# DIFFERENTIAL FORMS IN CARNOT GROUPS: A Г-CONVERGENCE APPROACH 

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#### Abstract

Carnot groups (connected simply connected nilpotent stratified Lie groups) can be endowed with a complex $\left(E_{0}^{*}, d_{c}\right)$ of "intrinsic" differential forms. In this paper we prove that, in a free Carnot group of step $\kappa$, intrinsic 1 -forms as well as their intrinsic differentials $d_{c}$ appear naturally as limits of usual "Riemannian" differentials $d_{\varepsilon}, \varepsilon>0$. More precisely, we show that $L^{2}$-energies associated with $\varepsilon^{-\kappa} d_{\varepsilon}$ on 1-forms $\Gamma$-converge, as $\varepsilon \rightarrow 0$, to the energy associated with $d_{c}$.


## 1. Introduction

In the last few years, sub-Riemannian structures have been largely studied in several respects, such as differential geometry, geometric measure theory, subelliptic differential equations, complex variables, optimal control theory, mathematical models in neurosciences, non-holonomic mechanics, robotics. Roughly speaking, a sub-Riemannian structure on a manifold $M$ is defined by a subbundle $H$ of the tangent bundle $T M$, that defines the "admissible" directions at any point of $M$ (typically, think of a mechanical system with non-holonomic constraints). Usually, $H$ is called the horizontal bundle. If we endow each fiber $H_{x}$ of $H$ with a scalar product, there is a naturally associated Carnot-Carathéodory (CC) distance $d$ on $M$, defined as the Riemannian length of the horizontal curves on $M$, i.e. of the curves $\gamma$ such that $\gamma^{\prime}(t) \in H_{\gamma(t)}$. In the spirit of the present paper, it is worth recalling that CC-distances can be seen as limits of "Riemannian" distances (see e.g. [17] and [23]). Basically, this is obtained by penalizing the directions of the tangent bundle that are orthogonal to the horizontal bundle $H$.

Among sub-Riemannian spaces, a privileged role is played by Carnot groups (see below for precise definition and [5] for a general survey), a role akin to that of Euclidean spaces versus Riemannian manifolds, acting in some sense as rigid "tangent" spaces to general sub-Riemannian spaces (rigid because they are invariant under left translations and group dilations). Roughly speaking, we can always think of a Carnot group $\mathbb{G}$ as of the Lie group ( $\mathbb{R}^{n}, \cdot$ ), where $\cdot$ is a (non-commutative) multiplication such that its Lie algebra $\mathfrak{g}$ is nilpotent and admits a step $\kappa$ stratification. This means

[^0]that there exist linear subspaces $V_{1}, \ldots, V_{\kappa}$ (the layers of the stratification) such that
$$
\mathfrak{g}=V_{1} \oplus \ldots \oplus V_{\kappa}, \quad\left[V_{1}, V_{i}\right]=V_{i+1}, \quad V_{\kappa} \neq\{0\}, \quad V_{i}=\{0\} \text { if } i>\kappa,
$$
where $\left[V_{1}, V_{i}\right]$ is the subspace of $\mathfrak{g}$ generated by the commutators $[X, Y]$ with $X \in V_{1}$ and $Y \in V_{i}$. We refer to the first layer $V_{1}$ as to the horizontal layer, which plays a key role in our theory, since it generates the all of $\mathfrak{g}$ by commutation.

The stratification of the Lie algebra induces a family of anisotropic dilations $\delta_{\lambda}(\lambda>0)$ on $\mathfrak{g}$ and therefore, through exponential map, on $\mathbb{G}$.

It is well known that the Lie algebra $\mathfrak{g}$ of $\mathbb{G}$ can be identified with the tangent space at the origin $e$ of $\mathbb{G}$, and hence the horizontal layer of $\mathfrak{g}$ can be identified with a subspace $H \mathbb{G}_{e}$ of $T \mathbb{G}_{e}$. By left translation, $H \mathbb{G}_{e}$ generates a subbundle $H \mathbb{G}$ of the tangent bundle $T \mathbb{G}$ and eventually a subRiemannian structure on $\mathbb{R}^{n}$. We stress that Carnot-Carathéodory geometry is not Riemannian at any scale (see [26]).

The first Heisenberg groups $\mathbb{H}^{1}$ provides the simplest example of noncommutative Carnot groups (of step $\kappa=2$ ). It can be identified with $\mathbb{R}^{3}$ with variables $(x, y, t)$. Set $X:=\partial_{x}-\frac{1}{2} y \partial_{t}, Y:=\partial_{y}+\frac{1}{2} x \partial_{t}, T:=\partial_{t}$. The stratification of the algebra $\mathfrak{g}$ is given by $\mathfrak{g}=V_{1} \oplus V_{2}$, where $V_{1}=\operatorname{span}\{X, Y\}$ and $V_{2}=\operatorname{span}\{T\}$.

From now on, we use the word "intrinsic" when we want to stress a privileged role played by the horizontal layer and by group translations and dilations.

Starting from de Rham complex $\left(\Omega^{*}, d\right)$ of differential forms in $\mathbb{R}^{n}$, we look for a complex of differential forms that has to be "intrinsic" for $\mathbb{G}$ in our sense. On one side, since the "intrinsic" vector fields are naturally sections of the horizontal bundle, and hence are vector fields of the first layer of $\mathfrak{g}$, "intrinsic" 1-forms should be their dual forms (for instance, if $\mathbb{G}=\mathbb{H}^{1}$, $d x$ and $d y$ are dual of $X$ and $Y$, respectively). On the other side, it is not so evident how to choose a class of "intrinsic" forms of degree 2 or higher, but, even more, the complex we are looking for can not be merely a subcomplex of de Rham complex. Indeed, already in $\mathbb{H}^{1}$, consider a smooth function $f$ : $\mathbb{H}^{1} \rightarrow \mathbb{R}$; as we have seen, a "natural" differential would be $d_{H} f:=(X f) d x+$ $(Y f) d y$. Clearly, this is no more de Rham differential $d f=(X f) d x+$ $(Y f) d y+(T f) \theta$ (here $\theta=d t+\frac{1}{2}(y d x-x d y)$ is the so-called contact form of $\mathbb{H}^{1}$ ). In addition, if we iterate this "differential", we get $d_{H}^{2} f=[X, Y] f d x \wedge$ $d y$, that does not vanish precisely because of the lack of commutativity of the group or, equivalently, of its Lie algebra. In other words, we do not have anymore the structure of a complex. In fact, we need a more sophisticated notion of "intrinsic" exterior differential to obtain a complex of differential forms that reflects the lack of commutativity of the group. It turns out that such a complex $\left(E_{0}^{*}, d_{c}\right)$, with $E_{0}^{*} \subset \Omega^{*}$, has been defined and studied by M. Rumin in [24] and [22] ([21] for contact structures). Rumin's theory needs a quite technical introduction that is sketched in Section 3 to make the paper self-consistent. For a more exhaustive presentation, we refer to original Rumin's papers, as well as to the presentation in [2]. The main properties of $\left(E_{0}^{*}, d_{c}\right)$ can be summarized in the following points:

- Intrinsic 1-forms are horizontal 1-forms, i.e. forms that are dual of horizontal vector fields, where by duality we mean that, if $v$ is a vector field in $\mathbb{R}^{n}$, then its dual form $v^{\natural}$ acts as $v^{\natural}(w)=\langle v, w\rangle$, for all $w \in \mathbb{R}^{n}$.
- The "intrinsic" exterior differential $d_{c}$ on a smooth function is its horizontal differential (that is dual operator of the gradient along a basis of the horizontal bundle).
- The complex $\left(E_{0}^{*}, d_{c}\right)$ is exact and self-dual under Hodge $*$-duality.

The first two properties above clearly fit our request for an "intrinsic" theory. Another evidence is provided by Theorem 3.16 in [15], that proves what we can call the "weak naturality" of the complex under homogeneous homomorphisms of the group $\mathbb{G}$. (notice homogeneous homomorphisms between Carnot groups appear naturally as Pansu differentials of maps between Carnot groups, see [19]). In fact, let $T$ be a homogeneous homomorphism of $\mathbb{G}$ (where homogeneous means that $T\left(\delta_{\lambda} x\right)=\delta_{\lambda}(T x)$ ). In exponential coordinates, $T$ can be identified with linear map $T: \mathbb{R}^{n} \rightarrow \mathbb{R}^{n}$. Suppose now that also ${ }^{t} T$ is a homogeneous homomorphism. Then the pull-back $T^{\#}$ maps $E_{0}^{*}$ into $E_{0}^{*}$ and the following diagram is commutative:


Since the class of homogeneous homomorphisms well reflects both the group structure and the stratification, the naturality of $d_{c}$ under homogeneous homomorphisms shows the intimate connection between the complex and the Carnot group. On the other hand, the "artificial assumption" on ${ }^{t} T$ is extensively discussed in Remarks 3.13 and 3.17 of [15], and is basically motivated by the fact that we are working with classes of "true differential forms" and not with quotient classes.

Recently, Rumin's theory has been fruitfully used for several questions in differential geometry, as well as in pde's theory in Carnot groups.

We stress now that a crucial property of $d_{c}$ relies on the fact that it is generally a non-homogeneous higher order differential operator. In this perspective, let us give a gist of how non-homogeneous higher order horizontal derivatives appear in $d_{c}$. We need now the notion of weight of vectors in $\mathfrak{g}$ and, by duality, of covectors. Elements of the $j$-th layer of $\mathfrak{g}$ are said to have (pure) weight $w=j$; by duality, a 1 -covector that is dual of a vector of (pure) weight $w=j$ will be said to have (pure) weight $w=j$.

This procedure can be extended to $h$-forms. Clearly, there are forms that have no pure weight, but we can decompose $E_{0}^{h}$ in the direct sum of orthogonal spaces of pure weight forms, and therefore we can find a basis of $E_{0}^{h}$ given by orthonormal forms of increasing pure weights. We refer to such a basis as to a basis adapted to the filtration of $E_{0}^{h}$ induced by the weight.

Then, once suitable adapted bases of $h$-forms and ( $h+1$ )-forms are chosen, $d_{c}$ can be seen as a matrix-valued operator such that, if $\alpha$ has weight $p$, then the component of weight $q$ of $d_{c} \alpha$ is given by an homogeneous differential
operator in the horizontal derivatives of order $q-p \geq 1$, acting on the components of $\alpha$.

In order to provide a concrete example of these phenomena, let us consider again the case $\mathbb{G}=\mathbb{H}^{1}$. We remind that $X^{\natural}=d x, Y^{\natural}=d y, T^{\natural}=\theta$. In this case

$$
\begin{aligned}
& E_{0}^{1}=\operatorname{span}\{d x, d y\} \\
& E_{0}^{2}=\operatorname{span}\{d x \wedge \theta, d y \wedge \theta\} \\
& E_{0}^{3}=\operatorname{span}\{d x \wedge d y \wedge \theta\} .
\end{aligned}
$$

The action of $d_{c}$ on $E_{0}^{1}$ is the following ([21], [13], [3]): let $\alpha=\alpha_{1} d x+\alpha_{2} d y \in$ $E_{0}^{1}$ be given. Then

$$
\begin{aligned}
d_{c} \alpha= & \left(X^{2} \alpha_{2}-2 X Y \alpha_{1}+Y X \alpha_{1}\right) d x \wedge \theta \\
& +\left(2 Y X \alpha_{2}-Y^{2} \alpha_{1}-X Y \alpha_{2}\right) d y \wedge \theta .
\end{aligned}
$$

We see that $d_{c}$ is a homogeneous operator of order 2 in the horizontal derivatives, since 2 -forms have weight 3 and 1 -forms have weight 1 .

In this paper we want to provide another evidence of the intrinsic character of Rumin's complex, in the spirit of the Riemannian approximation, like in [17] and [23]. More precisely, we want to show that the intrinsic differential $d_{c}$ is a limit of suitably weighted usual first order de Rham differentials $d_{\varepsilon}$. For this purpose, we notice preliminarily that the usual exterior differential $d$ acting on a form $\alpha$ of pure weight splits as

$$
d \alpha=d_{0} \alpha+d_{1} \alpha+\cdots+d_{\kappa} \alpha,
$$

where $d_{0} \alpha$ does not increase the weight, $d_{1} \alpha$ increases the weight by 1 , and, more generally, $d_{i} \alpha$ increases the weight by $i$ when $i=0,1, \ldots, \kappa$. Then, we define the $\varepsilon$-differential weighting the different terms of $d$ according to their different actions with respect to the stratification of the Lie algebra $\mathfrak{g}$. Therefore we set

$$
d_{\varepsilon}=d_{0}+\varepsilon d_{1}+\cdot+\varepsilon^{\kappa} d_{\kappa} .
$$

The issue now is to specify in what sense the $d_{\varepsilon}$ (that is a first order operator) converges to $d$, that is, in general, a higher order differential operator, as it has already been pointed out. Keep in mind somehow similar phenomena in elasticity theory, where, roughly speaking, the equations for vibrating plates (that are 4th-order differential equations) can be seen as limits of usual lower order equations for elastic materials. In these cases, the natural approach relies in the use of De Giorgi's $\Gamma$-convergence ( $[9],[8]$, and see also Section 4 below for precise definitions in our setting) for variational functionals (see, for instance $[7]$ and the references therein). Indeed, we are able to prove that the $L^{2}$-energies associated with $\varepsilon^{-\kappa} d_{\varepsilon}$ on 1 -forms $\Gamma$-converge, as $\varepsilon \rightarrow 0$, to the energy associated with $d_{c}$. We stress that intrinsic 1-forms in groups appear in several applications, like $H$-convergence of elliptic operators on groups ([3], [2]) and Maxwell's equations in Carnot groups ([4], [14], [15]).

More precisely, the main theorem of the present paper reads as follows. If we denote by $W^{\kappa, 2}\left(\mathbb{G}, \bigwedge^{1} \mathfrak{g}\right)$ the space of differential 1 -forms on $\mathbb{G}$ with coefficients belonging to the Folland-Stein space $W^{\kappa, 2}(\mathbb{G})$ (see Definition 2.2 ), we have:

Theorem 1.1. Let $\mathbb{G}$ be a free Carnot group of step $\kappa$. If $\omega \in W^{\kappa, 2}\left(\mathbb{G}, \bigwedge^{1} \mathfrak{g}\right)$, we set

$$
F_{\varepsilon}(\omega)=\frac{1}{\varepsilon^{2 \kappa}} \int_{\mathbb{G}}\left|d_{\varepsilon} \omega\right|^{2} d V
$$

where

$$
d_{\varepsilon}=d_{0}+\varepsilon d_{1}+\cdot+\varepsilon^{\kappa} d_{\kappa}
$$

Then $F_{\varepsilon}$ sequentially $\Gamma$-coverges to $F$ in the weak topology $W^{\kappa, 2}\left(\mathbb{G}, \Lambda^{1} \mathfrak{g}\right)$, as $\varepsilon \rightarrow 0$, where

$$
F(\omega)=\left\{\begin{array}{c}
\int_{\mathbb{G}}\left|d_{c} \omega\right|^{2} d V \quad \text { if } \omega \in W^{\kappa, 2}\left(\mathbb{G}, E_{0}^{1}\right) \\
+\infty \quad \text { otherwise }
\end{array}\right.
$$

We remind that the group $\mathbb{G}$ is said to be free if its Lie algebra is free, i.e. the commutators satisfy no linear relationships other than antisymmetry and the Jacobi identity. This is a large and relevant class of Carnot groups. We remind also that Carnot groups can always be "lifted" to free groups (see [20] and [5], Chapter 17). For our purposes, the main property of free Carnot groups relies on the fact that intrinsic 1 -forms and 2 -forms on free groups have all the same weight (see Theorem 3.10). This helps at several steps of the proofs. Unfortunately, the same assertion fails to hold for higher order forms (see Remark 3.11).

Finally, another point has to be put in evidence, i.e. the choice of the topology. Indeed, we prove a $\Gamma$-limit result with respect to the sequential weak convergence, and it is natural to ask whether we could get rid off the restriction "sequential". This would be possible if we had some kind of coercitivity of the functionals (see [8], Chapter 8). However, this is not the case, since our functionals contain only the $L^{2}$-norm of the differential and not of the codifferential, where it is well known that, already in the classical Euclidean setting, the differential alone does not control the $W^{1,2_{-}}$ norms (think of Gaffney's inequality: see e.g. [25], Corollary 2.1.6). For the same reasons our convergence result is not meant to derive a convergence of minima, but only to show in what sense Rumin's differential can be seen as the limit of "Riemannian" differentials.

## 2. Carnot groups

Let $(\mathbb{G}, \cdot)$ be a Carnot group of step $\kappa$ identified to $\mathbb{R}^{n}$ through exponential coordinates (see [5] for details). By definition, the Lie algebra $\mathfrak{g}$ has dimension $n$, and admits a step $\kappa$ stratification, i.e. there exist linear subspaces $V_{1}, \ldots, V_{\kappa}$ (the layers of the stratification) such that
(1) $\mathfrak{g}=V_{1} \oplus \ldots \oplus V_{\kappa}, \quad\left[V_{1}, V_{i}\right]=V_{i+1}, \quad V_{\kappa} \neq\{0\}, \quad V_{i}=\{0\}$ if $i>\kappa$,
where $\left[V_{1}, V_{i}\right.$ ] is the subspace of $\mathfrak{g}$ generated by the commutators $[X, Y]$ with $X \in V_{1}$ and $Y \in V_{i}$. Set $m_{i}=\operatorname{dim}\left(V_{i}\right)$, for $i=1, \ldots, \kappa$ and $h_{i}=m_{1}+\cdots+m_{i}$ with $h_{0}=0$. Clearly, $h_{\kappa}=n$. Choose now a basis $e_{1}, \ldots, e_{n}$ of $\mathfrak{g}$ adapted to the stratification, i.e. such that

$$
e_{h_{j-1}+1}, \ldots, e_{h_{j}} \text { is a basis of } V_{j} \text { for each } j=1, \ldots, \kappa
$$

We refer to the first layer $V_{1}$ as to the horizontal layer. It plays a key role in our theory, since it generates the all of $\mathfrak{g}$ by commutation.

Let $X=\left\{X_{1}, \ldots, X_{n}\right\}$ be the family of left invariant vector fields such that $X_{i}(0)=e_{i}$. Given (1), the subset $X_{1}, \ldots, X_{m_{1}}$ generates by commutations all the other vector fields; we will refer to $X_{1}, \ldots, X_{m_{1}}$ as to the generating vector fields of the algebra, or as to the horizontal derivatives of the group.

The Lie algebra $\mathfrak{g}$ can be endowed with a scalar product $\langle\cdot, \cdot\rangle$, making $\left\{X_{1}, \ldots, X_{n}\right\}$ be an orthonormal basis.

We can write the elements of $\mathbb{G}$ in exponential coordinates, identifying $p$ with the n-tuple $\left(p_{1}, \ldots, p_{n}\right) \in \mathbb{R}^{n}$ and we identify $\mathbb{G}$ with $\left(\mathbb{R}^{n}, \cdot\right)$, where the explicit expression of the group operation • is determined by the CampbellHausdorff formula.

For any $x \in \mathbb{G}$, the (left) translation $\tau_{x}: \mathbb{G} \rightarrow \mathbb{G}$ is defined as

$$
z \mapsto \tau_{x} z:=x \cdot z
$$

For any $\lambda>0$, the dilation $\delta_{\lambda}: \mathbb{G} \rightarrow \mathbb{G}$, is defined as

$$
\begin{equation*}
\delta_{\lambda}\left(x_{1}, \ldots, x_{n}\right)=\left(\lambda^{d_{1}} x_{1}, \ldots, \lambda^{d_{n}} x_{n}\right), \tag{2}
\end{equation*}
$$

where $d_{i} \in \mathbb{N}$ is called homogeneity of the variable $x_{i}$ in $\mathbb{G}$ (see [12] Chapter 1 ) and is defined as

$$
\begin{equation*}
d_{j}=i \quad \text { whenever } h_{i-1}+1 \leq j \leq h_{i}, \tag{3}
\end{equation*}
$$

hence $1=d_{1}=\ldots=d_{m_{1}}<d_{m_{1}+1}=2 \leq \ldots \leq d_{n}=\kappa$.
The Haar measure of $\mathbb{G}=\left(\mathbb{R}^{n}, \cdot\right)$ is the Lebesgue measure $\mathcal{L}^{n}$ in $\mathbb{R}^{n}$.
We denote also by $Q$ the homogeneous dimension of $\mathbb{G}$, i.e. we set

$$
Q:=\sum_{i=1}^{\kappa} i \operatorname{dim}\left(V_{i}\right) .
$$

The Euclidean space $\mathbb{R}^{n}$ endowed with the usual (commutative) sum of vectors provides the simplest example of Carnot group. It is a trivial example, since in this case the stratification of the algebra consists of only one layer, i.e. the Lie algebra reduces to the horizontal layer.

Definition 2.1. Let $m \geq 2$ and $\kappa \geq 1$ be fixed integers. We say that $\mathfrak{f}_{m, \kappa}$ is the free Lie algebra with $m$ generators $x_{1}, \ldots, x_{m}$ and nilpotent of step $\kappa$ if:
i) $\mathfrak{f}_{m, \kappa}$ is a Lie algebra generated by its elements $x_{1}, \ldots, x_{m}$, i.e. $\mathfrak{f}_{m, \kappa}=$ $\operatorname{Lie}\left(x_{1}, \ldots, x_{m}\right)$;
ii) $\mathfrak{f}_{m, \kappa}$ is nilpotent of step $\kappa$;
iii) for every Lie algebra $\mathfrak{n}$ nilpotent of step $\kappa$ and for every map $\phi$ from the set $\left\{x_{1}, \ldots, x_{m}\right\}$ to $\mathfrak{n}$, there exists a (unique) homomorphism of Lie algebras $\Phi$ from $\mathfrak{f}_{m, \kappa}$ to $\mathfrak{n}$ which extends $\phi$.
The Carnot group $\mathbb{G}$ is said free if its Lie algebra $\mathfrak{g}$ is isomorphic to a free Lie algebra.

When $\mathbb{G}$ is a free group, we can assume $\left\{X_{1}, \ldots, X_{n}\right\}$ a Grayson-GrossmanHall basis of $\mathfrak{g}$ (see [16], [5], Theorem 14.1.10). This makes several computations much simpler. In particular, $\left\{\left[X_{i}, X_{j}\right], X_{i}, X_{j} \in V_{1}, i<j\right\}$ provides an orthonormal basis of $V_{2}$.

From now on, following [12], we also adopt the following multi-index notation for higher-order derivatives. If $I=\left(i_{1}, \ldots, i_{n}\right)$ is a multi-index, we
set $X^{I}=X_{1}^{i_{1}} \cdots X_{n}^{i_{n}}$. By the Poincaré-Birkhoff-Witt theorem (see, e.g. [6], I.2.7), the differential operators $X^{I}$ form a basis for the algebra of left invariant differential operators in $\mathbb{G}$. Furthermore, we set $|I|:=i_{1}+\cdots+i_{n}$ the order of the differential operator $X^{I}$, and $d(I):=d_{1} i_{1}+\cdots+d_{n} i_{n}$ its degree of homogeneity with respect to group dilations. From the Poincaré-Birkhoff-Witt theorem, it follows, in particular, that any homogeneous linear differential operator in the horizontal derivatives can be expressed as a linear combination of the operators $X^{I}$ of the special form above.

Since here we are dealing only with integer order Folland-Stein function spaces, we can this simpler definition (for a general presentation, see e.g. [11]).
Definition 2.2. If $1<s<\infty$ and $m \in \mathbb{N}$, then the space $W_{\mathbb{G}}^{m, s}(\mathbb{G})$ is the space of all $u \in L^{s}(\mathbb{G})$ such that

$$
X^{I} u \in L^{s}(\mathbb{G}) \quad \text { for all multi-index } I \text { with } d(I)=m
$$

endowed with the natural norm.
We remind that
Proposition 2.3 ([11], Corollary 4.14). If $1<s<\infty$ and $m \geq 0$, then the space $W_{\mathbb{G}}^{m, s}(\mathbb{G})$ is independent of the choice of $X_{1}, \ldots, X_{m_{1}}$.

Proposition 2.4. If $1<s<\infty$ and $m \geq 0$, then $\mathcal{S}(\mathbb{G})$ and $\mathcal{D}(\mathbb{G})$ are dense subspaces of $W_{\mathbb{G}}^{m, s}(\mathbb{G})$.

The dual space of $\mathfrak{g}$ is denoted by $\bigwedge^{1} \mathfrak{g}$. The basis of $\bigwedge^{1} \mathfrak{g}$, dual of the basis $X_{1}, \cdots, X_{n}$, is the family of covectors $\left\{\theta_{1}, \cdots, \theta_{n}\right\}$. We indicate by $\langle\cdot, \cdot\rangle$ also the inner product in $\Lambda^{1} \mathfrak{g}$ that makes $\theta_{1}, \cdots, \theta_{n}$ an orthonormal basis. We point out that, except for the trivial case of the commutative group $\mathbb{R}^{n}$, the forms $\theta_{1}, \cdots, \theta_{n}$ may have polynomial (hence variable) coefficients.

Following Federer (see [10] 1.3), the exterior algebras of $\mathfrak{g}$ and of $\bigwedge^{1} \mathfrak{g}$ are the graded algebras indicated as $\bigwedge_{*} \mathfrak{g}=\bigoplus_{h=0}^{n} \bigwedge_{h} \mathfrak{g}$ and $\bigwedge^{*} \mathfrak{g}=\bigoplus_{h=0}^{n} \bigwedge^{h} \mathfrak{g}$ where $\bigwedge_{0} \mathfrak{g}=\Lambda^{0} \mathfrak{g}=\mathbb{R}$ and, for $1 \leq h \leq n$,

$$
\begin{aligned}
& \bigwedge_{h} \mathfrak{g}:=\operatorname{span}\left\{X_{i_{1}} \wedge \cdots \wedge X_{i_{h}}: 1 \leq i_{1}<\cdots<i_{h} \leq n\right\}, \\
& \bigwedge^{h} \mathfrak{g}:=\operatorname{span}\left\{\theta_{i_{1}} \wedge \cdots \wedge \theta_{i_{h}}: 1 \leq i_{1}<\cdots<i_{h} \leq n\right\}
\end{aligned}
$$

The elements of $\bigwedge_{h} \mathfrak{g}$ and $\bigwedge^{h} \mathfrak{g}$ are called $h$-vectors and $h$-covectors.
We denote by $\Theta^{h}$ the basis $\left\{\theta_{i_{1}} \wedge \cdots \wedge \theta_{i_{h}}: 1 \leq i_{1}<\cdots<i_{h} \leq n\right\}$ of $\bigwedge^{h} \mathfrak{g}$. We remind that $\operatorname{dim} \bigwedge^{h} \mathfrak{g}=\operatorname{dim} \bigwedge_{h} \mathfrak{g}=\binom{n}{h}$.

The dual space $\Lambda^{1}\left(\bigwedge_{h} \mathfrak{g}\right)$ of $\Lambda_{h} \mathfrak{g}$ can be naturally identified with $\Lambda^{h} \mathfrak{g}$. The action of a $h$-covector $\varphi$ on a $h$-vector $v$ is denoted as $\langle\varphi \mid v\rangle$.

The inner product $\langle\cdot, \cdot\rangle$ extends canonically to $\bigwedge_{h} \mathfrak{g}$ and to $\Lambda^{h} \mathfrak{g}$ making the bases $X_{i_{1}} \wedge \cdots \wedge X_{i_{h}}$ and $\theta_{i_{1}} \wedge \cdots \wedge \theta_{i_{h}}$ orthonormal.

Starting from $\bigwedge_{*} \mathfrak{g}$ and $\Lambda^{*} \mathfrak{g}$, by left translation, we can define now two families of vector bundles (still denoted by $\bigwedge_{*} \mathfrak{g}$ and $\bigwedge^{*} \mathfrak{g}$ ) over $\mathbb{G}$ (see [2] for details). Sections of these vector bundles are said respectively vector fields and differential forms.

Definition 2.5. If $0 \leq h \leq n, 1 \leq s \leq \infty$ and $m \geq 0$, we denote by $W_{\mathbb{G}}^{m, s}\left(\mathbb{G}, \Lambda^{h} \mathfrak{g}\right)$ the space of all sections of $\bigwedge^{h} \mathfrak{g}$ such that their components with respect to the basis $\Theta^{h}$ belong to $W_{\mathbb{G}}^{m, s}(\mathbb{G})$, endowed with its natural norm. Clearly, this definition is independent of the choice of the basis itself.

Sobolev spaces of vector fields are defined in the same way.
We conclude this section recalling the classical definition of Hodge duality: see [10] 1.7.8.

Definition 2.6. We define linear isomorphisms

$$
*: \bigwedge_{h} \mathfrak{g} \longleftrightarrow \bigwedge_{n-h} \mathfrak{g} \text { and } *: \bigwedge^{h} \mathfrak{g} \longleftrightarrow \bigwedge^{n-h} \mathfrak{g}
$$

for $1 \leq h \leq n$, putting, for $v=\sum_{I} v_{I} X_{I}$ and $\varphi=\sum_{I} \varphi_{I} \theta_{I}$,

$$
* v:=\sum_{I} v_{I}\left(* X_{I}\right) \quad \text { and } \quad * \varphi:=\sum_{I} \varphi_{I}\left(* \theta_{I}\right)
$$

where

$$
* X_{I}:=(-1)^{\sigma(I)} X_{I^{*}} \quad \text { and } \quad * \theta_{I}:=(-1)^{\sigma(I)} \theta_{I^{*}}
$$

with $I=\left\{i_{1}, \cdots, i_{h}\right\}, 1 \leq i_{1}<\cdots<i_{h} \leq n, X_{I}=X_{i_{1}} \wedge \cdots \wedge X_{i_{h}}$, $\theta_{I}=\theta_{i_{1}} \wedge \cdots \wedge \theta_{i_{h}}, I^{*}=\left\{i_{1}^{*}<\cdots<i_{n-h}^{*}\right\}=\{1, \cdots, n\} \backslash I$ and $\sigma(I)$ is the number of couples $\left(i_{h}, i_{\ell}^{*}\right)$ with $i_{h}>i_{\ell}^{*}$.

The following properties of the $*$ operator follow readily from the definition: $\forall v, w \in \bigwedge_{h} \mathfrak{g}$ and $\forall \varphi, \psi \in \bigwedge^{h} \mathfrak{g}$

$$
\begin{align*}
& * * v=(-1)^{h(n-h)} v, \quad * * \varphi=(-1)^{h(n-h)} \varphi, \\
& v \wedge * w=\langle v, w\rangle X_{\{1, \cdots, n\}}, \quad \varphi \wedge * \psi=\langle\varphi, \psi\rangle \theta_{\{1, \cdots, n\}},  \tag{4}\\
& \langle * \varphi \mid * v\rangle=\langle\varphi \mid v\rangle .
\end{align*}
$$

We refer to $d V:=\theta_{\{1, \cdots, n\}}$ as to the canonical volume form of $\mathbb{G}$.
If $v \in \Lambda_{h} \mathfrak{g}$ we define $v^{\natural} \in \Lambda^{h} \mathfrak{g}$ by the identity $\left\langle v^{\natural} \mid w\right\rangle:=\langle v, w\rangle$, for all $w \in \bigwedge_{h} \mathfrak{g}$, and analogously we define $\varphi^{\natural} \in \bigwedge_{h} \mathfrak{g}$ for $\varphi \in \bigwedge^{h} \mathfrak{g}$.

## 3. Differential forms in Carnot groups

The notion of intrinsic form in Carnot groups is due to M. Rumin ([24], [22]). A more extended presentation of the results of this section can be found in [2], [15].

The notion of weight of a differential form plays a key role.
Definition 3.1. If $\alpha \in \Lambda^{1} \mathfrak{g}, \alpha \neq 0$, we say that $\alpha$ has pure weight $p$, and we write $w(\alpha)=p$, if $\alpha^{\natural} \in V_{p}$. More generally, if $\alpha \in \Lambda^{h} \mathfrak{g}$, we say that $\alpha$ has pure weight $p$ if $\alpha$ is a linear combination of covectors $\theta_{i_{1}} \wedge \cdots \wedge \theta_{i_{h}}$ with $w\left(\theta_{i_{1}}\right)+\cdots+w\left(\theta_{i_{h}}\right)=p$.

In particular, the canonical volume form $d V$ has weight $Q$ (the homogeneous dimension of the group).
Remark 3.2. If $\alpha, \beta \in \bigwedge^{h} \mathfrak{g}$ and $w(\alpha) \neq w(\beta)$, then $\langle\alpha, \beta\rangle=0$. Indeed, it is enough to notice that, if $w\left(\theta_{i_{1}} \wedge \cdots \wedge \theta_{i_{h}}\right) \neq w\left(\theta_{j_{1}} \wedge \cdots \wedge \theta_{j_{h}}\right)$, with $i_{1}<i_{2}<\cdots<i_{h}$ and $j_{1}<j_{2}<\cdots<j_{h}$, then for at least one of the indices $\ell=1, \ldots, h, i_{\ell} \neq j_{\ell}$, and hence $\left\langle\theta_{i_{1}} \wedge \cdots \wedge \theta_{i_{h}}, \theta_{j_{1}} \wedge \cdots \wedge \theta_{j_{h}}\right\rangle=0$.

We have ([2], formula (16))

$$
\begin{equation*}
\bigwedge^{h} \mathfrak{g}=\bigoplus_{p=M_{h}^{\min }}^{M_{h}^{\max }} \bigwedge^{h, p} \mathfrak{g} \tag{5}
\end{equation*}
$$

where $\Lambda^{h, p} \mathfrak{g}$ is the linear span of the $h$-covectors of weight $p$ and $M_{h}^{\text {min }}$, $M_{h}^{\max }$ are respectively the smallest and the largest weight of left-invariant $h$-covectors.

Since the elements of the basis $\Theta^{h}$ have pure weights, a basis of $\bigwedge^{h, p} \mathfrak{g}$ is given by $\Theta^{h, p}:=\Theta^{h} \cap \bigwedge^{h, p} \mathfrak{g}$. In other words, the basis $\Theta^{h}=\cup_{p} \Theta^{h, p}$ is a basis adapted to the filtration of $\bigwedge^{h} \mathfrak{g}$ associated with (5).

We denote by $\Omega^{h, p}$ the vector space of all smooth $h$-forms in $\mathbb{G}$ of pure weight $p$, i.e. the space of all smooth sections of $\bigwedge^{h, p} \mathfrak{g}$. We have

$$
\begin{equation*}
\Omega^{h}=\bigoplus_{p=M_{h}^{\min }}^{M_{h}^{\max }} \Omega^{h, p} . \tag{6}
\end{equation*}
$$

The following crucial property of the weight follows from Cartan identintity: see [24], Section 2.1:
Lemma 3.3. We have $d\left(\bigwedge^{h, p} \mathfrak{g}\right) \subset \bigwedge^{h+1, p} \mathfrak{g}$, i.e., if $\alpha \in \bigwedge^{h, p} \mathfrak{g}$ is a left invariant $h$-form of weight $p$ with $d \alpha \neq 0$, then $w(d \alpha)=w(\alpha)$.

Definition 3.4. Let now $\alpha=\sum_{\theta_{i}^{h} \in \Theta^{h, p}} \alpha_{i} \theta_{i}^{h} \in \Omega^{h, p}$ be a (say) smooth form of pure weight $p$. Then we can write

$$
d \alpha=d_{0} \alpha+d_{1} \alpha+\cdots+d_{\kappa} \alpha
$$

where

$$
d_{0} \alpha=\sum_{\theta_{i}^{h} \in \Theta^{h, p}} \alpha_{i} d \theta_{i}^{h}
$$

does not increase the weight,

$$
d_{1} \alpha=\sum_{\theta_{i}^{h} \in \Theta^{h, p}} \sum_{j=1}^{m_{1}}\left(X_{j} \alpha_{i}\right) \theta_{j} \wedge \theta_{i}^{h}
$$

increases the weight of 1 , and, more generally,

$$
d_{i} \alpha=\sum_{\theta_{i}^{h} \in \Theta^{h, p}} \sum_{X_{j} \in V_{i}}\left(X_{j} \alpha_{i}\right) \theta_{j} \wedge \theta_{i}^{h}
$$

when $i=0,1, \ldots, \kappa$. In particular, $d_{0}$ is an algebraic operator.
Definition 3.5 (M. Rumin). If $0 \leq h \leq n$ we set

$$
E_{0}^{h}:=\operatorname{ker} d_{0} \cap \operatorname{ker} \delta_{0}=\operatorname{ker} d_{0} \cap\left(\operatorname{Im} d_{0}\right)^{\perp} \subset \Omega^{h}
$$

In the sequel, we refer to the elements of $E_{0}^{h}$ as to intrinsic $h$-forms on $\mathbb{G}$. Since the construction of $E_{0}^{h}$ is left invariant, this space of forms can be seen as the space of sections of a fiber subbundle of $\Lambda^{h} \mathfrak{g}$, generated by left translation and still denoted by $E_{0}^{h}$. In particular $E_{0}^{h}$ inherits from $\Lambda^{h} \mathfrak{g}$ the scalar product on the fibers.

Moreover, there exists a left invariant orthonormal basis $\Xi_{0}^{h}=\left\{\xi_{j}\right\}$ of $E_{0}^{h}$ that is adapted to the filtration (5).

Since it is easy to see that $E_{0}^{1}=\operatorname{span}\left\{\theta_{1}, \ldots, \theta_{m}\right\}$, where the $\theta_{i}$ 's are dual of the elements of the basis of $V_{1}$, without loss of generality, we can take $\xi_{j}=\theta_{j}$ for $j=1, \ldots, m$.

Finally, we denote by $N_{h}^{\min }$ and $N_{h}^{\max }$ respectively the lowest and highest weight of forms in $E_{0}^{h}$.

Definition 3.6. If $0 \leq h \leq n, 1 \leq s \leq \infty$ and $m \geq 0$, we denote by $W_{\mathbb{G}}^{m, s}\left(\mathbb{G}, E_{0}^{h}\right)$ the space of all sections of $E_{0}^{h}$ such that their components with respect to the basis $\Xi_{0}^{h}$ belong to $W_{\mathbb{G}}^{m, s}(\mathbb{G})$, endowed with its natural norm. Clearly, this definition is independent of the choice of the basis itself.

Moreover, as in Proposition 2.4, $\mathcal{D}\left(\mathbb{G}, E_{0}^{h}\right)$ and $\mathcal{S}\left(\mathbb{G}, E_{0}^{h}\right)$ are dense in $W_{\mathbb{G}}^{m, s}(\mathbb{G})$.

Lemma 3.7 ([2], Lemma 2.11). If $\beta \in \bigwedge^{h+1} \mathfrak{g}$, then there exists a unique $\alpha \in \bigwedge^{h} \mathfrak{g} \cap\left(\operatorname{ker} d_{0}\right)^{\perp}$ such that

$$
d_{0}^{*} d_{0} \alpha=d_{0}^{*} \beta . \quad \text { We set } \quad \alpha:=d_{0}^{-1} \beta
$$

Here $d_{0}^{*}: \Lambda^{h+1} \mathfrak{g} \rightarrow \bigwedge^{h} \mathfrak{g}$ is the adjoint of $d_{0}$ with respect to our fixed scalar product. In particular

$$
\alpha=d_{0}^{-1} \beta \quad \text { if and only if } \quad d_{0} \alpha-\beta \in \mathcal{R}\left(d_{0}\right)^{\perp} .
$$

In particular
i) $\left(\operatorname{ker} d_{0}\right)^{\perp}=\mathcal{R}\left(d_{0}^{-1}\right)$;
ii) $d_{0}^{-1} d_{0}=I d$ on $\left(\operatorname{ker} d_{0}\right)^{\perp}$;
iii) $d_{0} d_{0}^{-1}-I d: \bigwedge^{h+1} \mathfrak{g} \rightarrow \mathcal{R}\left(d_{0}\right)^{\perp}$.

The following theorem summarizes the construction of the intrinsic differential $d_{c}$ (for details, see [24] and [2], Section 2) .

Theorem 3.8. The de Rham complex $\left(\Omega^{*}, d\right)$ splits in the direct sum of two sub-complexes $\left(E^{*}, d\right)$ and $\left(F^{*}, d\right)$, with

$$
E:=\operatorname{ker} d_{0}^{-1} \cap \operatorname{ker}\left(d_{0}^{-1} d\right) \quad \text { and } \quad F:=\mathcal{R}\left(d_{0}^{-1}\right)+\mathcal{R}\left(d d_{0}^{-1}\right) .
$$

We have
i) Let $\Pi_{E}$ be the projection on $E$ along $F$ (that is not an orthogonal projection). Then for any $\alpha \in E_{0}^{h, p}$, if we denote by $\left(\Pi_{E} \alpha\right)_{j}$ the component of $\Pi_{E} \alpha$ of weight $j$, then

$$
\begin{aligned}
\left(\Pi_{E} \alpha\right)_{p} & =\alpha \\
\left(\Pi_{E} \alpha\right)_{p+k+1} & =-d_{0}^{-1}\left(\sum_{1 \leq \ell \leq k+1} d_{\ell}\left(\Pi_{E} \alpha\right)_{p+k+1-\ell}\right)
\end{aligned}
$$

Notice that $\alpha \rightarrow\left(\Pi_{E} \alpha\right)_{p+k+1}$ is an homogeneous differential operator of order $k+1$ in the horizontal derivatives.
ii) $\Pi_{E}$ is a chain map, i.e.

$$
d \Pi_{E}=\Pi_{E} d
$$

iii) Let $\Pi_{E_{0}}$ be the orthogonal projection from $\Omega^{*}$ on $E_{0}^{*}$, then

$$
\begin{equation*}
\Pi_{E_{0}}=I d-d_{0}^{-1} d_{0}-d_{0} d_{0}^{-1}, \quad \Pi_{E_{0}^{\perp}}=d_{0}^{-1} d_{0}+d_{0} d_{0}^{-1} \tag{8}
\end{equation*}
$$

Notice that, since $d_{0}$ and $d_{0}^{-1}$ are algebraic, then formulas (8) hold also for covectors.
iv) $\Pi_{E_{0}} \Pi_{E} \Pi_{E_{0}}=\Pi_{E_{0}}$ and $\Pi_{E} \Pi_{E_{0}} \Pi_{E}=\Pi_{E}$.

Set now

$$
d_{c}=\Pi_{E_{0}} d \Pi_{E}: E_{0}^{h} \rightarrow E_{0}^{h+1}, \quad h=0, \ldots, n-1 .
$$

We have:
v) $d_{c}^{2}=0$;
vi) the complex $E_{0}:=\left(E_{0}^{*}, d_{c}\right)$ is exact;
vii) with respect to the bases $\Xi_{0}^{*}$ the intrinsic differential $d_{c}$ can be seen as a matrix-valued operator such that, if $\alpha$ has weight $p$, then the component of weight $q$ of $d_{c} \alpha$ is given by an homogeneous differential operator in the horizontal derivatives of order $q-p \geq 1$, acting on the components of $\alpha$.

Remark 3.9. Let us give a gist of the construction of $E$. The map $d_{0}^{-1} d$ induces an isomorphism from $\mathcal{R}\left(d_{0}^{-1}\right)$ to itself. Thus, since $d_{0}^{-1} d_{0}=I d$ on $\mathcal{R}\left(d_{0}^{-1}\right)$, we can write $d_{0}^{-1} d=I d+D$, where $D$ is a differential operator that increases the weight. Clearly, $D: \mathcal{R}\left(d_{0}^{-1}\right) \rightarrow \mathcal{R}\left(d_{0}^{-1}\right)$. As a consequence of the nilpotency of $\mathbb{G}, D^{k}=0$ for $k$ large enough, and therefore the Neumann series of $d_{0}^{-1} d$ reduces to a finite sum on $\mathcal{R}\left(d_{0}^{-1}\right)$. Hence there exist a differential operator

$$
P=\sum_{k=1}^{N}(-1)^{k} D^{k}, \quad N \in \mathbb{N} \text { suitable }
$$

such that

$$
P d_{0}^{-1} d=d_{0}^{-1} d P=\operatorname{Id}_{\mathcal{R}\left(d_{0}^{-1}\right)} .
$$

We set $Q:=P d_{0}^{-1}$. Then $\Pi_{E}$ is given by

$$
\Pi_{E}=I d-Q d-d Q .
$$

From now on, we restrict ourselves to assume $\mathbb{G}$ is a free group of step $\kappa$ (see Definition 2.1 above). The technical reason for this choice relies in the following property.

Theorem 3.10 ([15], Theorem 5.9). Let $\mathbb{G}$ be a free group of step $\kappa$. Then all forms in $E_{0}^{1}$ have weight 1 and all forms in $E_{0}^{2}$ have weight $\kappa+1$.

In particular, the differential $d_{c}: E_{0}^{1} \rightarrow E_{0}^{2}$ can be identified, with respect to the adapted bases $\Xi_{0}^{1}$ and $\Xi_{0}^{2}$, with a homogeneous matrix-valued differential operator of degree $\kappa$ in the horizontal derivatives.

Moreover, if $\xi \in \bigwedge^{2, p} \mathfrak{g}$ with $p \neq \kappa+1$, then $\Pi_{E_{0}} \xi=0$. Indeed, $\Pi_{E_{0}} \xi$ has weight $p$, and therefore has to be zero, since $\Pi_{E_{0}} \xi \in \Lambda^{2, \kappa+1} \mathfrak{g}$.

Remark 3.11. Theorem 3.10 might suggest that in free groups all forms in $E_{0}^{*}$ have pure weight. Unfortunately, this assertion fails to hold, at least in this naïf form. Indeed, A. Ottazzi showed us a counterexample for $E_{0}^{3}$ in the free group of step 2 with 3 generators. Actually, this is a general phenomenon, due to the fact that, denoting as usual by $Q$ the homogeneous dimension of the group, in this case $n=6$ (even), so that $E_{0}^{3}=* E_{0}^{3}$, but $Q=9$ (odd), yielding a contradiction with $w(* \alpha)=w(\alpha)$ when $\alpha \in E_{0}^{3}$, since $w(* \alpha)=Q-w(\alpha)$. Clearly, this situation occurs whenever $n$ is even and $Q$ is odd.

Lemma 3.12. If $\mathbb{G}$ is a free group of step $\kappa \geq 2$, then

$$
\bigwedge^{2,2} \mathfrak{g} \subset d_{0}\left(\bigwedge^{1,2} \mathfrak{g}\right) \subset \mathcal{R}\left(d_{0}\right)
$$

or, equivalently,

$$
\Omega^{2,2} \subset d_{0}\left(\Omega^{1,2}\right) \subset d_{0}\left(\Omega^{1}\right)
$$

Proof. Since a basis of $\bigwedge^{2,2} \mathfrak{g}$ is given by covectors $\theta$ of the form $\theta=\theta_{i} \wedge \theta_{j}$, with $\theta_{i}=X_{i}^{\natural}, \theta_{j}=X_{j}^{\natural}, X_{i}, X_{j} \in V_{1}, i<j$, we need only to prove that

$$
d_{0}\left(\left[X_{i}, X_{j}\right]^{\mathrm{\natural}}\right)=-\theta_{i} \wedge \theta_{j} .
$$

Thus, if $X, Y \in \mathfrak{g}$, we are left to show that

$$
\left\langle d_{0}\left(\left[X_{i}, X_{j}\right]^{\natural}\right) \mid X \wedge Y\right\rangle=-\left\langle\theta_{i} \wedge \theta_{j} \mid X \wedge Y\right\rangle .
$$

Since $d_{0}$ preserves the weights, we may assume that $X \wedge Y$ has weight 2 . Therefore, without loss of generality, we can take $X=X_{k}, Y=X_{h}$, with $X_{k}, X_{h} \in V_{1}$. Therefore

$$
\begin{aligned}
& \left\langle d_{0}\left(\left[X_{i}, X_{j}\right]^{\natural}\right) \mid X_{k} \wedge X_{h}\right\rangle=\left\langle d\left(\left[X_{i}, X_{j}\right]^{\mathrm{h}}\right) \mid X_{k} \wedge X_{h}\right\rangle \\
& =-\left\langle\left[X_{i}, X_{j}\right]^{\mathrm{h}} \mid\left[X_{k}, X_{h}\right]\right\rangle=-\left\langle\left[X_{i}, X_{j}\right],\left[X_{k}, X_{h}\right]\right\rangle .
\end{aligned}
$$

On the other hand, as pointed out in Definition 2.1,

- $\left\langle\left[X_{i}, X_{j}\right],\left[X_{k}, X_{h}\right]\right\rangle=0$ if $\{i, j\} \neq\{k, h\}$,
- $\left\langle\left[X_{i}, X_{j}\right],\left[X_{k}, X_{h}\right]\right\rangle=1$ if $(i, j)=(k, h)$, and
- $\left\langle\left[X_{i}, X_{j}\right],\left[X_{k}, X_{h}\right]\right\rangle=-1$ if $(i, j)=(h, k)$,
whereas

$$
\left\langle\theta_{i} \wedge \theta_{j} \mid X_{k} \wedge X_{h}\right\rangle=\operatorname{det}\left(\begin{array}{cc}
\left\langle\theta_{i} \mid X_{k}\right\rangle & \left\langle\theta_{i} \mid X_{h}\right\rangle \\
\left\langle\theta_{j} \mid X_{k}\right\rangle & \left\langle\theta_{j} \mid X_{h}\right\rangle
\end{array}\right) .
$$

This achieves the proof of the lemma.

## 4. $\Gamma$-convergence

Definition 4.1. Let $X$ be a separated topological space, and let

$$
F_{\varepsilon}, F: X \longrightarrow[-\infty,+\infty]
$$

with $\varepsilon>0$ be functionals on $X$. We say that $\left\{F_{\varepsilon}\right\}_{\varepsilon>0}$ sequentially $\Gamma$ converges to $F$ on $X$ as $\varepsilon$ goes to zero if the following two conditions hold:

1) for every $u \in X$ and for every sequence $\left\{u_{\varepsilon_{k}}\right\}_{k \in \mathbb{N}}$ with $\varepsilon_{k} \rightarrow 0$ as $k \rightarrow \infty$, which converges to $u$ in $X$, there holds

$$
\begin{equation*}
\liminf _{k \rightarrow \infty} F_{\varepsilon_{k}}\left(u_{\varepsilon_{k}}\right) \geq F(u) \tag{9}
\end{equation*}
$$

2) for every $u \in X$ and for every sequence $\left\{\varepsilon_{k}\right\}_{k \in \mathbb{N}}$ with $\varepsilon_{k} \rightarrow 0$ as $k \rightarrow \infty$ there exists a subsequence (still denoted by $\left\{\varepsilon_{k}\right\}_{k \in \mathbb{N}}$ ) such that $\left\{u_{\varepsilon_{k}}\right\}_{k \in \mathbb{N}}$ converges to $u$ in $X$ and

$$
\limsup _{k \rightarrow \infty} F_{\varepsilon_{k}}\left(u_{\varepsilon_{k}}\right) \leq F(u)
$$

For a deep and detailed survey on $\Gamma$-convergence, we refer to the monograph [8].

We recall the following reduction Lemma. The proof is only a minor variant of the one given in [18], Lemma $I V$ (see also [1]), hence we shall omit such a proof.

Lemma 4.2. Let $X$ be a separated topological space, let $F_{h}, F: M \longrightarrow$ $[-\infty,+\infty]$ with $h \in \mathbb{N}$; consider $D \subset M$ and $x \in M$. Let us suppose that

1) for every $y \in D$ there exists a sequence $\left(y_{h}\right)_{h \in \mathbb{N}} \subset M$ such that $y_{h} \rightarrow y$ in $M$ and $\limsup _{h \rightarrow \infty} F_{h}\left(y_{h}\right) \leq F(y)$;
2) there exists a sequence $\left(x_{h}\right)_{h \in \mathbb{N}} \subset D$ such that $x_{h} \rightarrow x$ and $\limsup _{h \rightarrow \infty} F\left(x_{h}\right) \leq$ $F(x) ;$
then there exists a sequence $\left(\bar{x}_{h}\right)_{h \in \mathbb{N}} \subset M$ such that $\limsup _{h \rightarrow \infty} F_{h}\left(\bar{x}_{h}\right) \leq F(x)$.
To avoid cumbersome notations, from now on we write systematically $\lim _{\varepsilon \rightarrow 0}$ to mean a limit with $\varepsilon=\varepsilon_{k}$, where $\left\{\varepsilon_{k}\right\}_{k \in \mathbb{N}}$ is any sequence with $\varepsilon_{k} \rightarrow 0$ as $k \rightarrow \infty$.

## 5. Intrinsic differential as a $\Gamma$-Limit

Let $\varepsilon>0$ be given. If $\omega \in W^{\kappa, 2}\left(\mathbb{G}, \bigwedge^{1} \mathfrak{g}\right)$, we set

$$
F_{\varepsilon}(\omega)=\frac{1}{\varepsilon^{2 \kappa}} \int_{\mathbb{G}}\left|d_{\varepsilon} \omega\right|^{2} d V
$$

where

$$
d_{\varepsilon}=d_{0}+\varepsilon d_{1}+\cdot+\varepsilon^{\kappa} d_{\kappa} .
$$

We stress that $F_{\varepsilon}(\omega)$ is always finite, since the coefficients of $d_{i} \omega$ contain horizontal derivatives of order $i \leq \kappa$ of the coefficients of $\omega$.

Theorem 5.1. Let $\mathbb{G}$ be a free Carnot group of step $\kappa$. Then
$F_{\varepsilon}$ sequentially $\Gamma$-coverges to $F$ in the weak topology $W^{\kappa, 2}\left(\mathbb{G}, \bigwedge^{1} \mathfrak{g}\right)$, as $\varepsilon \rightarrow 0$, where

$$
F(\omega)=\left\{\begin{array}{c}
\int_{\mathbb{G}}\left|d_{c} \omega\right|^{2} d V \quad \text { if } \omega \in W^{\kappa, 2}\left(\mathbb{G}, E_{0}^{1}\right) \\
+\infty \quad \text { otherwise } .
\end{array}\right.
$$

Proof. Let $\omega^{\varepsilon} \rightarrow \omega$ as $\varepsilon \rightarrow 0$ weakly in $W^{\kappa, 2}\left(\mathbb{G}, \bigwedge^{1} \mathfrak{g}\right)$. We want to show that

$$
\begin{equation*}
F(\omega) \leq \liminf _{\varepsilon \rightarrow 0} F_{\varepsilon}\left(\omega^{\varepsilon}\right) \tag{11}
\end{equation*}
$$

In particular, it follows that $\omega \in W^{\kappa, 2}\left(\mathbb{G}, E_{0}^{1}\right)$ provided $\liminf _{\varepsilon \rightarrow 0} F_{\varepsilon}\left(\omega^{\varepsilon}\right)<$ $\infty$.

Keeping in mind (6), we write

$$
\omega^{\varepsilon}=\omega_{1}^{\varepsilon}+\cdots+\omega_{\kappa}^{\varepsilon}
$$

with $\omega_{i}^{\varepsilon} \in \Omega^{1, i}, i=1, \ldots, \kappa$. Reordering the terms of $d_{\varepsilon} \omega_{\varepsilon}$ according to their weights, as in (6), we have the following orthogonal decomposition:

$$
\begin{align*}
d_{\varepsilon} \omega^{\varepsilon}= & \left(d_{0} \omega_{2}^{\varepsilon}+\varepsilon d_{1} \omega_{1}^{\varepsilon}\right) \\
& +\left(d_{0} \omega_{3}^{\varepsilon}+\varepsilon d_{1} \omega_{2}^{\varepsilon}+\varepsilon^{2} d_{2} \omega_{1}^{\varepsilon}\right) \\
& +\cdots \\
& +\left(d_{0} \omega_{\kappa}^{\varepsilon}+\varepsilon d_{1} \omega_{\kappa-1}^{\varepsilon}+\cdots \varepsilon^{\kappa-1} d_{\kappa-1} \omega_{1}^{\varepsilon}\right) \\
& +\left(\varepsilon d_{1} \omega_{\kappa}^{\varepsilon}+\cdots+\varepsilon^{\kappa} d_{\kappa} \omega_{1}^{\varepsilon}\right) \\
& +\cdots \\
& +\left(\varepsilon^{\kappa-1} d_{\kappa-1} \omega_{\kappa}^{\varepsilon}+\varepsilon^{\kappa} d_{\kappa} \omega_{\kappa-1}^{\varepsilon}\right)  \tag{12}\\
& +\varepsilon^{\kappa} d_{\kappa} \omega_{\kappa}^{\varepsilon} \\
= & \sum_{2 \leq p \leq \kappa} \sum_{i=0}^{p-1} \varepsilon^{i} d_{i} \omega_{p-i}^{\varepsilon} \\
& +\left(\varepsilon d_{1} \omega_{\kappa}^{\varepsilon}+\cdots+\varepsilon^{\kappa} d_{\kappa} \omega_{1}^{\varepsilon}\right) \\
& +\sum_{\kappa+2 \leq p \leq 2 \kappa} \sum_{i=p-\kappa}^{\kappa} \varepsilon^{i} d_{i} \omega_{p-i}^{\varepsilon}
\end{align*}
$$

Therefore we can write

$$
\begin{align*}
F_{\varepsilon}\left(\omega_{\varepsilon}\right)=\varepsilon^{-2 \kappa} & \sum_{2 \leq p \leq \kappa} \int_{\mathbb{G}}\left\|\sum_{i=0}^{p-1} \varepsilon^{i} d_{i} \omega_{p-i}^{\varepsilon}\right\|^{2} d V \\
& +\varepsilon^{2(1-\kappa)} \int_{\mathbb{G}}\left\|d_{1} \omega_{\kappa}^{\varepsilon}+\cdots+\varepsilon^{\kappa-1} d_{\kappa} \omega_{1}^{\varepsilon}\right\|^{2} d V  \tag{13}\\
& +\sum_{\kappa+2 \leq p \leq 2 \kappa} \varepsilon^{2(p-2 \kappa)} \int_{\mathbb{G}}\left\|\sum_{i=p-\kappa}^{\kappa} \varepsilon^{i-p+\kappa} d_{i} \omega_{p-i}^{\varepsilon}\right\|^{2} d V
\end{align*}
$$

Without loss of generality, we may assume $\liminf _{\varepsilon \rightarrow 0} F_{\varepsilon}\left(\omega_{\varepsilon}\right)<\infty$. This implies that, if $2 \leq p \leq \kappa$, then, if $\varepsilon \in(0,1)$,

$$
\begin{equation*}
\left.\varepsilon^{-\kappa} \sum_{i=0}^{p-1} \varepsilon^{i} d_{i} \omega_{p-i}^{\varepsilon} \quad \text { is uniformly bounded in } L^{2}\left(\mathbb{G}, \bigwedge^{2} \mathfrak{g}\right)\right) \tag{14}
\end{equation*}
$$

Moreover, again if $\varepsilon \in(0,1)$,
(15) $\varepsilon^{1-\kappa}\left(d_{1} \omega_{\kappa}^{\varepsilon}+\cdots+\varepsilon^{\kappa-1} d_{\kappa} \omega_{1}^{\varepsilon}\right) \quad$ is uniformly bounded in $\left.L^{2}\left(\mathbb{G}, \bigwedge^{2} \mathfrak{g}\right)\right)$.

In particular,

$$
\begin{equation*}
\left.\sum_{i=0}^{p-1} \varepsilon^{i} d_{i} \omega_{p-i}^{\varepsilon} \longrightarrow 0 \quad \text { in } L^{2}\left(\mathbb{G}, \bigwedge^{2} \mathfrak{g}\right)\right) \tag{16}
\end{equation*}
$$

as $\varepsilon \rightarrow 0$, since we can write (16) as

$$
\begin{equation*}
d_{0} \omega_{p}^{\varepsilon}+\varepsilon \sum_{i=1}^{p-1} \varepsilon^{i-1} d_{i} \omega_{p-i}^{\varepsilon} \longrightarrow 0 \tag{17}
\end{equation*}
$$

as $\varepsilon \rightarrow 0$. By assumption, we know that $\omega_{p}^{\varepsilon} \rightarrow \omega_{p}$ weakly in $L^{2}\left(\mathbb{G}, \bigwedge^{1} \mathfrak{g}\right)$ for $p \geq 1$, and therefore

$$
\begin{equation*}
d_{0} \omega_{p}^{\varepsilon} \rightarrow d_{0} \omega_{p} \quad \text { in } L^{2}\left(\mathbb{G}, \bigwedge^{1} \mathfrak{g}\right) \tag{18}
\end{equation*}
$$

since $d_{0}$ is algebraic.
Combining (17) with the boundedness of $\left\{\omega^{\varepsilon}\right\}$ in $W^{\kappa, 2}\left(\mathbb{G}, \bigwedge^{1} \mathfrak{g}\right)$ and with (18), it follows that

$$
\begin{equation*}
d_{0} \omega_{p}=0 \quad \text { for } p=2, \ldots, \kappa \tag{19}
\end{equation*}
$$

(obviously, (19) holds also for $p=1$ since $\left.d_{0}\left(\bigwedge^{1,1} \mathfrak{g}\right)=\{0\}\right)$. Hence $\omega \in$ ker $d_{0}=E_{0}^{1}$, and therefore $\omega=\omega_{1}$.

Recall now that, by definition, $d_{c} \omega=\Pi_{E_{0}} d \Pi_{E} \omega$. But, by Theorem 3.10, $\Pi_{E_{0}}$ vanishes on all 2-forms of weight $p \neq \kappa+1$. Therefore, the full expression of $d_{c} \omega$ reduces to

$$
\begin{equation*}
d_{c}(\omega)=\Pi_{E_{0}}\left(\sum_{\ell=1}^{\kappa} d_{\ell}\left(\Pi_{E} \omega\right)_{\kappa+1-\ell}\right) . \tag{20}
\end{equation*}
$$

Let us show now that

$$
\begin{equation*}
d_{j}\left(\Pi_{E} \omega\right)_{\kappa+1-\ell}=\lim _{\varepsilon \rightarrow 0} \varepsilon^{\ell-\kappa} d_{j} \omega_{\kappa+1-\ell}^{\varepsilon}, \tag{21}
\end{equation*}
$$

in the sense of distributions for $\ell=1, \ldots, \kappa, j=0, \ldots, \ell$, i.e.

$$
\begin{equation*}
\int\left\langle\varepsilon^{\ell-\kappa} d_{j} \omega_{\kappa+1-\ell}^{\varepsilon}, \varphi\right\rangle d V \rightarrow \int\left\langle d_{j}\left(\Pi_{E} \omega\right)_{\kappa+1-\ell}, \varphi\right\rangle d V \tag{22}
\end{equation*}
$$

for $\ell=1, \ldots, \kappa, j=0, \ldots, \ell$, and for any $\varphi \in \mathcal{D}\left(\mathbb{G}, \bigwedge^{2} \mathfrak{g}\right)$. Notice $\left(\Pi_{E} \omega\right)_{\kappa+1-\ell} \in$ $W^{\ell, 2}\left(\mathbb{G}, \Lambda^{1} \mathfrak{g}\right)$, by Theorem 3.8 i ), so that (21) follows straightforwardly if we prove that

$$
\begin{equation*}
\varepsilon^{\ell-\kappa} \omega_{\kappa+1-\ell}^{\varepsilon} \xrightarrow{\varepsilon \rightarrow 0}\left(\Pi_{E} \omega\right)_{\kappa+1-\ell} \tag{23}
\end{equation*}
$$

in the sense of distributions for $\ell=1, \ldots, \kappa, j=0, \ldots, \ell$. Indeed, denoting by $\delta_{j}$ the formal adjoint of $d_{j}$ in $L^{2}\left(\mathbb{G}, \wedge^{*} \mathfrak{g}\right)$,

$$
\begin{align*}
& \int\left\langle\varepsilon^{\ell-\kappa} d_{j} \omega_{\kappa+1-\ell}^{\varepsilon}, \varphi\right\rangle d V=\int\left\langle\varepsilon^{\ell-\kappa} \omega_{\kappa+1-\ell}^{\varepsilon}, \delta_{j} \varphi\right\rangle d V  \tag{24}\\
& \stackrel{\varepsilon \rightarrow 0}{\longrightarrow}-\int\left\langle\left(\Pi_{E} \omega\right)_{\kappa+1-\ell}, \delta_{j} \varphi\right\rangle d V=-\int\left\langle d_{j}\left(\Pi_{E} \omega\right)_{\kappa+1-\ell}, \varphi\right\rangle d V .
\end{align*}
$$

In order to prove (23), we set $p:=\kappa+1-\ell, p=1, \ldots, \kappa$. Thus, formulae (23) and (21) become

$$
\begin{equation*}
\varepsilon^{1-p} \omega_{p}^{\varepsilon} \underset{15}{\stackrel{\varepsilon \rightarrow 0}{\longrightarrow}}\left(\Pi_{E} \omega\right)_{p} \tag{25}
\end{equation*}
$$

in the sense of distributions, and

$$
\begin{equation*}
\varepsilon^{1-p} d_{j} \omega_{p}^{\varepsilon} \xrightarrow{\varepsilon \rightarrow 0} d_{j}\left(\Pi_{E} \omega\right)_{p} \tag{26}
\end{equation*}
$$

in the sense of distributions, respectively.
To prove (25) and (26), we argue by iteration in $p$.
Step $1(\mathbf{p}=1)$ : the proof is trivial, since, by assumption,

$$
\begin{equation*}
\omega_{1}^{\varepsilon} \rightarrow \omega_{1}:=\left(\Pi_{E} \omega\right)_{1} \tag{27}
\end{equation*}
$$

weakly in $L^{2}\left(\mathbb{G}, \bigwedge^{2} \mathfrak{g}\right)$ as $\varepsilon \rightarrow 0$.
Step $2(\mathbf{p}=\mathbf{2})$ : by (14), we have

$$
\begin{equation*}
d_{0} \omega_{2}^{\varepsilon}+\varepsilon d_{1} \omega_{1}^{\varepsilon}=O\left(\varepsilon^{\kappa}\right) \quad \text { in } L^{2}\left(\mathbb{G}, \bigwedge^{2} \mathfrak{g}\right) \tag{28}
\end{equation*}
$$

Remember now that, by Lemma 3.7, ii), $d_{0}^{-1} d_{0} \omega_{2}^{\varepsilon}=\omega_{2}^{\varepsilon}$, since $\omega_{2}^{\varepsilon}$ has weight 2 and hence is orthogonal to $\Omega^{1,1}=\operatorname{ker} d_{0}$. Thus, keeping in mind that $d_{0}^{-1}$ is algebraic, it follows from (28) that

$$
\begin{equation*}
\frac{1}{\varepsilon} \omega_{2}^{\varepsilon}+d_{0}^{-1} d_{1} \omega_{1}^{\varepsilon}=O\left(\varepsilon^{\kappa-1}\right) \quad \text { in } L^{2}\left(\mathbb{G}, \bigwedge^{2} \mathfrak{g}\right) \tag{29}
\end{equation*}
$$

Moreover, since $d_{1}$ is an homogeneous differential operator in the horizontal derivatives of order $1 \leq \kappa$ and, again, $d_{0}^{-1}$ is algebraic, then

$$
\begin{equation*}
d_{0}^{-1} d_{1} \omega_{1}^{\varepsilon} \rightarrow d_{0}^{-1} d_{1} \omega_{1} \quad \text { weakly in } L^{2}\left(\mathbb{G}, \bigwedge^{2} \mathfrak{g}\right) . \tag{30}
\end{equation*}
$$

Combining (29) and (30), we obtain

$$
\begin{equation*}
\varepsilon^{-1} \omega_{2}^{\varepsilon} \rightarrow-d_{0}^{-1} d_{1} \omega_{1}=\left(\Pi_{E} \omega\right)_{2} \tag{31}
\end{equation*}
$$

weakly in $L^{2}\left(\mathbb{G}, \bigwedge^{2} \mathfrak{g}\right)$, as $\varepsilon \rightarrow 0$. This proves (25) for $p=2$.
Step 3: suppose (25) and (26) have been proved for $p=1, \ldots, q-1<\kappa$, with $q>1$. Let us prove the assertion for $p=q$. We argue as follows: using (14) with $p=q$, we get

$$
\begin{gather*}
\frac{1}{\varepsilon^{q-1}} d_{0} \omega_{q}^{\varepsilon}+\frac{1}{\varepsilon^{q-2}} d_{1} \omega_{q-1}^{\varepsilon}+\cdots+\frac{1}{\varepsilon} d_{q-2} \omega_{2}^{\varepsilon}+d_{q-1} \omega_{1}^{\varepsilon}  \tag{32}\\
=O\left(\varepsilon^{\kappa-q+1}\right)=o(1)
\end{gather*}
$$

as $\varepsilon \rightarrow 0$. Remember now that, by Lemma 3.7, ii), $d_{0}^{-1} d_{0} \omega_{q}^{\varepsilon}=\omega_{q}^{\varepsilon}$, since $\omega_{q}^{\varepsilon}$ has weight $q$ and hence is orthogonal to $\Omega^{1,1}=\operatorname{ker} d_{0}$. Thus, keeping in mind that $d_{0}^{-1}$ is algebraic, it follows from (32) that

$$
\begin{align*}
\frac{1}{\varepsilon^{q-1}} \omega_{q}^{\varepsilon}+ & \frac{1}{\varepsilon^{q-2}} d_{0}^{-1} d_{1} \omega_{q-1}^{\varepsilon}+\cdots+\frac{1}{\varepsilon} d_{0}^{-1} d_{q-2} \omega_{2}^{\varepsilon}+d_{0}^{-1} d_{q-1} \omega_{1}^{\varepsilon}  \tag{33}\\
& =o(1)
\end{align*}
$$

as $\varepsilon \rightarrow 0$. Hence, by the inductive hypothesis and keeping in mind (7),

$$
\begin{gather*}
\frac{1}{\varepsilon^{q-1}} \omega_{q}^{\varepsilon} \longrightarrow-d_{0}^{-1}\left(d_{1}\left(\Pi_{E} \omega\right)_{q-1}+\cdots+d_{q-2}\left(\Pi_{E} \omega\right)_{2}+d_{q-1}\left(\Pi_{E} \omega\right)_{1}\right)  \tag{34}\\
=\left(\Pi_{E} \omega\right)_{q}
\end{gather*}
$$

in the sense of distribution, as $\varepsilon \rightarrow 0$.
This achieves the proof of (23) and hence of (21).
So far, we have used the equiboundedness of the first sum in (13) for $\varepsilon$ close to zero. We proceed now to estimate the liminf of the second term in (13).

To this end, we take $j=\ell$ in (21)and we sum up for $\ell=1, \ldots, \kappa$. We obtain

$$
\begin{equation*}
\frac{1}{\varepsilon^{\kappa-1}}\left(d_{1} \omega_{\kappa}^{\varepsilon}+\cdots+\varepsilon^{\kappa-1} d_{\kappa} \omega_{1}^{\varepsilon}\right) \longrightarrow \sum_{\ell=1}^{\kappa} d_{\ell}\left(\Pi_{E} \omega\right)_{\kappa+1-\ell} \tag{35}
\end{equation*}
$$

as $\varepsilon \rightarrow 0$ in the sense of distributions. On the other hand, the limit $\sum_{\ell=1}^{\kappa} d_{\ell}\left(\Pi_{E} \omega\right)_{\kappa+1-\ell}$ belongs to $L^{2}\left(\mathbb{G}, \Lambda^{2} \mathfrak{g}\right)$ (since $d_{\ell}\left(\Pi_{E} \omega\right)_{\kappa+1-\ell}$ is an homogeneous differential operator in the horizontal derivatives of order $\kappa$, by Theorem 3.8, i) and Definition 3.4), and

$$
\begin{equation*}
\left\{\frac{1}{\varepsilon^{\kappa-1}}\left(d_{1} \omega_{\kappa}^{\varepsilon}+\cdots+\varepsilon^{\kappa-1} d_{\kappa} \omega_{1}^{\varepsilon}\right)\right\}_{\varepsilon>0} \quad \text { is equibounded in } L^{2}\left(\mathbb{G}, \bigwedge^{2} \mathfrak{g}\right) \tag{36}
\end{equation*}
$$

as $\varepsilon \rightarrow 0$, by (15). Combining (36) and (35) we obtain that the limit in (35) is in fact a weak limit in $L^{2}\left(\mathbb{G}, \bigwedge^{2} \mathfrak{g}\right)$. Thus, by (20), (13) and taking into account that $\Pi_{E_{0}}$ is an othogonal projection, we obtain eventually

$$
\begin{aligned}
F(\omega) & =\int_{\mathbb{G}}\left\|\Pi_{E_{0}}\left(\sum_{\ell=1}^{\kappa} d_{\ell}\left(\Pi_{E} \omega\right)_{\kappa+1-\ell}\right)\right\|^{2} d V \\
& \leq \int_{\mathbb{G}}\left\|\sum_{\ell=1}^{\kappa} d_{\ell}\left(\Pi_{E} \omega\right)_{\kappa+1-\ell}\right\|^{2} d V \\
& \leq \liminf _{\varepsilon \rightarrow 0} \varepsilon^{2(1-\kappa)} \int_{\mathbb{G}}\left\|d_{1} \omega_{\kappa}^{\varepsilon}+\cdots+\varepsilon^{\kappa-1} d_{\kappa} \omega_{1}^{\varepsilon}\right\|^{2} d V \leq \liminf _{\varepsilon \rightarrow 0} F_{\varepsilon}\left(\omega_{\varepsilon}\right) .
\end{aligned}
$$

This proves (11).
We prove now that, if $\omega \in W^{\kappa, 2}\left(\mathbb{G}, E_{0}^{1}\right)$, then there exists a sequence $\left(\omega_{\varepsilon}\right)_{\varepsilon>0}$ in $W^{\kappa, 2}\left(\mathbb{G}, \Lambda^{1} \mathfrak{g}\right)$ such that
i) $\omega_{\varepsilon} \rightarrow \omega$ weakly in $W^{\kappa, 2}\left(\mathbb{G}, \bigwedge^{1} \mathfrak{g}\right)$;
ii) $F_{\varepsilon}\left(\omega_{\varepsilon}\right) \rightarrow F(\omega)$ as $\varepsilon \rightarrow 0$.

By Lemma 4.2, without loss of generality we may assume $\omega \in \mathcal{D}\left(\mathbb{G}, E_{0}^{1}\right)$.
We choose

$$
\begin{equation*}
\omega_{\varepsilon}=\omega+\varepsilon\left(\Pi_{E} \omega\right)_{2}+\cdots+\varepsilon^{\kappa-1}\left(\Pi_{E} \omega\right)_{\kappa} \tag{37}
\end{equation*}
$$

I we write the identity $d^{2}=0$ gathering all terms of the same weight, we get

$$
0=\sum_{p=0}^{\kappa} \sum_{j=0}^{p} d_{p-j} d_{j} .
$$

and therefore

$$
\begin{equation*}
\sum_{j=0}^{p} d_{p-j} d_{j}=0 \quad \text { for } p=0, \ldots, \kappa \tag{38}
\end{equation*}
$$

since these terms are mutually orthogonal when applied to a form of pure weight. In particular,

$$
\begin{equation*}
d_{0}^{2}=0, \quad d_{0} d_{1}=-d_{1} d_{0}, \quad d_{0} d_{2}=-d_{2} d_{0}-d_{1}^{2}, \quad \cdots \tag{39}
\end{equation*}
$$

Thus,

$$
\begin{align*}
& F_{\varepsilon}\left(\omega_{\varepsilon}\right)= \frac{1}{\varepsilon^{2 \kappa}} \int_{\mathbb{G}}\left\|d_{\varepsilon}\left(\sum_{i=1}^{\kappa} \varepsilon^{i-1}\left(\Pi_{E} \omega\right)_{i}\right)\right\|^{2} d V \\
&=\frac{1}{\varepsilon^{2 \kappa}}\left(\int_{\mathbb{G}}\left\|\Pi_{E_{0}}\left(d_{\varepsilon}\left(\sum_{i=1}^{\kappa} \varepsilon^{i-1}\left(\Pi_{E} \omega\right)_{i}\right)\right)\right\|^{2} d V\right.  \tag{40}\\
&\left.+\int_{\mathbb{G}}\left\|\Pi_{E_{0}^{\perp}}\left(d_{\varepsilon}\left(\sum_{i=1}^{\kappa} \varepsilon^{i-1}\left(\Pi_{E} \omega\right)_{i}\right)\right)\right\|^{2} d V\right) .
\end{align*}
$$

Arguing as in (12), we can write

$$
\begin{aligned}
& d_{\varepsilon}\left(\sum_{i=1}^{\kappa} \varepsilon^{i-1}\left(\Pi_{E} \omega\right)_{i}\right) \\
&= \sum_{2 \leq p \leq \kappa} \varepsilon^{p-1} \sum_{i=0}^{p-1} d_{i}\left(\Pi_{E} \omega\right)_{p-i} \\
&+\varepsilon^{\kappa}\left(d_{1}\left(\Pi_{E} \omega\right)_{\kappa}+\cdots+d_{\kappa}\left(\Pi_{E} \omega\right)_{1}\right) \\
&+\sum_{\kappa+2 \leq p \leq 2 \kappa} \varepsilon^{p-1} \sum_{i=p-\kappa}^{\kappa} d_{i}\left(\Pi_{E} \omega\right)_{p-i} \\
&=I_{1}+I_{2}+I_{3} .
\end{aligned}
$$

Now, by Theorem 3.10,

$$
\Pi_{E_{0}} I_{1}=0 .
$$

Let us prove that

$$
\begin{equation*}
\Pi_{E_{0}}^{\perp} I_{1}=\Pi_{E_{0}}^{\perp} I_{2}=0 . \tag{41}
\end{equation*}
$$

To this end, we have but to prove that

$$
\begin{equation*}
\sum_{i=0}^{p-1} d_{i}\left(\Pi_{E} \omega\right)_{p-i} \in E_{0}^{2} \quad \text { for } 2 \leq p \leq \kappa+1 \tag{42}
\end{equation*}
$$

Recalling that $E_{0}^{2}=\operatorname{ker} d_{0} \cap \mathcal{R}\left(d_{0}\right)^{\perp}$, we prove by induction that
i) $d_{0}\left(d_{0}\left(\Pi_{E} \omega\right)_{p}+d_{1}\left(\Pi_{E} \omega\right)_{p-1}+\cdots+d_{p-1}\left(\Pi_{E} \omega\right)_{1}\right)=0$
ii) $d_{0}\left(\Pi_{E} \omega\right)_{p}+d_{1}\left(\Pi_{E} \omega\right)_{p-1}+\cdots+d_{p-1}\left(\Pi_{E} \omega\right)_{1} \in \mathcal{R}\left(d_{0}\right)^{\perp}$,
for $2 \leq p \leq \kappa+1$. Obviously, by our choice of $\omega \in E_{0}^{1}$, (42) holds also for $p=1$.
Step $1(\mathbf{p}=\mathbf{2})$ : by Lemma 3.12, $d_{0} d_{0}^{-1} d_{1}\left(\Pi_{E} \omega\right)_{1}=d_{1}\left(\Pi_{E} \omega\right)_{1}$, since $d_{1}\left(\Pi_{E} \omega\right)_{1} \in$ $\Omega^{2,2}$. Hence we can write

$$
d_{0}\left(\Pi_{E} \omega\right)_{2}+d_{1}\left(\Pi_{E} \omega\right)_{1}=d_{0}\left(\left(\Pi_{E} \omega\right)_{2}+d_{0}^{-1} d_{1}\left(\Pi_{E} \omega\right)_{1}\right)=0,
$$

by (7). Trivially, i) and ii) hold.
Step 2 (if $\mathbf{p}$ then $\mathbf{p}+\mathbf{1}$ ): first of all, we notice that $p<\kappa+1$ (since $p+1 \leq \kappa+1)$. Hence, by (42) and by Theorem 3.10, we get in particular

$$
\begin{equation*}
d_{0}\left(\Pi_{E} \omega\right)_{j}+d_{1}\left(\Pi_{E} \omega\right)_{j-1}+\cdots+d_{j-1}\left(\Pi_{E} \omega\right)_{1}=0 \quad \text { for } j \leq p<\kappa+1 . \tag{43}
\end{equation*}
$$

Now, using (39), we get

$$
\begin{aligned}
& d_{0}\left(d_{0}\left(\Pi_{E} \omega\right)_{p+1}+d_{1}\left(\Pi_{E} \omega\right)_{p}+\cdots+d_{p}\left(\Pi_{E} \omega\right)_{1}\right) \\
& \quad=-d_{1} d_{0}\left(\Pi_{E} \omega\right)_{p}-\left(d_{2} d_{0}\left(\Pi_{E} \omega\right)_{p-1}+d_{1}^{2}\left(\Pi_{E} \omega\right)_{p-1}\right) \\
& \quad-\cdots-\left(d_{p} d_{0}\left(\Pi_{E} \omega\right)_{1}+d_{p-1} d_{1}\left(\Pi_{E} \omega\right)_{1}+\cdots+d_{1} d_{p-1}\left(\Pi_{E} \omega\right)_{1}\right) \\
& \quad=-\sum_{i=1}^{p-1} d_{i} \sum_{j=0}^{p-i} d_{j}\left(\Pi_{E} \omega\right)_{p+1-i-j}=0
\end{aligned}
$$

by (43), and the assertion i) holds for $p+1$.
In order to prove ii), we remind that, by (7),

$$
\left(\Pi_{E} \omega\right)_{p}=-d_{0}^{-1}\left(d_{1}\left(\Pi_{E} \omega\right)_{p-1}+\cdots+d_{p-1}\left(\Pi_{E} \omega\right)_{1}\right),
$$

so that, by Lemma 3.7, iii),

$$
d_{0}\left(\Pi_{E} \omega\right)_{p}=-\left(d_{1}\left(\Pi_{E} \omega\right)_{p-1}+\cdots+d_{p-1}\left(\Pi_{E} \omega\right)_{1}\right)+\xi
$$

with $\xi \in \mathcal{R}\left(d_{0}\right)^{\perp}$. This proves ii) and eventually (42).
Coming back to (40) we get,

$$
\begin{aligned}
F_{\varepsilon}\left(\omega_{\varepsilon}\right) & =\frac{1}{\varepsilon^{2 \kappa}} \int_{\mathbb{G}}\left\|\Pi_{E_{0}} I_{2}\right\|^{2} d V+\frac{1}{\varepsilon^{2 \kappa}} \int_{\mathbb{G}}\left\|I_{3}\right\|^{2} d V \\
& =\int_{\mathbb{G}}\left\|\Pi_{E_{0}}\left(d_{1}\left(\Pi_{E} \omega\right)_{\kappa}+\cdots+d_{\kappa}\left(\Pi_{E} \omega\right)_{1}\right)\right\|^{2} d V \\
& +\frac{1}{\varepsilon^{2 \kappa}} \int_{\mathbb{G}}\left\|\sum_{\kappa+2 \leq p \leq 2 \kappa} \varepsilon^{p-1} \sum_{i=p-\kappa}^{\kappa} d_{i}\left(\Pi_{E} \omega\right)_{p-i}\right\|^{2} d V ;
\end{aligned}
$$

observing that the second term in previous expression goes to zero as $\varepsilon \rightarrow 0$, we get $\lim _{\varepsilon \rightarrow 0} F_{\varepsilon}\left(\omega_{\varepsilon}\right)=F(\omega)$ in $\mathcal{D}\left(\mathbb{G}, E_{0}^{1}\right)$. This achieves the proof of the theorem.

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