COMPACTNESS AND LOWER SEMICONTINUITY IN GSBD

ANTONIN CHAMBOLLE AND VITO CRISMALE

CMAP, École Polytechnique, CNRS, 91128 Palaiseau Cedex, France

Abstract. In this paper we prove a compactness and semicontinuity result in GSBD for sequences with bounded Griffith energy. This generalises classical results in (G)SBV by Ambrosio [1, 2, 3] and SBD by Bellettini-Coscia-Dal Maso [9]. As a result, the static problem in Francfort-Marigo’s variational approach to crack growth [30] admits (weak) solutions.

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1. Introduction

The variational approach to fracture was introduced by Francfort and Marigo in [30] in order to build crack evolutions in brittle materials, following Griffith’s laws [36], without a priori knowledge of the crack path (or surface in higher dimension). It relies on successive minimisations of the Griffith energy:

$$(u, K) \mapsto \int_{\Omega \setminus K} C(e(u)) : e(u) \, dx + \gamma \mathcal{H}^{n-1}(K)$$

where $\Omega \subset \mathbb{R}^n$ is a bounded open set, the reference configuration, $u : \Omega \to \mathbb{R}^n$ is an (infinitesimal) displacement, $e(u)$ its symmetrised gradient (the infinitesimal elastic strain) and $C$ the Cauchy stress tensor defining the Hooke’s law (in particular, $Ca : a$ defines a positive definite quadratic form of the $n \times n$ symmetric tensor $a$). The symmetrised gradient $e(u)$ is defined out of the crack set $K$, which is in the theory a compact $(n-1)$-dimensional set and is penalised by its surface (multiplied by a coefficient $\gamma$ called the toughness).

The minimisation of the energy is under the constraint that $K$ should contain a previously computed crack $K_0$, and that $u$ should satisfy a Dirichlet condition $u = u_0$ on a subset $\partial_D \Omega \setminus K$ of $\partial \Omega$, where $\partial_D \Omega$ is a regular part of the boundary and $u_0$ a sufficiently regular displacement. Hence an important question in the theory is whether the problem

$$\min_{\substack{u = u_0 \text{ on } \partial_D \Omega \setminus K \setminus K_0 \subset K \text{ compact}}} \int_{\Omega \setminus K} C(e(u)) : e(u) \, dx + \gamma \mathcal{H}^{n-1}(K)$$

has a solution.

\textit{E-mail address:} antonin.chambolle@cmap.polytechnique.fr, vito.crismale@polytechnique.edu.

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This problem however is not easy to analyse, since the energy controls very little of the function $u$: for instance if $K$ almost cuts out from $\partial D\Omega$ a connected component of $\Omega$, the function $u$ may have any (arbitrarily large) value in this component at small cost.

From a technical point of view, one cannot take truncations or compositions with bounded transformations to get an \textit{a priori} $L^\infty$ bound for minimisers. In fact, the integrability of $e(u)$ is in general lost by $e(\psi(u))$, unless $\psi(y) = y_0 + \lambda y$, for some $y_0 \in \mathbb{R}^n$, $\lambda \in \mathbb{R}$ (see e.g. the introduction of \cite{21}).

For this reason, most of the “sound” approaches to problem (1.1) consider additional assumptions. In particular, a global $L^\infty$ bound on the displacements ensures one may work in the class $SBD$ of \textit{Special functions with Bounded Deformation} \cite{5}, provided one considers a \textit{weak formulation} of the problem where $K$ is replaced with the intrinsic jump set $J_u$ of $u$ (which needs not to be closed anymore): in this space minimising sequences are shown to be compact \cite{3}, and the energy to be lower semicontinuous. Another possible assumption is, in dimension two

$$K: \text{for instance if } u \in GSBD, \text{ and let } h \to \infty,$$

which needs not to be closed anymore): in this space minimising sequences are shown to be compact \cite{3}, and the energy to be lower semicontinuous. Another possible assumption is, in

$$2d, \text{that the crack set } K \text{ be connected } \cite{25, 12}.$$

The natural space for studying (1.1), in fact, is not $SBD(\Omega)$ (which assumes that the symmetrised gradient of $u$ is a measure and hence $u$ is in $L^{n/(n-1)}(\Omega;\mathbb{R}^n)$) but the space $GSBD(\Omega)$, introduced by Dal Maso in \cite{21}. This space, defined by the slicing properties of the functions, is designed in order to contain “all” displacements $u$ for which the energy is finite. Even if \cite{21} proves compactness under very mild assumptions on the integrability of displacements, no compactness result was available in $GSBD$ for minimizing sequences of (the weak formulation of) (1.1) until very recently.

The first existence result without further constraint has been proven indeed in \cite{35}, in \textit{dimension two}. It relies on a delicate construction showing a \textit{piecewise Korn inequality}, in \cite{33} (for approximated Korn and Korn-Poincaré inequalities see also e.g. \cite{19, 13, 32}, for piecewise rigidity cf. \cite{15}).

In this paper, we prove the following general compactness result for sequences bounded in energy, in the space $GSBD(\Omega)$, in any dimension.

\textbf{Theorem 1.1.} Let $\phi: \mathbb{R}^+ \to \mathbb{R}^+$ be a non-decreasing function with

$$\lim_{t \to +\infty} \frac{\phi(t)}{t} = +\infty,$$

and let $(u_h)_h$ be a sequence in $GSBD(\Omega)$ such that

$$\int_{\Omega} \phi(|e(u_h)|) \, dx + \mathcal{H}^{n-1}(J_{u_h}) < M,$$

for some constant $M$ independent of $h$. Then there exists a subsequence, still denoted by $(u_h)_h$, such that

$$A := \{x \in \Omega: |u_h(x)| \to +\infty\}$$

has finite perimeter, and $u \in GSBD(\Omega)$ with $u = 0$ on $A$ for which

$$u_h \to u \text{ $L^n$-a.e. in } \Omega \setminus A,$$

$$e(u_h) \to e(u) \text{ in } L^1(\Omega \setminus A; \mathbb{R}^{n \times n}_{\text{sym}}),$$

$$\mathcal{H}^{n-1}(J_u \cup \partial^* A) \leq \liminf_{h \to \infty} \mathcal{H}^{n-1}(J_{u_h}).$$

The proof of this theorem is in our opinion simpler than \cite{35}, even if a fundamental tool is a quite technical Korn-Poincaré inequality for functions with small jump set, proved in \cite{15} and employed also in \cite{14, 15, 17}. We combine this inequality with arguments in the spirit of Rellich’s type compactness theorems.

Theorem 1.1 gives then the existence of minimisers for the Griffith energy with Dirichlet boundary conditions in the weak formulation (see Theorem 4.1), which by results in \cite{20, 15} satisfy the properties of strong solutions in the interior of $\Omega$. In the forthcoming paper \cite{16} we prove existence of solutions for the strong formulation (1.1) by extending the regularity theorems in \cite{20, 15} up to the boundary, when $\partial D\Omega$ is of class $C^1$ and $u_0$ is Lipschitz.
The major issue for establishing the compactness result of Theorem 1.1 comes from the lack of control on both the displacement and its full gradient, as is natural in the study of brittle fracture in small strain (linearised) elasticity [36].

A bound such as (1.3) for the full gradient in place of the symmetrised gradient is available for brittle fractures models in finite strain elasticity or in small strain elasticity in the simplified antiplane case (i.e. when the displacement is vertical and depends only on the horizontal components). In these cases, the energy is closely related to the Mumford-Shah functional in image reconstruction [39] (which however includes a fidelity term, artificial from a mechanical standpoint).

In this context, the original strategy of passing through a weak formulation in terms of $u$ was first proposed by De Giorgi and realised by Ambrosio [1] [2] [3] [4], for the existence of weak solutions, and De Giorgi, Carriero, Leaci in [27] (see also e.g. [11] [28]), for the regularity giving the improvement to strong solutions (an alternative approach, where the discontinuity set is the main variable, has been successfully employed in [24] [35]).

Ambrosio’s results are obtained in the space $GSBV$ [26], and have been extended to $GSBD$ by Dal Maso in [21]. In both cases, a control of the values is required to obtain compactness, guaranteeing that the set $A$ in Theorem 1.1 is empty. Without such a control, it is still relatively simple to obtain a $GSBV$ version of Theorem 1.1. For instance, in the scalar case one can consider as in [1] the sequences of truncated functions $u^N_k := \max\{-N, \min\{u_k, N\}\}$ for any integer $N \geq 1$, which are compact in $BV$ and converge up to subsequences. Then, by a diagonal argument, sending then $N$ to $+\infty$, one builds a subsequence $(u^N_k)_h$ which converges a.e. to some $u$, except on a possible set $A$ where it goes to $+\infty$ or $-\infty$. The scalar version of (1.5b) is obtained exactly as in [1] (see in particular [1] Prop. 4.4), considering perturbations $w \in L^1(\Omega)$ with $w = 0$ a.e. in $A$. A slicing argument then is used to show that $A$ has finite perimeter, whose measure is controlled by (1.5c). (Existence for (1.1) is then deduced by considering the limit of a minimising sequence and setting in $A$ the limit function equal to 0, or to any ground state of the elastic energy.)

This strategy however fails in our case since, as already mentioned, the space $GSBD$ is not stable by truncations. The way out to get compactness without any assumption on the displacements is to locally approximate $GSBD$ functions with piecewise infinitesimal rigid motions, by means of the Korn-Poincaré inequality in [13], and use that such motions belong to a finite dimensional space. We then obtain compactness with respect to the convergence in $L^1$-measure, but still, we can not exclude the existence of a set $A$ of points where the limit is not in $\mathbb{R}^n$. A slicing argument then is used to show that $A$ has finite perimeter, whose measure is controlled by (1.5c). (Existence for (1.1) is then deduced by considering the limit of a minimising sequence and setting in $A$ the limit function equal to 0, or to any ground state of the elastic energy.)

A more general (and difficult) approach, for $GSBV^p$, has been proposed by Friedrich in [33]: there, the set $A$ is a priori removed by a careful modification at the level of the minimising sequence, with a control of the energy. Friedrich and Solombrino also prove in [35] existence of quasistatic evolutions in dimension two, extending in that case the antiplane result by Francfort and Larsen in [29], (see [8] for the existence of strong quasistatic evolutions in dimension two, and e.g. [22] [23] for quasistatic evolutions for brittle fractures with finite strain elasticity).

2. NOTATION AND PRELIMINARIES

For every $x \in \mathbb{R}^n$ and $\rho > 0$ let $B_\rho(x)$ be the open ball with center $x$ and radius $\rho$. For $x, y \in \mathbb{R}^n$, we use the notation $x \cdot y$ for the scalar product and $|x|$ for the norm. We denote by $L^n$ and $\mathcal{H}^k$ the $n$-dimensional Lebesgue measure and the $k$-dimensional Hausdorff measure. For any locally compact subset $B$ of $\mathbb{R}^n$, the space of bounded $\mathbb{R}^m$-valued Radon measures on $B$ is denoted by $\mathcal{M}_b(B; \mathbb{R}^m)$. For $m = 1$ we write $\mathcal{M}_b(B)$ for $\mathcal{M}_b(B; \mathbb{R})$ and $\mathcal{M}^+_b(B)$ for the subspace of positive measures of $\mathcal{M}_b(B)$. For every $\mu \in \mathcal{M}_b(B; \mathbb{R}^m)$, its total variation is denoted by $|\mu|(B)$. We write $\chi_E$ for the indicator function of any $E \subset \mathbb{R}^n$, which is 1 on $E$ and 0 otherwise. We call infinitesimal rigid motion any affine function with skew-symmetric gradient. Let us also set $\mathbb{R} := \mathbb{R} \cup \{-\infty, +\infty\}$ and $\mathbb{R}^* := \mathbb{R} \setminus \{0\}$. 
Definition 2.1. Let \( E \subset \mathbb{R}^n, v: E \to \mathbb{R}^m \) an \( \mathcal{L}^n \)-measurable function, \( x \in \mathbb{R}^n \) such that
\[
\limsup_{\varepsilon \to 0^+} \frac{\mathcal{L}^n(E \cap B_\varepsilon(x))}{\varepsilon^n} > 0.
\]
A vector \( a \in \mathbb{R}^n \) is the approximate limit of \( v \) as \( y \) tends to \( x \) if for every \( \varepsilon > 0 \)
\[
\lim_{\varepsilon \to 0^+} \frac{\mathcal{L}^n(E \cap B_\varepsilon(x) \cap \{v-a\} \geq \varepsilon)}{\varepsilon^n} = 0,
\]
and then we write
\[
\operatorname{ap lim}_{y \to x} v(y) = a. \tag{2.1}
\]

Remark 2.2. Let \( E, v, x, \) and \( a \) be as in Definition 2.1 and let \( \psi \) be a homeomorphism between \( \mathbb{R}^m \) and a bounded open subset of \( \mathbb{R}^m \). Then (2.1) holds if and only if
\[
\lim_{\varepsilon \to 0^+} \frac{1}{\varepsilon^n} \int_{E \cap B_\varepsilon(x)} |\psi(v(y)) - \psi(a)| \, dy = 0.
\]

Definition 2.3. Let \( U \subset \mathbb{R}^n \) open, and \( v: U \to \mathbb{R}^m \) be \( \mathcal{L}^n \)-measurable. The approximate jump set \( J_v \) is the set of points \( x \in U \) for which there exist \( a, b \in \mathbb{R}^m \), with \( a \neq b \), and \( v \in \mathbb{S}^{n-1} \) such that
\[
\operatorname{ap lim}_{(y-x) \cdot v > 0, y \to x} v(y) = a \quad \text{and} \quad \operatorname{ap lim}_{(y-x) \cdot v < 0, y \to x} v(y) = b.
\]
The triplet \((a, b, v)\) is uniquely determined up to a permutation of \((a, b)\) and a change of sign of \( v \), and is denoted by \((v^+(x), v^-(x), \nu_v(x))\). The jump of \( v \) is the function defined by \([v](x) := v^+(x) - v^-(x)\) for every \( x \in J_v \). Moreover, we define
\[
J_v^1 := \{ x \in J_v : |[v](x)| \geq 1 \}. \tag{2.2}
\]

Remark 2.4. By Remark 2.2, \( J_v \) and \( J_v^1 \) are Borel sets and \([v]\) is a Borel function. By Lebesgue’s differentiation theorem, it follows that \( \mathcal{L}^n(J_v) = 0 \).

**BV and BD functions.** If \( U \subset \mathbb{R}^n \) open, a function \( v \in L^1(U) \) is a function of bounded variation on \( U \), and we write \( v \in BV(U) \), if \( D_i v \in \mathcal{M}_b(U) \) for \( i = 1, \ldots, n \), where \( D_i v = (D_1 v, \ldots, D_n v) \) is its distributional gradient. A vector-valued function \( v: U \to \mathbb{R}^m \) is in \( BV(U; \mathbb{R}^m) \) if \( v_j \in BV(U) \) for every \( j = 1, \ldots, m \). The space \( BV_{\text{loc}}(U) \) is the space of \( v \in L^1_{\text{loc}}(U) \) such that \( D_i v \in \mathcal{M}_b(U) \) for \( i = 1, \ldots, n \).

A \( \mathcal{L}^n \)-measurable bounded set \( E \subset \mathbb{R}^n \) is a set of finite perimeter if \( \chi_E \) is a function of bounded variation. The reduced boundary of \( E \), denoted by \( \partial^* E \), is the set of points \( x \) in \( \operatorname{supp} D\chi_E \) such that the limit \( \nu_E(x) := \lim_{x \to 0^+} \frac{D\chi_E(B_\varepsilon(x))}{\mathcal{L}^n(B_\varepsilon(x))} \) exists and satisfies \(|\nu_E(x)| = 1\). The reduced boundary is countably \((\mathcal{H}^{n-1}, n - 1)\) rectifiable, and the function \( \nu_E \) is called generalised inner normal to \( E \).

A function \( v \in L^1(U; \mathbb{R}^n) \) belongs to the space of functions of bounded deformation if its distributional symmetric gradient \( E v \) belongs to \( \mathcal{M}_b(U; \mathbb{R}^n) \). It is well known (see [5, 40]) that for \( v \in BD(U) \), \( J_v \) is countably \((\mathcal{H}^{n-1}, n - 1)\) rectifiable, and that
\[
E v = E^s v + E^c v + E^f v, \tag{2.3}
\]
where \( E^s v \) is absolutely continuous with respect to \( \mathcal{L}^n \), \( E^c v \) is singular with respect to \( \mathcal{L}^n \) and such that \(|E^c v|(B) = 0 \) if \( \mathcal{H}^{n-1}(B) < \infty \), while \( E^f v \) is concentrated on \( J_v \). The density of \( E^s v \) with respect to \( \mathcal{L}^n \) is denoted by \( e(v) \), and we have that (see [5, Theorem 4.3] and recall (2.1)) for \( \mathcal{L}^n\text{-a.e.} \ x \in U \)
\[
\operatorname{ap lim}_{y \to x} \frac{(v(y) - v(x) - e(v)(x)(y-x)) \cdot (y-x)}{|y-x|^2} = 0. \tag{2.4}
\]

The space \( SBD(U) \) is the subspace of all functions \( v \in BD(U) \) such that \( E^s v = 0 \), while for \( p \in (1, \infty) \)
\[
SBD^p(U) := \{ v \in SBD(U) : e(v) \in L^p(\Omega; M_{\text{sym}}^{n \times n}), \mathcal{H}^{n-1}(J_v) < \infty \}.
\]
Analogous properties hold for BV, as the countable rectifiability of the jump set and the decomposition of $Dv$, and the spaces $SBV(U;\mathbb{R}^m)$ and $SBV^p(U;\mathbb{R}^m)$ are defined similarly, with $\nabla v$, the density of $D^a v$, in place of $e(v)$. For a complete treatment of BV, SBV functions and BD, SBD functions, we refer to [6] and to [5, 7, 40], respectively.

**GBD functions.** We now recall the definition and the main properties of the space GBD of generalised functions of bounded deformation, introduced in [21], referring to that paper for a general treatment and more details. Since the definition of GBD is given by slicing (differently from the definition of GBV, cf. [26, 21]), we introduce before some notation.

Fixed $\xi \in \mathbb{S}^{n-1} := \{\xi \in \mathbb{R}^n : ||\xi|| = 1\}$, for any $y \in \mathbb{R}^n$ and $B \subset \mathbb{R}^n$ let

$$\Pi^\xi := \{y \in \mathbb{R}^n : y \cdot \xi = 0\}, \quad B^\xi_y := \{t \in \mathbb{R} : y + t\xi \in B\}.$$  

and for every function $v : B \to \mathbb{R}^n$ and $t \in B^\xi_y$ let

$$v^\xi_y(t) := v(y + t\xi), \quad \tilde{v}^\xi_y(t) := v^\xi_y(t) \cdot \xi .$$

**Definition 2.5 ([21]).** Let $\Omega \subset \mathbb{R}^n$ be bounded and open, and $v : \Omega \to \mathbb{R}^n$ be $\mathcal{L}^n$-measurable. Then $v \in GBD(\Omega)$ if there exists $\lambda_v \in \mathcal{M}_b^+(\Omega)$ such that one of the following conditions holds true for every $\xi \in \mathbb{S}^{n-1}$:

(a) for every $\tau \in C^1(\mathbb{R})$ with $-\frac{1}{2} \leq \tau \leq \frac{1}{2}$ and $0 \leq \tau' \leq 1$, the partial derivative $D\xi(\tau(v \cdot \xi)) = D(\tau(v \cdot \xi)) \cdot \xi$ belongs to $\mathcal{M}_b(\Omega)$, and for every Borel set $B \subset \Omega$

$$||D\xi(\tau(v \cdot \xi))||(B) \leq \lambda_v(B);$$

(b) $\tilde{v}^\xi \in BV_{loc}^\xi(\Omega)$ for $H^{n-1}$-a.e. $y \in \Pi^\xi$, and for every Borel set $B \subset \Omega$

$$\int_{\Pi^\xi} \left( ||D\xi^\xi_y||(B^\xi_y \setminus J^1_{\tilde{v}_y}(\xi)) + H^0(B^\xi_y \cap J^1_{\tilde{v}_y}(\xi)) \right) dH^{n-1}(y) \leq \lambda_v(B), \tag{2.5}$$

where $J^1_{\tilde{v}_y} := \{t \in J_{v^\xi_y} : ||\tilde{v}^\xi_y||(t) \geq 1\}$.

The function $v$ belongs to GSB(\Omega) if $v \in GBD(\Omega)$ and $\tilde{v}^\xi \in SBV_{loc}^\xi(\Omega)$ for every $\xi \in \mathbb{S}^{n-1}$ and for $H^{n-1}$-a.e. $y \in \Pi^\xi$.

**GBD(\Omega) and GSB(\Omega) are vector spaces, as stated in [21] Remark 4.6, and one has the inclusions BD(\Omega) \subset GBD(\Omega), SBD(\Omega) \subset GSB(\Omega), which are in general strict (see [21] Remark 4.5 and Example 12.3).** For every $v \in GBD(\Omega)$ the approximate jump set $J_v$ is still countably $(H^{n-1}, n-1)$-rectifiable (cf. [21] Theorem 6.2) and can be reconstructed from the jump of the slices $\tilde{v}^\xi$ ([21] Theorem 8.1). Indeed, for every $C^1$ manifold $M \subset \mathbb{R}^n$ with unit normal $\nu$, it holds that for $H^{n-1}$-a.e. $x \in M$ there exist the traces $v^\pm_M(x), v^\pm_M(x) \in \mathbb{R}^n$ such that

$$\text{ap lim}_{\pm(y-x) \nu(x) > 0, y \to x} v^\pm_M(x) = v^\pm_M(x), \tag{2.6}$$

and they can be reconstructed from the traces of the one-dimensional slices (see [21] Theorem 5.2). Every $v \in GBD(\Omega)$ has an approximate symmetric gradient $e(v) \in L^1(\Omega; M_{sym}^{n \times n})$, characterised by (2.4) and such that for every $\xi \in \mathbb{S}^{n-1}$ and $H^{n-1}$-a.e. $y \in \Pi^\xi$

$$e(v)^\xi \cdot \xi = \nabla \tilde{v}^\xi \cdot \xi \quad \text{in} \quad L^1(\Omega). \tag{2.7}$$

By these properties of slices it follows that, if $v \in GSB(\Omega)$ with $e(v) \in L^1(\Omega; M_{sym}^{n \times n})$ and $H^{n-1}(J_v) < +\infty$, then for every Borel set $B \subset \Omega$

$$H^{n-1}(J_v \cap B) = (2\omega_{n-1})^{-1} \int_{\mathbb{S}^{n-1}} \left( \int_{\Pi^\xi} H^0(J_{\tilde{v}_y} \cap B^\xi_y) dH^{n-1}(y) \right) d\mathcal{H}^{n-1}(\xi) \tag{2.8}$$

and the two conditions in the definition of GSB for $v$ hold for $\lambda_v \in \mathcal{M}_b^+(\Omega)$ such that

$$\lambda_v(B) \leq \int_B |e(v)| dx + H^{n-1}(J_v \cap B), \tag{2.9}$$

for every Borel set $B \subset \Omega$ (cf. also [31] Theorem 1 and [37] Remark 2).
We now recall the following result, proven in \cite[Proposition 2]{13}. Notice that the proposition is therein stated in $SB$, but the proof, which is based on the Fundamental Theorem of Calculus along lines, still holds for $GSBD$, with small adaptations.

**Proposition 2.6** (\cite{13}). Let $Q_r = (-r, r)^n$, $v \in GSBD(Q_r)$, $p \in [1, \infty)$. Then there exist a Borel set $\omega \subset Q_r$ and an affine function $a: \mathbb{R}^n \to \mathbb{R}^n$ with $e(a) = 0$ such that

$$
\mathcal{L}^n(\omega) \leq c r \mathcal{H}^{n-1}(J_v) \quad \text{and} \quad \int_{Q_r \setminus \omega} |v - a|^p \, dx \leq c r^p \int_{Q_r} |e(v)|^p \, dx.
$$

(2.10)

The constant $c$ depends only on $p$ and $n$.

We conclude the section with a technical lemma.

**Lemma 2.7.** Let $E \subset \mathbb{R}^n$ Borel, $v_h: E \to \mathbb{R}^n$ for every $h$, and consider the $n$ sequences $(v_h \cdot e_1)_h$, obtained by taking every component of $v_h$ with respect to the canonical basis of $\mathbb{R}^n$ 

\begin{align*}
&\text{Assume that every $(v_h \cdot e_1)_h$ converges pointwise } \mathcal{L}^n\text{-a.e.} \text{ to a } v: E \to \mathbb{R}, \text{ and that for } \mathcal{L}^n\text{-a.e. } x \in E \text{ there is } i \in \{1, \ldots, n\} \text{ for which } v_i(x) \in [-\infty, \infty]. \text{ Then for } \mathcal{H}^{n-1}\text{-a.e. } \xi \in S^{n-1}
&\quad \{v_h \cdot \xi \to +\infty \} \text{ is } \mathcal{L}^n\text{-a.e. in } E.
\end{align*}

(2.11)

**Proof.** On the sets

$$E_i := \{v_h \cdot e_i \to +\infty\} \cap \left(\limsup_{h \to \infty} (v_h \cdot e_i) < +\infty\right),$$

we have that (2.11) holds for every $\xi \in \{\xi \in S^{n-1}: \xi \neq 0\}$, which is of full $\mathcal{H}^{n-1}$ measure in $S^{n-1}$.

Let us consider the case when there are $m$ components of $v_h$, with $1 < m \leq n$, that we may assume up to a permutation $v_h \cdot e_1, \ldots, v_h \cdot e_m$, such that $\frac{v_h \cdot e_i}{v_h \cdot e_j} \to \xi_{i,j} \in \mathbb{R}^n$ for $1 \leq i < j \leq m$ and $|\frac{v_h \cdot e_i}{v_h \cdot e_j}| \to +\infty$ for $i \in \{1, \ldots, m\}$ and $j \in \{m+1, \ldots, n\}$ (if $m < n$). In this case (2.11) does not hold only for

$$S^{n-1} \cap (1, \xi_{1,2}, \ldots, \xi_{1,m}, 0, \ldots, 0)^\perp,$$

which has dimension $n - 2$. Notice now that for every $m$ for which $m$ components go faster to infinity than the other ones, there is an at most countable collection of $(\xi_{1,2}, \ldots, \xi_{1,m}) \in (\mathbb{R}^n)^{m-1}$ for which $\frac{v_h \cdot e_i}{v_h \cdot e_j} \to \xi_{i,j}$ for $j \in \{2, \ldots, m\}$ on a subset of $E$ of positive $\mathcal{L}^n$ measure. Thus (2.11) holds for every $\xi$ except on an at most countable union of $\mathcal{H}^{n-1}$-negligible sets of $S^{n-1}$. \hfill \Box

3. THE MAIN COMPACTNESS AND LOWER SEMICONtinuity RESULT

In this section we prove Theorem 1.1 the main result of the paper.

**Proof of Theorem 1.1** We divide the proof into three parts: compactness (with respect to the convergence in measure, by means of approximation through piecewise infinitesimal rigid motions), lower semicontinuity, and closure (in $GSBD$).

**Compactness.** For every $k \in \mathbb{N}$ and $z \in (2k^{-1})\mathbb{Z}^n$ we consider the cubes of center $z$

$$q_{k,z} := z + (-k^{-1}, k^{-1})^n.$$

Then $\Omega_k := \Omega \setminus \bigcup_{q_{k,z} \subset \Omega} \overline{q_{k,z}}$ is essentially the union of the cubes which are contained in $\Omega$.

We apply Proposition 2.6 with $p = 1$ in any $q_{k,z} \subset \Omega$, so for $r = k^{-1}$. Then there exist sets $\omega_{k,z}^h \subset q_{k,z}$ with

$$\mathcal{L}^n(\omega_{k,z}^h) \leq ck^{-1}\mathcal{H}^{n-1}(J_{u_h} \cap q_{k,z})$$

(3.1)
and affine functions $a_{k,z}^h : \mathbb{R}^n \to \mathbb{R}^n$, with $\epsilon(a_{k,z}^h) = 0$, such that

$$\int_{q_{k,z}\setminus \omega_{k,z}} |u_k - a_{k,z}^h| \, dx \leq c k^{-1} \int_{q_{k,z}} |\epsilon(u_k)| \, dx. \quad (3.2)$$

The functions $(a_{k,z}^h)_{h \geq 1}$ belong to the finite dimensional space of affine functions. For any sequence of the $i$-th component $(a_{k,z}^h \cdot e_i)_h$, $i = 1, \ldots, n$, we have the following cases:

- it is bounded, and then converges uniformly (up to a subsequence) to an affine function;
- it is unbounded, and then one of the two alternative possibilities below occurs:
  - it converges globally, up to a subsequence, to $+\infty$ or $-\infty$;
  - there is a hyperplane $\{x \cdot \nu = t\}$ ($\nu \in \mathbb{R}^n$, $t \in \mathbb{R}$) and a subsequence such that $a_{k,z}^h(x) \cdot e_i \to +\infty$ if $x \cdot \nu > t$ and $a_{k,z}^h(x) \cdot e_i \to -\infty$ if $x \cdot \nu < t$.

(To see this, consider the bounded sequence $a_{k,z}^h \cdot e_i \|a_{k,z}^h\|$ for any norm $\| \|$ on the space of affine functions, which has converging subsequences.)

Let $\tau$ denote the function tanh (or any smooth, 1-Lipschitz increasing function from $-1$ to 1 with $\tau(0) = 0$). As a consequence we obtain that, up to a subsequence, the function

$$a_k^h(x) := \sum_{q_{z,k} \subset \Omega} a_{k,z}^h(x) \chi_{q_{z,k}}(x)$$

is such that $(\tau(a_k^h \cdot e_i))_h$ converges to some function in $L^1(\Omega_k)$, for any $i = 1, \ldots, n$. Indeed, we have

$$\tau(a_k^h \cdot e_i)(x) = \sum_{q_{z,k} \subset \Omega} \tau(a_{k,z}^h \cdot e_i)(x) \chi_{q_{z,k}}(x),$$

and in any cube $q_{k,z}$ the sequence $(\tau(a_{k,z}^h \cdot e_i))_h$ converges uniformly either to a function valued in $(-1, 1)$, if $(a_{k,z}^h \cdot e_i)_h$ is bounded, or to a function with values $-1$ and 1, attained where the limit of $(a_{k,z}^h \cdot e_i)_h$ is $+\infty$ or $-\infty$, respectively (notice that at this stage $k$ is fixed).

Clearly the subsequences could be extracted from a previous subsequence built at the stage $k - 1$, hence by a diagonal argument, we may assume that for any $k$, $(\tau(a_k^h \cdot e_i))_h$ converges for all $i = 1, \ldots, n$, in $L^1(\Omega_k)$.

We have that for each $i = 1, \ldots, n$, $k \geq 1$, and $l, m \geq 1$,

$$\int_{\Omega} |\tau(u_m \cdot e_i) - \tau(u_l \cdot e_i)| \, dx \leq 2|\Omega \setminus \Omega_k| + \int_{\Omega_k} |\tau(u_m \cdot e_i) - \tau(u_l \cdot e_i)| \, dx$$

$$+ \int_{\Omega_k} |\tau(a_k^m \cdot e_i) - \tau(a_k^l \cdot e_i)| \, dx + \int_{\Omega_k} |\tau(u_l \cdot e_i) - \tau(a_k^l \cdot e_i)| \, dx. \quad (3.3)$$

By construction,

$$\lim_{l,m \to +\infty} \int_{\Omega_k} |\tau(a_k^m \cdot e_i) - \tau(a_k^l \cdot e_i)| \, dx = 0.$$  

On the other hand,

$$\int_{\Omega_k} |\tau(u_m \cdot e_i) - \tau(a_k^m \cdot e_i)| \, dx = \sum_{q_{z,k} \subset \Omega} \int_{q_{z,k} \setminus \omega_{k,z}} |\tau(u_m \cdot e_i) - \tau(a_{k,z}^m \cdot e_i)| \, dx$$

$$\leq \sum_{q_{z,k} \subset \Omega} \left(2|\omega_{k,z}^m| + \int_{q_{z,k} \setminus \omega_{k,z}^m} |u_m - a_{k,z}^m| \, dx\right)$$

$$\leq \frac{2C}{k} \left(H^{n-1}(\mathcal{J}_{u_m}) + \int_{\Omega_k} |\epsilon(u_m)| \, dx\right) \leq \frac{C}{k}.$$

Using that $|\Omega \setminus \Omega_k| \to 0$ as $k \to \infty$, we deduce from (3.3) that $(\tau(u_k \cdot e_i))_h$ is a Cauchy sequence (for each $i$) and therefore converges in $L^1(\Omega)$ to some limit which we denote $\tilde{\tau}$. Up to a further subsequence, we may assume that the convergence occurs almost everywhere.
and, by (1.2) and (1.3), that \((e(u_h))_h\) converges weakly in \(L^1(\Omega; M_{sym}^{n \times n})\). This determines the (sub)sequence \((u_h)_h\) for which we are going to prove the result, fixed from now on. First notice that the set \(A\) defined in (1.4) (in correspondence to the subsequence) is such that \((u_h)_h\) converges pointwise \(L^n\)-a.e. in \(\Omega \setminus A\) to a function with finite values (that is in \(\mathbb{R}^n\)).

We define \(\tilde{u} : \Omega \to (\mathbb{R})^n\) and \(u : \Omega \to \mathbb{R}^n\) such that

\[
\tilde{u} := (\hat{u}^1, \ldots, \hat{u}^n), \quad \text{where } \hat{u}^i = \tau^{-1}(\tilde{\tau}_i); \quad u := \tilde{u} \chi_{\Omega \setminus A},
\]

(3.4)

with the convention that \(\tau^{-1}(\pm 1) = \pm \infty\).

The set \(A\), which coincides with \(\{ x \in \Omega : \tilde{\tau}^i(x) \in \{-\infty, +\infty\} \text{ for some } i \in \{1, \ldots, n\} \}\), is measurable, since \(\hat{u}^i(x) \in \mathbb{R}\) if and only if \(|\tau(\hat{u}^i)| < 1\) and the functions \(\tilde{\tau}_i : \Omega \to [-1, 1]\) are measurable. Since \((u_h)_h\) converges pointwise \(L^n\)-a.e. in \(\Omega \setminus A\) to \(u\) we have that for every \(\xi \in \mathbb{S}^{n-1}\)

\[
u_h \cdot \xi \to u \cdot \xi \quad L^n\text{-a.e. in } \Omega \setminus A.
\]

(3.5)

Notice that we have not extracted further subsequences depending on \(\xi\), and that the limit function \(u\) (equal to \(\tilde{u}\) since we are in \(\Omega \setminus A\)) does not depend on \(\xi\). Eventually, by Lemma 2.7 we have that for \(H^{n-1}\text{-a.e. } \xi \in \mathbb{S}^{n-1}\)

\[|\nu_h \cdot \xi| \to +\infty \quad L^n\text{-a.e. in } A.
\]

(3.6)

**Lower semicontinuity.** Here we prove first (1.5c), which is specific of our approach due to the description of \(A\), and then (1.5b), which follows the lines of [9] Theorem 1.1.

As in [9] Theorem 1.1 (see also [21] Theorem 11.3), we introduce

\[I_y^\xi(u_h) := \int_{\Omega} \phi(||(\hat{u}_h)_y^\xi||) \, d\Omega,
\]

(3.7)

where \((\hat{u}_h)_y^\xi\) is the density of the absolutely continuous part of \(D(\hat{u}_h)_y^\xi\), the distributional derivative of \((\hat{u}_h)_y^\xi ((\hat{u}_h)_y^\xi \in SBV_{loc}(\mathbb{R})^n)\) for every \(\xi \in \mathbb{S}^{n-1}\) and for \(H^{n-1}\text{-a.e. } y \in \mathbb{R}^n\), since \(u_h \in GSBD(\Omega)\). Thus for any \(\xi \in \mathbb{S}^{n-1}\) it holds that

\[
\int_{\mathbb{R}^n} I_y^\xi(u_h) \, dH^{n-1}(y) = \int_{\Omega} \phi(|e(u_h)(x) \xi|) \, dx \leq \int_{\Omega} \phi(|e(u_h)|) \, dx \leq M,
\]

(3.8)

by Fubini-Tonelli’s theorem and (1.3), recalling that \(\phi\) is non-decreasing. Moreover, since \(u_h \in GSBD(\Omega), D\xi(\tau(u_h) \cdot \xi) \in M_{sym}^n(\Omega)\) for every \(\xi \in \mathbb{S}^{n-1}\) and

\[
\int_{\mathbb{R}^n} |D(\tau(u_h) \cdot \xi)_y^\xi||(\Omega)^n_y \, dH^{n-1}(y) = |D_\xi(\tau(u_h) \cdot \xi)||(\Omega) \leq M,
\]

(3.9)

by (2.9) and (1.3). We denote

\[\Pi_y^\xi(u_h) := |D(\tau(u_h) \cdot \xi)|(\Omega)^n_y.
\]

(3.10)

Let \((u_k)_k = (u_h)_h\) be a subsequence of \((u_h)_h\) such that

\[
\lim_{k \to \infty} H^{n-1}(J_{u_k}) = \liminf_{k \to \infty} H^{n-1}(J_{u_h}) < +\infty,
\]

(3.11)

so that, by (2.8), (3.8), and Fatou’s lemma, we have that for \(H^{n-1}\text{-a.e. } \xi \in \mathbb{S}^{n-1}\)

\[
\liminf_{k \to \infty} \int_{\mathbb{R}^n} \left[ \mathcal{H}^0(J_{(\hat{u}_k)_y}^\xi) + \varepsilon(\Pi_y^\xi(u_k) + \Pi_y^\xi(u_h)) \right] \, dH^{n-1}(y) < +\infty,
\]

(3.12)

for a fixed \(\varepsilon \in (0, 1)\). Let us fix \(\xi \in \mathbb{S}^{n-1}\) such that (3.6) and (3.12) hold. Then there is a subsequence \((u_m)_m = (u_{k_m})_m\) of \((u_k)_k\), depending on \(\varepsilon\) and \(\xi\), such that

\[
\lim_{m \to \infty} \int_{\mathbb{R}^n} \left[ \mathcal{H}^0(J_{(\hat{u}_m)_y}^\xi) + \varepsilon(\Pi_y^\xi(u_m) + \Pi_y^\xi(u_m)) \right] \, dH^{n-1}(y) = \liminf_{k \to \infty} \int_{\mathbb{R}^n} \left[ \mathcal{H}^0(J_{(\hat{u}_k)_y}^\xi) + \varepsilon(\Pi_y^\xi(u_k) + \Pi_y^\xi(u_k)) \right] \, dH^{n-1}(y).
\]

(3.13)
Therefore, by (3.13), (3.5), and (3.6), employing Fatou’s lemma, we have that for $H^{n-1}$-a.e. $y \in \Pi^\xi$

$$
\liminf_{m \to \infty} \left[ H^0(J(\check{u}_m)^\xi_y) + \varepsilon (I^\xi_y(u_m) + \Pi^\xi_y(u_m)) \right] < +\infty,
$$

(3.14)

and

$$(\check{u}_m)^\xi_y \to \check{u}_y^\xi \quad \text{L}^1\text{-a.e. in } (\Omega \setminus A)^\xi_y \quad |(\check{u}_m)^\xi_y| \to \infty, \quad \text{L}^1\text{-a.e. in } A_y^\xi,$n

(3.15)

and

$$
\tau(u_m \cdot \xi)^\xi_y \to \tilde{\tau}_y^\xi \quad \text{in } L^1(\Omega_y^\xi),
$$

for a suitable $\tilde{\tau}_y^\xi \in L^1(\Omega_y^\xi)$. Now we employ (3.5), (3.6), and (3.15), (3.16) to get

$$
\left\{ \begin{array}{ll}
\tilde{\tau}_y^\xi = \tau(u \cdot \xi)^\xi_y \quad &\text{L}^1\text{-a.e. in } (\Omega \setminus A_y^\xi) \\
|\tilde{\tau}_y^\xi| = 1 \quad &\text{L}^1\text{-a.e. in } A_y^\xi.
\end{array} \right.
$$

(3.17)

Fixed $y \in \Pi^\xi$ satisfying (3.14) and (3.15), and such that $(\check{u}_m)^\xi_y \in SBV_{loc}(\Omega_y^\xi)$ for every $m$, we extract a subsequence $(u_j) = (u_m)_j$ from $(u_m)_m$, depending also on $y$, for which

$$
\lim_{j \to \infty} \left[ H^0(J(\check{u}_j)^\xi_y) + \varepsilon (I^\xi_y(u_j) + \Pi^\xi_y(u_j)) \right] = \liminf_{m \to \infty} \left[ H^0(J(\check{u}_m)^\xi_y) + \varepsilon (I^\xi_y(u_m) + \Pi^\xi_y(u_m)) \right].
$$

(3.18)

Then by (3.16) we have that

$$
\tau(u_j \cdot \xi)^\xi_y \rightharpoonup \tilde{\tau}_y^\xi \quad \text{in } SBV(\Omega_y^\xi).
$$

(3.19)

In order to describe the set $A$, we consider its slices $A_y^\xi$ and prove that for $H^{n-1}$-a.e. $y \in \Pi^\xi$ $A_y^\xi$ is a finite union of intervals where $\tilde{\tau}_y^\xi$ has either the value 1 or $-1$,

(3.20)

and

$$
\partial A_y^\xi \subset J_{\tilde{\tau}_y^\xi}.
$$

(3.21)

Recalling that $|\tilde{\tau}_y^\xi| < 1$ in $(\Omega \setminus A_y^\xi)$, by (3.17), the property above states that there is a jump each time one passes from values of $\tilde{\tau}_y^\xi$ with absolute value less than 1 to $A_y^\xi$, that is the set where $|\tilde{\tau}_y^\xi| = 1$. In terms of the slices of $u$, one passes from finite to infinite values.

Let us show the claimed properties. Up to considering a subsequence of $(\check{u}_j)_y$, we may assume that for every $j$

$$
H^0(J(\check{u}_j)^\xi_y) = N_y \in \mathbb{N},
$$

namely there is a fixed number $N_y$ of jump points. These points tend to $M_y \leq N_y$ points

$$
\ell_1, \ldots, \ell_{M_y}.
$$

Then (recall that $\Pi^\xi_y(u_j)$ is equibounded in $j$ by (3.18)) for every $l = 1, \ldots, M_y - 1$

$$
\tau(u_j \cdot \xi)^\xi_y \rightharpoonup \tilde{\tau}_y^\xi \quad \text{in } W^{1,1}_{loc}(\ell_l, \ell_{l+1}),
$$

and the convergence above is locally uniform (for the precise representatives). Moreover, since $\Pi^\xi_y(u_j)$ is equibounded again by (3.18), it follows that $x \mapsto (\check{u}_j)^\xi_y(x) - (\check{u}_j)^\xi_y(\bar{x})$ is locally uniformly bounded in $(\ell_l, \ell_{l+1})$, for any choice of $\bar{x} \in (\ell_l, \ell_{l+1})$ (by the Fundamental Theorem of Calculus). Hence for any $l$ we have two alternative possibilities:

- there is $\bar{x} \in (\ell_l, \ell_{l+1})$ such that

$$
\lim_{j \to \infty} (\check{u}_j)^\xi_y(\bar{x}) = \check{u}_y^\xi(\bar{x}) \in \mathbb{R}
$$

(that is $\bar{x} \notin A_y^\xi$), and then $(\check{u}_j)^\xi_y$ converge locally uniformly in $(\ell_l, \ell_{l+1})$ to $\check{u}_y^\xi$;

- for $\mathcal{L}^1$-a.e. $x \in (\ell_l, \ell_{l+1})$,

$$
\lim_{j \to \infty} |(\check{u}_j)^\xi_y(x)| = \infty,
$$

that is $(\ell_l, \ell_{l+1}) \subset A_y^\xi$.
Therefore any \((t_l, t_{l+1})\) is contained either in \((\Omega \setminus A)_y^j\) or in \(A_y^j\). Moreover, in the first case we have that \(\tilde{u}_y^j \in W^{1,1}(t_l, t_{l+1}) \subset L^\infty(t_l, t_{l+1})\). In particular, in this case there is \(\eta \in (0, 1)\) such that

\[
\tilde{x}_y^j(t_l, t_{l+1}) \subset [-1 + \eta, 1 - \eta]. \tag{3.22}
\]

This implies \((3.20)\) and \((3.21)\).

By \((3.18)\), \((3.19)\), \((3.21)\), and since the jump sets of \(\tau(u_j, \xi)_y^j\) and \((\tilde{u}_j)_y^j\) coincide, we deduce, by lower semicontinuity for \(SBV\) functions defined in one-dimensional domains (see \([11\text{ Proposition 4.2}])\), that

\[
\mathcal{H}^0(J_{\tilde{u}_y^j} \cap (\Omega \setminus A)_y^j) + \mathcal{H}^0(\partial A_y^j) \leq \liminf_{n \to \infty} \left[ \mathcal{H}^0(J_{\tilde{u}_m}^j) + \varepsilon(I_y^j(u_m) + H_y^j(u_m)) \right]. \tag{3.23}
\]

We now integrate over \(y \in \Pi^x\) and use Fatou’s lemma with \((3.13)\) to get

\[
\int_{\Pi^x} \left[ \mathcal{H}^0(J_{\tilde{u}_y^j} \cap (\Omega \setminus A)_y^j) + \mathcal{H}^0(\partial A_y^j) \right] d\mathcal{H}^{n-1}(y) \leq \liminf_{k \to \infty} \int_{\Pi^x} \left[ \mathcal{H}^0(J_{\tilde{u}_y}^j) + \varepsilon(I_y^j(u_k) + H_y^j(u_k)) \right] d\mathcal{H}^{n-1}(y) \tag{3.24}
\]

for \(\mathcal{H}^{n-1}\)-a.e. \(\xi \in \mathbb{S}^{n-1}\). In particular we deduce that \(A\) has finite perimeter (cf. \([6\text{ Remark 3.104}]).

We integrate \((3.24)\) over \(\xi \in \mathbb{S}^{n-1}\); by \((2.8)\), \((3.8)\), \((3.9)\), and \((3.11)\) we get

\[
\mathcal{H}^{n-1}(J_u \cap (\Omega \setminus A) + \mathcal{H}^{n-1}(\partial^* A) \leq C M \varepsilon + \liminf_{h \to \infty} \mathcal{H}^{n-1}(J_{u_h}), \tag{3.25}
\]

for a universal constant \(C\). By the arbitrariness of \(\varepsilon\) and the definition of \(u\) we obtain \((1.5c)\).

The property \((1.5b)\) follows by an adaptation of the arguments in \([9\text{ Theorem 1.1}]\) as in \([11\text{ Theorem 11.3}]\) (which employ Ambrosio-Dal Maso’s \([11\text{ Prop. 4.4}])\). We report the proof for the reader’s convenience.

Fatou’s lemma and \((2.8)\) give that for \(\mathcal{H}^{n-1}\)-a.e. \(\xi \in \mathbb{S}^{n-1}\)

\[
\liminf_{h \to \infty} \int_{\Pi^x} \mathcal{H}^0(J_{(\tilde{u}_h)_y}^j \cap \Omega_y^j) d\mathcal{H}^{n-1}(y) < +\infty. \tag{3.26}
\]

In particular there is a basis \(\{\xi_1, \ldots, \xi_n\}\) of \(\mathbb{R}^n\) such that this holds for every \(\xi\) of the form \(\xi = \xi_i + \xi_j\), \(i, j = 1, \ldots, n\). We fix a \(\xi\) of this type, and we find a subsequence \((u_k)_k = (u_{h_k})_k\) of \((u_h)_h\), depending on \(\xi\), such that

\[
\lim_{k \to \infty} \int_{\Pi^x} \mathcal{H}^0(J_{(\tilde{u}_h)_y}^j \cap \Omega_y^j) d\mathcal{H}^{n-1}(y) = \liminf_{h \to \infty} \int_{\Pi^x} \mathcal{H}^0(J_{(\tilde{u}_h)_y}^j \cap \Omega_y^j) d\mathcal{H}^{n-1}(y). \tag{3.27}
\]

For a given \(w \in L^1(\Omega)\) let (recall \((3.7)\) for the definition of \((\hat{u}_k)_y^j\))

\[
III_y^j(u_k, w) := \int_{(\Omega \setminus A)_y^j} |(\hat{u}_k)_y^j - w| \, d\mathcal{H}^{n-1}(y). \tag{3.28}
\]

By \((2.7)\), \((3.3)\) (the sequence \((u_h)_h\) has been fixed before \((3.4)\)), and Fubini-Tonelli’s theorem there is a subsequence \((u_{l_k}) = (u_{h_{l_k}})_k\) of \((u_h)_h\) such that

\[
\lim_{l \to \infty} \int_{\Pi^x} III_y^j(u_l, w) d\mathcal{H}^{n-1}(y) = \liminf_{k \to \infty} \int_{\Omega \setminus A} |\varepsilon(u_l)\xi \cdot \xi - w| \, dx < +\infty. \tag{3.28}
\]
Let us also fix $\varepsilon \in (0,1)$. Again by Fubini-Tonelli’s theorem, there is a subsequence $(u_m)_m = (u_{m,j})_j$ of $(u_{m})_m$, depending on $\xi$, $w$, $\varepsilon$, such that (3.15) holds for $\mathcal{H}^{n-1}$-a.e. $y \in \Pi^\varepsilon$ and

$$
\lim_{m \to \infty} \int_{\Pi^\varepsilon} \Pi^\varepsilon_y(u_m, w) + \varepsilon \left[ \mathcal{H}^0(J(u_m)_y^\varepsilon) + \Pi^\varepsilon_y(u_m) \right] \, d\mathcal{H}^{n-1}(y) = \lim_{i \to \infty} \int_{\Pi^\varepsilon} \Pi^\varepsilon_y(u_i, w) + \varepsilon \left[ \mathcal{H}^0(J(u_i)_y^\varepsilon) + \Pi^\varepsilon_y(u_i) \right] \, d\mathcal{H}^{n-1}(y) .
$$

(3.29)

By (3.8), (3.11), (3.27), (3.28), and Fatou’s lemma, for $\mathcal{H}^{n-1}$-a.e. $y \in \Pi^\varepsilon$

$$
\liminf_{m \to \infty} \left[ \Pi^\varepsilon_y(u_m, w) + \varepsilon \left[ \mathcal{H}^0(J(u_m)_y^\varepsilon) + \Pi^\varepsilon_y(u_m) \right] \right] < +\infty .
$$

(3.30)

Let $y \in \Pi^\varepsilon$ be such that (3.15) and (3.30) hold, and $(\tilde{u}_m)_m \in SBV_{loc}(\Omega^\varepsilon_y)$ for every $m$. We find a subsequence $(u_j)_j = (u_{m,j})_j$ of $(u_{m})_m$, depending also on $y$, for which

$$
\lim_{j \to \infty} \left[ \Pi^\varepsilon_y(u_j, w) + \varepsilon \left[ \mathcal{H}^0(J(u_j)_y^\varepsilon) + \Pi^\varepsilon_y(u_j) \right] \right] = \liminf_{m \to \infty} \left[ \Pi^\varepsilon_y(u_m, w) + \varepsilon \left[ \mathcal{H}^0(J(u_m)_y^\varepsilon) + \Pi^\varepsilon_y(u_m) \right] \right] .
$$

(3.31)

Recalling the form of $A^\varepsilon_y$ and (3.22), we deduce that $(\tilde{u}_j)_j$ converge to $\tilde{u}^\varepsilon_y$ weakly* in $BV(I)$ for any $I$ compactly contained in $(\Omega \setminus A)^\varepsilon_y$, and then $(\tilde{u}_j)_j \rightharpoonup \tilde{u}^\varepsilon_y$ in $L^1((\Omega \setminus A)^\varepsilon_y)$, by (1.2).

Together with (3.31) this gives

$$
\Pi^\varepsilon_y(u, w) \leq \liminf_{m \to \infty} \Pi^\varepsilon_y(u_m, w) \leq \liminf_{m \to \infty} \left[ \Pi^\varepsilon_y(u_m, w) + \varepsilon \left[ \mathcal{H}^0(J(u_m)_y^\varepsilon) + \Pi^\varepsilon_y(u_m) \right] \right] .
$$

Integrating with respect to $y \in \Pi^\varepsilon$, by Fatou’s lemma and (3.28), (3.29) plus the bounds (3.8), (3.9), (3.12), we get

$$
\int_{\Omega \setminus A} |e(u)\xi - w| \leq \liminf_{k \to \infty} \int_{\Omega \setminus A} |e(u_k)\xi - w| \, dx + \varepsilon \left( C M + \liminf_{h \to h_0} \int_{\Pi^\varepsilon} \mathcal{H}^0(J(u_h)_y^\varepsilon \cap \Omega^\varepsilon_y) \, d\mathcal{H}^{n-1}(y) \right) .
$$

By (3.26) and the arbitrariness of $\varepsilon$, we deduce that for all $w \in L^1(\Omega)$,

$$
\int_{\Omega \setminus A} |e(u)\xi - w| \leq \liminf_{k \to \infty} \int_{\Omega \setminus A} |e(u_k)\xi - w| \, dx .
$$

Since the sequence $(e(u_k))_k$ weakly converges in $L^1(\Omega \setminus A; M^{n,n}_{sym})$, then [1], Proposition 4.4] gives

$$
e(u_k)\xi \cdot \xi - w \to e(u)\xi \cdot \xi \quad \text{in} \quad L^1(\Omega \setminus A) ,$$

and by the arbitrariness of $\xi = \xi_2 + \xi_3$ we deduce (1.5).

**Closure.** We now show that the limit function $u$, defined in (3.4), is in $GSBD(\Omega)$.

Employing (2.9) and recalling (1.3), we have that there exist $\lambda u_h \in M^+_b(\Omega)$ such that

$$\lambda u_h(\Omega) \leq M ,$$

and for every $\xi \in S^{n-1}$ and every Borel set $B \subset \Omega$

$$|D_\xi (\tau(u_h \cdot \xi))|(B) \leq \lambda u_h(B) .$$

Let $\tilde{\lambda} \in M^+_b(\Omega)$ be a weak* limit of a subsequence of $(\lambda u_h)_h$, so that $\tilde{\lambda}(\Omega) \leq M$. Notice that

$$D_\xi \tau(u \cdot \xi) \in M^+_b(\Omega) \quad \text{for every} \quad \xi \in S^{n-1}
$$

(3.32)

and

$$|D_\xi \tau(\tilde{u} \cdot \xi)|(B) \leq \tilde{\lambda}(B) =: \lambda u(B) .
$$

(3.33)

for every Borel set $B \subset \Omega$, where $\tilde{\lambda}$ has been defined above. This follows by a slicing procedure and the use of Fatou’s lemma for every $\xi$, to reconstruct at the end $|D_\xi (\tau(u \cdot \xi))|(\Omega)$ from
Existence of weak solutions. Assume \( \Omega \subset \mathbb{R}^n \) be an open, bounded domain for which
\[
\partial \Omega = \partial_D \Omega \cup \partial_N \Omega \cup N,
\]
with \( \partial_D \Omega \) and \( \partial_N \Omega \) relatively open, \( \partial_D \Omega \cap \partial_N \Omega = \emptyset \). \( H^{n-1}(N) = 0 \), \( \partial_D \Omega \neq \emptyset \), and \( \partial(\partial_D \Omega) = \partial(\partial_N \Omega) \). Let \( u_0 \in W^{1,p}(\mathbb{R}^n; \mathbb{R}^n) \) and \( W: \mathbb{R}^{n \times n} \rightarrow [0, \infty) \) be convex, with \( W(0) = 0 \) and
\[
W(\xi) \geq \phi(\|\xi\|) \quad \text{for } \xi \in \mathbb{R}^{n \times n},
(4.1)
\]
where \( \phi \) satisfies \( (1.2) \).

Let \( K_0 \subset \Omega \cup \partial_D \Omega \) be \( (n-1) \)-countably rectifiable with \( H^{n-1}(K_0) < +\infty \), and consider the minimisation problem:
\[
\min_{v \in GSBD(\tilde{\Omega})} \left\{ \int_{\tilde{\Omega}} W(e(v)) \, dx + H^{n-1}(J_v \cup (\partial_D \Omega \cap \{\text{tr}_\Omega v \neq \text{tr}_\Omega u_0\}) \setminus K_0) \right\}.
(4.2)
\]
Notice that, defining \( \tilde{\Omega} := \Omega \cup U \), where \( U \) is an open bounded set with \( U \cap \partial_D \Omega = \partial_D \Omega \), we can recast the problem as
\[
\min_{v \in GSBD(\tilde{\Omega})} \left\{ \int_{\tilde{\Omega}} W(e(v)) \, dx + H^{n-1}(J_v \setminus K_0): v = u_0 \text{ in } \tilde{\Omega} \setminus (\Omega \cup \partial_D \Omega) \right\}.
(4.3)
\]
Then we have the following existence result.

**Theorem 4.1.** Problem (4.3) admits solutions.

**Proof.** Let \( u_h \in GSBD(\tilde{\Omega}) \) with \( u_h = u_0 \) in \( \tilde{\Omega} \setminus (\Omega \cup \partial_D \Omega) \) be the elements of a minimising sequence for (4.3). Observe that the infimum of problem (4.3) is finite, since the functional is nonnegative and \( u_0 \) is an admissible competitor.

Assume for the moment that \( K_0 \) is compact. By \( (4.1) \) the functions \( u_h \) satisfy the hypotheses of Theorem 1.1 with \( \tilde{\Omega} = \Omega \setminus K_0 \), so that there exist \( A \subset \tilde{\Omega} \setminus K_0 \) with finite perimeter and a measurable function \( u: \tilde{\Omega} \setminus K_0 \rightarrow \mathbb{R}^n \) with \( u = 0 \) in \( A \) such that (up to a subsequence)
\[
A = \{ x \in \tilde{\Omega} \setminus K_0 : |u_h(x)| \rightarrow \infty \}, \quad u_h \rightharpoonup u \quad \text{L}^n \text{-a.e. in } \Omega \setminus (K_0 \cup A)
(4.4)
\]
(since \( \mathcal{L}^n(K_0) = 0 \) we could consider just \( \tilde{\Omega} \) above, but we keep \( \tilde{\Omega} \setminus K_0 \) to indicate the set where we apply Theorem 1.1 and
\[
\int_{\tilde{\Omega}} W(e(u)) \, dx + H^{n-1}(J_u \setminus K_0) \leq \liminf_{h \rightarrow \infty} \int_{\tilde{\Omega}} W(e(u_h)) \, dx + H^{n-1}(J_{u_h} \setminus K_0),
\]
Moreover, by \( (4.4) \) and the admissibility condition for \( u_h \) it follows that \( u = u_0 \) in \( \tilde{\Omega} \setminus (\Omega \cup \partial_D \Omega) \), and in particular \( A \) does not intersect \( (\tilde{\Omega} \setminus (\Omega \cup \partial_D \Omega)) \). Since \( W \) is convex, we have lower semicontinuity for the bulk term, and \( u \) solves (4.3). This proves the theorem if \( K_0 \) is compact. Notice that this holds for any other function \( v \) which coincides with \( u \) in \( \tilde{\Omega} \setminus A \) and is set equal to any fixed infinitesimal rigid motion in \( A \), since the energy of \( v \) in \( A \) is null, and then by (1.5) the Griffith energy of \( v \) is less than the the liminf of the energies of \( u_h \).
If $K_0$ is not compact, for any $\varepsilon > 0$ consider $\overline{K}_0 \subset K_0$ with $\mathcal{H}^{n-1}(K_0 \setminus \overline{K}_0) < \varepsilon$. Then, arguing as above for the open set $\Omega \setminus \overline{K}_0 \supset \overline{\Omega} \setminus K_0$, we get still
\[
\int_{\overline{\Omega}} W(e(u)) \, dx \leq \liminf_{h \to \infty} \int_{\overline{\Omega}} W(e(u_h)) \, dx,
\]
and
\[
\mathcal{H}^{n-1}(J_u \setminus K_0) \leq \mathcal{H}^{n-1}(J_u \setminus \overline{K}_0) \leq \liminf_{h \to \infty} \mathcal{H}^{n-1}(J_{u_h} \setminus \overline{K}_0)
\]
\[
\leq \liminf_{h \to \infty} \mathcal{H}^{n-1}(J_{u_h} \setminus K_0) + \varepsilon \leq \liminf_{h \to \infty} \mathcal{H}^{n-1}(J_{u_h} \setminus K_0) + \varepsilon,
\]
since $J_u \setminus K_0 \subset J_{u_h} \setminus \overline{K}_0$ and $J_{u_h} \setminus \overline{K}_0 \subset (J_{u_h} \setminus K_0) \cup (K_0 \setminus \overline{K}_0)$ (cf. also [35, Theorem 2.5]). We conclude since $\varepsilon > 0$ is arbitrary. \hfill $\square$

Remark 4.2. Since, as observed in the proof, a family of minimisers is obtained by adding any fixed infinitesimal rigid motion in $A$ to a given minimiser, we conclude that
\[
\mathcal{H}^{n-1}(\partial^* A \cap \{\text{tr} u = a\}) = 0
\]
for every infinitesimal rigid motion $a$ ($a(x) = a \cdot x + b$, $a + a^T = 0$), where tr denotes here the trace of $u$ on $\partial^* A$ (which is $(n-1)$-countably rectifiable) from $\Omega \setminus A$.

Existence of strong solutions. In recent works [20, 15], Chambolle, Conti, Focardi, and Iurlano have shown more regularity for the possible minimisers of (4.3) (or (4.2)) if $W(\xi) = C e(\xi) : e(\xi)$ (in [15]), or $n = 2$ and
\[
W(\xi) = f_\mu(\xi) := \frac{1}{p} \left((C\xi : \xi + \mu)^{p/2} - \mu^{p/2}\right)
\]
(4.5)
in [20], requiring that $C: M^{n \times n}_{sym} \to M^{n \times n}_{sym}$ is a symmetric linear map with
\[
C(\xi - \xi^T) = 0 \text{ and } C\xi : \xi \geq c_0|\xi + \xi^T|^2 \text{ for all } \xi \in M^{n \times n}_{sym}.
\]
More precisely, the essential closedness of the jump set is established:

Theorem 4.3. Let $K_0 \subset \Omega \cup \partial D \Omega$ closed, with $\mathcal{H}^{n-1}(K_0) < +\infty$, and $u \in GSBD^2(\Omega \setminus K_0)$ (or $u \in GSBD^p(\Omega \setminus K_0)$, if $\Omega \subset \mathbb{R}^2$) be a minimiser of
\[
\int_{\Omega} C e(v) : e(v) \, dx + \mathcal{H}^{n-1}(J_u \cup (\partial D \Omega \cap \{\text{tr} v \neq \text{tr} u \}) \setminus K_0)
\]
(4.6)
(a minimiser of (4.3) with (4.5), respectively). Then
\[
\mathcal{H}^{n-1}((\Omega \setminus K_0) \cap (J_u \setminus J_a)) = 0, \quad u \in C^1(\Omega \setminus (K_0 \cup J_u))
\]
In [16], this is extended to $\Omega \cup \partial D \Omega$, yielding the following result (see [8] for the SBV case):

Theorem 4.4. Let $\partial D \Omega$ be of class $C^1$, $u_0 \in W^{1,\infty}(\mathbb{R}^n; \mathbb{R}^n)$, and $u \in GSBD^2(\Omega \setminus K_0)$, be a minimiser of (4.6). Then
\[
\mathcal{H}^{n-1}(((\Omega \cup \partial D \Omega) \setminus K_0) \cap (J_u \setminus J_a)) = 0, \quad u - u_0 \in C^1((\Omega \cup \partial D \Omega) \setminus (K_0 \cup J_u))
\]

Another consequence of Theorem 4.1 is a compactness result for phase-field approximations of (1.1), which are used for the numerical simulations of evolutions in brittle fracture (such as in [10]), see [17] for details.

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REFERENCES


