

1 **Maximal monotone operator in non-reflexive Banach space and the**
2 **application to thin film equation in epitaxial growth on vicinal surface**

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7 **Keywords** Fourth-order degenerate parabolic equation, Radon measure, global strong solution,
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10 **Abstract** In this work we consider

$$w_t = [(w_{hh} + c_0)^{-3}]_{hh}, \quad w(0) = w^0, \quad (1) \quad \{\text{abs}\}$$

11 which is derived from a thin film equation for epitaxial growth on vicinal surface. We establish a general
12 framework for abstract evolution equations with maximal monotone operators in non-reflexive Banach
13 space which can be used in a wide class of degenerate parabolic equations. Following this framework,
14 we formulate a global strong solution to (1) which allows a Radon measure occurring. Then we prove
15 the existence of such a global strong solution and obtain the almost everywhere positivity of $w_{hh} + c_0$.

16 **1 Introduction**

17 **1.1 Background and motivation.**

18 Below the roughening transition temperature, the crystal surface is not smooth and forms steps, terraces
19 and adatoms on the substrate, which form solid films. Adatoms detach from steps, diffuse on the

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terraces until they meet other steps and reattach again, which lead to a step flow on the crystal surface. The evolution of individual steps is described mathematically by the Burton-Cabrera-Frank (BCF) type discrete models [5]. Although discrete models do have the advantage of reflecting physical principle directly, when we study the evolution of crystal growth from macroscopic view, continuum approximation for the discrete models involves fewer variables than discrete models and can briefly show the evolution of step flow. Many interesting continuum models can be found in the literature on surface morphological evolution; see [20, 23, 7, 25, 26, 22, 18, 6, 11] for one dimensional models and [17, 27] for two dimensional models. KOHN clarified the evolution of surface height from the thermodynamic viewpoint in the book [14]. He considered the classical surface energy, which dates back to the pioneering work of Mullins [19], given by

$$F(h) := \int_{\Omega} (\beta_1 |\nabla h| + \beta_3 |\nabla h|^3) dx, \quad (2)$$

where Ω is the “step locations area” we are concerned with. Then, by conservation of mass, we have the equation for surface height h

$$\begin{aligned} h_t &= -\nabla \left(M(\nabla h) \nabla \frac{\delta F}{\delta h} \right) \\ &= -\nabla \left(M(\nabla h) \nabla \left(\nabla \cdot \left(\beta_1 \frac{\nabla h}{|\nabla h|} + \beta_3 |\nabla h| \nabla h \right) \right) \right), \end{aligned} \quad (3)$$

where $M(\nabla h)$ is a suitable “mobility” term depending on the dominating process of surface motion. Often two limit cases are considered. For diffusion-limited (DL) case, the dominated dynamics is diffusion across the terraces, we have $M(\nabla h) = 1$; while for attachment-detachment-limited (ADL) case, the dominating processes are the attachment and detachment of atoms at steps edges, and $M(\nabla h) = \frac{1}{|\nabla h|}$. In the DL regime, [13] obtained a fully understanding of the evolution and proved the finite-time flattening. However, in the ADL regime, due to the difficulty brought by mobility term $M(\nabla h) = \frac{1}{|\nabla h|}$, the dynamics of the solution to surface height equation (3), with both $\beta_1 = 0$ or $\beta_1 \neq 0$, is still an open question (see for instance [14]).

Although a general surface may have peaks and valleys, the analysis of step motion on the level of continuous PDE is complicated and we focus on a simpler situation first: a monotone one-dimensional step train. In this simpler case, $\beta_1 = 0$, and by taking $\beta_3 = \frac{1}{2}$, (3) becomes

$$h_t = - \left[\frac{1}{h_x} (3h_x h_{xx}) \right]_x. \quad (4)$$

OZDEMIR, ZANGWILL [20] and AL HAJJ SHEHADEH, KOHN AND WEARE [1] realized that using the step slope as a new variable is a convenient way to study the continuum PDE model, i.e.,

$$u_t = -u^2 (u^3)_{hhh}, \quad u(0) = u_0, \quad (5)$$

where u , considered as a $[0, 1)$ -periodic function of the step height h , is the step slope of the surface. [11] provided a method to rigorously obtain the convergence rate of discrete model to its corresponding continuum approximation.

Two questions then arise. One is how to formulate a proper solution to (5) and prove the well-posedness of its solution. The other one is the positivity of the solution. More explicitly, we want to know whether the sign of the solution u to (5) is persistent. Our goal in this work is to validate the continuum slope PDE (5) by answering the above two questions. The equation (5) is a degenerate equation and we cannot prevent u from touching zero, where singularity arise. We observe that we are able to rewrite (5) as an abstract evolution equation with maximal monotone operator using $\frac{1}{u}$. However, the main difficulty is that we have to work in a non-reflexive Banach space L^1 , which does not possess weak compactness, so the classical theorem for maximal monotone operators in reflexive Banach space

cannot be applied directly. In fact, due to the loss of weak compactness it is natural to allow a Radon measure being our solution $\frac{1}{u}$ and we do observe the singularity when u approaches zero in numerical simulations [10]. Also see [15] for an example where a measure appears in the case of an exponential nonlinearity. Therefore, we devote ourselves to the establishment of a general abstract framework for problems associated with nonlinear monotone operators in non-reflexive Banach spaces and to solve our problem (5) by the abstract framework. Furthermore, the established abstract framework can be applied to a wide class of degenerate parabolic equations which can be recast as an abstract evolution equation with maximal monotone operator in some non-reflexive Banach space. We refer [2] for deep study of monotone operators in metric space and the existence of curves of maximal slope. However, we do not have λ -convex functionals in our problem. Let us introduce our framework in the next section.

1.2 Formal observations and abstract setup

Denote by $\varphi(h, t)$ as the step location when considered as a function of surface height h . Formally, we have

$$u(h, t) = h_x(\varphi(h, t), t) = \frac{1}{\varphi_h(h, t)},$$

and the u -equation (5) can be rewritten as φ -equation

$$\varphi_t = \left(\frac{1}{\varphi_h^3} \right)_{hhh}; \quad (6) \quad \{\text{eq:phin}\}$$

for further details we refer to the appendix of [12].

Motivated by the φ -equation, we want to recast (5) as an abstract evolution equation. If u has a positive lower-bound $u \geq \alpha > 0$, then (5) can be rewritten as

$$\left(\frac{1}{u} \right)_t = (u^3)_{hhh}, \quad u(0) = u_0. \quad (7) \quad \{\text{nueq}\}$$

Formally, if we take $w_{hh} = \frac{1}{u}$, then we have

$$w_t = (w_{hh}^{-3})_{hh}.$$

Since our problem (7) is in 1-periodic setting, i.e., one period $[0, 1)$, we also want w to be periodic. Denote by \mathbb{T} the $[0, 1)$ -torus. For measure space, we can define periodic distributions as distributions on \mathbb{T} , i.e., linear functionals on $C(\mathbb{T})$. Let the $[0, 1)$ -periodic function w be the solution of the Laplace equation

$$w_{hh} = \frac{1}{u} - c_0, \quad \int_{\mathbb{T}} w \, dh = 0,$$

with compatibility condition

$$\int_{\mathbb{T}} w_{hh} \, dh = \int_{\mathbb{T}} \frac{1}{u} - c_0 \, dh = 0.$$

If (7) holds a.e., then we have

$$\int_{\mathbb{T}} \frac{1}{u} \, dh \equiv \int_{\mathbb{T}} \frac{1}{u_0} \, dh =: c_0$$

due to the periodicity of u . However, we cannot show that (7) holds almost everywhere. Actually, the possible existence of singular part for w_{hh} or $\frac{1}{u}$ is intrinsic, since the equation (5) becomes degenerated when u approaches zero. We cannot prevent u from touching zero, and can only show $w_{hh} = \frac{1}{u} - c_0 \in$

$\mathcal{M}(\mathbb{T})$, where \mathcal{M} is the set of finitely additive, finite, signed Radon measures. Hence the compatibility condition becomes

$$\int_{\mathbb{T}} d\left(\frac{1}{u} - c_0\right) = 0.$$

Therefore, in this paper we consider the parabolic evolution equation

$$w_t = [(w_{hh} + c_0)^{-3}]_{hh}, \quad w(0) = w^0, \quad (8) \quad \{\text{maineq}\}$$

under the assumption w is periodic with period \mathbb{T} and has mean value zero in one period, i.e., $\int_{\mathbb{T}} w \, dh = 0$.

For $1 \leq p < \infty$, $k \in \mathbb{Z}$, set

$$\begin{aligned} W_{\mathbb{T}_0}^{k,p}(\mathbb{T}) &:= \{u \in W^{k,p}(\mathbb{T}); u(h) = u(h+1), \text{ a.e. and } u \text{ has mean value zero in one period}\}, \\ L_{\mathbb{T}_0}^p(\mathbb{T}) &:= \{u \in L^p(\mathbb{T}); u(h) = u(h+1), \text{ a.e. and } u \text{ has mean value zero in one period}\}. \end{aligned}$$

Standard notations for Sobolev spaces are assumed above. If $k < 0$ and $(1/p) + (1/q) = 1$, $1 \leq p < +\infty$, then it can be shown that $W_{\mathbb{T}_0}^{k,q}(\mathbb{T})$ is the dual of $W_{\mathbb{T}_0}^{-k,p}(\mathbb{T})$.

Our main functional spaces will be

$$V := \{v \in W_{\mathbb{T}_0}^{2,1}(\mathbb{T})\}, \quad (9) \quad \{\text{space V}\}$$

and

$$\tilde{V} := \{u \in W_{\mathbb{T}_0}^{1,2}(\mathbb{T}); u_{hh} \in \mathcal{M}(\mathbb{T})\}. \quad (10)$$

Define also

$$U := \{v \in L_{\mathbb{T}_0}^2(\mathbb{T})\}. \quad (11) \quad \{\text{space U}\}$$

Endow U and V with the norms $\|u\|_U := \|u\|_{L^2(\mathbb{T})}$ and $\|v\|_V := \|v_{hh}\|_{L^1(\mathbb{T})}$ respectively. Note that the zero-mean conditions for functions of V give the equivalence between $\|\cdot\|_V$ and $\|\cdot\|_{W^{2,1}(\mathbb{T})}$. Note also that the embeddings $V \hookrightarrow U \hookrightarrow V'$ are all dense and continuous.

The space \tilde{V} .

Note also that any \mathbb{T} -periodic function w who has mean value zero such that w_{hh} is a finite Radon measure will belong to $W_{\mathbb{T}_0}^{1,2}(\mathbb{T})$, since the first derivative w_h is a BV function (the total variation of w_h is exactly the total mass of w_{hh}). Thus we can endow the space \tilde{V} with the norm

$$\|w\|_{W^{1,2}(\mathbb{T})} + \|w_{hh}\|_{\mathcal{M}(\mathbb{T})}.$$

Since w is periodic and has mean value zero, we have

$$\|w_h\|_{L^2(\mathbb{T})} \leq \|w_{hh}\|_{\mathcal{M}(\mathbb{T})}, \quad \|w\|_{L^2(\mathbb{T})} \leq \|w_h\|_{L^2(\mathbb{T})}.$$

So we can use the equivalent norm

$$\|w\|_{\tilde{V}} := \|w_{hh}\|_{\mathcal{M}(\mathbb{T})} = \sup_{f \in C(\mathbb{T}), |f| \leq 1} \int_{\mathbb{T}} f \, dw_{hh}.$$

The weak $*$ -convergence on \tilde{V} is then characterized as: a sequence w^n converges weakly- $*$ to w in \tilde{V} if w^n converges weakly to w in $W^{1,2}(\mathbb{T})$, and w_{hh}^n converges weakly $*$ - to w_{hh} in $\mathcal{M}(\mathbb{T})$, i.e.

$$\int_{\mathbb{T}} f \, dw_{hh}^n \rightarrow \int_{\mathbb{T}} f \, dw_{hh} \quad \text{for any } f \in C(\mathbb{T}).$$

83 Relations between V and \tilde{V} .

84 Since V is not reflexive, we first present a characterization of the bidual space V'' . For any $v \in V$, we
85 have $v_{hh} \in L^1_{\mathbb{T}_0}(\mathbb{T})$. Since also $C(\mathbb{T}) \subseteq L^\infty(\mathbb{T})$, we have:

- 86 (i) the dual space $V' = \{u \in (W_{\mathbb{T}_0}^{2,1}(\mathbb{T}))'\} = W_{\mathbb{T}_0}^{-2,\infty}(\mathbb{T})$;
87 (ii) for any $\xi \in V'$, $\eta \in V$, from the Riesz representation, there exists $\bar{\xi} \in L^\infty(\mathbb{T})$ such that

$$\langle \xi, \eta \rangle_{V',V} := \int_{\mathbb{T}} \bar{\xi} \eta_{hh} \, dh,$$

87 and we denote $\bar{\xi}_{hh}$ as ξ without risk of confusion;

- 88 (iii) the bidual space V'' is a subspace of \tilde{V} .

Thus

$$V \subseteq V'' \subseteq \tilde{V} \subseteq W_{\mathbb{T}_0}^{1,2}(\mathbb{T}) \subseteq U.$$

89 Therefore, we conclude that the canonical embedding $V \hookrightarrow V'' \hookrightarrow \tilde{V} \hookrightarrow W_{\mathbb{T}_0}^{1,2}(\mathbb{T}) \hookrightarrow U$ is continuous
90 and each one is a dense subset of the next, since V is dense in U .

91 Observation 1.

92 From (8), one formal observation is that if we set

$$\phi(w) := \frac{1}{2} \int_{\mathbb{T}} (w_{hh} + c_0)^{-2} \, dh, \quad (12)$$

then

$$w_t = -\frac{\delta\phi}{\delta w} = [(w_{hh} + c_0)^{-3}]_{hh}$$

93 is the first variation of ϕ ; see exact definition and calculation in (18) and Proposition 15. Hence we have

$$\frac{d\phi}{dt} = \int_{\mathbb{T}} \frac{\delta\phi}{\delta w} w_t \, dh = \int_{\mathbb{T}} -w_t^2 \, dh = \int_{\mathbb{T}} [(w_{hh} + c_0)^{-3}]_{hh}^2 \, dh \leq 0. \quad (13)$$

94 Besides, we also notice that $\phi(w) = \frac{1}{2} \int_{\mathbb{T}} (w_{hh} + c_0)^{-2} \, dh$ is a convex functional. Recall that the sub-
95 differential of a proper, convex, lower-semicontinuous function is a maximal monotone operator (see
96 for instance [3]), which gives us the idea of using maximal monotone operator to formally rewrite our
97 problem (8), i.e.,

$$w_t = -\partial\phi(w). \quad (14)$$

98 Observation 2.

99 Set also

$$E(w) := \frac{1}{2} \int_{\mathbb{T}} [(w_{hh} + c_0)^{-3}]_{hh}^2 \, dh = \int_{\mathbb{T}} w_t^2 \, dh; \quad (15)$$

100 see exact definition in Definition 3. Taking the derivative ∂_{hh} on the both side of (8), we have

$$[w_{hh} + c_0]_t = [(w_{hh} + c_0)^{-3}]_{hhhh}. \quad (16) \quad \{\text{whh}\}$$

Then another formal observation is that

$$\begin{aligned} \frac{dE(w)}{dt} &= \int_{\mathbb{T}} [(w_{hh} + c_0)^{-3}]_{hh} [(w_{hh} + c_0)^{-3}]_{hht} dh \\ &= \int_{\mathbb{T}} [(w_{hh} + c_0)^{-3}]_{hhhh} [(w_{hh} + c_0)^{-3}]_t dh = \int_{\mathbb{T}} [w_{hh} + c_0]_t [(w_{hh} + c_0)^{-3}]_t dh \\ &= \int_{\mathbb{T}} -3 \frac{[(w_{hh} + c_0)_t]^2}{(w_{hh} + c_0)^4} dh \leq 0. \end{aligned}$$

101 We point out the dissipation of $E(w)$ is important for the proof of existence result.

102 Observation 3.

Moreover, to ensure the surjectivity of the maximal monotone operator $\partial\phi$, we need to find a proper invariant ball. Another formal observation from (16) is that

$$\frac{d}{dt} \int_{\mathbb{T}} (w_{hh} + c_0) dh = \int_{\mathbb{T}} [(w_{hh} + c_0)^{-3}]_{hhhh} dh = 0.$$

So for a constant C depending only on the initial value w_0 , $\{\|w\|_V \leq C\}$ could be an invariant ball provided $w_{hh} + c_0 > 0$ almost everywhere. But note that V is not a reflexive space and that bounded sets in $L^1(\mathbb{T})$ do not have any compactness property. Actually we only obtain

$$w_{hh} + c_0 > 0, \text{ a.e. } (t, h) \in [0, T] \times \mathbb{T},$$

$$\int_{\mathbb{T}} d(w_{hh} + c_0) \leq C, \text{ for any } t \geq 0,$$

103 and choose $\{\|w\|_{\tilde{V}} \leq C\}$ to be the invariant ball. That is consistent with the prediction that $w_{hh} = \frac{1}{u} - c_0$
104 is possible to be a Randon measure.

105 After those formal observations, in order to rewrite our problem as an abstract problem precisely,
106 we introduce the following definition.

107 **Definition 1** For any $w \in \tilde{V}$, from [8, p.42], we have the decomposition

$$\{\text{decom}\} \quad w_{hh} = \eta + \nu \quad (17)$$

108 with respect to the Lebesgue measure, where $\eta \in L^1(\mathbb{T})$ is the absolutely continuous part of w_{hh} and ν
109 is the singular part, i.e., the support of ν has Lebesgue measure zero. Recall c_0 is a constant in (8).

110 Denote $g := \eta + c_0$. Then $w_{hh} + c_0 = g + \nu$ and $g \in L^1(\mathbb{T})$ is the absolutely continuous part of $w_{hh} + c_0$.

111 Define the proper, convex functional

$$\phi : \tilde{V} \longrightarrow \mathbb{R} \cup \{+\infty\}, \quad \phi(w) := \begin{cases} \int_0^1 \Phi(g) dh & \text{if } w_{hh} + c \in \mathcal{M}^+(\mathbb{T}), \\ +\infty & \text{otherwise,} \end{cases} \quad \Phi(x) := \begin{cases} +\infty & \text{if } x \leq 0, \\ x^{-2}/2 & \text{if } x > 0, \end{cases} \quad (18)$$

\{\text{defphi}\}

112 where $g \in L^1(\mathbb{T})$ is the absolutely continuous part of $w_{hh} + c_0$. For some constant $C > 0$ large enough,
113 define the proper, convex functional

$$\{\text{defpsi}\} \quad \psi : \tilde{V} \longrightarrow \{0, +\infty\}, \quad \psi(w) := \begin{cases} 0 & \text{if } \|w\|_{\tilde{V}} \leq C, \\ +\infty & \text{if } \|w\|_{\tilde{V}} > C. \end{cases} \quad (19)$$

The domain of $\phi + \psi$ is

$$D(\phi + \psi) := \{w \in \tilde{V}; (\phi + \psi)(w) < +\infty\} \subseteq \tilde{V} \cap \{\|w\|_{\tilde{V}} \leq C\}.$$

114 Note that $\mathcal{K} := \{w \in \tilde{V} : \|w\|_{\tilde{V}} \leq C\}$ is closed and convex, hence its indicator (i.e., ψ) is convex,
115 lower-semicontinuous and proper. Later, we will determine the constant C by initial data w_0 .

116 Now we can state two definitions of solutions we study in this work.

Definition 2 Given ϕ, ψ defined in Definition 1, for any $T > 0$, we call the function

$$w \in L^\infty([0, T]; \tilde{V}) \cap C^0([0, T]; U), \quad w_t \in L^\infty([0, T]; U)$$

117 a variational inequality solution to (8) if it satisfies

$$\langle w_t, v - w \rangle_{U', U} + (\phi + \psi)(v) - (\phi + \psi)(w) \geq 0 \quad (20) \quad \{\text{vi}\}$$

118 for a.e. $t \in [0, T]$ and all $v \in \tilde{V}$.

119 **Definition 3** For any $T > 0$, let $\eta \in L^1(\mathbb{T})$ be the absolutely continuous part of w_{hh} in (17). Define

$$E(t) := \frac{1}{2} \int_{\mathbb{T}} [((\eta + c_0)^{-3})_{hh}]^2 dh. \quad (21)$$

We call the function

$$w \in L^\infty([0, T]; \tilde{V}) \cap C^0([0, T]; U), \quad w_t \in L^\infty([0, T]; U)$$

120 a strong solution to (8) if

121 (i) it satisfies

$$w_t = [(\eta + c_0)^{-3}]_{hh} \quad (22) \quad \{\text{s1u}\}$$

122 for a.e. $(t, h) \in [0, T] \times \mathbb{T}$;

123 (ii) we have $((\eta + c_0)^{-3})_{hh} \in L^\infty([0, T]; U)$ and the dissipation inequality

$$E(t) = \frac{1}{2} \int_{\mathbb{T}} [((\eta + c_0)^{-3})_{hh}]^2 dh \leq E(0). \quad (23)$$

124 The main result in this work is to prove existence of the variational inequality solution and strong
125 solution to (8), which is stated in Theorem 14 and Theorem 16 separately.

126 1.3 Overview of our method and related method

127 The key of our method is to rewrite the original problem as an abstract evolution equation $w_t = -\tilde{B}w$,
128 where \tilde{B} is the sub-differential of a proper, convex, lower semi-continuous function, i.e. $\tilde{B} = \partial(\phi + \psi)$.
129 \tilde{B} is a maximal monotone operator by classical results (see for instance [3]). ψ is the indicator of the
130 invariant ball \mathcal{K} in (19). By constructing the proper invariant ball \mathcal{K} , we also obtain the restriction of \tilde{B}
131 to $L^2(\mathbb{T})$ is also a maximal monotone operator; see Lemma 11. Notice the definition of the functional
132 ϕ involves only the absolutely continuous part of w_{hh} , so we need to prove that it is still lower semi-
133 continuous on \tilde{V} ; see details in Proposition 7. Then by standard theorem for m-accretive operator (see
134 Definition 6) in [3], we can prove the variational inequality solution to (8) in Theorem 14. Another key
135 point is to prove the multi-valued operator $\partial(\phi + \psi)$ is actually single valued, which concludes that
136 the variational inequality solution is also the strong solution defined in Definition 3. In order to prove
137 $\partial(\phi + \psi)$ is single valued, we need to first determine $\partial\phi$ in Proposition 15, then make two *a-priori*

assumptions and finally verify them using those dissipation observations in Section 1.2; see details in Theorem 16.

Actually, our definitions for variational inequality solution and strong solution in Definitions 2 and 3 hide a Radon measure in it. As we said before, this kind of fourth order degenerate equation has the intrinsic property of singular measure. We want to mention that [9] also used maximal monotone operator method for diffusion limited (DL) case. However, due to the mobility for DL model is $M = 1$ rather than $M = \frac{1}{h_x}$ in our case, DL model can be recast as an abstract evolution equation with maximal monotone operator using the anti-derivative of h . The coercivity of the this maximal monotone operator in DL case is natural and hence the operated space is a reflexive space. It is much easier than our case and singular part will not appear.

Recently, [12,16] also analyzed the positivity and the weak solution to the same equation (5) separately. They all considered this nonlinear fourth order parabolic equation, which comes from the same step flow model on vicinal surface. The aim is to answer the two questions in Section 1.1, which also are stated as open questions in [14]. The nonlinear structure of this equation, the key for both previous and current works, is important for the positivity of solution because we know that the sign changing is a general property for solutions to linear fourth order parabolic equations. For one dimensional case, following the regularized method in [4], [12] defined the weak solution on a subset, which has full measure, of $[0, 1]$ and proved positivity and existence. Using the method of approximating solutions, based on the implicit time-discretization scheme and carefully chosen regularization, [16] expanded the result in [12] to high dimensional case. Our results is consistent with theirs, but uses a totally different approach. Paper [16]'s method is delicate and subtle while our method seems to be more natural and we obtain the variational inequality solution to (8).

We point out that our method establishes a general structure for this kind of equation whose invariant ball exists in a non-reflexive Banach space. We believe this method can be applied to many similar degenerated problems as long as they can be reduced to an abstract evolution equation with maximal monotone operator which unfortunately is a non-reflexive Banach space.

The rest of this work is devoted to first recall some useful definitions in Section 1.4. Then in Section 2, we rigorously study the sub-differential $\partial(\phi + \psi)$ and prove it is m-accretive on U , which leads to the existence result for variational inequality solution. In Section 3, we calculate the exact value of $\partial(\phi + \psi)$ and prove the variational inequality solution is actually a strong solution.

1.4 Preliminaries

In this section, we first recall the following classical definitions (see for instance [3]).

Definition 4 Given a Banach space X with the duality pairing $\langle \cdot, \cdot \rangle_{X', X}$, an element $x \in X$, a functional $f : X \rightarrow \mathbb{R} \cup \{+\infty\}$, the sub-differential of f at x is the set defined as

$$\partial f(x) := \{x' \in X' : f(y) - f(x) \geq \langle x', y - x \rangle_{X', X} \quad \forall y \in X\}.$$

We denote the domain of ∂f as usual by $D(\partial f)$, i.e. the set of all $x \in X$ such that $\partial f(x) \neq \emptyset$.

Definition 5 Given a Banach space X with the duality pairing $\langle \cdot, \cdot \rangle_{X', X}$, a multivalued operator $A : X \rightarrow X'$ identified with its graph $\Gamma_A := \{(u, Au) : u \in X\} \subseteq X \times X'$ is:

1. **monotone** if for any pair $(u, u'), (v, v') \in \Gamma_A$, it holds

$$\langle u' - v', u - v \rangle_{X', X} \geq 0;$$

2. **maximal monotone** if the graph Γ_A is not a proper subset of any monotone set;

174 **Definition 6** Given a Hilbert space X , a multivalued operator $B : X \rightarrow X$, with graph $\Gamma_B(X) :=$
 175 $\{(x, Bx) : x \in X \text{ such that } Bx \in X\}$. Denote $J_X : X \rightarrow X'$ as the canonical isomorphism of X to X' .
 176 B is

- 177 1. **accretive** if for any pair $(x, Bx), (y, By)$, there exists an element $z \in J_X(x - y)$ such that $\langle z, Bx -$
 178 $By \rangle_{X', X} \geq 0$;
 179 2. **m-accretive** if it is accretive and $R(I + B) = X$, where $R(I + B)$ denotes the range of $(I + B)$;

180 *Remark:* For general Banach space, J_X is the duality mapping of X ; see details in [3, Section 1.1].
 181 In our case, $X = U$, so J_X is the identity operator I in U .

182 2 Existence result for variational inequality solution

183 This section is devoted to obtain a variational inequality solution to (8). By restricting the operator in
 184 the non-reflexive Banach space \tilde{V} to U , we want to apply the classical result for m-accretive operator
 185 in U . However, since we do not have weak compactness for sequences in V , and a Radon measure may
 186 appear when taking the limit, we need to first prove weak-* lower semi-continuity for functional ϕ in
 187 \tilde{V} .

188 2.1 Weak-* lower semi-continuity for functional in \tilde{V} .

189 Since for any $w \in \tilde{V}$, ϕ defined only on its absolutely continuous part, we need the following proposition
 190 to guarantee ϕ is lower-semi-continuous with respect to the weakly-* convergence in \tilde{V} .

Proposition 7 *The function ϕ defined in Definition 1 is lower semi-continuous with respect to the weakly-* convergence in \tilde{V} , i.e., if $w_n \xrightarrow{*} w$ in \tilde{V} , we have*

$$\liminf_{n \rightarrow +\infty} \phi(w_n) \geq \phi(w).$$

191 For any $\mu \in \mathcal{M}(\mathbb{T})$, we denote $\mu \ll \mathcal{L}^1$ if μ is absolutely continuous with respect to Lebesgue measure
 192 and denote $\bar{\mu} := \frac{d\mu}{d\mathcal{L}^1}$ as the density of μ . For notational simplification, denote $\mu_{||}$ (resp. μ_{\perp}) as the
 193 absolutely continuous part (resp. singular part) of μ with respect to Lebesgue measure. Before proving
 194 Proposition 7, we first state some lemmas. The following Lemma comes from the weak-* compactness
 195 of L^∞ .

Lemma 8 *For any $N \geq 0$, given a sequence of measures μ_n in $\mathcal{M}(\mathbb{T})$, such that each $\mu_n \ll \mathcal{L}^1$ for any n , and the densities satisfy*

$$\sup_n \|\bar{\mu}_n\|_{L^\infty(\mathbb{T})} \leq N,$$

196 *then there exists a measure $\mu \ll \mathcal{L}^1$, $\|\bar{\mu}\|_{L^\infty(\mathbb{T})} \leq N$ and a subsequence μ_{n_k} such that $\mu_{n_k} \xrightarrow{*} \mu$ in $\mathcal{M}(\mathbb{T})$.*

197 From now on, we identify μ_n with its density $\bar{\mu}_n := \frac{d\mu_n}{d\mathcal{L}^1}$ and do not distinguish them for brevity.
 198 Given a sequence of measures μ_n such that $\mu_n \ll \mathcal{L}^1$, $\bar{\mu}_n \geq 0$ and $N > 0$, observe that

$$\mu_n = \min\{\mu_n, N\} + \max\{\mu_n, N\} - N. \quad (24) \quad \{b\}$$

199 From Lemma 8 we know, upon subsequence, $\min\{\mu_n, N\} \xrightarrow{*} \mu_-$ for some measure μ_- satisfying $\mu_- \ll \mathcal{L}^1$
 200 and $N \geq \mu_- \geq 0$. We also need the following useful Lemma to clarify the relation between μ_- and the
 201 weak-* limit of μ_n .

202 **Lemma 9** Given a sequence of measures μ_n such that $\mu_n \ll \mathcal{L}^1$, $\bar{\mu}_n \geq 0$, we assume moreover that
 203 $\mu_n \xrightarrow{*} \mu$, for some measure $\mu \geq 0$. Then for any $N > 0$, there exist $\mu_-, \mu_+ \in \mathcal{M}(\mathbb{T})$, such that

$$\min\{\mu_n, N\} \xrightarrow{*} \mu_-, \quad \mu_- \ll \mathcal{L}^1, \quad \mu_- \leq \mu_{\parallel}, \quad (25) \quad \{\star\text{lem1}\}$$

204

$$\max\{\mu_n, N\} \xrightarrow{*} \mu_+, \quad \mu_{+\parallel} \geq N, \quad (26) \quad \{\star\text{lem2}\}$$

205 where μ_{\parallel} (resp. μ_{\perp}) is the absolutely continuous part (resp. singular part) of μ . Moreover, for the
 206 function Φ defined in (18), we have

$$\{\text{PHI}\} \quad \Phi(\mu_{\parallel}) \leq \Phi(\mu_-). \quad (27)$$

Proof From Lemma 8 we know, upon subsequence, $\min\{\mu_n, N\} \xrightarrow{*} \mu_-$ for some measure μ_- satisfying
 $\mu_- \ll \mathcal{L}^1$ and $N \geq \mu_- \geq 0$. By Lebesgue decomposition theorem, there exist unique measures $\mu_{\parallel} \ll \mathcal{L}^1$
 and $\mu_{\perp} \perp \mathcal{L}^1$ such that $\mu = \mu_{\parallel} + \mu_{\perp}$. The decomposition (24) then gives

$$0 \leq \mu_n - \min\{\mu_n, N\} = \max\{\mu_n, N\} - N \xrightarrow{*} \mu - \mu_-.$$

207 Taking $\mu_+ := \mu - \mu_- + N$, as the sequence $\max\{\mu_n, N\} - N \geq 0$, we obtain $\max\{\mu_n, N\} \xrightarrow{*} \mu_+$ and
 208 $(\mu - \mu_-)_{\parallel} = \mu_{+\parallel} - N \geq 0$. Besides, since $\Phi(\mu_{\parallel})$ is decreasing with respect to μ_{\parallel} , we obtain (27). \square

209

Now we can start to prove Proposition 7.

Proof (Proof of Proposition 7) Without loss of generality we assume $\sup_{n \rightarrow +\infty} \phi(w_n) < +\infty$. This
 implies immediately that all $w_{nhh} + c_0$ are positive measures. Assume $w_n \xrightarrow{*} w$ in \tilde{V} , thus we have
 $w_{nhh} \xrightarrow{*} w_{hh}$ in \mathcal{M} . Denote $f_n := w_{nhh} + c_0$ and $f := w_{hh} + c_0$. Since $\Phi(f_{\parallel})$ is decreasing with respect
 to f_{\parallel} , we only concern the case $f_{n\parallel}$ may weakly-* converge to a singular measure. Thus without loss
 of generality, we assume $f_n \ll \mathcal{L}^1$, i.e., $f_{n\perp} = 0$. For any $M > 0$ large enough, denote $\phi_M(w_n) :=$
 $\int_{\mathbb{T}} \Phi(\min\{f_n, M\})$. From the definition of Φ in (18), the truncated measures $\min\{f_n, M\}$ satisfy

$$\begin{aligned} \phi_M(w_n) &= \int_{\mathbb{T}} \Phi(\min\{f_n, M\}) dh \\ &= \int_{\{f_n \leq M\}} \Phi(\min\{f_n, M\}) dh + \frac{1}{2M^2} \mathcal{L}^1(\{f_n > M\}) \\ &\geq \int_{\{f_n \leq M\}} \Phi(f_n) dh + \int_{\{f_n > M\}} \Phi(f_n) dh = \phi(w_n). \end{aligned}$$

The second equality also shows

$$\begin{aligned} \phi_M(w_n) - \frac{1}{2M^2} \mathcal{L}^1(\{f_n > M\}) &= \int_{\{f_n \leq M\}} \Phi(\min\{f_n, M\}) dh \\ &\leq \int_{\mathbb{T}} \Phi(f_n) dh = \phi(w_n). \end{aligned}$$

210 Hence we obtain

$$\{0104_28o\} \quad |\phi(w_n) - \phi_M(w_n)| \leq \frac{1}{2M^2} \mathcal{L}^1(\{f_n > M\}) \leq \frac{1}{2M^2}. \quad (28)$$

211

From Lemma 8 and Lemma 9, we know the truncated sequence $\min\{f_n, M\}$ satisfies

$$\{\star 0105_1\} \quad \min\{f_n, M\} \xrightarrow{*} f_-, \quad f_- \ll \mathcal{L}^1, \quad \Phi(f_{\parallel}) \leq \Phi(f_-). \quad (29)$$

212

Hence by the convexity and lower semi-continuity of ϕ on V , we infer

$$\{\text{conv}\} \quad \liminf_{n \rightarrow +\infty} \int_{\mathbb{T}} \Phi(\min\{f_n, M\}) dh \geq \int_{\mathbb{T}} \Phi(f_-) dh \geq \int_{\mathbb{T}} \Phi(f_{\parallel}) dh = \phi(w). \quad (30)$$

213 Combining this with (28), we obtain

$$\begin{aligned}
\liminf_{n \rightarrow +\infty} \phi(w_n) &\geq \liminf_{n \rightarrow +\infty} \phi_M(w_n) - \frac{1}{2M^2} \\
&= \liminf_{n \rightarrow +\infty} \int_{\mathbb{T}} \Phi(\min\{f_n, M\}) dh - \frac{1}{2M^2} \\
&\geq \phi(w) - \frac{1}{2M^2},
\end{aligned} \tag{31}$$

214 and thus we complete the proof of Proposition 7 by the arbitrariness of M . \square

215 2.2 Maximal monotone and m-accretive operator in U .

216 In this section, we first define the sub-differential of $\phi + \psi$ and then obtain some useful lemmas to ensure
217 $\partial(\phi + \psi)$ is also a maximal monotone operator when restricted to U .

218 Let $\tilde{B} := \partial(\phi + \psi) : \tilde{V} \rightarrow \tilde{V}'$ be the sub-differential of $\phi + \psi$. Let us consider the operator B as the
219 restriction of \tilde{B} from U to U' .

Definition 10 Define the operator $B : D(B) \subseteq U \rightarrow U'$ such that

$$Bw = \tilde{B}w, \text{ for any } w \in D(B) = \{w \in \tilde{V}; \tilde{B}w \in U\}.$$

220 We first prove B is maximal monotone in $U \times U'$, which is important to prove the existence result.

221 **Lemma 11** The operator $B : D(B) \subseteq U \rightarrow U'$ in Definition 10 is maximal monotone in $U \times U'$.

222 *Proof* It suffices to prove that $\phi + \psi$ is (i) proper, i.e. $D(\phi + \psi) \neq \emptyset$, (ii) convex and (iii) lower semi-
223 continuous when considered as a functional from $(U, \|\cdot\|_U)$ to $\mathbb{R} \cup \{+\infty\}$. (i) First it is clear that $\phi + \psi$
224 is proper.

(ii) *Convexity.* Let $u_1, u_2 \in U$ be arbitrarily given, and we need to show

$$(1-t)(\phi + \psi)(u_1) + t(\phi + \psi)(u_2) \geq (\phi + \psi)((1-t)u_1 + tu_2).$$

If u_1 or u_2 does not belong to $D(\phi + \psi)$, then the left-hand side term is $+\infty$. If both u_1 and u_2 belong
to $D(\phi + \psi)$, then $(1-t)u_1 + tu_2$ also belongs to $\tilde{V} \cap \{\|\cdot\|_{\tilde{V}} \leq C\}$, hence

$$\psi(u_1) = \psi(u_2) = \psi((1-t)u_1 + tu_2) = 0.$$

Notice the convexity of ϕ , and the fact that the absolutely continuous part of $((1-t)u_1 + tu_2)_{hh}$ is
 $(1-t)((u_1)_{hh})_{\parallel} + t((u_2)_{hh})_{\parallel}$, where $((u_i)_{hh})_{\parallel}$ are notations representing the absolutely continuous parts
of $(u_i)_{hh}$ separately. Then we obtain

$$(1-t)\phi(u_1) + t\phi(u_2) \geq \phi((1-t)u_1 + tu_2).$$

225 Thus $\phi + \psi$ is convex.

(iii) *Lower-semicontinuity.* Note that the lower-semicontinuity is here intended as with respect to
the strong convergence in U (less restrictive than the convergence in \tilde{V}). Consider an arbitrary sequence
 $u^n \subseteq U$ converging to $u \in U$. We need to prove

$$(\phi + \psi)(u) \leq \liminf_{n \rightarrow +\infty} (\phi + \psi)(u^n).$$

If $\liminf_{n \rightarrow +\infty} (\phi + \psi)(u^n) = +\infty$ then the thesis is trivial. Thus assume (upon subsequence)

$$\liminf_{n \rightarrow +\infty} (\phi + \psi)(u^n) = \lim_{n \rightarrow +\infty} (\phi + \psi)(u^n) \leq D < +\infty.$$

Without loss of generality we can further assume $(u^n)_n \subseteq \tilde{V} \cap \{\|\cdot\|_{\tilde{V}} \leq C\}$. This implies $\|u_{hh}^n\|_{\mathcal{M}(\mathbb{T})} \leq C$, so

$$\psi(u^n) = 0, \quad \forall n$$

and there exists $\xi \in \mathcal{M}(\mathbb{T})$ such that $u_{hh}^n \xrightarrow{*} \xi$ in $\mathcal{M}(\mathbb{T})$. Since from Proposition 7, $\phi + \psi$ is convex and lower-semicontinuous in \tilde{V} , we infer

$$\phi(\xi) \leq \liminf_{n \rightarrow +\infty} \phi(u^n). \quad (32)$$

The uniform boundedness of $\|u_{hh}^n\|_{\mathcal{M}(\mathbb{T})}$ also implies $(u^n)_n$ is bounded in $W^{1,\infty}(\mathbb{T})$ (and hence in $W^{1,p}(\mathbb{T})$ for any $p < +\infty$). Thus $(u^n)_n$ is (upon subsequence) weakly convergent in $W^{1,p}(\mathbb{T})$, and strongly convergent in $L^2(\mathbb{T})$ to u . Thus $u_{hh} = \xi$, and (32) now becomes

$$\phi(u_{hh}) = \phi(u) = \phi(\xi) \leq \liminf_{n \rightarrow +\infty} \phi(u^n).$$

Therefore $\phi + \psi : U \rightarrow \mathbb{R} \cup \{+\infty\}$ is proper, convex and lower-semicontinuous. Then by [3, Theorem 2.8] we have that $\partial(\phi + \psi) : U \rightarrow U'$ is maximal monotone in $U \times U'$. \square

Notice $U' = U$. From Lemma 11 and the Definition 6 we deduce

Proposition 12 *The operator $B : D(B) \subseteq U \rightarrow U'$ in Definition 10 is m -accretive from $D(B) \subseteq U$ to U .*

2.3 Existence result for variational inequality solution.

After those preliminary results, we can apply [3, Theorem 4.5] to obtain the existence of variational inequality solution to (8).

First let us recall [3, Theorem 4.5].

Theorem 13 ([3, Theorem 4.5]) *For any $T > 0$, let U be a Hilbert space and let B be a m -accretive operator from $D(B) \subseteq U$ to U . Then for each $y_0 \in D(B)$, the cauchy problem*

$$\begin{cases} \frac{dy}{dt}(t) + By(t) \ni 0, & t \in [0, T], \\ y(0) = y_0, \end{cases}$$

has a unique strong solution $y \in W^{1,\infty}([0, T]; U)$ in the sense that

$$-\frac{dy}{dt}(t) \in By(t), \quad a.e. t \in (0, T), \quad y(0) = y_0.$$

Moreover, y satisfies the estimate

$$\|y_t\|_U \leq |-By_0|_*, \quad (33)$$

where $|-By_0|_* = \inf\{\|u\|_U; u \in -By_0\}$.

Proposition 12 shows that B defined in Definition 10 is m -accretive from $D(B) \subseteq U$ to U . Hence we can apply Theorem 13 to obtain

Theorem 14 *Let $B : D(B) \subseteq U \rightarrow U$ be the operator defined in Definition 10. Given $T > 0$, initial datum $w^0 \in D(B)$, then*

243 (i) there exists a unique function $w \in W^{1,\infty}([0, T]; U)$ such that

$$\{slu1\} \quad -w_t(t) \in Bw(t), \quad w(0) = w^0, \quad \text{for a.e. } t \in [0, T]. \quad (34)$$

244 (ii) w is also a variational inequality solution to (8). Moreover,

$$\{\text{positive}\} \quad \begin{aligned} w_{hh} + c_0 &\in \mathcal{M}^+(\mathbb{T}), \text{ for a.e. } (t, h) \in [0, T] \times \mathbb{T}, \\ \eta + c_0 &> 0, \text{ for a.e. } (t, h) \in [0, T] \times \mathbb{T}, \end{aligned} \quad (35)$$

245 with respect to the Lebesgue measure, where η is the absolutely continuous part of w_{hh} in (17),
246 and $\mathcal{M}^+(\mathbb{T})$ denotes the set of positive Radon measures.

247 *Proof* Proof of (i). From Proposition 12, we know B is a m -accretive operator in $U \times U$. So (34) follows
248 from Theorem 13 and $w \in W^{1,\infty}([0, T]; U)$. From (33), we also have

$$\|w_t\|_U \leq |-Bw_0|_\star, \quad (36) \quad \{\text{es1107}_2\}$$

249 where $|-Bw_0|_\star = \inf\{\|u\|_U; u \in -Bw_0\}$.

250 Proof of (ii). Since $-w_t(t) \in Bw(t) = \partial(\phi + \psi)(w)$ for a.e. $t \in [0, T]$ and $w \in W^{1,\infty}([0, T]; U)$, by
251 Definition 4, we have

$$\langle w_t, v - w \rangle_{U', U} + (\phi + \psi)(v) - (\phi + \psi)(w) \geq 0, \quad \text{a.e. } t \in [0, T], \quad (37) \quad \{\text{v12}\}$$

for all $v \in U$, and

$$w \in C^0([0, T]; U), \quad w_t \in L^\infty([0, T]; U).$$

252 Let $v \in U$ such that $(\phi + \psi)(v) \leq 1$. Then from (37), we also have

$$(\phi + \psi)(w) \leq \|w_t\|_U \|w - v\|_U + 1 \quad \text{for a.e. } t \in [0, T]. \quad (38) \quad \{\text{45}\}$$

This implies

$$w \in L^\infty([0, T]; \tilde{V})$$

and with respect to the Lebesgue measure,

$$\eta + c_0 > 0 \text{ for a.e. } (t, h) \in [0, T] \times \mathbb{T},$$

$$w_{hh} + c_0 \in \mathcal{M}^+(\mathbb{T}) \text{ for a.e. } (t, h) \in [0, T] \times \mathbb{T},$$

253 where η is the absolutely continuous part of w_{hh} in (17). Therefore we obtain the variational inequality
254 solution to (8) and w satisfies the positivity property (35). \square

255 3 Existence result for strong solution

256 Although we obtained a unique variational inequality solution in Theorem 14, we do not know whether
257 B is single-valued and which element belongs to B . We will determine B in Section 3.1 and then prove
258 the variational inequality solution is actually a strong solution to (8) in Section 3.2.

3.1 Sub-differential of ϕ is single-valued.

For any $f \in C(\mathbb{T})$, f_{hh} induces a continuous linear form on \tilde{V} given by

$$f_{hh}(z) := \int_{\mathbb{T}} f dz_{hh} < \infty \quad (39) \quad \{\text{n39}\}$$

for any $z \in \tilde{V}$. Hence we can write $f_{hh} \in \tilde{V}'$. In the following proposition, we calculate the sub-differential $\partial\phi(w)$, for $w \in D(\partial\phi) \cap \{w \in \tilde{V}; g^{-3} \in C(\mathbb{T})\}$.

Proposition 15 *Let $g \in L^1(\mathbb{T})$ be the absolutely continuous part of $w_{hh} + c_0$ from Definition 1. Let ϕ be the proper, convex, lower-semicontinuous functional defined in (18). Let*

$$\partial\phi : D(\partial\phi) \subseteq \tilde{V} \longrightarrow \tilde{V}'$$

be the sub-differential of ϕ . For any $w \in D(\partial\phi) \cap \{w \in \tilde{V}; g^{-3} \in C(\mathbb{T})\}$, we have

$$\partial\phi(w) = \{-(g^{-3})_{hh}\} \in \tilde{V}' \subseteq V', \quad (40)$$

and

$$\langle \partial\phi(w), v \rangle_{V', V} := \int_{\mathbb{T}} -g^{-3} v_{hh} dh$$

for any $v \in V$.

Proof Step 1. Let

$$w \in D(\partial\phi) \cap \{w \in \tilde{V}; g^{-3} \in C(\mathbb{T})\}$$

be given. We first prove that

$$\{-(g^{-3})_{hh}\} \subseteq \partial\phi(w). \quad (41)$$

Since $w \in D(\partial\phi)$, we know $\partial\phi(w) \neq \emptyset$ and $w \in D(\partial\phi) \subseteq D(\phi)$ by [3, Proposition 1.6]. Thus we can assume $\phi(w) < +\infty$, so $g > 0$ a.e. and $g^{-2} \in L^1(\mathbb{T})$.

Let $z \in V$ be given. If $\phi(w + \varepsilon z) = +\infty$, then the inequality

$$\phi(w + \varepsilon z) - \phi(w) \geq \langle -(g^{-3})_{hh}, \varepsilon z \rangle_{V', V}$$

is trivial.

If $\phi(w + \varepsilon z) < +\infty$, then $g + \varepsilon z_{hh} > 0$ a.e.. Using the convexity of Φ , one can still have

$$\Phi(g + \varepsilon z_{hh}) - \Phi(g) \geq \Phi'(g)\varepsilon z_{hh} = -g^{-3}\varepsilon z_{hh} \quad \text{for a.e. } h. \quad (42)$$

Since $g^{-3} \in C(\mathbb{T})$, the right-hand side of (42) is integrable, i.e.,

$$\int_{\mathbb{T}} |g^{-3}\varepsilon z_{hh}| dh < +\infty. \quad (43)$$

Then by the definition for ϕ , from (42), we have

$$\phi(w + \varepsilon z) - \phi(w) = \int_{\mathbb{T}} [\Phi(g + \varepsilon z_{hh}) - \Phi(g)] dh \geq \int_{\mathbb{T}} -g^{-3}\varepsilon z_{hh} dh \quad (44)$$

for any $z \in V$.

Since $g^{-3} \in C(\mathbb{T})$, from (39), we know

$$\{(g^{-3})_{hh}\} \subseteq \tilde{V}' \subseteq V', \quad (45)$$

which implies that for any $v \in V$ we have

$$\int_{\mathbb{T}} (g^{-3})_{hh} v \, dh = \int_{\mathbb{T}} g^{-3} v_{hh} \, dh.$$

274 Hence $g^{-3} \in (L^1(\mathbb{T}))' \subseteq L^2(\mathbb{T})$.

275 Now in both cases we have

$$\{Gvep\} \quad \phi(w + \varepsilon z) - \phi(w) \geq \int_{\mathbb{T}} -g^{-3} \varepsilon z_{hh} \, dh = \langle -(g^{-3})_{hh}, \varepsilon z \rangle_{V', V} \quad (46)$$

276 for any $z \in V$.

277 Note that for any $\xi \in \tilde{V}'$, $z \in \tilde{V}$, since V is dense in \tilde{V} , there exists a sequence $z^n \in V$ such that

$$\langle \xi, z \rangle_{\tilde{V}', \tilde{V}} = \lim_{n \rightarrow \infty} \langle \xi, z^n \rangle_{\tilde{V}', \tilde{V}} = \lim_{n \rightarrow \infty} \langle \xi, z^n \rangle_{V', V}. \quad (47) \quad \{G2\}$$

278 Taking $\xi = -(g^{-3})_{hh}$ in (47), together with (46), we obtain

$$\phi(w + \varepsilon z) - \phi(w) \geq \langle -(g^{-3})_{hh}, \varepsilon z \rangle_{\tilde{V}', \tilde{V}}, \quad (48)$$

279 and by the Definition 4, we have (41).

Step 2. Let

$$w \in D(\partial\phi) \cap \{w \in \tilde{V}; g^{-3} \in C(\mathbb{T})\}$$

280 be given. We want to prove $\partial\phi$ is single-valued.

First we give a density argument. Assume for some w , there are two elements $\eta_1, \eta_2 \in \partial\phi(w)$. Assume the Gâteaux-derivative of ϕ exists for all z in some dense set G of \tilde{V} . Then by the definition of sub-differential we have

$$\begin{aligned} \lim_{\varepsilon \rightarrow 0} \frac{\phi(w + \varepsilon z) - \phi(w)}{\varepsilon} &\geq \langle \eta_1, z \rangle_{\tilde{V}', \tilde{V}}, & \lim_{\varepsilon \rightarrow 0} \frac{\phi(w - \varepsilon z) - \phi(w)}{\varepsilon} &\geq -\langle \eta_1, z \rangle_{\tilde{V}', \tilde{V}}, \\ \lim_{\varepsilon \rightarrow 0} \frac{\phi(w - \varepsilon z) - \phi(w)}{\varepsilon} &\geq -\langle \eta_2, z \rangle_{\tilde{V}', \tilde{V}}, & \lim_{\varepsilon \rightarrow 0} \frac{\phi(w + \varepsilon z) - \phi(w)}{\varepsilon} &\geq \langle \eta_2, z \rangle_{\tilde{V}', \tilde{V}} \end{aligned}$$

for $z \in G$. Hence we have

$$\langle \eta_1 - \eta_2, z \rangle_{\tilde{V}', \tilde{V}} \leq 0, \quad \langle \eta_2 - \eta_1, z \rangle_{\tilde{V}', \tilde{V}} \leq 0, \quad (49)$$

281 and then

$$\langle \eta_1, z \rangle_{\tilde{V}', \tilde{V}} = \langle \eta_2, z \rangle_{\tilde{V}', \tilde{V}} \quad (50) \quad \{dense21_1\}$$

282 for $z \in G$. Since G is dense in \tilde{V} , we know (50) holds for any $z \in \tilde{V}$, which shows $\partial\phi$ is single valued.

283 Then it is sufficient to prove for all functions z belonging to some dense subset G of V , which is a
284 dense subset of \tilde{V} , the Gâteaux-derivative of ϕ exists.

Second, we find the dense subset G . For any $c \geq 0$, define

$$D_{w,c} := \{\tilde{w} \in L^2(0, 1); |\tilde{w}/g| \leq c \text{ a.e.}\}, \quad D_w := \bigcup_{c \geq 0} D_{w,c}.$$

285 The above sets are well-defined since $g > 0$ almost everywhere.

We claim that for any given g , such that $g > 0$, the set D_w is dense in $L^2(\mathbb{T}) \cap L^\infty(\mathbb{T})$. Indeed, for any $f \in L^2(\mathbb{T}) \cap L^\infty(\mathbb{T})$, we consider the sequence in D_w given by

$$f_n := \begin{cases} f & \text{if } |f/g| \leq n, \\ 0 & \text{otherwise.} \end{cases}$$

Then, note that, for any $p \in [1, \infty)$,

$$\|f_n - f\|_{L^p(\mathbb{T})} = \left[\int_{\{|f/g| > n\}} |f - f_n|^p dh \right]^{1/p},$$

and $\left| \frac{f}{g} \right| \leq \frac{\|f\|_{L^\infty(\mathbb{T})}}{g}$. Hence

$$\left\{ \left| \frac{f}{g} \right| > n \right\} \subseteq \left\{ g < \frac{\|f\|_{L^\infty(\mathbb{T})}}{n} \right\},$$

which yields

$$\begin{aligned} \left[\int_{\left\{ \left| \frac{f}{g} \right| > n \right\}} |f - f_n|^p dh \right]^{1/p} &\leq \left[\int_{\left\{ g < \frac{\|f\|_{L^\infty(\mathbb{T})}}{n} \right\}} |f - f_n|^p dh \right]^{1/p} \\ &\leq \|f\|_{L^\infty(\mathbb{T})} \left| \left\{ g < \frac{\|f\|_{L^\infty(\mathbb{T})}}{n} \right\} \right|^{1/p} \xrightarrow{n \rightarrow +\infty} 0. \end{aligned}$$

286 Thus we have proven that D_w is dense in $L^2(\mathbb{T}) \cap L^\infty(\mathbb{T})$, and hence dense in $L^2(\mathbb{T})$. Now define
287 $G := \{z \in V; z_{hh} \in D_w\}$, which is a dense subset of V .

288 Third, we calculate the Gâteaux-derivative of ϕ . Consider a test function $z \in G$ so that $z_{hh} \in D_w$.

289 Since $D_w = \bigcup_{c \geq 0} D_{w,c}$, there exists c such that $z_{hh} \in D_{w,c}$.

Direct computation gives

$$\phi(w + \varepsilon z) - \phi(w) = \int_{\mathbb{T}} \left(\frac{1}{2(g + \varepsilon z_{hh})^2} - \frac{1}{2g^2} \right) dh = \int_{\mathbb{T}} \frac{-2\varepsilon g z_{hh} - \varepsilon^2 z_{hh}^2}{2g^2(g + \varepsilon z_{hh})^2} dh =: I_1 + I_2,$$

where

$$\{I\} \quad I_1 := -\varepsilon \int_{\mathbb{T}} \frac{z_{hh}}{g(g + \varepsilon z_{hh})^2} dh, \quad I_2 := -\varepsilon^2 \int_{\mathbb{T}} \frac{z_{hh}^2}{2g^2(g + \varepsilon z_{hh})^2} dh. \quad (51)$$

290 *Estimate for I_2 .* Since we are interested in the behavior of I_2 in the limit $\varepsilon \rightarrow 0$, we will consider
291 only $\varepsilon < 1/(2c)$. Thus

$$\{\text{w below}\} \quad g + \varepsilon z_{hh} \geq g - \varepsilon |z_{hh}/g| \geq g(1 - c\varepsilon) \geq g/2 \quad \text{for a.e. } h, \quad (52)$$

and

$$\begin{aligned} |I_2| &= \left| \varepsilon^2 \int_{\mathbb{T}} \frac{z_{hh}^2}{2g^2(g + \varepsilon z_{hh})^2} dh \right| \\ &\leq \varepsilon^2 \frac{c^2}{2} \int_{\mathbb{T}} (g + \varepsilon z_{hh})^{-2} dh \quad (\text{due to } z_{hh} \in D_{w,c}) \\ &\leq 2\varepsilon^2 c^2 \int_{\mathbb{T}} g^{-2} dh = O(\varepsilon^2), \end{aligned}$$

292 where we also used (52) and $g^{-2} \in L^1(\mathbb{T})$.

Estimate for I_1 . Again, since we are interested in the behavior of I_1 in the limit $\varepsilon \rightarrow 0$, we will consider only $\varepsilon < 1/(2c)$. Then

$$I_1/\varepsilon = - \int_{\mathbb{T}} \frac{z_{hh}}{g(g + \varepsilon z_{hh})^2} dh,$$

and

$$\frac{z_{hh}}{g(g + \varepsilon z_{hh})^2} \rightarrow \frac{z_{hh}}{g^3} \quad \text{for a.e. } h.$$

Using (52), we have

$$\left| \frac{z_{hh}}{g(g + \varepsilon z_{hh})^2} \right| \leq \left| \frac{4z_{hh}}{g^3} \right| \in L^1(\mathbb{T}),$$

with the last inclusion due to

$$\int_{\mathbb{T}} \left| \frac{z_{hh}}{g^3} \right| dh \leq \|z_{hh}\|_{L^2(\mathbb{T})} \|g^{-3}\|_{L^2(\mathbb{T})}.$$

Thus by Lebesgue's dominated convergence we infer

$$I_1/\varepsilon = - \int_{\mathbb{T}} \frac{z_{hh}}{g(g + \varepsilon z_{hh})^2} dh \xrightarrow{\varepsilon \rightarrow 0} - \int_{\mathbb{T}} \frac{z_{hh}}{g^3} dh = \langle -(g^{-3})_{hh}, z \rangle_{V', V}.$$

Since $(g^{-3})_{hh} \in \tilde{V}'$, $z \in \tilde{V}$, we have

$$\langle -(g^{-3})_{hh}, z \rangle_{V', V} = \langle -(g^{-3})_{hh}, z \rangle_{\tilde{V}', \tilde{V}}.$$

293 Therefore ,

$$\lim_{\varepsilon \rightarrow 0} \frac{\phi(w + \varepsilon z) - \phi(w)}{\varepsilon} = \langle -(g^{-3})_{hh}, z \rangle_{\tilde{V}', \tilde{V}} \quad (53)$$

294 for any $z \in G$. This together with the density argument (50), gives (40) and completes the proof. \square

295 3.2 Existence result for strong solution.

296 Now combining Proposition 15 and Theorem 14, we can state existence result for strong solution as
297 follows.

298 **Theorem 16** *Given $T > 0$, initial data $w^0 \in D(B)$, then the variational inequality solution w obtained
299 in Theorem 14 is also a strong solution to (8), i.e.,*

$$w_t = (g^{-3})_{hh} = ((\eta + c_0)^{-3})_{hh}, \quad (54) \quad \{\text{main21_3}\}$$

300 for a.e. $(t, h) \in [0, T] \times \mathbb{T}$. Besides, we have $((\eta + c_0)^{-3})_{hh} \in L^\infty([0, T]; U)$ and the dissipation inequality

$$E(t) := \frac{1}{2} \int_{\mathbb{T}} [((\eta + c_0)^{-3})_{hh}]^2 dh \leq E(0), \quad (55) \quad \{\text{dissiE}\}$$

301 where η is the absolutely continuous part of w_{hh} in (17).

302 *Proof* First we want to prove that

$$\partial\psi(w) = \{0\}, \quad (56) \quad \{\text{pd1}\}$$

303 and

$$g^{-3} \in C(\mathbb{T}) \text{ for a.e. } t \in [0, T]. \quad (57) \quad \{\text{pd2}\}$$

Indeed, if both (56) and (57) hold, Proposition 15 gives that

$$Bw = -(g^{-3})_{hh} = -((\eta + c_0)^{-3})_{hh}$$

304 is single-valued, which together with (34) implies (54).

305 Hence it is sufficient to prove

$$w \in \text{Int } \mathcal{K}, \quad (58) \quad \{\text{int}\}$$

and (57). We choose the constant C in Definition 1 as

$$C := \|w_0\|_{\tilde{V}} + 2c_0 + 3.$$

306 We need to use two *a priori* assumptions

$$\|w\|_{\tilde{V}} \leq C - 1 < C, \quad (59) \quad \{\text{apri}\}$$

307

$$\|(g^{-3})_{hh}\|_U \leq \sqrt{2E(0)} + 1, \quad (60) \quad \{\text{apri1}\}$$

308 which will be verified later.

309 Under (59), we have $w \in \text{Int } \mathcal{K}$ and $\partial\psi(w) = 0$. Under (60), we know for any $h_1, h_2 \in \mathbb{T}$,

$$\{cc\} \quad |g^{-3}(h_1) - g^{-3}(h_2)| \leq |h_1 - h_2| \|(g^{-3})_{hh}\|_U. \quad (61)$$

310 Notice that from (38) we also have $\int_{\mathbb{T}} \frac{1}{g^2} dh < +\infty$. Hence

$$\left(\frac{1}{g^3}\right)_{\min}^{2/3} = \left(\frac{1}{g^2}\right)_{\min} \leq \int_{\mathbb{T}} \frac{1}{g^2} dh < +\infty. \quad (62)$$

311 This, together with (61), gives $g^{-3} \in C(\mathbb{T})$. Therefore, Proposition 15 shows that $Bw = -(g^{-3})_{hh}$.

312 Then from assertion (i) in Theorem 14 we know

$$\{50\} \quad w_t = (g^{-3})_{hh} = ((\eta + c_0)^{-3})_{hh}, \quad (63)$$

313 for a.e. $(t, h) \in [0, T] \times \mathbb{T}$.

Next we verify the *a priori* assumption (59). Let J_δ be the standard $C^\infty(\mathbb{T})$ mollifier. Then we have

$$J_{\delta hh} \star w_t = J_{\delta hh} \star ((\eta + c_0)^{-3})_{hh}$$

314 for a.e. $(t, h) \in [0, T] \times \mathbb{T}$. Integrating on \mathbb{T} shows that

$$\frac{d}{dt} \int_{\mathbb{T}} J_\delta \star w_{hh} dh = \int_{\mathbb{T}} [J_\delta \star ((\eta + c_0)^{-3})_{hh}]_{hh} dh = 0. \quad (64)$$

Thus we obtain

$$\int_{\mathbb{T}} J_\delta \star w_{hh} dh \leq \int_{\mathbb{T}} J_\delta \star w_{0hh} dh \leq \|w_0\|_{\tilde{V}},$$

and

$$\int_{\mathbb{T}} J_\delta \star (w_{hh} + c_0) dh = \int_{\mathbb{T}} J_\delta \star w_{hh} dh + c_0 \leq \|w_0\|_{\tilde{V}} + c_0.$$

315 Since we also have $w \in L^\infty([0, T]; \tilde{V})$, we know that for δ small enough,

$$\begin{aligned} \|w\|_{\tilde{V}} &\leq \|w_{hh} + c_0\|_{\mathcal{M}(\mathbb{T})} + c_0 \\ &\leq \int_{\mathbb{T}} J_\delta \star (w_{hh} + c_0) dh + 1 + c_0 \\ &\leq \|w_0\|_{\tilde{V}} + 2c_0 + 1 \\ &= C - 2 < C - 1, \end{aligned} \quad (65)$$

316 which verifies the *a priori* assumption (59).

317 Now we turn to verify (60). Combining (63) and (36), we have the dissipation law

$$\{tE\} \quad E(t) = \frac{1}{2} \|w_t\|_U^2 = \frac{1}{2} \|((\eta + c_0)^{-3})_{hh}\|_U^2 \leq \frac{1}{2} \|Bw_0\|_U^2 = E(0) \quad (66)$$

318 for $E(t) = \frac{1}{2} \int_{\mathbb{T}} [((\eta + c_0)^{-3})_{hh}]^2 dh$. Hence we know

$$\|(g^{-3})_{hh}\|_U = \|((\eta + c_0)^{-3})_{hh}\|_U \leq \sqrt{2E(0)} < \sqrt{2E(0)} + 1, \quad (67)$$

319 which verifies the *a priori* assumption (60).

320 Therefore, we know (56) and (57), so B is single-valued and (54) holds. The dissipation inequality
321 (55) comes from (66) and we completes the proof of (iii). \square

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324 References

- 325 1. H. Al Hajj Shehadeh, R. V. Kohn and J. Weare, The evolution of a crystal surface: Analysis of a one-dimensional
326 step train connecting two facets in the adl regime, *Physica D: Nonlinear Phenomena* 240 (2011), no. 21, 1771–1784.
- 327 2. L. Ambrosio, N. Gigli and G. Savaré, *Gradient flows in metric spaces and in the space of probability measures*,
328 Birkhäuser Verlag, 2005.
- 329 3. V. Barbu, *Nonlinear differential equations of monotone types in banach spaces*, Springer, New York, 2010.
- 330 4. F. Bernis and A. Friedman, Higher order nonlinear degenerate parabolic equations, *Journal of Differential Equations*
331 83 (1990), no. 1, 179–206.
- 332 5. W. K. Burton, N. Cabrera and F. C. Frank, The growth of crystals and the equilibrium structure of their surfaces,
333 *Philosophical Transactions of the Royal Society of London A: Mathematical, Physical and Engineering Sciences* 243
334 (1951), no. 866, 299–358.
- 335 6. G. Dal Maso, I. Fonseca and G. Leoni, Analytical validation of a continuum model for epitaxial growth with elasticity
336 on vicinal surfaces, *Archive for Rational Mechanics and Analysis* 212 (2014), no. 3, 1037–1064.
- 337 7. W. E and N. K. Yip, Continuum theory of epitaxial crystal growth. I, *Journal Statistical Physics* 104 (2001), no. 1-2,
338 221–253.
- 339 8. L.C. EVANS, R.F. GARIEPY, *Measure Theory and Fine Properties of Functions*, CRC Press, 1992.
- 340 9. I. Fonseca, G. Leoni and X. Y. Lu, Regularity in time for weak solutions of a continuum model for epitaxial growth
341 with elasticity on vicinal surfaces, *Communications in Partial Differential Equations* 40 (2015), no. 10, 1942–1957.
- 342 10. Y. Gao, H. Ji, J.-G. Liu and T. P. Witelski, A vicinal surface model for epitaxial growth with logarithmic entropy,
343 submitted.
- 344 11. Y. Gao, J.-G. Liu and J. Lu, Continuum limit of a mesoscopic model of step motion on vicinal surfaces, *J Nonlinear*
345 *Sci* (2016). doi:10.1007/s00332-016-9354-1.
- 346 12. Y. Gao, J.-G. Liu and J. Lu, Weak solution of a continuum model for vicinal surface in the attachment-detachment-
347 limited regime, accepted by *SIAM J. Math. Anal.*
- 348 13. Y. Giga, R.V. Kohn, Scale-invariant extinction time estimates for some singular diffusion equations. Hokkaido Uni-
349 versity Preprint Series in Mathematics (2010), no. 963.
- 350 14. R. V. Kohn, "Surface relaxation below the roughening temperature: Some recent progress and open questions,"
351 *Nonlinear partial differential equations: The abel symposium 2010*, H. Holden and H. K. Karlsen (Editors), Springer
352 Berlin Heidelberg, Berlin, Heidelberg, 2012, pp. 207–221.
- 353 15. J.-G. Liu and X. Xu, *Existence theorems for a multi-dimensional crystal surface model*, *SIAM J. Math. Anal.* 48
354 (2016), no. 6, 3667–3687.
- 355 16. J.-G. Liu and X. Xu, Analytical validation of a continuum model for the evolution of a crystal surface in multiple
356 space dimensions, submitted.
- 357 17. D. Margetis and R. V. Kohn, Continuum relaxation of interacting steps on crystal surfaces in 2 + 1 dimensions,
358 *Multiscale Modeling & Simulation* 5 (2006), no. 3, 729–758.
- 359 18. D. Margetis, K. Nakamura, From crystal steps to continuum laws: Behavior near large facets in one dimension,
360 *Physica D*, 240 (2011), 1100–1110.
- 361 19. W.W. Mullins, Theory of Thermal Grooving. *Journal of Applied Physics* 28 (1957), 333–339.
- 362 20. M. Ozdemir and A. Zangwill, Morphological equilibration of a corrugated crystalline surface, *Physical Review B* 42
363 (1990), no. 8, 5013–5024.
- 364 21. A. Pimpinelli and J. Villain, *Physics of Crystal Growth*, Cambridge University Press, New York, 1998.
- 365 22. V. Shenoy and L. Freund, A continuum description of the energetics and evolution of stepped surfaces in strained
366 nanostructures, *Journal of the Mechanics and Physics of Solids* 50 (2002), no. 9, 1817–1841.
- 367 23. L.-H. Tang, Flattening of grooves: From step dynamics to continuum theory, *Dynamics of crystal surfaces and inter-*
368 *faces* (1997), 169.
- 369 24. J. D. Weeks and G. H. Gilmer, Dynamics of Crystal Growth, in *Advances in Chemical Physics*, Vol. 40, edited by I.
370 Prigogine and S. A. Rice (John Wiley, New York, 1979), pp. 157–228.

- 371 25. Y. Xiang, Derivation of a continuum model for epitaxial growth with elasticity on vicinal surface, SIAM Journal on
372 Applied Mathematics 63 (2002), no. 1, 241–258.
- 373 26. Y. Xiang and W. E, Misfit elastic energy and a continuum model for epitaxial growth with elasticity on vicinal
374 surfaces, Physical Review B 69 (2004), no. 3, 035409.
- 375 27. H. Xu and Y. Xiang, Derivation of a continuum model for the long-range elastic interaction on stepped epitaxial
376 surfaces in $2 + 1$ dimensions, SIAM Journal on Applied Mathematics 69 (2009), no. 5, 1393–1414.