ell)\exp(aS_1)$$

By the Campbell-Hausdorff formula and the Jacobi identity they get the following equality :

$$C_2(a; S_1, S_2) = \exp(a^2[S_1, S_2] + \sum_{d(I)>2} c_I a^{d(I)} S_{[I]})$$

Iterating:

$$C_\ell(a; S_1, S_2, \dots, S_\ell) = \exp(a^\ell S_{[1, \dots, \ell]} + \sum_{d(I)>\ell} c_I a^{d(I)} S_{[I]})$$

Here $d(I)$ represents the length of the multiindex I , so the iteration will end after r steps, where r is the step of the group.

3. HAMILTON-JACOBI EQUATION

In order to define viscosity solutions we must first identify the first order jets adapted to our framework. Motivated by the Taylor expansion [5] consider a differentiable function $u: G \mapsto \mathbb{R}$ at the point ξ_0 . We have

$$u(\xi) = u(\xi_0) + \langle D_h u(\xi_0), \overline{\xi_0^{-1} \cdot \xi} \rangle + o(|\xi_0^{-1} \cdot \xi|_G),$$

where $D_h u = X_1 u X_1 + \dots + X_m u X_m$ is the horizontal gradient of u and $\bar{\xi}$ is the horizontal projection.

Definition 1. *A function $u: G \mapsto \mathbb{R}$ is of class C^1 if the horizontal derivatives $X_1 u, \dots, X_m u$ are continuous.*

If a function u is not necessarily smooth but merely upper semicontinuous, the collection of vectors $p \in \mathbb{R}^m$ such that

$$(3.1) \quad u(\xi) \leq u(\xi_0) + \langle p, \overline{\xi_0^{-1} \cdot \xi} \rangle + o(|\xi_0^{-1} \cdot \xi|_G)$$

is denoted by $J_u^{1,+}(\xi_0)$ and called the first order superjet of u at the point ξ_0 . Analogously we define $J_v^{1,-}(\xi_0)$, the first order subjet of a lower semicontinuous function v at ξ_0 as the set of vectors $q \in \mathbb{R}^m$ such that

$$(3.2) \quad v(\xi) \geq v(\xi_0) + \langle q, \overline{\xi_0^{-1} \cdot \xi} \rangle + o(|\xi_0^{-1} \cdot \xi|_G).$$

Jets can also be characterized by test functions ψ as follows, see [17] in the Heisenberg Group case.

Let u defined on a neighborhood of a point ξ_0 . Suppose that ψ is a C^1 function touching u from above at ξ_0 ,

$$u(\xi_0) = \psi(\xi_0)$$

and

$$u(\xi) \leq \psi(\xi)$$

in a neighborhood of ξ_0 . Then the vector

$$(D_h \psi(\xi_0)) \in J_u^{1,+}(\xi_0).$$

Moreover, every vector

$$\eta \in J_u^{1,+}(\xi_0)$$

is of the form

$$(D_h \psi(\xi_0))$$

for some C^1 function ψ that touches u from above at ξ_0 .

A similar statement holds for $J_u^{1,-}(\xi_0)$ replacing “touching from above” by “touching from below”.

An upper semicontinuous function u is a viscosity subsolution of a first order equation of the form

$$(3.3) \quad F(\xi, u, D_h u) = 0$$

where F is continuous in all variables and increasing in u if, for any ξ_0 and every $p \in J_u^{1,+}(\xi_0)$, the inequality holds:

$$F(\xi_0, u(\xi_0), p) \leq 0$$

A lower semicontinuous function v is a viscosity supersolution of (3.3) if for any ξ_0 and every $q \in J_v^{1,-}(\xi_0)$, then:

$$F(\xi_0, v(\xi_0), q) \geq 0$$

A continuous function u that is both a viscosity subsolution and a viscosity supersolution is called a viscosity solution. Uniform limits of viscosity solutions are viscosity solutions.

Let us start from a simple example of first order equation: the eikonal equation. The smooth gauge $|\cdot|_G$ is a viscosity subsolution of $|D_h u| = 1$ in $B_R(\xi) - \{\xi\}$, with the boundary conditions $u = R$ on ∂B_R and $u(\xi) = 0$, see [13]. Let us consider a more general eikonal equation: $|D_h u| = C$ and prove that every subsolution is Lipschitz continuous with respect to the smooth gauge. By Pansu Theorem, it is also almost everywhere differentiable. The Lie bracket generating property transports the Lipschitz continuity from the horizontal layer to the entire group. In fact, control theory guarantees the continuity only for horizontally connected points. For arbitrary points, the Lie geometry plays the major role.

Theorem 3. *Any bounded viscosity solution of $|D_h u| \leq C$ satisfies:*

$$|u(\xi \cdot \eta) - u(\xi)| \leq C|\eta|$$

Proof. First step: Without loss of generality we can assume that ξ is equal to the origin $\underline{0}$. We shall prove the inequality for any horizontal $\eta \in G_1$. This is done following the same proof of [6], using Theorem 5.21 of [3]. We write the norm of the horizontal gradient:

$$|D_h u| = |\sigma(\xi) \nabla u| = \sup_{\eta \in G_1, |\eta| \leq 1} \sigma^T \eta \cdot \nabla u$$

where σ is the $m \times N$ matrix with entries a_{ij} , $X_i = \sum_j a_{ij} \frac{\partial}{\partial x_j}$. We are considering the evolution of a horizontal curve starting from the origin:

$$(3.4) \quad \begin{cases} \gamma'(t) &= \sum_{i=1}^m \eta_i X_i(\gamma(t)) \\ \gamma(0) &= \underline{0} \end{cases}$$

where $\eta \in \mathbb{R}^m$. The solution $y(s, \eta)$ can be parametrized using the dilations: $y(s, \eta) = \delta_s(\eta)$. Geometrically, it represents a horizontal curve joining the origin and $\eta \in G_1$.

The function u is a viscosity subsolution of the Hamilton-Jacobi-Bellmann equation:

$$\sup_{\eta \in G_1, |\eta| \leq 1} \{\langle \eta, D_h u \rangle - C\} = 0$$

Hence u satisfies:

$$u(y(s, \eta)) - u(y(t, \eta)) \leq (t - s)$$

for $s \leq t$. Changing η in $-\eta$ we get the assertion for horizontally connected points. The general case is achieved using the Lie bracket generating property. We write the point η as direct sum of projection onto G_1, G_2, \dots, G_r , namely $\eta = h_1 \oplus h_2 \oplus \dots \oplus h_r$. We write the difference in a telescopic sum:

$$\begin{aligned} u(\underline{0}) - u(\eta) &= u(\underline{0}) - u(h_1, 0, \dots, 0) + u(h_1, 0, \dots, 0) - u(h_1, h_2, \dots, 0) \\ &\quad + \dots + u(h_1, h_2, \dots, h_{r-1}, 0) - u(h_1, h_2, \dots, h_r) \end{aligned}$$

The first difference is controlled by step 1. To estimate the second one, we take $a = |h_2|^{\frac{1}{2}}$ and apply C_2 to the point $(h_1, 0, \dots, 0)$. Iterating the construction, after a finite number of steps, we get the all gauge.

The telescoping sum is controlled by :

$$|h_1| + |h_2|^{\frac{1}{2}} + \dots + |h_r|^{\frac{1}{r}}$$

Therefore, we reach the assertion , having a control in terms of the smooth gauge. \square

From now on, we will always consider Lipschitz continuity with respect to the Carnot smooth gauge.

Our Hamilton-Jacobi equations is of the form:

$$(3.5) \quad \gamma u + H(\xi, D_h u) = 0$$

with $\gamma > 0$. The Hamiltonian H verifies:

$$\begin{aligned} |H(\xi', p) - H(\xi, p)| &\leq m(|\xi^{-1} \cdot \xi'|_G (1 + |p|)) Q(\xi, \xi', p) \\ |H(\xi, 0)| &\leq C_3 \\ \lim_{|p| \rightarrow \infty} H(\xi, p) &= +\infty \quad \text{uniformly in } \xi \end{aligned}$$

where $m(t) \rightarrow 0$ when $t \rightarrow 0$ and $Q(\xi, \xi', p) = \max\{\Phi(H(\xi, p)), \Phi(\xi', p)\}$ and Φ is a continuous function from \mathbb{R} into \mathbb{R}^+ . Examples of Hamiltonians that satisfy our hypotheses are: $H(x, p) = \Psi(H_o(x, p))$ where H_o is a continuous function on \mathbb{R}^{2N} satisfying: $p \rightarrow H_o(x, p)$ convex , $H_o(x, \lambda p) = \lambda H_o(x, p)$, $H_o(x, p) \geq 0$, $|H_o(x, p) - H_o(y, p)| \leq \omega(|x - y|(1 + |p|))$, with ω modulus of continuity such that $\lim_{s \rightarrow 0^+} \omega(s) = 0$, for all $\lambda > 0$, $x, y \in \mathbb{R}^N$. We assume also that H_o is degenerate coercive in the sense that for some $\varepsilon > 0$ $\sigma(x)([-\varepsilon, \varepsilon]^m) \subset \partial H_o(x, 0)$. Here ∂H_o is the subdifferential of the convex function H_o and $\sigma(x)$ is the $m \times N$ matrix as in the introduction. The function Ψ is convex, increasing , $\Phi(0) = 0$. The function H_o plays the role of a generalized gauge in G , for example, one can take $H_o(\xi) = |\xi|_{CC}$.

Remark 1. *The hypotheses on the Hamiltonian H ensure that $-\gamma C_3$ and γC_3 are respectively subsolution and supersolution of the equation (3.5). In particular, they are also Euclidean viscosity subsolution and supersolution. Perron's method for viscosity solutions ([12]) proves the existence of an Euclidean viscosity solution, that is also subelliptic viscosity solution. Also, a solution, as a consequence of the coercivity of H , verifies an equation of the form $|D_h u| \leq M$, for some $M > 0$, so it is (subelliptic) Lipschitz continuous, see Theorem 3.*

It will remain to prove the uniqueness Theorem. The proof for second order subelliptic equations has been done in [19]. The author uses the Euclidean Theorem to get the jets and twist them into the subelliptic setting. For first order equations, we can proceed directly.

Theorem 4. *Let u a bounded subsolution and v a bounded supersolution of equation (3.5) . Then we have $u \leq v$ on G . As a consequence, the solution is unique*

Proof. Suppose, by contradiction, that $\sup(u - v) = \delta > 0$. We define the functions:

$$A(\xi, \eta) = |\xi^{-1} \cdot \eta|_G^r \quad \rho(\xi) = |\xi|_G^r$$

We have the following inequalities :

$$|D_{h,\xi}A(\xi, \eta)| \leq CrA^{\frac{r-1}{r}}(\xi, \eta), |D_{h,\eta}A(\xi, \eta)| \leq CrA^{\frac{r-1}{r}}(\xi, \eta)$$

Consider for τ and α positive:

$$\psi(\xi, \eta) = u(\eta) - v(\xi) - \tau A(\xi, \eta) - \alpha \rho(\xi)$$

For every τ and α , since u and v are bounded there is a maximum of ψ attained in a point $(\xi_{\tau,\alpha}, \eta_{\tau,\alpha})$. In particular, the point $\xi_{\tau,\alpha}$ is a minimum for $\xi \rightarrow v(\xi) + \tau A(\xi, \eta_{\tau,\alpha}) + \alpha \rho(\xi)$ and v is a supersolution, we get that $p_{\tau,\alpha} = -\tau D_{h,\xi_{\tau,\alpha}}A(\xi_{\tau,\alpha}, \eta_{\tau,\alpha}) - \alpha D_h \rho(\xi_{\tau,\alpha}) \in J_v^{1,-}(\xi_{\tau,\alpha})$ and plugging it into the equation:

$$(3.6) \quad \gamma v(\xi_{\tau,\alpha}) + F(\xi_{\tau,\alpha}, p_{\tau,\alpha}) \geq 0$$

From the other hand, the point $\eta_{\tau,\alpha}$ is a maximum for the function: $\eta \rightarrow u(\eta) - \tau A(\xi_{\tau,\alpha}, \eta)$ and u is a viscosity subsolution:

$$(3.7) \quad \gamma u(\eta_{\tau,\alpha}) + F(\eta_{\tau,\alpha}, q_{\tau,\alpha}) \leq 0$$

where $q_{\tau,\alpha} = \tau D_{h,\eta_{\tau,\alpha}}A(\xi_{\tau,\alpha}, \eta_{\tau,\alpha}) \in J_u^{1,+}(\eta_{\tau,\alpha})$.

Now the coercitivity condition on H and the viscosity subsolution property of u guarantee that u is in fact a viscosity subsolution of an eikonal equation $|D_h u| \leq M$ for some $M > 0$. We then apply the Lipschitz continuity to get:

$$\tau A(\xi_{\tau,\alpha}, \eta_{\tau,\alpha}) \leq u(\eta_{\tau,\alpha}) - u(\xi_{\tau,\alpha}) \leq M |\xi_{\tau,\alpha}^{-1} \cdot \eta_{\tau,\alpha}|_G \leq MA(\xi_{\tau,\alpha}, \eta_{\tau,\alpha})^{\frac{1}{r}}$$

which implies

$$(3.8) \quad A(\xi_{\tau,\alpha}, \eta_{\tau,\alpha}) \leq \left(\frac{M}{\tau}\right)^{\frac{r}{r-1}}$$

From the other hand, writing $\psi(\xi_{\tau,\alpha}, \eta_{\tau,\alpha}) \geq \psi(\xi_{\tau,\alpha}, \xi_{\tau,\alpha})$, we get:

$$\tau A(\xi_{\tau,\alpha}, \eta_{\tau,\alpha}) + \alpha \rho(\xi_{\tau,\alpha}) \leq C$$

We get

$$\begin{aligned} |\xi_{\tau,\alpha}^{-1} \eta_{\tau,\alpha}|_G &\rightarrow 0 & (\tau, \alpha) &\rightarrow (+\infty, 0) \\ \alpha \rho(\xi_{\tau,\alpha}) &\rightarrow 0 & (\tau, \alpha) &\rightarrow (+\infty, 0) \\ M_{\tau,\alpha} &\rightarrow M & (\tau, \alpha) &\rightarrow (+\infty, 0) \end{aligned}$$

Taking into account the hypothese on H :

$$\begin{aligned} \gamma(u(\eta_{\tau,\alpha}) - v(\xi_{\tau,\alpha})) &\leq H(\eta_{\tau,\alpha}, q_{\tau,\alpha}) - H(\xi_{\tau,\alpha}, p_{\tau,\alpha}) \\ H(\eta_{\tau,\alpha}, q_{\tau,\alpha}) - H(\xi_{\tau,\alpha}, q_{\tau,\alpha}) &+ H(\xi_{\tau,\alpha}, q_{\tau,\alpha}) - H(\xi_{\tau,\alpha}, p_{\tau,\alpha}) \\ &\leq m(|\xi_{\tau,\alpha}^{-1} \eta_{\tau,\alpha}|_G (1 + |q_{\tau,\alpha}|)) + n_\tau(\alpha |\xi_{\tau,\alpha}|) \end{aligned}$$

where $n(\cdot)$ is a modulus of continuity of H on the set $G \times \bar{B}_{s_\tau}$ and $s_\tau = \max\{\|u\|, \|v\|, 1 + \tau\}$. To conclude, we fix τ and let α goes to zero, then we let τ goes to $+\infty$. \square

4. HOMOGENIZATION PROBLEMS

Consider the cube $Q = [-\frac{1}{2}, \frac{1}{2}]^N$ and construct a tiling of G moving only from horizontal points of the form $(\underline{k}, 0)$, $\underline{k} \in \mathbb{Z}^m$, where $(\underline{k}, 0) \in G_1$, $Q_k = (\underline{k}, 0) \cdot Q$. The bracket generating property yields the following property:

Lemma 1. *For every point $\xi \in G$ there exists a point $\xi_0 \in Q$ and a finite number of group actions generated by elements of the form $(\underline{k}, 0)$, $\underline{k} \in \mathbb{Z}^m$ that applied to ξ_0 give ξ .*

Proof. We put $k_i = [x_i + \frac{1}{2}]$, for $i = 1 \cdots m$, where $[\cdot]$ stand for the integer part; hence $\tilde{x}_i = x_i - k_i \in [-\frac{1}{2}, \frac{1}{2}]^m$. We consider $\bar{k} = (k_1, \cdots, k_m) \in \mathbb{Z}^m$ and compute

$$(\bar{k}, 0, \cdots, 0) \cdot \tilde{\xi}$$

The group operation gives a point with the first m coordinates equal to the first m coordinates of ξ and the $m+1$ equal to $\tilde{\xi}_{m+1} + Q_2$, where Q_2 is an homogeneous polynomial of degree 2 in the horizontal variables. Using the Lie bracket generating property towards the mappings C_ℓ we could solve a system of equations in terms of the components of ξ_0 . □

We will adopt this latter one as an intrinsic pavage and therefore the corresponding notion of periodicity, we will call it G -periodicity.

Definition 2. *A function f defined on a Carnot group G is said G -periodic iff*

$$f((\underline{k}, 0) \cdot \xi) = f(\xi) \quad \forall \underline{k} \in \mathbb{Z}^m, \forall \xi \in Q.$$

Consider now the action of the dilatations of the group:

$$Q_k^\varepsilon = \delta_\varepsilon(Q_k) = \delta_\varepsilon(\underline{k}, 0) \cdot Q$$

We notice that $\{Q_k^\varepsilon\}$ is still a tiling of G . Furthermore, if f is G -periodic, then $f(\delta_{\frac{1}{\varepsilon}} \xi)$ is εG -periodic.

In this context, homogenizations problems are:

$$(4.1) \quad u^\varepsilon + H(\xi, \delta_{\frac{1}{\varepsilon}} \xi, D_h u^\varepsilon(\xi)) = 0$$

with $\eta \rightarrow H(\xi, \eta, p)$ G -periodic for each $(\xi, p) \in G \times \mathbb{R}^m$. The function H is defined on $\mathbb{R}^{2N} \times \mathbb{R}^m$ and satisfies for all $\xi, \xi', \eta, \eta' \in \mathbb{R}^N$ and all $p, p' \in \mathbb{R}^m$:

$$|H(\xi', \eta', p') - H(\xi, \eta, p)| \leq C(|\xi^{-1} \cdot \xi'|_G + |\eta^{-1} \cdot \eta'|_G + |p - p'|)$$

$$\nu|p| - C_3 \leq H(\xi, \eta, p) \leq \nu|p| + C_3.$$

Remark 2. *From the previous section we do know that it exists a sequence of solutions of the (4.1); moreover this sequence is equibounded and equilipschitz. In fact, $-C_3$ and C_3 are, respectively, subsolution and supersolution of equations (4.1) and, therefore, by comparison $\|u_\varepsilon\|_\infty \leq C$. In addition, using the coercivity of H , we get that $|D_h u_\varepsilon| \leq C$ in the viscosity sense, that implies, by Theorem 3, the equilipschitz bound. So, up to a subsequence, it converges uniformly, as ε tends to zero, to a continuous function u , the solution of the homogenized problem, see Theorem 6.*

The nonlinear eigenvalue problem for the effective Hamiltonian is:

$$(4.2) \quad H(\xi, \eta, p + D_h v(\eta)) = \lambda$$

and then the effective Hamiltonian \bar{H} is

$$\bar{H}(\xi, p) = \lambda$$

for every $(\xi, p) \in G \times \mathbb{R}^m$. The uniqueness of λ depending on (ξ, p) is a consequence of the comparison theorem. As in the Euclidean case, (see [15],[8]), one consider a sequence of approximated problems:

$$\alpha v_\alpha(\eta) + H(\xi, \eta, p + D_h v_\alpha(\eta)) = 0$$

and study the asymptotic behaviour as α tends to zero. We already know that there exists a solution v_α bounded and Lipschitz continuous, that is also periodic. It can be viewed indeed (see [15]) that the limit of $\alpha v_\alpha(\eta)$ does not depend on η and that a solution (λ, v) of (4.2) is given by

$$\lambda = \lim_{\alpha \rightarrow 0} \alpha v_\alpha \quad v(\eta) = \lim_{\alpha \rightarrow 0} (v_\alpha - \inf_Q v_\alpha)$$

Theorem 5. *For each p there exists a unique λ such that there exists v viscosity solution of (4.2).*

Proof. From Theorem 4 we know that there exists unique v_α that is also Lipschitz and, from the uniqueness, it follows that is periodic. Moreover

$$(4.3) \quad -\sup_\eta H(\xi, \eta, p) \leq \alpha v_\alpha(\eta) \leq -\inf_\eta H(\xi, \eta, p)$$

and from the coercivity of H , we also get:

$$\|D_h v_\alpha\|_\infty \leq C$$

for some C independent of α . Set $\tilde{v}_\alpha = v_\alpha - \inf_Q v_\alpha$ and observe that \tilde{v}_α is periodic and equiLipschitz. Up to a subsequence, we may assume that $(-\alpha v_\alpha, \tilde{v}_\alpha)$ converges uniformly to some (λ, v) . From the stability property of viscosity solutions, v is a solution of (4.2). Suppose now that there exists another couple (μ, w) , solution of the cell problem, with $\lambda \neq \mu$. We may assume that $\lambda < \mu$ is the case. Since if w is a solution of

$$H(\xi, \eta, p + D_h w(\eta)) = \mu$$

then so is $w + c$; therefore, we can also assume that $w > v$. Choose α so that $\lambda + \alpha v \leq \mu + \alpha w$. We notice that v, w are respectively the unique viscosity solutions of

$$(4.4) \quad H(\xi, \eta, p + D_h v(\eta)) + \alpha v(\eta) = \lambda + \alpha v(\eta)$$

$$(4.5) \quad H(\xi, \eta, p + D_h w(\eta)) + \alpha w(\eta) = \mu + \alpha w(\eta)$$

Then, in particular, w is a super solution of (4.4), and by comparison, $v \leq w$, that is a contradiction. □

We now conclude by proving the uniform convergence of u^ε towards the solution of the homogenized problem, using the perturbed test functions method adapted to Hamilton-Jacobi equations.

Theorem 6. *The sequence u^ε converges uniformly to the solution of the homogenized equation:*

$$(4.6) \quad u + \bar{H}(\xi, D_h u) = 0$$

Proof. We consider a point ξ_0 and a test function ϕ such that $u - \phi$ has a strict local maximum at ξ_0 , we need to prove that

$$\phi(\xi_0) + \bar{H}(\xi_0, D_h \phi(\xi_0)) \leq 0$$

We start by solving the cell problem (4.2) at the point ξ_0 and with $p = D_h \phi(\xi_0)$

$$(4.7) \quad H(\xi_0, y, D_h \phi(\xi_0) + D_h \psi(y)) = \bar{H}(\xi_0, D_h \phi(\xi_0))$$

The function ψ is G -periodic and defined up to a constant, we normalize it, by requiring that $\int_Q \psi = 1$. Since ψ is a.e. differentiable, we take a test function χ touching ψ from above at the point $\delta_{\frac{1}{\varepsilon}}(\xi_0)$. Consider the perturbed test function:

$$\phi^\varepsilon(\xi) = \phi(\xi) + \varepsilon \chi(\delta_{\frac{1}{\varepsilon}}(\xi))$$

Since $u - \phi$ has a strict local maximum at ξ_0 and, up to a subsequence, $u^\varepsilon \rightarrow u$, $\phi^\varepsilon \rightarrow \phi$ uniformly, we see that for each ε , $u^\varepsilon - \phi^\varepsilon$ has a local maximum at some point ξ_ε , with $\xi_\varepsilon \rightarrow \xi_0$ as $\varepsilon \rightarrow 0$. By using that u^ε is a solution of (4.1), we get:

$$(4.8) \quad \phi^\varepsilon(\xi_\varepsilon) + H(\xi_\varepsilon, \delta_{\frac{1}{\varepsilon}}(\xi_\varepsilon), D_h \phi^\varepsilon(\xi_\varepsilon)) \leq 0$$

We compute:

$$D_h \phi^\varepsilon(\xi_\varepsilon) = D_h \phi(\xi_\varepsilon) + D_h \chi(\delta_{\frac{1}{\varepsilon}}(\xi_\varepsilon)) = D_h \phi(\xi_0) + D_h \chi(\delta_{\frac{1}{\varepsilon}}(\xi_\varepsilon)) + o(1)$$

since the horizontal gradient commutes with dilations and ξ_ε converges to ξ_0 .

By inserting $D_h \phi^\varepsilon(\xi_\varepsilon)$ in (4.8) and use the uniform convergence of ϕ^ε to ϕ and the continuity of H :

$$\phi(\xi_0) + H(\xi_0, \delta_{\frac{1}{\varepsilon}}(\xi_\varepsilon), D_h \phi(\xi_0) + D_h \chi(\delta_{\frac{1}{\varepsilon}}(\xi_\varepsilon))) \leq o(1)$$

We now set $y = \delta_{\frac{1}{\varepsilon}}(\xi_\varepsilon)$ in (4.7) and substitute with \bar{H} , to conclude as $\varepsilon \rightarrow 0$. Analogously, starting from a strict local minimum for $u - \phi$, we can show the opposite inequality. \square

5. APPENDIX

We have defined the function:

$$A(\xi, \eta) = \left\{ \sum_{j=1}^m (x_{1,j} - y_{1,j})^2 \right\}^{r!} + [(x_{m+1} - y_{m+1} + Q_2)^2]^{\frac{r!}{2}} +$$

$$+ \cdots + [(x_{m+r-1} - y_{m+r-1} + Q_{r-1})^2]^{\frac{r!}{r-1}} + [(x_{m+r} - y_{m+r} + Q_r)^2]^{(r-1)!} \Big\}^{\frac{r}{2r!}}$$

Let us develop the calculations for:

$$|D_{h,\xi} A(\xi, \eta)| \leq C r A^{\frac{r-1}{r}}(\xi, \eta)$$

Denote by B the argument of A inside the graph parenthesis:

$$B = \left[\sum_{j=1}^m (x_{1,j} - y_{1,j})^2 \right]^{r!} + \cdots + [(x_{m+r-1} - y_{m+r-1} + Q_{r-1})^2]^{\frac{r!}{r-1}} + [(x_{m+r} - y_{m+r} + Q_r)^2]^{(r-1)!}$$

$$= I_1^{r!} + I_2^{\frac{r!}{2}} + \cdots + I_{r-1}^{\frac{r!}{r-1}} + I_r^{(r-1)!}$$

$$\begin{aligned} X_{j,\xi}A(\xi, \eta) &= X_{j,\xi}\{B(\xi, \eta)\}^{\frac{r}{2r!}} \\ &= \frac{r}{2r!}\{B\}^{\frac{r}{2r!}-1}X_{j,\xi}B \end{aligned}$$

Observe that the exponent $\frac{r}{2r!} - 1$ is negative. Recall now the formula (2.1) in a simplified version accordingly to our notations:

$$X_j(g) = \frac{\partial}{\partial x_j} + \sum_{s=m+1}^{m+r} b_{s,j}(x_1(g), \dots, x_{s-1}(g)) \frac{\partial}{\partial x_s}, \quad j = 1, \dots, m$$

where each $b_{s,j}$ is a homogeneous polynomial of weighted degree $s - m$ and observe that Q_j are homogeneous polynomial of weighted degree j . We get:

$$\begin{aligned} &[C_1 I_1^{r!-1} I_1^{\frac{1}{2}} + \dots + C_{r-1} I_{r-1}^{\frac{r!}{r-1} + \frac{1}{2}} R_{r-2}] + C_r I_r^{(r-1)! + \frac{1}{2}} R_{r-1} \\ &+ b_{m+1,j} \{C_2 I_1^{\frac{r!}{2}-1} I_2^{\frac{1}{2}} + \dots + C_{r-1} I_{r-1}^{\frac{r!}{r-1} + \frac{1}{2}} R_{r-3}\} + C_r I_r^{(r-1)! + \frac{1}{2}} R_{r-2} + \dots + b_{m+r,j} C_r I_r^{(r-1)! + \frac{1}{2}} \end{aligned}$$

where R_t are homogeneous polynomial of weighted degree t . To conclude, observe that the polynomial R_j are equivalent to the polynomial Q_j that appear in the group operation.

REFERENCES

- [1] Alvarez, O; and Bardi M., *Viscosity solutions methods for singular perturbations in deterministic and stochastic control*, SIAM J. Control and Optim., 40,(2001), 1159-1188.
- [2] Arisawa,M., *Quasi-periodic homogenizations for second-order Hamilton-Jacobi-Bellmann equations*, Adv. Sci. Appl. (1), 11, (2001), 465-480.
- [3] Bardi,M., and Capuzzo Dolcetta, I. *Optimal control and viscosity solutions of Hamilton-Jacobi -Bellman equations* , Birkhäuser Boston,Boston, MA, 1997.
- [4] Barles, G., *Solutions de viscosité des équations de Hamilton-Jacobi*, Mathematiques et Applications 17, Springer-Verlag 1994.
- [5] Bellaïche, André and Risler, Jean-Jacques, ed., *Sub-Riemannian Geometry*, Progress in Mathematics, Vol. 144, Birkhäuser, 1996.
- [6] Birindelli, I. and Wigniolle, J. *Homogenization of Hamilton-Jacobi equations in the Heisenberg Group* , Commun. Pure Appl. Anal., 2, (2003), 461-479.
- [7] Biroli, M.,Mosco, U. and Tchou, N.*Homogenization in the Heisenberg Group*, Adv. Math. Sci. Appl., (2), 7, (1997), 809-831.
- [8] Capuzzo Dolcetta, I. and Ishii, H. *On the rate of convergence in homogenization of Hamilton-Jacobi equations*, Indiana University Math. J. , 50, (2001), 1113-1129.
- [9] Evans, L.C. *The perturbed test function method for viscosity solutions of nonlinear PDE*, Proc. of the Royal Soc. of Edinburgh, 11A, (1989), 359-375.
- [10] Evans, L.C. *Periodic homogenization of certain fully nonlinear PDE*, Proc. of the Royal Soc. of Edinburgh, 120, (1992), 245-265.
- [11] Folland, G. B. and Stein, Elias M. *Hardy spaces on homogeneous groups.*, Mathematical Notes, 28. Princeton University Press, Princeton, N.J.; University of Tokyo Press, Tokyo, 1982
- [12] Ishii, H., *Perron's method for Hamilton-Jacobi equations*, Duke Math. J., 55, (1987), 369-384.
- [13] Juutinen,P.,Lu, G., Manfredi, J. and Stroffolini, B., *Convex functions on Carnot Groups* , submitted.
- [14] Lu, G., Manfredi, J. and Stroffolini, B., *Convex functions on the Heisenberg Group* , Calc. Var. Partial Differential Equations 19 (2004), no. 1,
- [15] Lions, P. L., Papanicolau,G. and Varadhan, R.S.*Homogenization of Hamilton-Jacobi equations*, preprint 1986.
- [16] Manfredi, J. , *Nonlinear subelliptic equations on Carnot Groups*, notes of a course at the School on Analysis and Geometry, Trento, 2003.
- [17] Manfredi, J. and Stroffolini B., *A Version of the Hopf-Lax Formula in the Heisenberg Group*, Comm. in Partial Differential Equations, 27, (2002), 1139-1159.

- [18] Montgomery, R., *A Tour of Subriemannian Geometries, their geodesics and applications*, Mathematical Surveys and Monographs, 91, American Mathematical Society, Providence, RI, 2002.
- [19] Morbidelli, D. *Fractional Sobolev norms and structure of Carnot-Carathéodory balls for Hörmander vector fields*, *Studia Math.*, 139, (2000), 213–244.
- [MSc] Monti, R. and Serra Cassano, F., *Surface measures in Carnot-Carathéodory spaces*, *Calc. of Variations*, 13, (2001), 339-376.
- [20] Nagel, A., Stein, E.M. and Wainger, S., *Balls and metrics defined by vector fields I: basic properties*, *Acta Math.*, 137, (1976), 247–320.

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