

Graphs of $W^{1,1}$ -maps with values into S^1 : relaxed energies, minimal connections and lifting

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The aim of this paper is to link the analytic results of [6] [7] [19] relative to $W^{1,1}$ -mappings from B^n into S^1 to the measure theoretical geometric results in [12] [15]. The paper also contains a few remarks about mappings in $W^{1,p}$, $p \geq 2$, with values into S^2 .

1 The relaxed energy of $W^{1,1}$ -maps

Let Ω be a simply connected, smooth, n -dimensional domain. For simplicity, we take $\Omega = B^n$, the n -dimensional unit ball, and we let \tilde{B}^n be the unit ball of radius 2, so that $B^n \subset\subset \tilde{B}^n$. Also, let $S^1 \subset \mathbb{R}^2 \simeq \mathbb{C}$ be the unit sphere. For any non-negative integer k and for $U = \tilde{B}^n$ or $\tilde{B}^n \times S^1$, we will denote by $\mathcal{D}^k(U)$ the class of smooth compactly supported k -forms in U and by $\mathcal{D}_k(U)$ the usual class of k -dimensional currents in U , i.e., the dual of $\mathcal{D}^k(U)$. Moreover, $\mathcal{R}_k(U)$ denotes the subclass of k -dimensional integer multiplicity (say i.m.) rectifiable currents in U , compare [8] [10] [13]. We set

$$W^{1,1}(\tilde{B}^n, S^1) := \{u \in W^{1,1}(\tilde{B}^n, \mathbb{R}^2) : |u(x)| = 1 \text{ for a.e. } x \in \tilde{B}^n\}$$

and, in the sequel, $\varphi : \tilde{B}^n \rightarrow S^1$ being a given smooth $W^{1,1}$ -map, we set

$$\begin{aligned} W_\varphi^{1,1}(\tilde{B}^n, S^1) &:= \{u \in W^{1,1}(\tilde{B}^n, S^1) : u = \varphi \text{ in } \tilde{B}^n \setminus \overline{B}^n\}. \\ C_\varphi^1(\tilde{B}^n, S^1) &:= \{u \in C^1(\tilde{B}^n, S^1) : u = \varphi \text{ in } \tilde{B}^n \setminus \overline{B}^n\}. \end{aligned}$$

Also, $\pi : B^n \times S^1 \rightarrow B^n$ and $\hat{\pi} : B^n \times S^1 \rightarrow S^1$ will denote the projections onto the first and second factor, respectively. Finally, we denote by ω_{S^1} the *volume 1-form on $S^1 \subset \mathbb{R}^2$*

$$\omega_{S^1} := y^1 dy^2 - y^2 dy^1.$$

GRAPHS OF $W^{1,1}$ -MAPS. We recall from [13] that the i.m. rectifiable current $G_u \in \mathcal{R}_n(\tilde{B}^n \times S^1)$ associated to the "graph" of a function $u \in W^{1,1}(\tilde{B}^n, S^1)$ is defined in an *approximate sense* by

$$G_u := (Id \bowtie u)_\# [\tilde{B}^n], \tag{1.1}$$

where $(Id \bowtie u)(x) := (x, u(x))$, u standing for the restriction of u to the set of *approximate differentiability* of u , i.e.

$$G_u(\omega) := \int_{\tilde{B}^n} (Id \bowtie u)_\# \omega, \quad \omega \in \mathcal{D}^n(\tilde{B}^n \times S^1).$$

THE SINGULAR SET. Following [13, Vol. II], for any $u \in W_\varphi^{1,1}(\tilde{B}^n, S^1)$ we define the $(n-2)$ -current $\mathbb{P}(u) \in \mathcal{D}_{n-2}(\tilde{B}^n)$ by $2\pi \cdot \mathbb{P}(u) := -\pi_\#(\partial G_u) \llcorner \hat{\pi}^\# \omega_{S^1}$, so that for every $\xi \in \mathcal{D}^{n-2}(\tilde{B}^n)$

$$\mathbb{P}(u)(\xi) = -\frac{1}{2\pi} \partial G_u(\hat{\pi}^\# \omega_{S^1} \wedge \pi^\# \xi) = \frac{1}{2\pi} \int_{\tilde{B}^n} u^\# \omega_{S^1} \wedge d\xi. \tag{1.2}$$

Since $u = \varphi$ outside \overline{B}^n , and φ is smooth, we infer that $\mathbb{P}(u)$ is a boundaryless current supported in the closure of B^n ,

$$\partial \mathbb{P}(u) = 0, \quad \text{spt } \mathbb{P}(u) \subset \overline{B}^n.$$

Define now the $(n-1)$ -current $\mathbb{D}(u) \in \mathcal{D}_{n-1}(\tilde{B}^n)$ by $2\pi \cdot \mathbb{D}(u) := \pi_{\#}(G_u \lrcorner \hat{\pi}^{\#} \omega_{S^1})$, so that for every $\gamma \in \mathcal{D}^{n-1}(\tilde{B}^n)$

$$\mathbb{D}(u)(\gamma) := \frac{1}{2\pi} G_u(\hat{\pi}^{\#} \omega_{S^1} \wedge \pi^{\#} \gamma) = \frac{1}{2\pi} \int_{\tilde{B}^n} u^{\#} \omega_{S^1} \wedge \gamma.$$

Since $u = \varphi$ on $\tilde{B}^n \setminus \bar{B}^n$, we have $\text{spt}(\mathbb{D}(u) - \mathbb{D}(\varphi)) \subset \bar{B}^n$. Moreover,

$$\mathbb{P}(u) = \partial(\mathbb{D}(u) - \mathbb{D}(\varphi)). \quad (1.3)$$

In fact, since φ is smooth, $d\varphi^{\#} \omega_{S^1} = \varphi^{\#} d\omega_{S^1} = 0$ and hence $d(\varphi^{\#} \omega_{S^1} \wedge \xi) = -\varphi^{\#} \omega_{S^1} \wedge d\xi$, which yields

$$\begin{aligned} 2\pi \mathbb{P}(u)(\xi) &= \int_{\bar{B}^n} u^{\#} \omega_{S^1} \wedge d\xi + \int_{\tilde{B}^n \setminus \bar{B}^n} \varphi^{\#} \omega_{S^1} \wedge d\xi \\ &= \int_{B^n} u^{\#} \omega_{S^1} \wedge d\xi - \int_{\tilde{B}^n \setminus \bar{B}^n} d(\varphi^{\#} \omega_{S^1} \wedge \xi) \\ &= \int_{B^n} u^{\#} \omega_{S^1} \wedge d\xi + \int_{\partial B^n} \varphi^{\#} \omega_{S^1} \wedge \xi \end{aligned} \quad (1.4)$$

for every $\xi \in \mathcal{D}^{n-2}(\tilde{B}^n)$. On the other hand, we compute

$$\begin{aligned} 2\pi \partial(\mathbb{D}(u) - \mathbb{D}(\varphi))(\xi) &:= 2\pi (\mathbb{D}(u) - \mathbb{D}(\varphi))(d\xi) \\ &= \int_{\bar{B}^n} u^{\#} \omega_{S^1} \wedge d\xi - \int_{\bar{B}^n} \varphi^{\#} \omega_{S^1} \wedge d\xi \\ &= \int_{B^n} u^{\#} \omega_{S^1} \wedge d\xi + \int_{B^n} d(\varphi^{\#} \omega_{S^1} \wedge \xi) \\ &= \int_{B^n} u^{\#} \omega_{S^1} \wedge d\xi + \int_{\partial B^n} \varphi^{\#} \omega_{S^1} \wedge \xi. \end{aligned}$$

REAL AND INTEGRAL MASSES. We now recall the following

Definition 1.1 Let $0 \leq k \leq n-2$. For every k -dimensional current $\Gamma \in \mathcal{D}_k(\tilde{B}^n)$, with support $\text{spt} \Gamma \subset \bar{B}^n$, we denote by

$$m_i(\Gamma) := \inf\{\mathbf{M}(L) \mid L \in \mathcal{R}_{k+1}(\tilde{B}^n), \text{ spt } L \subset \bar{B}^n, \partial L = \Gamma\}$$

the integral mass of Γ and by

$$m_r(\Gamma) := \inf\{\mathbf{M}(D) \mid D \in \mathcal{D}_{k+1}(\tilde{B}^n), \text{ spt } D \subset \bar{B}^n, \partial D = \Gamma\}$$

the real mass of Γ . Moreover, in case $m_i(\Gamma) < \infty$, we say that an i.m. rectifiable current $L \in \mathcal{R}_{k+1}(\tilde{B}^n)$ is an integral minimal connection of Γ if $\partial L = \Gamma$, $\text{spt } L \subset \bar{B}^n$, and $\mathbf{M}(L) = m_i(\Gamma)$.

Of course, $m_r(\Gamma) < \infty$ if $\mathbf{M}(\Gamma) < \infty$ and $\partial\Gamma = 0$. Moreover, in the definition of integral and real mass, respectively, the infimum is actually a minimum, provided that the set on the right-hand side is non-empty. Now, if $k = n-2$ and $\Gamma = \mathbb{P}(u)$ for some $u \in W_{\varphi}^{1,1}(\tilde{B}^n, S^1)$, by (1.3) we infer that the real mass is finite,

$$m_r(\mathbb{P}(u)) \leq \mathbf{M}(\mathbb{D}(u) - \mathbb{D}(\varphi)) < \infty. \quad (1.5)$$

We recall that in general

$$m_r(\Gamma) \leq m_i(\Gamma), \quad (1.6)$$

and the strict inequality may occur if $1 \leq k \leq n-3$, compare [20] [27]. However, as shown by Federer [9], and by Hardt-Pitts [17], equality holds in (1.6) if Γ has dimension zero or if $k = n-2$. In particular, we obtain

$$m_r(\mathbb{P}(u)) = m_i(\mathbb{P}(u)) \quad \forall u \in W_{\varphi}^{1,1}(\tilde{B}^n, S^1) \quad (1.7)$$

and hence that *the integral minimal connection of the singularities of any map $u \in W_{\varphi}^{1,1}(\tilde{B}^n, S^1)$ is finite* provided we are able to show that $\mathbb{P}(u)$ is an *integral flat chain*, i.e., $\mathbb{P}(u)$ is the boundary of an i.m. rectifiable current with finite mass.

THE RELAXED $W^{1,1}$ -ENERGY. For $\Omega = \widetilde{B}^n$ or B^n , denote

$$\mathcal{E}_{1,1}(u, \Omega) := \int_{\Omega} |\nabla u| dx, \quad u \in W_{\varphi}^{1,1}(\widetilde{B}^n, S^1),$$

and consider the *lower semicontinuous envelope* of the functional

$$\overline{\mathcal{E}}_{1,1}(u, \Omega) := \begin{cases} \int_{\Omega} |\nabla u| dx & \text{if } u \in C_{\varphi}^1(\widetilde{B}^n, S^1) \\ +\infty & \text{elsewhere in } W_{\varphi}^{1,1}(\widetilde{B}^n, S^1). \end{cases} \quad (1.8)$$

More precisely, we define the *relaxed $W^{1,1}$ -energy* $u \mapsto \widetilde{\mathcal{E}}_{1,1}(u, \Omega)$ as the greatest functional on $W_{\varphi}^{1,1}(\widetilde{B}^n, S^1)$ which is lower than or equal to $u \mapsto \overline{\mathcal{E}}_{1,1}(u, \Omega)$ and is lower semicontinuous with respect to the strong L^1 -convergence. Of course, for every $u \in W_{\varphi}^{1,1}(\widetilde{B}^n, S^1)$ we have

$$\widetilde{\mathcal{E}}_{1,1}(u, \Omega) = \inf \left\{ \liminf_{h \rightarrow \infty} \int_{\Omega} |\nabla u_h| dx : \{u_h\} \subset C_{\varphi}^1(\widetilde{B}^n, S^1), \quad u_h \rightarrow u \quad \text{a.e.} \right\}.$$

It is well-known that

$$\mathcal{E}_{1,1}(u, \Omega) \leq \widetilde{\mathcal{E}}_{1,1}(u, \Omega) \quad \forall u \in W_{\varphi}^{1,1}(\widetilde{B}^n, S^1). \quad (1.9)$$

However, since $\pi_1(S^1) \neq 0$, in general the strict inequality " $<$ " may hold in (1.9), see e.g. [13].

CARTESIAN CURRENTS. In order to analyze the structure property of the relaxed $W^{1,1}$ -energy, we recall from [13] some facts concerning the class of *Cartesian currents* $\text{cart}(\widetilde{B}^n \times S^1)$.

Definition 1.2 *We let*

$$\text{cart}(\widetilde{B}^n \times S^1) := \{T \in \text{cart}(\widetilde{B}^n \times \mathbb{R}^2) \mid \text{spt } T \subset \widetilde{B}^n \times S^1\}$$

and

$$\text{cart}_{\varphi}(\widetilde{B}^n \times S^1) := \{T \in \text{cart}(\widetilde{B}^n \times S^1) \mid T \llcorner (\widetilde{B}^n \setminus \overline{B}^n) \times S^1 = G_{\varphi} \llcorner (\widetilde{B}^n \setminus \overline{B}^n) \times S^1\},$$

where $\text{cart}(\widetilde{B}^n \times \mathbb{R}^N)$ is defined as the class of i.m. rectifiable currents T in $\mathcal{R}_n(\widetilde{B}^n \times \mathbb{R}^N)$ which have no inner boundary, $\partial T \llcorner \widetilde{B}^n \times \mathbb{R}^N = 0$, have finite mass, $\mathbf{M}(T) < \infty$, and are such that $\|T\|_1 < \infty$, $\pi_{\#}T = \llbracket \widetilde{B}^n \rrbracket$ and $T^{\overline{0}} \geq 0$, where

$$\|T\|_1 := \sup\{T(\varphi(x, y)|y| dx) \mid \varphi \in C_c^0(\widetilde{B}^n \times \mathbb{R}^N), \|\varphi\| \leq 1\},$$

and $T^{\overline{0}}$ is the Radon measure in $\widetilde{B}^n \times \mathbb{R}^N$ given by

$$T^{\overline{0}}(\varphi(x, y)) := T(\varphi(x, y) dx) \quad \forall \varphi \in C_c^0(\widetilde{B}^n \times \mathbb{R}^N).$$

Then, to any $T \in \text{cart}_{\varphi}(\widetilde{B}^n \times S^1)$ we may associate a function $u_T \in BV_{\varphi}(\widetilde{B}^n, S^1)$, where

$$BV_{\varphi}(\widetilde{B}^n, S^1) := \{u \in BV(\widetilde{B}^n, \mathbb{R}^2) : |u(x)| = 1 \text{ a.e. and } u = \varphi \text{ in } \widetilde{B}^n \setminus \overline{B}^n\},$$

such that

$$T(\phi(x, y) dx) = \int_{\widetilde{B}^n} \phi(x, u_T(x)) dx \quad (1.10)$$

for all $\phi \in C^0(\widetilde{B}^n \times \mathbb{R}^2)$ such that $|\phi(x, y)| \leq C(1 + |y|)$, and for any $\psi \in C_c^1(\widetilde{B}^n)$ and for $j = 1, 2$

$$(-1)^{n-i} T(\psi(x) \widehat{dx}^i \wedge dy^j) = \langle D_i u_T^j, \psi \rangle := - \int_{\widetilde{B}^n} u_T^j(x) \cdot D_i \psi(x) dx, \quad (1.11)$$

where

$$\widehat{dx}^i := dx^1 \wedge \cdots \wedge dx^{i-1} \wedge dx^{i+1} \wedge \cdots \wedge dx^n.$$

Moreover, if the corresponding BV -function u_T belongs to the Sobolev class $W_\varphi^{1,1}(\tilde{B}^n, S^1)$, we have

$$T = G_{u_T} + L \times \llbracket S^1 \rrbracket, \quad (1.12)$$

where $\llbracket S^1 \rrbracket$ is the 1-current integration of 1-forms on S^1 , with respect to the counterclockwise orientation, and L is an $(n-1)$ -dimensional i.m. rectifiable current in $\mathcal{R}_{n-1}(\tilde{B}^n)$ with *finite mass*, $\mathbf{M}(L) < \infty$, and support in \tilde{B}^n , $\text{spt } L \subset \tilde{B}^n$. Now, since T satisfies the null-boundary condition

$$\partial T \llcorner \tilde{B}^n \times S^1 = 0,$$

i.e., $T(d\omega) = 0$ for all $\omega \in \mathcal{D}^{n-1}(\tilde{B}^n \times S^1)$, if $\omega = \pi^\# \xi \wedge \hat{\pi}^\# \omega_{S^1}$ for some $\xi \in \mathcal{D}^{n-2}(\tilde{B}^n)$, since $d\omega = \pi^\# d\xi \wedge \hat{\pi}^\# \omega_{S^1} = (-1)^{n-1} \hat{\pi}^\# \omega_{S^1} \wedge \pi^\# d\xi$, whereas

$$(L \times \llbracket S^1 \rrbracket)(\pi^\# d\xi \wedge \hat{\pi}^\# \omega_{S^1}) = L(d\xi) \cdot \llbracket S^1 \rrbracket(\omega_{S^1}) = 2\pi \partial L(\xi),$$

we infer that

$$\partial L(\xi) = (-1)^n \frac{1}{2\pi} G_{u_T}(\hat{\pi}^\# \omega_{S^1} \wedge \pi^\# d\xi) = (-1)^n \mathbb{P}(u_T)(\xi).$$

In conclusion, L satisfies the boundary condition

$$\partial L = (-1)^n \mathbb{P}(u_T). \quad (1.13)$$

On the other hand, if L is an i.m. rectifiable current in $\mathcal{R}_{n-2}(\tilde{B}^n)$, hence with finite mass, such that $\text{spt } L \subset \tilde{B}^n$, satisfying (1.13), the corresponding current (1.12) belongs to the class $\text{cart}_\varphi(\tilde{B}^n \times S^1)$.

Denote now for any $u \in W_\varphi^{1,1}(\tilde{B}^n, S^1)$ by

$$\mathcal{T}_u := \{T \in \text{cart}_\varphi(\tilde{B}^n \times S^1) \mid u_T = u\} \quad (1.14)$$

the class of Cartesian currents T in $\text{cart}_\varphi(\tilde{B}^n \times S^1)$ such that the underlying BV -function u_T is equal to u . By the previous discussion we have

$$\mathcal{T}_u = \{G_u + L \times \llbracket S^1 \rrbracket : L \in \mathcal{R}_{n-1}(\tilde{B}^n), \text{spt } L \subset \tilde{B}^n, \partial L = (-1)^n \mathbb{P}(u)\}. \quad (1.15)$$

THE DENSITY RESULT OF BETHUEL. If $n \geq 2$, we denote by $R_{1,\varphi}^\infty(\tilde{B}^n, S^1)$ the set of all the maps $u \in W_\varphi^{1,1}(\tilde{B}^n, S^1)$ which are smooth except on a singular set $\Sigma(u)$ of the type

$$\Sigma(u) = \bigcup_{i=1}^r \Sigma_i, \quad r \in \mathbb{N}, \quad (1.16)$$

where Σ_i is a smooth $(n-2)$ -dimensional subset of \tilde{B}^n with smooth boundary, if $n \geq 3$, and Σ_i is a point if $n = 2$. The following density result appears in [4].

Theorem 1.3 $R_{1,\varphi}^\infty(\tilde{B}^n, S^1)$ is strongly dense in $W_\varphi^{1,1}(\tilde{B}^n, S^1)$.

Theorem 1.3 yields, in particular, see [11], that $\mathbb{P}(u)$ is an integral flat chain. Thus by (1.7) we readily obtain the following

Corollary 1.4 For any $u \in W_\varphi^{1,1}(\tilde{B}^n, S^1)$ the class \mathcal{T}_u is non-empty.

Moreover, the following density result holds true, see e.g. [15].

Proposition 1.5 Let $n \geq 2$ and let $T \in \mathcal{T}_u$ satisfy (1.12) for some $u \in W_\varphi^{1,1}(\tilde{B}^n, S^1)$. There exists a smooth sequence $\{u_h\} \subset C_\varphi^1(\tilde{B}^n, S^1)$ such that $G_{u_h} \rightharpoonup T$ as $h \rightarrow \infty$ weakly in $\mathcal{D}_n(\tilde{B}^n \times S^1)$, i.e., $G_{u_h}(\omega) \rightarrow T(\omega)$ for all $\omega \in \mathcal{D}^n(\tilde{B}^n \times S^1)$, and

$$\lim_{h \rightarrow \infty} \int_{\tilde{B}^n} |\nabla u_h| dx = \mathcal{E}_{1,1}(T) := \int_{\tilde{B}^n} |\nabla u| dx + 2\pi \mathbf{M}(L).$$

We now recall that the weak convergence $G_{u_h} \rightharpoonup T$ yields the convergence of u_h to u_T weakly in the BV -sense, that the energy $T \mapsto \mathcal{E}_{1,1}(T)$ is lower semicontinuous with respect to the weak convergence as currents, and that the class $\text{cart}_\varphi(\widetilde{B}^n \times S^1)$ is closed. As a consequence, compare [15], from the above density result we readily obtain the following

Proposition 1.6 *For any $n \geq 2$ and $u \in W_\varphi^{1,1}(\widetilde{B}^n, S^1)$ we have $\widetilde{\mathcal{E}}_{1,1}(u, \widetilde{B}^n) < \infty$. Moreover, the relaxed energy of u is given by*

$$\begin{aligned} \widetilde{\mathcal{E}}_{1,1}(u, \widetilde{B}^n) &= \inf\{\mathcal{E}_{1,1}(T) \mid T \in \mathcal{T}_u\} \\ &= \int_{\widetilde{B}^n} |\nabla u| dx + 2\pi m_i(\mathbb{P}(u)) \\ &= \int_{\widetilde{B}^n} |\nabla u| dx + 2\pi \inf\{\mathbf{M}(L) : L \in \mathcal{R}_{n-1}(\widetilde{B}^n), \text{spt } L \subset \overline{B}^n, \partial L = (-1)^n \mathbb{P}(u)\}. \end{aligned}$$

2 Flat norm and minimal connections

In this section we write more explicitly the action of the current $\mathbb{P}(u)$ associated to the singular set of a Sobolev map $u \in W_\varphi^{1,1}(\widetilde{B}^n, S^1)$, and recover in the case $n = 2$ some results from [6].

THE SINGULAR SET AS A DISTRIBUTION. We will denote by $\text{Lip}(B^n, \Lambda^k TB^n)$ the class of k -forms in B^n with coefficients in $\text{Lip}(B^n)$, for every $k = 0, \dots, n$. Every $(n-2)$ -form $\zeta \in \text{Lip}(B^n, \Lambda^{n-2} TB^n)$ will be written as

$$\zeta = \sum_{1 \leq i < j \leq n} \zeta^{i,j} \widehat{dx^{i,j}}, \quad (2.1)$$

where $\zeta^{i,j} \in \text{Lip}(B^n; \mathbb{R})$ and

$$\widehat{dx^{i,j}} := dx^1 \wedge \dots \wedge dx^{i-1} \wedge dx^{i+1} \wedge \dots \wedge dx^{j-1} \wedge dx^{j+1} \wedge \dots \wedge dx^n.$$

Let $g \in W^{1,1}(B^n, \mathbb{R}^2) \cap L^\infty$. For any $i < j$ introduce the distribution $T_{i,j}(g) \in \mathcal{D}'(B^n, \mathbb{R})$ given by

$$T_{i,j}(g) := -(g \times g_{x_i})_{x_j} + (g \times g_{x_j})_{x_i} \quad (2.2)$$

where for every i

$$g \times g_{x_i} := g^1 g_{x_i}^2 - g^2 g_{x_i}^1,$$

that is,

$$\langle T_{i,j}(g), \zeta^{i,j} \rangle = \int_{B^n} ((g \times g_{x_i})_{x_j} \zeta_{x_j}^{i,j} - (g \times g_{x_j})_{x_i} \zeta_{x_i}^{i,j}) dx \quad \forall \zeta^{i,j} \in \text{Lip}(B^n, \mathbb{R}).$$

Definition 2.1 *Let $n \geq 2$. To any $g \in W^{1,1}(B^n, \mathbb{R}^2) \cap L^\infty$ we associate the distribution $T(g)$ of order $(n-2)$ defined by*

$$\langle T(g), \zeta \rangle := \sum_{1 \leq i < j \leq n} (-1)^{i+j-1} \langle T_{i,j}(g), \zeta^{i,j} \rangle \quad \forall \zeta \in \text{Lip}(B^n, \Lambda^{n-2} TB^n),$$

where ζ is decomposed as in (2.1).

Viewing the target space S^1 as a subspace of \mathbb{C} , from Definition 2.1 we readily obtain

Proposition 2.2 *Let $n \geq 2$. We have:*

- (a) $T(g) = -T(\bar{g})$ for all $g \in W^{1,1}(B^n, \mathbb{R}^2) \cap L^\infty$;
- (b) $T(gh) = T(g) + T(h)$ for all $g, h \in W^{1,1}(B^n, S^1)$.

PROOF: Property (a) is trivial, (b) follows arguing as in [6]. \square

For any $n \geq 2$ and $g \in W^{1,1}(B^n, \mathbb{R}^2) \cap L^\infty$, defining

$$g \times \nabla g := (g \times g_{x_1}, \dots, g \times g_{x_n}),$$

we can write

$$\langle T(g), \zeta \rangle = \int_{B^n} \left(\sum_{i=1}^n (g \times g_{x_i}) F^i(\zeta) \right) d\mathcal{H}^n \quad \forall \zeta \in \text{Lip}(B^n, \Lambda^{n-2}TB^n),$$

where for every fixed i

$$F^i(\zeta) := \sum_{1 \leq h < i} (-1)^{i+h} \zeta_{x_h}^{h,i} - \sum_{i < h \leq n} (-1)^{i+h} \zeta_{x_h}^{i,h},$$

ζ being decomposed as in (2.1). Therefore, setting $F(\zeta) := (F^1(\zeta), \dots, F^n(\zeta))$, we have

$$\langle T(g), \zeta \rangle = \int_{B^n} (g \times \nabla g) \cdot F(\zeta) dx \quad \forall \zeta \in \text{Lip}(B^n, \Lambda^{n-2}TB^n).$$

Notice that, if $n = 2$,

$$F(\zeta) = \nabla^\perp \zeta := (\zeta_{x_2}, -\zeta_{x_1}),$$

so that, as in [6], we have

$$\langle T(g), \zeta \rangle := \int_{B^2} (g \times \nabla g) \cdot \nabla^\perp \zeta dx \quad \forall \zeta \in \text{Lip}(B^2, \mathbb{R}).$$

THE LINK BETWEEN $T(u)$ AND $\mathbb{P}(u)$. Suppose now that $u \in W^{1,1}(B^n, S^1)$. By the dominated convergence theorem, the action of the current G_u , see (1.1), extends e.g. to forms $\omega := \widehat{\pi}^\# \omega_{S^1} \wedge \pi^\# d\zeta$, where $\zeta \in \text{Lip}(B^n, \Lambda^{n-2}TB^n)$, so that

$$G_u(\widehat{\pi}^\# \omega_{S^1} \wedge \pi^\# d\zeta) = \int_{B^n} u^\# \omega_{S^1} \wedge d\zeta.$$

If ζ is given by (2.1), since for every $i < j$

$$dx^i \wedge \widehat{dx^{i,j}} = (-1)^{i-1} \widehat{dx^j} \quad \text{and} \quad dx^j \wedge \widehat{dx^{i,j}} = (-1)^j \widehat{dx^i},$$

we have

$$d\zeta = \sum_{1 \leq i < j \leq n} ((-1)^{i-1} \zeta_{x_i}^{i,j} \widehat{dx^j} + (-1)^j \zeta_{x_j}^{i,j} \widehat{dx^i}).$$

Consequently, since $dx^h \wedge \widehat{dx^h} = (-1)^{h-1} dx^1 \wedge \dots \wedge dx^n$, we have

$$\begin{aligned} u^\# \omega_{S^1} \wedge d\zeta &= (u^1 du^2 - u^2 du^1) \wedge d\zeta \\ &= \sum_{1 \leq i < j \leq n} (-1)^{i+j-1} ((u \times u_{x_i}) \zeta_{x_j}^{i,j} - (u \times u_{x_j}) \zeta_{x_i}^{i,j}) dx^1 \wedge \dots \wedge dx^n. \end{aligned}$$

Therefore, for every $\zeta \in \text{Lip}(B^n, \Lambda^{n-2}TB^n)$ satisfying (2.1) we have

$$G_u(\widehat{\pi}^\# \omega_{S^1} \wedge \pi^\# d\zeta) = \sum_{1 \leq i < j \leq n} (-1)^{i+j-1} \langle T_{i,j}(u), \zeta^{i,j} \rangle.$$

On account of Definition 2.1, we conclude that

$$\frac{1}{2\pi} \langle T(u), \zeta \rangle = \mathbb{P}(u)(\zeta) \quad \forall \zeta \in \text{Lip}(B^n, \Lambda^{n-2}TB^n). \quad (2.3)$$

THE FLAT NORM. Let $\Gamma \in \mathcal{D}_k(\widetilde{B}^n)$ with $\text{spt} \Gamma \subset \overline{B}^n$, and suppose that Γ is the boundary of a $(k+1)$ -dimensional current $D \in \mathcal{D}_{k+1}(\widetilde{B}^n)$, with $\text{spt} D \subset \overline{B}^n$. The *flat norm* of Γ is defined by

$$F_{\overline{B}^n}(\Gamma) := \sup\{\Gamma(\xi) \mid \xi \in \mathcal{D}^k(\widetilde{B}^n), \max\{\|\xi\|, \|d\xi\|\} \leq 1 \text{ in } \overline{B}^n\}.$$

Taking $k = n - 2$, we now define for any $n \geq 2$

$$L(u) := F_{\overline{B}^n}(\mathbb{P}(u)), \quad u \in W_{\varphi}^{1,1}(\widetilde{B}^n, S^1), \quad (2.4)$$

so that for every $u \in W_{\varphi}^{1,1}(\widetilde{B}^n, S^1)$ we have

$$\begin{aligned} L(u) &:= \sup\{\mathbb{P}(u)(\xi) \mid \xi \in \mathcal{D}^{n-2}(\widetilde{B}^n), \max\{\|\xi\|, \|d\xi\|\} \leq 1 \text{ in } \overline{B}^n\} \\ &= \frac{1}{2\pi} \sup\left\{\int_{\widetilde{B}^n} u^{\#}\omega_{S^1} \wedge d\xi \mid \xi \in \mathcal{D}^{n-2}(\widetilde{B}^n), \max\{\|\xi\|, \|d\xi\|\} \leq 1 \text{ in } \overline{B}^n\right\}. \end{aligned}$$

We now observe that by (1.4) and (2.3) we have

$$\mathbb{P}(u)(\xi) = \langle T(u|_{B^n}), \xi \rangle - \int_{\partial B^n} \varphi^{\#}\omega_{S^1} \wedge \xi \quad (2.5)$$

where $u|_{B^n}$ is the restriction of u to B^n . In order to write explicitly the boundary term, we notice that

$$\begin{aligned} j > i &\implies dx^j \wedge \widehat{dx^{i,j}} = (-1)^j \widehat{dx^i} \\ h < i &\implies dx^h \wedge \widehat{dx^{h,i}} = (-1)^{h+1} \widehat{dx^i}. \end{aligned}$$

Therefore, if

$$\xi = \sum_{1 \leq i < j \leq n} \xi^{i,j} \widehat{dx^{i,j}}, \quad \xi^{i,j} \in C_c^{\infty}(\widetilde{B}^n),$$

we have

$$\varphi^{\#}\omega_{S^1} \wedge \xi = (\varphi^1 d\varphi^2 - \varphi^2 d\varphi^1) \wedge \xi = \sum_{i=1}^n A_i(x) \widehat{dx^i},$$

where for every i

$$A_i(x) = \sum_{j>i} (-1)^j \xi^{i,j} (\varphi \times \varphi_{x_j}) + \sum_{h<i} (-1)^{h+1} \xi^{h,i} (\varphi \times \varphi_{x_h}).$$

We then obtain

$$\int_{\partial B^n} \varphi^{\#}\omega_{S^1} \wedge \xi = \int_{\partial B^n} \sum_{i=1}^n A_i(x) \widehat{dx^i} = \int_{\partial B^n} \langle F, \nu \rangle d\mathcal{H}^{n-1},$$

where ν is the outward unit normal to ∂B^n and $F(x) = (F^1, \dots, F^n)(x)$ is the vector field of components

$$F^i(x) := (-1)^{i-1} A_i(x) = \sum_{j>i} (-1)^{i+j-1} \xi^{i,j} (\varphi \times \varphi_{x_j}) + \sum_{h<i} (-1)^{i+h} \xi^{h,i} (\varphi \times \varphi_{x_h}). \quad (2.6)$$

Remark 2.3 Notice that the above boundary term depends only on the tangential components of the derivatives of φ , i.e., it can be expressed in terms of $(\varphi \times \varphi_{\tau_i}) := (\varphi^1 \varphi_{\tau_i}^2 - \varphi^2 \varphi_{\tau_i}^1)$, where $\tau_1, \dots, \tau_{n-1}$ is any orthonormal frame tangent to ∂B^n such that $(\tau_1, \dots, \tau_{n-1}, \nu)$ is a positively oriented orthonormal frame to \mathbb{R}^n . In particular, it is zero if e.g. the boundary datum φ is constant on ∂B^n . Moreover, in the simpler case $n = 2$ we have

$$F^1(x) = \xi (\varphi \times \varphi_{x_2}), \quad F^2(x) = -\xi (\varphi \times \varphi_{x_1}),$$

so that for every $\xi \in C_c^{\infty}(\widetilde{B}^2)$

$$\int_{\partial B^2} \varphi^{\#}\omega_{S^1} \wedge \xi = \int_{\partial B^2} \xi D(\varphi) \cdot \nu d\mathcal{H}^1,$$

where $D(\varphi) := (\varphi \times \varphi_{x_2}, -\varphi \times \varphi_{x_1})$. Equivalently, if (τ, ν) is positively oriented, we may write

$$\int_{\partial B^2} \varphi^{\#}\omega_{S^1} \wedge \xi = - \int_{\partial B^2} \xi (\varphi \times \varphi_{\tau}) d\mathcal{H}^1.$$

In conclusion, we infer that for every $u \in W_\varphi^{1,1}(\tilde{B}^n, S^1)$

$$L(u) = \frac{1}{2\pi} \max \left\{ \langle T(u|_{B^n}), \zeta \rangle - \int_{\partial B^n} \langle F, \nu \rangle d\mathcal{H}^{n-1} \mid \zeta \in \text{Lip}(\tilde{B}^n, \Lambda^{n-2}T\tilde{B}^n), \right. \\ \left. \max\{\|\zeta\|_\infty, \|\nabla\zeta\|_\infty\} \leq 1 \text{ in } \overline{B}^n \right\}$$

where $F = F(\zeta, \varphi)$ is the vector field given by (2.6), with $\xi = \zeta$, and

$$\|\zeta\|_\infty := \sup_{i < j} \|\zeta^{i,j}\|_\infty, \quad \|\nabla\zeta\|_\infty := \sup_{i < j} \|\nabla\zeta^{i,j}\|_\infty$$

if ζ is decomposed as in (2.1). Moreover, in the case $n = 2$ the above formula simplifies to

$$L(u) = \frac{1}{2\pi} \max \left\{ \langle T(u|_{B^2}), \zeta \rangle + \int_{\partial B^2} \zeta(\varphi \times \varphi_\tau) d\mathcal{H}^1 \mid \zeta \in \text{Lip}(\tilde{B}^2), \max\{\|\zeta\|_\infty, \|\nabla\zeta\|_\infty\} \leq 1 \text{ in } \overline{B}^2 \right\}.$$

If $g \in W^{1,1}(S^n, S^1)$, with a similar computation we obtain

$$L(g) = \frac{1}{2\pi} \max \left\{ \langle T(g), \zeta \rangle \mid \zeta \in \text{Lip}(S^n, \Lambda^{n-2}TS^n), \|\nabla\zeta\|_\infty \leq 1 \right\}.$$

This last formula has been used in [6], in dimension $n = 2$, as definition of minimal connection of the singularities of g . Of course, the formula for $L(u)$ makes sense for any function $u \in W_\varphi^{1,1}(\tilde{B}^n, \mathbb{R}^2) \cap L^\infty$. Moreover, by the definition of $T(u)$ we readily obtain the following

Proposition 2.4 *Let $n \geq 2$. We have:*

$$(c) \quad L(u) \leq \frac{1}{2\pi} (\|u\|_{W^{1,1}(B^n)} \|u\|_{L^\infty(B^n)} + \|\nabla\varphi\|_{L^1(\partial B^n)}) \text{ for all } u \in W_\varphi^{1,1}(\tilde{B}^n, \mathbb{R}^2) \cap L^\infty;$$

$$(d) \quad \text{if } u_h, u \in W_\varphi^{1,1}(\tilde{B}^n, \mathbb{R}^2) \cap L^\infty \text{ are such that } u_h \rightarrow u \text{ in } W^{1,1} \text{ and } \|u_h\|_\infty \leq C, \text{ then } L(u_h) \rightarrow L(u).$$

PROOF: Property (c) is trivial whereas, since

$$|\langle T(u_h|_{B^n}), \zeta \rangle - \langle T(u|_{B^n}), \zeta \rangle| \leq C \left(\int_{B^n} |u_h| |\nabla(u_h - u)| |\nabla\zeta| dx + \int_{B^n} |u_h - u| |\nabla u| |\nabla\zeta| dx \right),$$

where $C > 0$ is an absolute constant, and the boundary term in (2.5) does not depend on u_h or u , we have

$$|L(u_h) - L(u)| \leq C (\|u_h - u\|_{W^{1,1}(B^n)} + \|(u_h - u) \nabla u\|_{L^1(B^n)})$$

and (d) follows from the dominated convergence theorem. \square

FLAT NORM AND REAL MASS. The following property goes back to Federer [9].

Proposition 2.5 *Let $\Gamma \in \mathcal{D}_k(\tilde{B}^n)$ and $D \in \mathcal{D}_{k+1}(\tilde{B}^n)$ be such that $\text{spt } \Gamma \subset \overline{B}^n$, $\text{spt } D \subset \overline{B}^n$, and $\partial D = \Gamma$, with $\mathbf{M}(D) < \infty$. We have*

$$F_{\overline{B}^n}(\Gamma) = m_r(\Gamma).$$

PROOF: Trivially

$$F_{\overline{B}^n}(\Gamma) \leq \sup\{D(d\xi) \mid \xi \in \mathcal{D}^k(\tilde{B}^n), \|d\xi\| \leq 1 \text{ in } \overline{B}^n\} \leq \mathbf{M}(D),$$

hence $F_{\overline{B}^n}(\Gamma) \leq m_r(\Gamma)$. Conversely, consider the seminorm $\nu(\eta) := \sup_{\overline{B}^n} \|\eta\|$, for $\eta \in \mathcal{D}^{k+1}(\tilde{B}^n)$. By means of the linear map $\xi \mapsto d\xi$ from $\mathcal{D}^k(\tilde{B}^n)$ to $\mathcal{D}^{k+1}(\tilde{B}^n)$, we may regard Γ as a linear functional $\tilde{\Gamma}$ on the subspace of exact forms in $\mathcal{D}^{k+1}(\tilde{B}^n)$, endowed with the sup-norm. Of course, $\tilde{\Gamma}(d\xi) \leq F_{\overline{B}^n}(\Gamma) \nu(d\xi)$ for every $\xi \in \mathcal{D}^k(\tilde{B}^n)$. Therefore, by Hahn-Banach theorem, we can extend $\tilde{\Gamma}$ to a linear functional S on

$\mathcal{D}^{k+1}(\tilde{B}^n)$ such that $S = \tilde{\Gamma}$ on exact forms in $\mathcal{D}^{k+1}(\tilde{B}^n)$, and such that $S(\eta) \leq F_{\tilde{B}^n}(\Gamma) \nu(\eta)$ for every $\eta \in \mathcal{D}^{k+1}(\tilde{B}^n)$. Since $\partial S = \Gamma$ and $\text{spt } S \subset \overline{B}^n$, we obtain the assertion. \square

Now, taking $k = n - 2$, $\Gamma = \mathbb{P}(u)$ and $D = (\mathbb{D}(u) - \mathbb{D}(\varphi))$, see (1.3), Proposition 2.5 yields

$$L(u) = m_r(\mathbb{P}(u)) \quad \forall u \in W_\varphi^{1,1}(\tilde{B}^n, S^1).$$

Moreover, by Definition 1.1 we have

$$m_r(\mathbb{P}(u)) \leq \mathbf{M}(\mathbb{D}(u) - \mathbb{D}(\varphi)).$$

Since $(\mathbb{D}(u) - \mathbb{D}(\varphi)) \llcorner \overline{B}^n = (\mathbb{D}(u) - \mathbb{D}(\varphi)) \llcorner \overline{B}^n$, and by the definition of mass, for $v = u, \varphi$,

$$\begin{aligned} \mathbf{M}(\mathbb{D}(v) \llcorner \overline{B}^n) &= \frac{1}{2\pi} \sup \left\{ \int_{\overline{B}^n} v^\# \omega_{S^1} \wedge \gamma : \gamma \in \mathcal{D}^{n-1}(B^n), \|\gamma\| \leq 1 \right\} \\ &\leq \frac{1}{2\pi} \int_{B^n} |\nabla v| dx, \end{aligned}$$

we conclude that

$$L(u) = m_r(\mathbb{P}(u)) \leq \frac{1}{2\pi} \left(\int_{B^n} |\nabla u| dx + \int_{B^n} |\nabla \varphi| dx \right) \quad \forall u \in W_\varphi^{1,1}(\tilde{B}^n, S^1). \quad (2.7)$$

FLAT NORM AND INTEGRAL MASS. By (1.7), on account of Definition 1.1, we also infer that

$$L(u) = m_i(\mathbb{P}(u)) := \inf \{ \mathbf{M}(L) : L \in \mathcal{R}_{n-1}(\tilde{B}^n), \text{spt } L \subset \overline{B}^n, \partial L = (-1)^n \mathbb{P}(u) \} \quad (2.8)$$

for every $u \in W_\varphi^{1,1}(\tilde{B}^n, S^1)$. We thus obtain that for every $n \geq 2$ and $u \in W_\varphi^{1,1}(\tilde{B}^n, S^1)$ there exists an i.m. rectifiable current $L \in \mathcal{R}_{n-1}(\tilde{B}^n)$, hence with finite mass, such that $\text{spt } L \subset \overline{B}^n$ and $(-1)^n \partial L = 2\pi T(u)$, compare [6, Thm. 3] for the case $n = 2$.

We notice that in [1] it is proved that the converse holds true. More precisely, for any i.m. rectifiable current $L \in \mathcal{R}_{n-1}(\tilde{B}^n)$, with $\text{spt } L \subset \overline{B}^n$, there exists a function $u \in W_\varphi^{1,1}(\tilde{B}^n, S^1)$ such that $(-1)^n \partial L = 2\pi T(u)$.

We also recall that by the boundary rectifiability theorem, see e.g. [25, Thm. 30.3], if $L \in \mathcal{R}_{n-1}(\tilde{B}^n)$ has boundary of finite mass, $\mathbf{M}(\partial L) < \infty$, then ∂L is an i.m. rectifiable current in $\mathcal{R}_{n-2}(\tilde{B}^n)$. Due to (2.3) we thus obtain

Corollary 2.6 *Let $n \geq 2$ and $u \in W_\varphi^{1,1}(\tilde{B}^n, S^1)$. The distribution $T(u)$ is a measure of finite mass if and only if the current $\mathbb{P}(u)$ is i.m. rectifiable in $\mathcal{R}_{n-2}(\tilde{B}^n)$.*

In the case $n = 2$, this yields the representation

$$T(u) = 2\pi \sum_{h=1}^M (\delta_{P_h} - \delta_{N_h}),$$

where δ_P is the unit Dirac mass at P and the sum is finite.

Finally, as a consequence of Proposition 1.6, by (2.8) we obtain the following link between the relaxed energy $\widetilde{\mathcal{E}}_{1,1}(u)$ and the minimal connection $L(u)$.

Corollary 2.7 *Let $n \geq 2$. For every $u \in W_\varphi^{1,1}(\tilde{B}^n, S^1)$ and for $\Omega = \tilde{B}^n$ or B^n we have*

$$\widetilde{\mathcal{E}}_{1,1}(u, \Omega) - \int_{\Omega} |\nabla u| dx = 2\pi \cdot L(u).$$

Therefore, by (2.7) we conclude with the following:

Corollary 2.8 *Let $n \geq 2$. For every $u \in W_{\varphi}^{1,1}(\widetilde{B}^n, S^1)$ and for $\Omega = \widetilde{B}^n$ or B^n we have*

$$\widetilde{\mathcal{E}}_{1,1}(u, \Omega) \leq \int_{\Omega} |\nabla u| dx + \int_{B^n} |\nabla u| dx + \int_{B^n} |\nabla \varphi| dx.$$

In particular, if φ is constant on ∂B^n we have

$$\widetilde{\mathcal{E}}_{1,1}(u, B^n) \leq 2 \int_{B^n} |\nabla u| dx.$$

THE CASE WITH NO BOUNDARY DATA. In a similar way, we argue as follows. Let $\Gamma \in \mathcal{D}_k(B^n)$, and suppose that Γ is the boundary in B^n of a $(k+1)$ -dimensional current $D \in \mathcal{D}_{k+1}(B^n)$, i.e., $(\partial D) \llcorner B^n = \Gamma$, with $\mathbf{M}(D) < \infty$. The *flat norm* of Γ is defined by

$$F_{B^n}(\Gamma) := \sup\{\Gamma(\xi) \mid \xi \in \mathcal{D}^k(B^n), \|\mathbf{d}\xi\| \leq 1\}.$$

Moreover, we denote respectively by

$$\begin{aligned} m_{i,B^n}(\Gamma) &:= \inf\{\mathbf{M}(L) \mid L \in \mathcal{R}_{k+1}(B^n), (\partial L) \llcorner B^n = \Gamma\} \\ m_{r,B^n}(\Gamma) &:= \inf\{\mathbf{M}(D) \mid D \in \mathcal{D}_{k+1}(B^n), (\partial D) \llcorner B^n = \Gamma\} \end{aligned}$$

the *integral* and *real mass* of Γ in B^n . Also, in case $m_{i,B^n}(\Gamma) < \infty$, we say that an i.m. rectifiable current $L \in \mathcal{R}_{k+1}(B^n)$ is an *integral minimal connection* of Γ *allowing connections to the boundary* if $(\partial L) \llcorner B^n = \Gamma$ and $\mathbf{M}(L) = m_{i,B^n}(\Gamma)$. Similarly to Proposition 2.5, we have

$$F_{B^n}(\Gamma) = m_{r,B^n}(\Gamma).$$

Taking $k = n - 2$, we now define for any $n \geq 2$

$$L(u) := F_{B^n}(\mathbb{P}(u)), \quad u \in W^{1,1}(B^n, S^1),$$

so that we obtain

$$L(u) = \frac{1}{2\pi} \max\{\langle T(u), \zeta \rangle \mid \zeta \in \text{Lip}(B^n, \Lambda^{n-2}TB^n), \|\nabla \zeta\|_{\infty} \leq 1\}.$$

Setting now

$$\widetilde{\mathcal{E}}_{1,1}(u) := \inf \left\{ \liminf_{h \rightarrow \infty} \int_{B^n} |\nabla u_h| dx : \{u_h\} \subset C^1(B^n, S^1), u_h \rightarrow u \text{ a.e.} \right\},$$

we obtain for every $u \in W^{1,1}(B^n, S^1)$

$$\begin{aligned} \widetilde{\mathcal{E}}_{1,1}(u) &= \int_{B^n} |\nabla u| dx + 2\pi m_{i,B^n}(\mathbb{P}(u)) \\ &= \int_{B^n} |\nabla u| dx + 2\pi \inf\{\mathbf{M}(L) : L \in \mathcal{R}_{n-1}(B^n), (\partial L) \llcorner B^n = (-1)^n \mathbb{P}(u)\}. \end{aligned}$$

Moreover, since

$$L(u) = m_{r,B^n}(\mathbb{P}(u)) = m_{i,B^n}(\mathbb{P}(u)) \leq \frac{1}{2\pi} \int_{B^n} |\nabla u| dx,$$

we obtain that

$$\widetilde{\mathcal{E}}_{1,1}(u) \leq 2 \int_{B^n} |\nabla u| dx \quad \forall u \in W^{1,1}(B^n, S^1).$$

Finally, a statement analogous to Proposition 3.7 below holds true. In particular,

$$\begin{aligned} \widetilde{\mathcal{E}}_{1,1}(u) = \mathcal{E}_{1,1}(u) &\iff G_u \in \text{cart}(B^n \times S^1) \iff T(u) = 0 \iff L(u) = 0 \\ &\iff u \text{ belongs to the strong } W^{1,1}\text{-closure of } C^{\infty}(B^n, S^1). \end{aligned}$$

DISTRIBUTIONAL MINORS. Let $G = (G_i^h)$ be a $(2 \times n)$ -matrix, e.g., $G = \nabla u$ for some $u \in W^{1,1}(B^n, \mathbb{R}^2)$. For $1 \leq i < j \leq n$, we denote by $G_{i,j}$ the (2×2) -submatrix obtained by selecting the columns by i and j . We will also denote by $M_{i,j}(G)$ its determinant

$$M_{i,j}(G) := \det G_{i,j}.$$

We recall that the *matrix of the adjoints* of $G_{i,j}$ is defined by the formula

$$\begin{aligned} (\operatorname{adj} G_{i,j})_i^1 &:= G_j^2, & (\operatorname{adj} G_{i,j})_j^1 &:= -G_i^2, \\ (\operatorname{adj} G_{i,j})_i^2 &:= -G_j^1, & (\operatorname{adj} G_{i,j})_j^2 &:= G_i^1. \end{aligned} \quad (2.9)$$

Definition 2.9 Let $u \in W^{1,1}(B^n, \mathbb{R}^2) \cap L^\infty$. The *distributional minor of indices* $1 \leq i < j \leq n$ of ∇u is defined by

$$\operatorname{Det} \nabla_{i,j} u := \frac{1}{2} \sum_{h=1}^2 \left(\frac{\partial}{\partial x_i} (u^h(x) ((\operatorname{adj} \nabla u)_{i,j})_i^h) + \frac{\partial}{\partial x_j} (u^h(x) ((\operatorname{adj} \nabla u)_{i,j})_j^h) \right).$$

More explicitly, since $G_i^h = \nabla_i u^h = u_{x_i}^h$ if $G = \nabla u$, by (2.9) we have

$$\operatorname{Det} \nabla_{i,j} u := \frac{1}{2} \left(\frac{\partial}{\partial x_i} (u^1 u_{x_j}^2 - u^2 u_{x_j}^1) - \frac{\partial}{\partial x_j} (u^1 u_{x_i}^2 - u^2 u_{x_i}^1) \right)$$

i.e., for every $\zeta \in \operatorname{Lip}(B^n)$,

$$\langle \operatorname{Det} \nabla_{i,j} u, \zeta \rangle := -\frac{1}{2} \left((u \times u_{x_j}) D_i \zeta - (u \times u_{x_i}) D_j \zeta \right). \quad (2.10)$$

In particular, if $n = 2$ we infer that $\operatorname{Det} \nabla_{1,2} u = \operatorname{Det} \nabla u$, the *distributional determinant* of ∇u . Moreover, by (2.2) we also have that

$$T_{i,j}(u) = 2 \operatorname{Det} \nabla_{i,j} u.$$

Notice that if $u \in W^{1,1}(B^n, \mathbb{R}^2) \cap L^\infty$ is smooth, then $\operatorname{Det} \nabla_{i,j} u$ coincides with the pointwise determinant $M_{i,j}(\nabla u)$. In fact by (2.9) we have $\frac{\partial}{\partial x_i} ((\operatorname{adj} \nabla u)_{i,j})_i^h + \frac{\partial}{\partial x_j} ((\operatorname{adj} \nabla u)_{i,j})_j^h = 0$, so that Laplace's formulas for $h = 1, 2$ yield

$$\begin{aligned} \frac{\partial}{\partial x_i} (u^h(x) ((\operatorname{adj} \nabla u)_{i,j})_i^h) + \frac{\partial}{\partial x_j} (u^h(x) ((\operatorname{adj} \nabla u)_{i,j})_j^h) &= \frac{\partial u^h}{\partial x_i} ((\operatorname{adj} \nabla u)_{i,j})_i^h + \frac{\partial u^h}{\partial x_j} ((\operatorname{adj} \nabla u)_{i,j})_j^h \\ &+ u^h \left(\frac{\partial}{\partial x_i} ((\operatorname{adj} \nabla u)_{i,j})_i^h + \frac{\partial}{\partial x_j} ((\operatorname{adj} \nabla u)_{i,j})_j^h \right) \\ &= M_{i,j}(\nabla u). \end{aligned}$$

Of course, if $u \in \operatorname{Lip}(B^n, S^1)$, the *area formula* yields that $M_{i,j}(\nabla u) = 0$.

Let now $u \in W^{1,1}(B^n, \mathbb{R}^2) \cap L^\infty$ be such $M_{i,j}(\nabla u) \in L^1(B^n)$ for every $i < j$. Suppose in addition that the boundary of the graph ∂G_u has finite mass in $B^n \times \mathbb{R}^2$, i.e., $T(u)$ is a bounded measure, compare Corollary 2.6. With these hypotheses, in [21] it is shown that for every $i < j$ the distributional minor $\operatorname{Det} \nabla_{i,j} u$ is a signed Radon measure with finite total variation, the density of its absolute continuous part is equal to the pointwise determinant $M_{i,j}(\nabla u)$

$$\operatorname{Det} \nabla_{i,j} u = M_{i,j}(\nabla u) \cdot d\mathcal{L}^n + (\operatorname{Det} \nabla_{i,j} u)^s, \quad (\operatorname{Det} \nabla_{i,j} u)^s \perp \mathcal{L}^n; \quad (2.11)$$

moreover, the singular part $(\operatorname{Det} \nabla_{i,j} u)^s$ is supported on a countably \mathcal{H}^{n-2} -rectifiable set, possibly with unbounded \mathcal{H}^{n-2} -measure. In particular, $(\operatorname{Det} \nabla_{i,j} u)^s$ does not contain any *Cantor type* mass and, in dimension $n = 2$, we have

$$(\operatorname{Det} \nabla u)^s = \sum_h c_h \delta_{x_h}, \quad c_h \in \mathbb{R},$$

where the sum is possibly infinite, but satisfying $\sum_h |c_h| < \infty$. Finally, if $n = 2$, notice that if the boundary ∂G_u has infinite mass, it may happen that the singular part of the distributional determinant is supported on a Cantor-type set of Hausdorff dimension $d \in]0, 2[$, compare [13] [22].

3 Lifting

In this section we extend to any dimension $n \geq 2$ a result from [6], proved in dimension $n = 2$, about the interpretation of the minimal connection $L(u)$ in terms of the L^1 -distance of the vector field $u \times \nabla u$ to the class of gradient maps, where

$$u \times \nabla u := (u \times u_{x_1}, \dots, u \times u_{x_n}). \quad (3.1)$$

More precisely, we will prove in any dimension $n \geq 2$ the following

Theorem 3.1 *For any $u \in W_\varphi^{1,1}(\tilde{B}^n, S^1)$ we have*

$$L(u) = \frac{1}{2\pi} \min_{\psi \in BV(\tilde{B}^n, \mathbb{R})} |u \times \nabla u - D\psi|(B^n).$$

In order to prove Theorem 3.1, we recall some results from [12] about the existence of a lifting of currents in $\text{cart}_\varphi(\tilde{B}^n \times S^1)$. To this aim we first recall, see [13], that the current *subgraph* of an L^1 -function $\psi \in L^1(\tilde{B}^n, \mathbb{R})$ is the $(n+1)$ -dimensional current in $\mathcal{D}_{n+1}(\tilde{B}^n \times \mathbb{R})$ defined by

$$SG_\psi(\phi(x, t) dx \wedge dt) := \int_{\tilde{B}^n} \left(\int_0^{\psi(x)} \phi(x, t) dt \right) dx, \quad \phi \in C_c^\infty(\tilde{B}^n \times \mathbb{R}). \quad (3.2)$$

Moreover, in the sequel we will denote by $i : \tilde{B}^n \times \mathbb{R} \rightarrow \tilde{B}^n \times S^1$ the map

$$i(x, t) := (x, \cos t, \sin t),$$

and by G_{q_0} the current in $\mathcal{D}_n(\tilde{B}^n \times S^1)$ integration over the graph of the constant map $q_0(x) \equiv (1, 0)$. In [12], see also [13, Vol. II, Sec. 6.2.2], the following is proved:

Proposition 3.2 *Let $T \in \text{cart}_\varphi(\tilde{B}^n \times S^1)$. The following facts hold:*

i) *There exists a current $\Sigma \in \mathcal{D}_{n+1}(\tilde{B}^n \times S^1)$ such that*

$$T - G_{q_0} = (-1)^n \partial \Sigma. \quad (3.3)$$

ii) *There exists a function $\psi \in BV(\tilde{B}^n, \mathbb{R})$ such that $\Sigma = i_\# SG_\psi$, i.e.,*

$$T - G_{q_0} = (-1)^n i_\# \partial SG_\psi. \quad (3.4)$$

In particular, $\mathbf{M}(\partial SG_\psi) = \mathbf{M}(T) + \mathcal{L}^n(\tilde{B}^n) < \infty$.

iii) *If $u_T \in BV_\varphi(\tilde{B}^n, S^1)$ is the BV-function corresponding to T , then*

$$u_T = e^{i\psi} \quad \mathcal{L}^n\text{-a.e. on } \tilde{B}^n. \quad (3.5)$$

Remark 3.3 In [15] it is shown that for every $u \in BV_\varphi(\tilde{B}^n, S^1)$ there exists a current $T \in \text{cart}_\varphi(\tilde{B}^n \times S^1)$ such that $u_T = u$. As a consequence, from Proposition 3.2 we infer that *every BV-function $u \in BV_\varphi(\tilde{B}^n, S^1)$ has a lift ψ in $BV(\tilde{B}^n, \mathbb{R})$* . Notice that in general the lift of a $W^{1,1}$ -function in $W_\varphi^{1,1}(\tilde{B}^n, S^1)$ is not a Sobolev function in $W^{1,1}(\tilde{B}^n, \mathbb{R})$, but only a BV-function. However, the existence of a lifting in the sense of (3.5) is not very useful, even if u_T belongs to $W_\varphi^{1,1}(\tilde{B}^n, S^1)$. In fact, to recover homological and topological properties from the lifting, the right condition is (3.4). As we shall see in Proposition 3.7 below, nice properties are recovered if the lifting ψ is a Sobolev map $\psi \in W^{1,1}(\tilde{B}^n, \mathbb{R})$.

Remark 3.4 From (3.4) it readily follows that *the area of the graph of ψ is equal to the mass of T ,*

$$\int_{\tilde{B}^n} \sqrt{1 + |\nabla \psi|^2} dx + |D^J \psi|(\tilde{B}^n) + |D^C \psi|(\tilde{B}^n) = \mathbf{M}(T),$$

where $\nabla\psi$, $D^J\psi$ and $D^C\psi$ denote the approximate gradient, the jump part and the Cantor part of the distributional derivative of ψ , see e.g. [3] [13]. In particular, if u_T belongs to the Sobolev class $W_\varphi^{1,1}(\tilde{B}^n, S^1)$, and $L \in \mathcal{R}_{n-1}(\tilde{B}^n)$ is given by (1.12), we have

$$|D\psi|(\tilde{B}^n) = \int_{\tilde{B}^n} |\nabla u_T| dx + 2\pi \mathbf{M}(L).$$

Remark 3.5 The formula (3.4) clearly yields that if $u_T \in W_\varphi^{1,1}(\tilde{B}^n, S^1)$, the derivative $D\psi$ of the lifting ψ has a *null Cantor part*. However, in general, the lifting ψ of a function $u \in W_\varphi^{1,1}(\tilde{B}^n, S^1)$ is not a Sobolev function in $W_\varphi^{1,1}(\tilde{B}^n, \mathbb{R})$, think for instance of $n = 2$ and $u(x) := x/|x|$. However, if the graph of u has no inner boundary, i.e., if G_u belongs to $\text{cart}_\varphi(\tilde{B}^n \times S^1)$, the existence of a lifting $\psi \in W^{1,1}(\tilde{B}^n, \mathbb{R})$ satisfying (3.4) with $T = G_u$ is provided. In fact, we have

Corollary 3.6 *Let $n \geq 2$ and let $u \in W_\varphi^{1,1}(\tilde{B}^n, S^1)$. Suppose that $\partial G_u = 0$ on $\mathcal{D}^n(\tilde{B}^n \times S^1)$, i.e., $G_u \in \text{cart}_\varphi(\tilde{B}^n \times S^1)$. There exists a function $\psi \in W^{1,1}(\tilde{B}^n, \mathbb{R})$ such that*

$$G_u - G_{q_0} = (-1)^n i_{\neq} \partial S G_\psi. \quad (3.6)$$

Moreover, $u = e^{i\psi}$ a.e. on \tilde{B}^n and

$$D\psi = \nabla\psi d\mathcal{L}^n = (u \times \nabla u) d\mathcal{L}^n,$$

where $u \times \nabla u$ is the L^1 -vector field given by (3.1).

Proposition 3.7 *Let $n \geq 2$ and let $u \in W_\varphi^{1,1}(\tilde{B}^n, S^1)$. The following facts are equivalent:*

- (a) $G_u \in \text{cart}_\varphi(\tilde{B}^n \times S^1)$;
- (b) $T(u) = 0$;
- (c) $L(u) = 0$;
- (d) there exists a function $\psi \in W^{1,1}(\tilde{B}^n, \mathbb{R})$ such that (3.6) holds and $u = e^{i\psi}$ a.e. on \tilde{B}^n ;
- (e) u belongs to the strong $W^{1,1}$ -closure of smooth maps in $C_\varphi^\infty(\tilde{B}^n, S^1)$.

The equivalence (b) \iff (d) \iff (e) was first proved in [7] for Sobolev maps in $W^{1,p}$, see also [5].

OPTIMAL LIFTING. Following [6], we finally consider the energy

$$\widehat{\mathcal{E}}_{1,1}(u, \Omega) := \inf\{|D\psi|(\Omega) : \psi \in BV(\Omega, \mathbb{R}), u = e^{i\psi} \text{ a.e. on } \Omega\},$$

where $\Omega = B^n$ or \tilde{B}^n . Since Ω is simply connected, arguing exactly as in [6, Prop. 2] we obtain that

$$\widehat{\mathcal{E}}_{1,1}(u, \Omega) = \widetilde{\mathcal{E}}_{1,1}(u, \Omega) \quad \forall u \in W_\varphi^{1,1}(\tilde{B}^n, S^1).$$

Therefore, by Corollary 2.7 we obtain that for every $n \geq 2$ and $u \in W_\varphi^{1,1}(\tilde{B}^n, S^1)$

$$\widehat{\mathcal{E}}_{1,1}(u, \Omega) - \int_\Omega |\nabla u| dx = 2\pi \cdot L(u).$$

Moreover, if φ is constant on ∂B^n , by Corollary 2.8 we have

$$\widehat{\mathcal{E}}_{1,1}(u, B^n) \leq 2 \int_{B^n} |\nabla u| dx.$$

In particular, since $L(u) = 0 \iff T(u) = 0$, by Proposition 3.7 we infer that

$$\forall u \in W_\varphi^{1,1}(\tilde{B}^n, S^1), \quad \widehat{\mathcal{E}}_{1,1}(u, \Omega) = \int_\Omega |\nabla u| dx \iff G_u \in \text{cart}_\varphi(\tilde{B}^n \times S^1).$$

PROOF OF THEOREM 3.1: We recall that for any $u \in W_\varphi^{1,1}(\tilde{B}^n, S^1)$ we have (1.15), see (1.14). This yields for the integral mass

$$m_i(u) = \inf\{\mathbf{M}(L_T) \mid G_u + L_T \times \llbracket S^1 \rrbracket \in \mathcal{T}_u\}, \quad (3.7)$$

compare (2.8). Moreover, by Proposition 3.2, to any $T \in \mathcal{T}_u$ it corresponds a function $\psi_T \in BV(\tilde{B}^n, \mathbb{R})$ such that

$$G_u + L_T \times \llbracket S^1 \rrbracket - G_{q_0} = (-1)^n i_\# \partial SG_{\psi_T} \quad \text{on } \mathcal{D}^n(\tilde{B}^n \times S^1). \quad (3.8)$$

Let $\omega \in \mathcal{D}^n(\tilde{B}^n \times S^1)$ be given by $\omega = \pi^\# \omega_\phi \wedge \hat{\pi}^\# \omega_{S^1}$, where $\omega_\phi \in \mathcal{D}^{n-1}(\tilde{B}^n)$ is given by

$$\omega_\phi := \sum_{i=1}^n (-1)^{i-1} \phi^i(x) \widehat{dx}^i, \quad \phi = (\phi^1, \dots, \phi^n) \in C_c^\infty(\tilde{B}^n, \mathbb{R}). \quad (3.9)$$

In the sequel we omit to write the action of the projection maps π and $\hat{\pi}$. Since

$$\begin{aligned} \omega_\phi \wedge u^\# \omega_{S^1} &= \sum_{i=1}^n (-1)^{i-1} \phi^i \widehat{dx}^i \wedge (u^1 du^2 - u^2 du^1) \\ &= (-1)^{n-1} \sum_{i=1}^n \phi^i \cdot (u \times u_{x_i}) dx, \end{aligned}$$

we have

$$G_u(\omega_\phi \wedge \omega_{S^1}) = \int_{\tilde{B}^n} \omega_\phi \wedge u^\# \omega_{S^1} = (-1)^{n-1} \int_{\tilde{B}^n} \langle u \times \nabla u, \phi \rangle dx, \quad (3.10)$$

whereas

$$L_T \times \llbracket S^1 \rrbracket (\omega_\phi \wedge \omega_{S^1}) = L_T(\omega_\phi) \cdot \llbracket S^1 \rrbracket (\omega_{S^1}) = 2\pi L_T(\omega_\phi)$$

and

$$G_{q_0}(\omega_\phi \wedge \omega_{S^1}) = 0. \quad (3.11)$$

Moreover, since $d\omega_\phi = \text{div} \phi dx$ and

$$i^\# d(\omega_\phi \wedge \omega_{S^1}) = i^\# (d\omega_\phi \wedge \omega_{S^1}) = i^\# (\text{div} \phi dx \wedge \omega_{S^1}) = \text{div} \phi dx \wedge dt,$$

on account of (3.2) we have

$$\begin{aligned} i_\# \partial SG_{\psi_T}(\omega_\phi \wedge \omega_{S^1}) &= i_\# SG_{\psi_T}(\text{div} \phi(x) dx \wedge \omega_{S^1}) \\ &= SG_{\psi_T}(\text{div} \phi(x) dx \wedge dt) \\ &= \int_{\tilde{B}^n} \text{div} \phi(x) (\psi_T(x) - 0) dx = -\langle D\psi_T, \phi \rangle. \end{aligned} \quad (3.12)$$

By (3.8) we have thus obtained

$$(-1)^n 2\pi L_T(\omega_\phi) = \int_{\tilde{B}^n} \langle u \times \nabla u, \phi \rangle dx - \langle D\psi_T, \phi \rangle \quad \forall \phi \in C_c^\infty(\tilde{B}^n, \mathbb{R}^n).$$

Moreover, since $T = G_\varphi$ on $(\tilde{B}^n \setminus \bar{B}^n) \times S^1$, we have $D\psi_T = \varphi \times \nabla \varphi$ on $\tilde{B}^n \setminus \bar{B}^n$ and hence

$$\mathbf{M}(L_T) := \sup\{L_T(\omega_\phi) \mid \|\phi\|_\infty \leq 1\} = \frac{1}{2\pi} |u \times \nabla u - D\psi_T|(B^n).$$

In conclusion, by (3.7) and (2.8) we obtain the assertion. \square

PROOF OF COROLLARY 3.6: By Proposition 3.2 we find the existence of a function $\psi \in BV(\tilde{B}^n, \mathbb{R})$ such that

$$G_u - G_{q_0} = (-1)^n i_{\#} \partial S G_{\psi},$$

see (3.4). Taking $\omega = \pi^{\#} \omega_{\phi} \wedge \hat{\pi}^{\#} \omega_{S^1} \in \mathcal{D}^n(\tilde{B}^n \times S^1)$, where ω_{ϕ} is given by (3.9), using (3.10), (3.11), and (3.12) we obtain that

$$\int_{\tilde{B}^n} \langle u \times \nabla u, \phi \rangle dx = \langle D\psi, \phi \rangle \quad \forall \phi \in C_c^{\infty}(\tilde{B}^n, \mathbb{R}^n).$$

Therefore, since $u \times \nabla u \in L^1(\tilde{B}^n, \mathbb{R}^n)$, we obtain $D\psi = (u \times \nabla u) d\mathcal{L}^n$ and hence the assertion. \square

PROOF OF PROPOSITION 3.7: If (a) holds, then $\partial G_u = 0$, hence by (1.2) we have $\mathbb{P}(u) = 0$, which yields $T(u) = 0$, by (2.3). Conversely, if $T(u) = 0$, we have $\partial G_u = 0$, and hence (a) \iff (b). The equivalence (a) \iff (c) is a trivial consequence of definition (2.4). If (a) holds, we obtain (d) by Corollary 3.6. If (d) holds, and $\{\psi_h\}_h \subset C^{\infty}(\tilde{B}^n, \mathbb{R})$ is such that $\psi_h \rightarrow \psi$ strongly in $W^{1,1}$, with $\varphi = e^{i\psi_h}$ on $\tilde{B}^n \setminus \bar{B}^n$ for any h , setting $u_h := e^{i\psi_h}$ we clearly have $\{u_h\} \subset C_{\varphi}^{\infty}(\tilde{B}^n, S^1)$ and $u_h \rightarrow u$ strongly in $W^{1,1}(\tilde{B}^n, \mathbb{R}^2)$. Finally, if (d) holds, since by Stokes theorem $\partial G_{u_h} = 0$ on $\mathcal{D}^{n-1}(\tilde{B}^n \times S^1)$ if u_h is smooth, and the strong $W^{1,1}$ -convergence yields the weak convergence $G_{u_h} \rightharpoonup G_u$ in the sense of currents in $\mathcal{D}_n(\tilde{B}^n \times S^1)$, we find that $\partial G_u = 0$ on $\mathcal{D}^{n-1}(\tilde{B}^n \times S^1)$, i.e., $G_u \in \text{cart}_{\varphi}(\tilde{B}^n \times S^1)$. \square

4 Examples

THE CASE $n = 2$. Let $a_{\pm} := (0, \pm 1/2)$ and consider the $W^{1,1}$ -maps

$$u_+(x) := \frac{x - a_+}{|x - a_+|}, \quad u_-(x) := \psi \left(\frac{x - a_-}{|x - a_-|} \right),$$

where $\psi : S^1 \rightarrow S^1$ is given by $\psi(y_1, y_2) := (y_1, -y_2)$. Since u_{\pm} has degree ± 1 , we may and do find a smooth $W^{1,1}$ -map $\phi : (\bar{B}^2 \setminus (B(a_+, 1/4) \cup B(a_-, 1/4))) \rightarrow S^1$ satisfying

$$\phi|_{\partial B^2} \equiv (1, 0), \quad \phi|_{\partial B(a_+, 1/4)} = u_+, \quad \phi|_{\partial B(a_-, 1/4)} = u_-.$$

Define $u : \bar{B}^2 \rightarrow S^1$ by

$$u(x) := \begin{cases} u_+(x) & \text{if } |x - a_+| < 1/4 \\ u_-(x) & \text{if } |x - a_-| < 1/4 \\ \phi(x) & \text{elsewhere on } B^2 \end{cases}$$

and set $u = \varphi \equiv (1, 0)$ in $\tilde{B}^2 \setminus B^2$, so that $u \in W_{\varphi}^{1,2}(\tilde{B}^2, S^1)$.

Remark 4.1 For future use, we also may and do define ϕ so that

$$u(x_1, -x_2) = u(x_1, x_2) \quad \forall (x_1, x_2) \in \bar{B}^2.$$

Following Sec. 3.2.2 in [13, Vol. I], we have

$$\partial G_u \llcorner \tilde{B}^2 \times S^1 = (\delta_{a_-} - \delta_{a_+}) \times \llbracket S^1 \rrbracket.$$

This yields that

$$\widetilde{\mathcal{E}}_{1,1}(u, \tilde{B}^2) = \int_{B^2} |Du| dx + 2\pi |a_+ - a_-|.$$

In fact, the current T of minimal mass in $\text{cart}_{\varphi}(\tilde{B}^2 \times S^1)$ satisfying $u_T = u$ is given by

$$T := G_u + \llbracket a_-, a_+ \rrbracket \times \llbracket S^1 \rrbracket,$$

where $\llbracket a_-, a_+ \rrbracket$ is the 1-current integration on the positively oriented segment connecting a_- to a_+ , so that $\partial \llbracket a_-, a_+ \rrbracket = \delta_{a_+} - \delta_{a_-}$. Moreover, see (1.13), we have

$$\mathbb{P}(u) = \delta_{a_+} - \delta_{a_-}, \quad T(u) = 2\pi(\delta_{a_+} - \delta_{a_-}), \quad L(u) = |a_+ - a_-|.$$

THE CASE $n \geq 3$. If $n = 3$, define $u : \overline{B}^3 \rightarrow S^1$ as the $W^{1,1}$ -map given by the rotation on the x_1 -axis of the map defined as in the case $n = 2$ on the 2-disk $\overline{B}^2 \simeq \overline{B}^3 \cap \{x_3 = 0\}$. By induction on the dimension n , define $u : \overline{B}^n \rightarrow S^1$ as the $W^{1,1}$ -map given by the rotation on the x_1 -axis of the map defined as in the case $n - 1$ on the $(n - 1)$ -disk $\overline{B}^{n-1} \simeq \overline{B}^n \cap \{x_n = 0\}$. By Remark 4.1, in the case $n = 3$, and by the inductive argument, we infer that u is smooth outside the $(n - 2)$ -sphere $\Delta := \{x \in B^n \mid x_1 = 0, |x| = 1/2\}$. Moreover, setting again $u = \varphi \equiv (1, 0)$ in $\widetilde{B}^n \setminus B^n$, this time we have

$$\partial G_u \llcorner \widetilde{B}^n \times S^1 = -\llbracket \Delta \rrbracket \times \llbracket S^1 \rrbracket,$$

where $\llbracket \Delta \rrbracket \in \mathcal{R}_{n-2}(B^n)$ is the $(n - 2)$ -current integration on the $(n - 2)$ -sphere Δ , oriented in the natural way. Notice that $\partial \llbracket \Delta \rrbracket = 0$. Moreover, we have

$$\widetilde{\mathcal{E}}_{1,1}(u, \widetilde{B}^n) = \int_{B^n} |Du| dx + 2\pi \mathcal{H}^{n-1}(D),$$

where $D := \{x \in B^n \mid x_1 = 0, |x| \leq 1/2\}$. In fact, the current T of minimal mass in $\text{cart}_\varphi(\widetilde{B}^n \times S^1)$ satisfying $u_T = u$ is given by

$$T := G_u + \llbracket D \rrbracket \times \llbracket S^1 \rrbracket,$$

where $\llbracket D \rrbracket$ is the $(n - 1)$ -current integration on the positively oriented $(n - 1)$ -disk D , so that $\partial \llbracket D \rrbracket = \llbracket \Delta \rrbracket$. Finally, according to (1.13), we have

$$\mathbb{P}(u) = (-1)^n \llbracket \Delta \rrbracket, \quad T(u) = (-1)^n 2\pi \llbracket \Delta \rrbracket, \quad L(u) = \mathcal{H}^{n-1}(D).$$

5 The relaxed energy of $W^{1,p}$ -maps into S^1

Let $p > 1$ and $u \in W_\varphi^{1,p}(\widetilde{B}^n, S^1)$, where

$$W_\varphi^{1,p}(\widetilde{B}^n, S^1) := \{u \in W^{1,p}(\widetilde{B}^n, \mathbb{R}^2) : |u(x)| = 1 \text{ a.e. in } \widetilde{B}^n, u = \varphi \text{ in } \widetilde{B}^n \setminus \overline{B}^n\}.$$

We now briefly discuss the relaxed $W^{1,p}$ -energy, defined for $u \in W_\varphi^{1,p}(\widetilde{B}^n, S^1)$ by

$$\widetilde{\mathcal{E}}_{1,p}(u) =: \inf \left\{ \liminf_{h \rightarrow \infty} \int_{\widetilde{B}^n} |\nabla u_h|^p dx : \{u_h\} \subset C_\varphi^1(\widetilde{B}^n, S^1), u_h \rightarrow u \text{ a.e.} \right\}.$$

It is well-known that

$$p \geq 2 \quad \Rightarrow \quad \widetilde{\mathcal{E}}_{1,p}(u) = \int_{\widetilde{B}^n} |\nabla u|^p dx \quad \forall u \in W_\varphi^{1,p}(\widetilde{B}^n, S^1).$$

This property follows from standard argument for $p > n$ and by Schoen-Uhlenbeck density theorem [24] in the critical case $p = n$. Since the higher order homotopy groups of the 1-sphere are all trivial, $\pi_i(S^1) = 0$ for all $i \geq 2$, this property follows from Bethuel's theorem [4] in the case $2 \leq p < n$.

We now prove the following

Theorem 5.1 *Let $1 < p < 2$ and $u \in W_\varphi^{1,p}(\widetilde{B}^n, S^1)$. Then*

$$\widetilde{\mathcal{E}}_{1,p}(u) = \begin{cases} \int_{\widetilde{B}^n} |\nabla u|^p dx & \text{if } T(u) = 0 \\ +\infty & \text{if } T(u) \neq 0 \end{cases}$$

where $T(u)$ is given by Definition 2.1.

This answers the Open Problems 1 and 2 stated in [6].

Notice that in [7] it is proved that for $1 \leq p < 2$ a Sobolev map $u \in W_\varphi^{1,p}(\tilde{B}^n, S^1)$ can be strongly approximated in $W^{1,p}$ by a smooth sequence in $C_\varphi^1(\tilde{B}^n, S^1)$ if and only if $T(u) = 0$. Therefore, we obtain that for every $p > 1$ a map in $W_\varphi^{1,p}(\tilde{B}^n, S^1)$ belongs to the sequential weak $W^{1,p}$ -closure of smooth maps, i.e., $\widetilde{\mathcal{E}}_{1,p}(u) < \infty$, if and only if it belongs to the strong $W^{1,p}$ -closure of smooth maps from \tilde{B}^n into S^1 .

This is false in the case $p = 1$. In fact, by Proposition 1.6 we know that $\widetilde{\mathcal{E}}_{1,1}(u) < \infty$ for every $u \in W_\varphi^{1,1}(\tilde{B}^n, S^1)$, whereas u belongs to the strong $W^{1,1}$ -closure of smooth maps from \tilde{B}^n into S^1 if and only if $T(u) = 0$.

PROOF OF THEOREM 5.1: Assume that $\widetilde{\mathcal{E}}_{1,p}(u) < \infty$ and let $\{u_h\} \subset C_\varphi^1(\tilde{B}^n, S^1)$ be a smooth sequence satisfying $\sup_h \|u_h\|_{W^{1,p}} < \infty$ and $u_h \rightarrow u \in W_\varphi^{1,p}(\tilde{B}^n, S^1)$ a.e.. Possibly passing to a subsequence, by [13] we infer that G_{u_h} weakly converges in $\mathcal{D}_n(\tilde{B}^n \times S^1)$ to some current $T \in \text{cart}_\varphi(\tilde{B}^n \times S^1)$ satisfying (1.12). Let \mathcal{L} be a compact subset of $\text{set}(L)$, the set of points of L which have non-zero density, \mathcal{L} with positive \mathcal{H}^{n-1} -measure. For any $x \in \mathcal{L}$ we denote by I_x the intersection with \tilde{B}^n of the straight line containing x and orthogonal to the approximate tangent $(n-1)$ -space to \mathcal{L} at x . Since \mathcal{L} is compact, for \mathcal{H}^{n-1} -a.e. $x \in \mathcal{L}$, the 1-dimensional restriction of G_{u_h} to I_x is a sequence of graphs of smooth functions $u_h|_{I_x}$ with equibounded $W^{1,p}$ -energies. For any $x_1, x_2 \in I_x$, let $[x_1, x_2]$ denote the line segment with end points x_1, x_2 . Now, in dimension one, if $p > 1$ the Hölder inequality yields

$$\int_{[x_1, x_2]} |\nabla u_h|_{I_x}| d\mathcal{H}^1 \leq \left(\int_{[x_1, x_2]} |\nabla u_h|_{I_x}|^p d\mathcal{H}^1 \right)^{1/p} \cdot |x_1 - x_2|^{1-1/p} \leq \tilde{C} |x_1 - x_2|^{1-1/p},$$

where \tilde{C} is an absolute constant, depending on the uniform upper bound for the energies $\int_{I_x} |\nabla u_h|_{I_x}|^p d\mathcal{H}^1$. This is in contradiction to the fact that, by a slicing argument, the 1-dimensional currents $G_{u_h|_{I_x}}$ have to converge "near" the point x to the graph $G_{u|_{I_x}}$ of the restriction $u|_{I_x}$ plus a vertical part of the type $\delta_x \times \llbracket S^1 \rrbracket$. In conclusion, we have shown that if $T(u) \neq 0$, then $\widetilde{\mathcal{E}}_{1,p}(u) = +\infty$. The assertion follows from [7], see also [5]. \square

6 The relaxed energy of $W^{1,2}$ -maps into S^2

In this section we collect a few remarks about the Dirichlet energy of $W^{1,2}$ -maps with values into S^2 , the unit sphere in \mathbb{R}^3 . Let $\varphi : \tilde{B}^n \rightarrow S^2$ be a given smooth $W^{1,2}$ -function. For $X := W^{1,2}$, $W^{1,p}$, C^1 , or BV , we set

$$X_\varphi(\tilde{B}^n, S^2) := \{u \in X(\tilde{B}^n, \mathbb{R}^3) : |u(x)| = 1 \text{ a.e. in } \tilde{B}^n, u = \varphi \text{ in } \tilde{B}^n \setminus \bar{B}^n\}.$$

Similarly to Sec. 1, if $u \in W_\varphi^{1,2}(\tilde{B}^n, S^2)$, the i.m. rectifiable current $G_u \in \mathcal{R}_n(\tilde{B}^n \times S^2)$ is defined in an approximate sense by (1.1). Moreover, if $n \geq 3$ we define the $(n-3)$ -current $\mathbb{P}(u) \in \mathcal{D}_{n-3}(\tilde{B}^n)$ by

$$\mathbb{P}(u)(\phi) := \frac{1}{4\pi} \partial G_u(\widehat{\pi}^\# \omega_{S^2} \wedge \pi^\# \phi) = \frac{1}{4\pi} \int_{\tilde{B}^n} u^\# \omega_{S^2} \wedge d\phi \quad (6.1)$$

for every $\phi \in \mathcal{D}^{n-3}(\tilde{B}^n)$, where ω_{S^2} is the volume 2-form on $S^2 \subset \mathbb{R}^3$,

$$\omega_{S^2} := y^1 dy^2 \wedge dy^3 + y^2 dy^3 \wedge dy^1 + y^3 dy^1 \wedge dy^2,$$

and the $(n-2)$ -current $\mathbb{D}(u) \in \mathcal{D}_{n-2}(\tilde{B}^n)$ by

$$\mathbb{D}(u)(\gamma) := \frac{1}{4\pi} G_u(\widehat{\pi}^\# \omega_{S^2} \wedge \pi^\# \gamma) = \frac{1}{4\pi} \int_{\tilde{B}^n} u^\# \omega_{S^2} \wedge \gamma$$

for every $\gamma \in \mathcal{D}^{n-2}(\tilde{B}^n)$. Again we have that $\text{spt } \mathbb{P}(u) \subset \bar{B}^n$, $\partial \mathbb{P}(u) = 0$ and (1.3) holds true. For $\Omega = \tilde{B}^n$ or B^n , denote by

$$\mathbf{D}(u, \Omega) := \frac{1}{2} \int_{\Omega} |\nabla u|^2 dx, \quad u \in W_\varphi^{1,2}(\tilde{B}^n, S^2),$$

the *Dirichlet energy* of u and consider the *relaxed $W^{1,2}$ -energy*

$$\tilde{\mathbf{D}}(u, \Omega) := \inf \left\{ \liminf_{h \rightarrow \infty} \mathbf{D}(u_h, \Omega) : \{u_h\} \subset C_\varphi^1(\tilde{B}^n, S^2), \quad u_h \rightarrow u \quad \text{a.e.} \right\}.$$

Moreover, we let

$$\text{cart}_\varphi(\tilde{B}^n \times S^2) := \{T \in \text{cart}(\tilde{B}^n \times \mathbb{R}^3) \mid \text{spt } T \subset \tilde{B}^n \times S^2, T = G_\varphi \text{ in } (\tilde{B}^n \setminus \bar{B}^n) \times S^2\}.$$

As in Sec. 1, to any $T \in \text{cart}_\varphi(\tilde{B}^n \times S^2)$ we can associate a function $u_T \in BV_\varphi(\tilde{B}^n, S^2)$ such that (1.10) and (1.11) holds true. In particular, if the corresponding BV -function u belongs to $W_\varphi^{1,2}(\tilde{B}^n, S^2)$, we infer that

$$T = G_u + L \times \llbracket S^2 \rrbracket, \quad (6.2)$$

where this time $L \in \mathcal{R}_{n-2}(\tilde{B}^n)$, with $\mathbf{M}(L) < \infty$ and $\text{spt } L \subset \bar{B}^n$. Moreover, the null-boundary condition $\partial T = 0$ applied to the forms $\omega = \pi^\# \phi \wedge \hat{\pi}^\# \omega_{S^2}$ for some $\phi \in \mathcal{D}^{n-3}(\tilde{B}^n)$, yields that

$$\partial L(\phi) = -\frac{1}{4\pi} G_u(\hat{\pi}^\# \omega_{S^2} \wedge \pi^\# d\phi) = -\mathbb{P}(u)(\phi).$$

Therefore, L satisfies the boundary condition

$$\partial L = -\mathbb{P}(u) \quad (6.3)$$

and, conversely, if L is an i.m. rectifiable current in $\mathcal{R}_{n-2}(\tilde{B}^n)$ with $\text{spt } L \subset \bar{B}^n$, hence with finite mass, satisfying (6.3), the corresponding current (6.2) belongs to the class $\text{cart}_\varphi(\tilde{B}^n \times S^2)$. Finally, denote

$$\mathcal{T}_u := \{T \in \text{cart}_\varphi(\tilde{B}^n \times S^2) \mid u_T = u\}, \quad u \in W_\varphi^{1,2}(\tilde{B}^n, S^2), \quad (6.4)$$

so that we have

$$\mathcal{T}_u = \{G_u + L \times \llbracket S^2 \rrbracket : L \in \mathcal{R}_{n-2}(\tilde{B}^n), \text{spt } L \subset \bar{B}^n, \partial L = -\mathbb{P}(u)\}.$$

The following density result was proved in [14].

Theorem 6.1 *Let $n \geq 3$ and let $T \in \text{cart}_\varphi(\tilde{B}^n \times S^2)$ satisfy (6.2) for some $u \in W_\varphi^{1,2}(\tilde{B}^n, S^2)$. There exists a smooth sequence $\{u_h\} \subset C_\varphi^1(\tilde{B}^n, S^2)$ such that $G_{u_h} \rightarrow T$ as $h \rightarrow \infty$ weakly in $\mathcal{D}_n(\tilde{B}^n \times S^2)$ and*

$$\lim_{h \rightarrow \infty} \mathbf{D}(u_h, \tilde{B}^n) = \mathbf{D}(T) := \mathbf{D}(u, \tilde{B}^n) + 4\pi \mathbf{M}(L).$$

Remark 6.2 We notice that by (1.3) we have $m_r(\mathbb{P}(u)) < \infty$, more precisely

$$m_r(\mathbb{P}(u)) \leq C \left(\int_{B^n} |\nabla u|^2 dx + \int_{B^n} |\nabla \varphi|^2 dx \right) \quad \forall u \in W_\varphi^{1,2}(\tilde{B}^n, S^2),$$

where $C > 0$ is an absolute constant, compare (2.7). Also, in the case of dimension $n = 3$, as a consequence of [4] one sees that $\mathbb{P}(u)$ is an integral flat chain, thus by (1.7) we conclude again that *the integral mass $m_i(\mathbb{P}(u))$ is finite for every $u \in W_\varphi^{1,2}(\tilde{B}^3, S^2)$* . This is not trivial, if $n \geq 4$. In fact, in [1] it is proved that for any i.m. rectifiable current $L \in \mathcal{R}_{n-2}(\tilde{B}^n)$, with $\text{spt } L \subset \bar{B}^n$, there exists a function $u \in W_\varphi^{1,2}(\tilde{B}^n, S^2)$ such that $\partial L = -4\pi \mathbb{P}(u)$. Therefore, on account of the counterexamples from [20] and [27], we infer that *if $n \geq 4$, there exist Sobolev maps $u \in W_\varphi^{1,2}(\tilde{B}^n, S^2)$ such that*

$$m_r(\mathbb{P}(u)) < m_i(\mathbb{P}(u)).$$

However, we have

Proposition 6.3 *For any $n \geq 3$ and $u \in W_\varphi^{1,2}(\tilde{B}^n, S^2)$ the class \mathcal{T}_u is non-empty. As a consequence, the integral mass $m_i(\mathbb{P}(u)) < \infty$ and the relaxed energy $\tilde{\mathbf{D}}(u, \Omega) < \infty$.*

PROOF: Denote by $R_{2,\varphi}^\infty(\tilde{B}^n, S^2)$ the set of all the maps $u \in W_\varphi^{1,2}(\tilde{B}^n, S^2)$ which are smooth except on a singular set $\Sigma(u)$ of the type (1.16), where this time Σ_i is a smooth $(n-3)$ -dimensional subset of \bar{B}^n with smooth boundary, if $n \geq 4$, and Σ_i is a point if $n = 3$. Let $\{u_h\} \subset R_{2,\varphi}^\infty(\tilde{B}^n, S^2)$ be such that $u_h \rightarrow u$ strongly in $W^{1,2}$, see [4]. By the "smoothness" of u_h , arguing as in [26], see [2] for the case $n = 3$, we obtain that the integral mass

$$m_i(\mathbb{P}(u_h)) \leq \frac{1}{4\pi}(\mathbf{D}(u_h, B^n) + \mathbf{D}(\varphi, B^n)) < \infty.$$

Let now $T_h := G_{u_h} + L_h \times \llbracket S^2 \rrbracket$, where $L_h \in \mathcal{R}_{n-2}(\tilde{B}^n)$ is such that $\text{spt } L_h \subset \bar{B}^n$, $\partial L_h = -\mathbb{P}(u_h)$, and $\mathbf{M}(L_h) \leq m_i(\mathbb{P}(u_h)) + 1/h$. The sequence $\{T_h\}$ belongs to $\text{cart}_\varphi(\tilde{B}^n \times S^2)$ and has equibounded Dirichlet energies, $\sup_h \mathbf{D}(T_h) < \infty$. By closure-compactness, and since the Dirichlet energy $T \mapsto \mathbf{D}(T)$ is lower semicontinuous with respect to the weak convergence as currents, possibly passing to a subsequence we obtain that $T_h \rightharpoonup T$ weakly in $\mathcal{D}_n(\tilde{B}^n \times S^2)$ to some $T \in \text{cart}_\varphi(\tilde{B}^n \times S^2)$ with $\mathbf{D}(T) < \infty$. As $u_h \rightarrow u$ strongly in $W^{1,2}$, we infer that $T \in \mathcal{T}_u$, whence $m_i(\mathbb{P}(u)) < \infty$. Finally, since \mathcal{T}_u is nonempty, on account of Theorem 6.1 we readily conclude that $\tilde{\mathbf{D}}(u, \Omega) < \infty$. \square

In particular, for every $n \geq 3$ the class $W_\varphi^{1,2}(\tilde{B}^n, S^2)$ agrees with the weak sequential closure of smooth maps in $W_\varphi^{1,2}(\tilde{B}^n, S^2)$, see [23] for a more general result.

Remark 6.4 We point out that the strong $W^{1,2}$ -convergence of "smooth" sequences $\{u_h\} \subset R_{2,\varphi}^\infty(\tilde{B}^n, S^2)$ yields that $\mathbf{M}(\mathbb{D}(u_h) - \mathbb{D}(u)) \rightarrow 0$ and hence, by (1.3), that the real mass $m_r(\mathbb{P}(u_h) - \mathbb{P}(u)) \rightarrow 0$. However, see Remark 6.2, differently to what happens for maps in $W_\varphi^{1,1}(\tilde{B}^n, S^1)$, a part from the easier case $n = 3$, this does not yield that the integral mass $m_i(\mathbb{P}(u_h) - \mathbb{P}(u)) \rightarrow 0$. This is one of the crucial points in the proof of Theorem 6.1.

As a consequence, from the above results we readily obtain the following representation.

Proposition 6.5 *For any $n \geq 3$ and $u \in W_\varphi^{1,2}(\tilde{B}^n, S^2)$, the relaxed energy is given by*

$$\begin{aligned} \tilde{\mathbf{D}}(u, \Omega) &= \inf\{\mathbf{D}(T) \mid T \in \mathcal{T}_u\} \\ &= \mathbf{D}(u, \Omega) + 4\pi \inf\{\mathbf{M}(L) : L \in \mathcal{R}_{n-2}(\tilde{B}^n), \text{spt } L \subset \bar{B}^n, \partial L = -\mathbb{P}(u)\} \\ &= \mathbf{D}(u, \Omega) + 4\pi m_i(\mathbb{P}(u)). \end{aligned}$$

Arguing as in [26], we finally obtain

Corollary 6.6 *For any $n \geq 3$ and $u \in W_\varphi^{1,2}(\tilde{B}^n, S^2)$, we have*

$$\tilde{\mathbf{D}}(u, \Omega) \leq \mathbf{D}(u, \Omega) + \mathbf{D}(u, B^n) + \mathbf{D}(\varphi, B^n).$$

In particular, if φ is constant on ∂B^n we have

$$\tilde{\mathbf{D}}(u, B^n) \leq 2\mathbf{D}(u, B^n).$$

THE RELAXED $W^{1,p}$ -ENERGY. Let us finally consider the case $p > 2$, and introduce the *relaxed $W^{1,p}$ -energy* of maps u in $W_\varphi^{1,p}(\tilde{B}^n, S^2)$, given by

$$\widetilde{\mathcal{E}}_{1,p}(u) := \inf \left\{ \liminf_{h \rightarrow \infty} \int_{\tilde{B}^n} |\nabla u_h|^p dx : \{u_h\} \subset C_\varphi^1(\tilde{B}^n, S^2), \quad u_h \rightarrow u \quad \text{a.e.} \right\}.$$

In [16] it is shown that if p is not an integer, $p \notin \mathbb{Z}$, a map in $W_\varphi^{1,p}(\tilde{B}^n, S^2)$ belongs to the sequential weak $W^{1,p}$ -closure of smooth maps if and only if it belongs to the strong $W^{1,p}$ -closure of smooth maps from \tilde{B}^n into S^2 . Moreover, if $2 < p < 3$ and $u \in W_\varphi^{1,p}(\tilde{B}^n, S^2)$, arguing as in Theorem 5.1 we infer that

$$\widetilde{\mathcal{E}}_{1,p}(u) = \begin{cases} \int_{\tilde{B}^n} |\nabla u|^p dx & \text{if } \mathbb{P}(u) = 0 \\ +\infty & \text{if } \mathbb{P}(u) \neq 0 \end{cases}$$

where $\mathbb{P}(u)$ is given by (6.1).

However, if $p = 3$, the situation is totally different from the one about the relaxed energy of Sobolev maps into S^1 , see the previous section. This is due to the fact that the higher order homotopy groups of the 2-sphere are not all trivial, e.g., $\pi_3(S^2) = \mathbb{Z}$. If $h : S^3 \rightarrow S^2$ is the Hopf map, namely the one that generates the third homotopy group of S^2 , and $u : B^4 \rightarrow S^2$ is given by $u(x) := h(x/|x|)$, then u belongs to $W^{1,p}(B^4, S^2)$ for every $p < 4$ and *the graph of u has no interior boundary*, i.e., $\partial G_u \lrcorner B^4 \times S^2 = 0$. In particular, $\mathbb{P}(u) = 0$. However, the *topological singularity* at the origin is relevant, even if it cannot be treated by means of a homological theory as above, compare [18].

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