

## Scuola Normale Superiore

Classe di Scienze Matematiche Fisiche e Naturali

Ph.D. Thesis

# Optimization problems for solutions of elliptic equations and stability issues 

Berardo Ruffini

Ph.D. Advisor
Prof. Giuseppe Buttazzo

## Contents

Introduction ..... i
I Optimization problems for shapes and for solutions of elliptic equations ..... 1
1 Optimization problems related to Schrödinger operators ..... 3
1.1 Introduction ..... 3
1.2 Capacitary measures and $\gamma$-convergence ..... 6
1.3 Existence of optimal potentials in $L^{p}(\Omega)$ ..... 9
1.4 Existence of optimal potentials for unbounded constraints ..... 13
1.4.1 Optimal potentials for the Dirichlet Energy and the first eigen- value of the Dirichlet Laplacian ..... 14
1.5 Maximization problems in $L^{p}$ concerning the Dirichlet Energy functional ..... 16
1.6 Optimization problems in unbounded domains ..... 22
1.6.1 Optimal potentials in $L^{p}\left(\mathbb{R}^{d}\right)$ ..... 23
1.6.2 Optimal potentials in $\mathbb{R}^{d}$ with unbounded constraint ..... 24
1.6.3 Further remarks ..... 29
2 Optimization problems for metric graphs ..... 31
2.1 Introduction ..... 31
2.2 Sobolev space and Dirichlet Energy of a rectifiable set ..... 32
2.2.1 Optimization problem for the Dirichlet Energy on the class of connected sets ..... 35
2.3 Sobolev space and Dirichlet Energy of a metric graph ..... 39
2.3.1 Optimization problem for the Dirichlet Energy on the class of metric graphs ..... 41
2.4 Some examples of optimal metric graphs ..... 48
2.5 Complements and further results ..... 57
3 A non-local isoperimetric problem ..... 59
3.1 Introduction ..... 59
3.2 The Riesz potential energy ..... 61
3.3 Relaxation of the problem ..... 68
3.4 Existence and characterization of minimizers under a regularity condition ..... 72
3.5 Stability of the ball ..... 75
II Quantitative stability problems ..... 83
4 Weighted isoperimetric inequalities in quantitative form ..... 85
4.1 Introduction ..... 85
4.2 Weighted isoperimetric inequality for the Lebesgue measure ..... 87
4.2.1 Nearly spherical ellipsoids ..... 92
4.3 Weighted isoperimetric inequalities for exponential measures ..... 94
5 Stability for the first Stekloff-Laplacian eigenvalue ..... 99
5.1 Introduction ..... 99
5.2 Spectral optimization for Stekloff eigenvalues ..... 101
5.3 The stability issue ..... 104
5.4 Sharpness of the quantitative Brock-Weinstock inequality ..... 106
5.4.1 Step 1: setting of the construction and basic properties ..... 106
5.4.2 Step 2: improving the decay rate ..... 112
5.4.3 Step 3: nearness estimates ..... 115
5.4.4 Step 4: conclusion ..... 121
6 A reduction theorem for the stability of Gagliardo-Nirenberg-Sobolev inequalities ..... 123
6.1 Introduction ..... 123
6.2 Continuity of $\lambda$ with respect to $\delta$ via a compactness theorem ..... 126
6.3 Reduction to $d$-symmetric functions ..... 131
6.4 Reduction inequalities ..... 134
7 Estimate of the dimension of the singular set of the MS functional: a short proof ..... 151
7.1 Introduction ..... 151
7.2 Caccioppoli partitions ..... 153
7.3 Proof of the main result ..... 154
Bibliography ..... 157

## Introduction

## Preamble

In this thesis we address some problems related to two topics in the Calculus of Variations which have attracted a growing interest in recent decades. In a general simplification we can refer to those two fields as optimization problems for shapes and for solutions of elliptic equations and quantitative stability problems for geometric and functional inequalities. It is worth bearing in mind that the historical and mathematical development of these two classes of problems have merged and looking at them as separate fields may not be the best approach to adopt. On the other hand a division of the works presented may simplify the reading, for this reason the thesis is divided into two main parts. In the first, Part I, containing three chapters, we deal with optimization problems related to the shape optimization field. In Part II, we address the quantitative stability of three problems, the first one regarding a class of isoperimetric inequalities, the second one about a spectral optimization problem and the third one concerning a class of functional inequalities.

The introduction, divided as well into two parts, is aimed to offer the basic background to get the reader into the two main fields taken into consideration, and to describe the organization and scopes of each chapter. More precisely in the first part we present a short introduction of the concept of optimization problems for shapes and for solutions of elliptic equations and recall some important examples, then we briefly describe the original contributions related to these subjects contained in the first three chapter of the thesis.

In the second part of the introduction, we describe the general concept of quantitative stability of a functional inequality, and then we delve into the state of the art of this subject by means of three important examples. After that, we describe the original contributions of the thesis in this topic, contained in Chapters 4, 5 and 6 .

Eventually, we briefly describe the results of the last chapter, which is unrelated to the rest of the thesis. For the sake of brevity this introduction does not contain all the technical preliminaries needed in the thesis, which are given, when necessary, within each chapter. All the results contained in this thesis are part of papers published, submitted or in preparation, and are the outcome of various collaborations which are cited both in this introduction and at the beginning of each chapter.

## Part I: Optimization problems for shapes and solutions of elliptic equations

An optimization problem in its general formulation takes the form

$$
\begin{equation*}
\min \{F(x): x \in X\} \tag{0.1}
\end{equation*}
$$

for a given set $X$ and a cost functional $F: X \rightarrow \mathbb{R}$. The general definition of a shape optimization problem is simply the case where $X$ is a subset of the powerset of $\mathbb{R}^{d}$, $\mathcal{P}\left(\mathbb{R}^{d}\right)$. So it is worth formulating problem (0.1) in the shape optimization's sense:

$$
\begin{equation*}
\min \{F(E): E \subset \mathcal{A}\}, \quad \mathcal{A} \subset \mathcal{P}\left(\mathbb{R}^{d}\right) \tag{0.2}
\end{equation*}
$$

The main issues one usually addresses regarding problem (0.2) are

- existence and uniqueness of a minimizer;
- qualitative properties of a minimizer as regularity, symmetry or more specific properties (e.g. convexity).

This class contains another sub-class of problems, that is when $F(E)=G\left(u_{E}\right)$ where $G$ is a suitable real function and $u_{E}$ is the solution of a given PDE on the domain $E$. For the sake of clearness, and since the two problems are strictly connected in many cases, we consider in this introduction mainly the case of shape optimization problems, keeping in mind that most of the forthcoming assertions are valid for both the cases.

In any case it is surely a good starting point to wonder how to prove the existence of a minimizer of the problem (0.2). This is usually done by means of the Direct Method in the Calculus of Variation (or more briefly: Direct Method). Namely, we consider a minimizing sequence of sets $E_{h} \subset \mathcal{A}$, that is a sequence such that

$$
F\left(E_{h}\right) \rightarrow \inf _{\mathcal{A}} F,
$$

and try to construct a set $E$ as limit point of the sets $E_{h}$ 's (this request is often abbreviated as a compactness request) such that $F$ is lower semicontinuous at least along this sequence, that is $\underline{\lim }_{h} F\left(E_{h}\right) \geq F(E)^{1}$. So the real problem in applying this classical and simple method is the choice of the right topology. Clearly the two requests needed to apply the Direct Method are in competition: to gain compactness we are led to ask the topology to be coarse enough, while to get semicontinuity, we would like it to be as fine as possible. Not surprisingly, it turns out that to attack different categories of problems we need different kinds of topologies. Immediately following we describe some of the main examples of problems in shape optimization which entail different kinds of topologies, with a twofold scope: to penetrate more deeply inside the world of the optimization problems for shapes and for solutions of elliptic equations, and to introduce some of the main mathematical tools and objects we deal with in Part I (but also in the rest of the thesis).

[^0]
## The isoperimetric problem

One of the most natural topologies on the class of sets is the $L^{1}$ topology: we say that $E_{h} \rightarrow E$ in $L^{1}$ if the Lebesgue measure of their symmetric difference $\left|E \Delta E_{h}\right|$ converges to 0 as $h \rightarrow \infty$. A very classical problem making use of this topology is the isoperimetric problem. This highly fascinating and ancient problem is the godmother of a huge class of problems in Geometric Measure Theory, in Shape Optimization and about geometric and functional inequalities. Indeed it is present, explicitly or not, in several topics of this thesis (actually most of them are present in Part II of the thesis).

A rough formulation of the problem is the following: find the container with minimal surface area which encompasses a fixed given volume. To translate it into precise mathematical language, we introduce the (proper!) object needed to measure the surface area of a subset of $\mathbb{R}^{d}$ : the perimeter. For a given set $E$ we recall the definition of the perimeter of $E$ relative to an open set $A$

$$
\begin{equation*}
P(E ; A)=\sup \left\{\int_{E} \operatorname{div}(\phi) d x: \phi \in C_{c}^{1}\left(\mathbb{R}^{d}, A\right), \quad|\phi| \leq 1\right\} \tag{0.3}
\end{equation*}
$$

first formulated by E. De Giorgi in [47]. If $A=\mathbb{R}^{d}$, we will denote the perimeter of $E$ as $P(E)$. We do not enter into the details of this topic and we refer the interested reader to the (beautiful) theory of sets of finite perimeter to [2] and to the recent book [92], but we limit ourselves to state the properties sufficient to solve the isoperimetric problem. It is not difficult to prove that the functional $P$ is lower semicontinuous with respect to the $L^{1}$-convergence, that is,

$$
P(E) \leq \underline{\lim }_{h \rightarrow \infty} P\left(E_{h}\right) \quad \text { if } \quad\left|E \Delta E_{h}\right| \rightarrow 0
$$

where $\left|E \Delta E_{h}\right|$ is the $d$-dimensional Lebesgue measure of the symmetric difference between $E$ and $E_{h}: E \Delta E_{h}:=\left(E \backslash E_{h}\right) \cup\left(E_{h} \backslash E\right)$. Moreover one can prove that the class of sets of equi-bounded perimeter, contained in a big fixed ball, is compact ${ }^{2}$. Then, by means of the Direct Method we get existence for the following version of the isoperimetric problem

$$
\begin{equation*}
\min \{P(E): E \subset B(0, R),|E|=c\} \tag{0.4}
\end{equation*}
$$

for two given positive constants $R$ and $c$. Although we will not enter into further detail, we mention the fact that the bound constraint on the sets is not necessary. Notice that in the definition of this problem we implicitly defined the topology we impose on the set: the $L^{1}$ topology (this notation is due to the fact that, if $\mathbf{1}_{A}$ and $\mathbf{1}_{B}$ are the characteristic functions of the Borel sets $A$ and $B$, then $\left.|A \Delta B|=\left\|\mathbf{1}_{A}-\mathbf{1}_{B}\right\|_{L^{1}\left(\mathbb{R}^{d}\right)}\right)$. Moreover we also know what the explicit solutions are: the balls (not just $a$ ball, since the problem is translation invariant). A complete proof of this fact has been given for the first time

[^1]in the celebrated paper of E. De Giorgi [47] by means of symmetrization techniques and followed later on by many other proofs.
An obvious but important property of the $L^{1}$ topology is that it is not sensible if we modify a set $E$ by adding (or subtracting) a set with zero $L^{1}$ mass. This turns out to be a problem when one considers the other main class of problems related to elliptic operators described in the following example.

## Spectral and Energy Function problems

A good guide to optimization problems related to spectral and Energy functionals can be the books [75], [77] and [25]. Here we introduce a couple of basic examples which are the main objects, together with the perimeter, we deal with in the first two chapters of Part I, and which engrafted a rich and interesting theory. Consider an open bounded Lipschitz set $\Omega$ and a function $f \in L^{2}(\Omega)$. Then we say that $u$ is a solution of the Dirichlet problem:

$$
\begin{cases}-\Delta u=f & \text { in } \Omega  \tag{0.5}\\ u=0 & \text { on } \partial \Omega\end{cases}
$$

if $u$ is the (unique) solution of the variational problem

$$
\left\{\begin{array}{l}
\int_{\Omega}\langle\nabla u, \nabla v\rangle-f v d x=0 \quad \forall v \in H_{0}^{1}(\Omega)  \tag{0.6}\\
u \in H_{0}^{1}(\Omega)
\end{array}\right.
$$

By means of the Poincaré inequality, and the Lax-Milgram Theorem (see for instance [57]) we know that there exists a solution $u_{\Omega}$ for problem (0.5) characterized as the unique minimum of the following variational problem:

$$
\begin{equation*}
\min \left\{\int_{\Omega}|\nabla u|^{2} d x-2 \int_{\Omega} u f d x: u \in H_{0}^{1}(\Omega)\right\} . \tag{0.7}
\end{equation*}
$$

Notice that by multiplying in equation (0.5) by the solution $u_{\Omega}$ and integrating by parts, we get that the quantity in (0.7) is equivalent to

$$
\begin{equation*}
\mathcal{E}_{f}(\Omega)=-\frac{1}{2} \int_{\Omega} u_{\Omega} f d x \tag{0.8}
\end{equation*}
$$

Before introducing the shape optimization problem we have in mind about the functional $\mathcal{E}_{f}$, we introduce the other class of functionals we are principally interested in: spectral functionals. To introduce them we recall that the resolvent of the DirichletLaplacian $-\Delta$ on $\Omega$, that is the functional that associates to every function $f \in L^{2}(\Omega)$ the solution of (0.5), is a compact, symmetric and (thanks to the minus sign) positive operator on $L^{2}(\Omega)$. Thus by classical results in functional analysis (see for instance [19]) it admits a positive spectrum, consisting in a sequence $\Lambda_{k}$ of eigenvalues, converging to 0 . As a consequence, there is a sequence of positive numbers $\lambda_{k}:=\lambda_{k}(\Omega)=1 / \Lambda_{k}$
(eigenvalues) accumulating at infinity, and of functions $u_{k} \in H_{0}^{1}(\Omega)$ (eigenfunctions) such that

$$
-\Delta u_{k}=\lambda_{k} u_{k}
$$

As in the case of the Energy Function, we would like to have a variational version to compute the eigenvalues. This is offered us by the Courant-Fischer Formulae, see [42]:

$$
\begin{equation*}
\lambda_{k}(\Omega)=\min \left\{\mathcal{R}_{\Omega}[u]:=\frac{\int_{\Omega}|\nabla u|^{2} d x}{\int_{\Omega} u^{2} d x}: \int_{\Omega} u u_{i} d x=0, i=0, \ldots, k-1\right\} \tag{0.9}
\end{equation*}
$$

The functional $\mathcal{R}_{\Omega}[u]$ is often referred to as Rayleigh quotient of $u$. With these examples in mind we can formulate optimization problems in the sense of $(0.2)$ where $F(\Omega)$ is, for instance, $\lambda_{k}$ or the Energy function $\mathcal{E}_{f}$. Other more general cost functionals $F$ that can be considered are spectral functionals that is $F(\Omega)=\Phi(\lambda(\Omega))$ with $\Phi: \mathbb{R}^{\mathbb{N}} \rightarrow \mathbb{R}$ is a given function and where $\lambda(\Omega)=\left(\lambda_{1}(\Omega), \lambda_{2}(\Omega), \ldots\right)$ is the spectrum of $-\Delta$ on $\Omega$, and integral functionals, where $F$ takes the form

$$
F(\Omega)=\int_{\Omega} j\left(x, u_{\Omega}(x), \nabla u_{\Omega}(x)\right) d x
$$

where $u_{\Omega}$ is the solution of (0.7) and $j$ is a suitable integrand which is usually assumed to be convex in the gradient variable and bounded from below.

The $L^{1}$ topology is not the right topology to afford the existence problem for the functionals above since they are sensible under variation of zero $L^{1}$ volume (think for example to the case of a ball in $\mathbb{R}^{2}$ and the same ball without a diameter). Indeed the crucial quantity to look at while affording shape optimization problems related to elliptic equations, is the capacity of a set instead of its Lebesgue measure. Again, for the sake of brevity we will not enter into the details of this interesting theory, which will be however introduced together with its main properties, in Section 1.2 of Chapter 1, and we refer the reader to the book [25]. Concerning the aforementioned shape optimization problems, we cite [30] where a general theorem is proved which entails existence for a large class of functionals, under the constraint of equiboundedness of the sets. As in the case of the perimeter such bound is not necessary, but the proof of this fact is quite involved and very recent (see [96] and [24]).

## Results of Part I of the Thesis

We describe hereafter the main results of the first three chapters of the thesis.

## Optimization problems related to Schrödinger operators

In the first chapter, based on the joint work [31] together with Giuseppe Buttazzo, Augusto Gerolin and Bozhidar Velichkov, we consider a problem of the form

$$
\begin{equation*}
\min \{F(V): V \in \mathcal{V}\} \tag{0.10}
\end{equation*}
$$

where $V: \mathbb{R}^{d} \rightarrow \mathbb{R}$ is a non-negative real function and $\mathcal{V}$ a suitable constraint which can take different forms depending on the functional we deal with. For every admissible potential $V$ we consider the Schrödinger equation formally written as

$$
-\Delta u+V u=f \quad u \in H_{0}^{1}(\Omega)
$$

and whose precise meaning has to be given in the weak form

$$
\left\{\begin{array}{l}
\int_{\Omega}\langle\nabla u, \nabla \varphi\rangle d x+\int_{\Omega} V u \varphi d x=\int_{\Omega} f \varphi d x \quad \forall \varphi \in H_{0}^{1}(\Omega) \\
u \in H_{0}^{1}(\Omega) \cap L^{2}(V d x)
\end{array}\right.
$$

One of the main tools we exploit, to deal with this problem, is the $\gamma$-convergence for capacitary measures, introduced, together with its main properties, in Subsection 1.2. Thanks to this, we are able to prove two existence results in the case where $\Omega$ is bounded for the problem (0.10), valid for a large class of cost functions $F$, where the constraint is taken to be a weakly compact set in $L^{p}(\Omega)$. Then we consider a volume constraint for $V$ (not necessarily bounded in $\left.L^{p}(\Omega)\right)$ of the form

$$
\mathcal{V}=\left\{V: \int_{\Omega} \Psi(V) d x \leq c\right\}
$$

where $\Psi:[0, \infty] \rightarrow[0, \infty]$ is a given function and, under some mild hypotheses on $F$, we prove existence for problem (0.10) for a class of constraints where the main example is $\Psi(t)=t^{-p}, p>0$. Then, we afford the existence issue for the particular case of the Dirichlet Energy function under a convex $L^{p}$ constraint for $V$ and with $\Omega$ bounded. In this case we are able to prove existence and to characterize the solutions in terms of the solutions of a particular PDE in $\mathbb{R}^{d}$. Later on we turn our attention to the case of $\Omega$ unbounded. In this case it is more difficult to get existence (in particular, it is more delicate to get compactness for a minimizing sequence). So we focus to the particular cases of the minimization and the maximization (under suitable volume constraints) of the first eigenvalue and the Dirichlet Energy function related to the Schrödinger operator $-\Delta+V$. In the particular case where $\Omega=\mathbb{R}^{d}$ and $F=\lambda_{1}$, after solving the existence problem, we study some qualitative properties of the solutions, such as their symmetry. Moreover we state a relation between the eigenfunction $u_{V}$ related the optimal potentials $V$ and the solutions of a class of celebrated functional inequalities: the Gagliardo-Nirenberg-Sobolev inequalities (GNS). This lead us to study, as in the case of the GNS inequalities (see for instance [104]), the support of the optimal potentials. Eventually, thanks to the characterization of the supports of such functions, we are able to describe qualitatively the shape of the potentials $V$ which minimize (under suitable constraints) the first and the second eigenvalue of the Schrödinger operator in $\mathbb{R}^{d}$.

## Optimization problems for metric graphs

In the second chapter, based on the joint work [32] with Giuseppe Buttazzo and Bozhidar Velichkov we study the problem of minimizing the first eigenvalue of the Laplacian and its Energy function in the case where $\Omega$ is a graph. By a graph we will initially mean a connected one-dimensional set; thanks to classical theory about such sets, see for instance [4], these sets are rectifiable so that quantities as the Dirichlet Energy or as the first eigenvalue of the Dirichlet Laplacian are well defined. After having developed the main tools needed to state and afford the problem in such setting, we prove by means of a counterexample that in general an existence result without additional assumptions cannot be expected. Thus we consider the larger class of metric graphs (see Section 2.3) where we are able to prove a general existence theorem. Then we develop a symmetrization technique, which may be seen as analogous to the Pólya-Szëgo inequality in this particular setting, and by means of such technique, we describe some specific examples of optimal graph for the Energy Function.

## A non-local isoperimetric problem

In the third chapter, based on a work in preparation with Michael Goldman and Matteo Novaga [70], we deal with a non-local isoperimetric problem. More precisely, we consider problem (0.2) where $\mathcal{A}$ is the class of sets with fixed volume, say $c$, and $F$ is given by

$$
\begin{equation*}
F(E)=P(E)+\mathcal{Q}_{\alpha}(E) \tag{0.11}
\end{equation*}
$$

where

$$
\mathcal{Q}_{\alpha}(E)=\min \left\{\int_{\mathbb{R}^{d} \times \mathbb{R}^{d}} \frac{d \mu(x) d \mu(y)}{|x-y|^{\alpha}}: \mu(E)=Q\right\}
$$

for $d \geq 2$ and $\alpha \in(0, d-1)$. Functionals of this kind are well known in the literature since they are related to the physical problem of finding the optimal shape of a liquid droplet in void, once it is provided with an electric charge $Q$. Similar problems have recently been addressed in several works, see for instance [80], [81] and the references therein. But the main case afforded is that where the charge is prescribed to be uniform on the set. In this chapter we prove that problem (0.11) is actually ill-posed: minimizing $F$ is equivalent to minimizing separately the two functionals $P$ and $\mathcal{Q}_{\alpha}$. This is not surprising if we notice that these two functionals are actually governed by different kind of measures: the first is invariant under $L^{1}$-negligible perturbations, while the second may be (depending on the optimal measure $\mu$ ) a capacitary functional (that is sensible up to positive capacitary perturbations). However, we are able to prove that under further regularity conditions, the existence problem is solvable. Moreover (at least) in the harmonic case $\alpha=d-2 \geq 1$, if we restrict the set of competitors just to sets which are small $C^{1,1}$ perturbation of the ball, the ball is the unique minimum.

## Part II: Stability problems

In recent years there has been renewed interest in the research of a quantitative stability version for several classes of inequalities. We can find several examples of quantitative versions of geometric inequalities see e.g [62], [58], [73], [72], [60], [61], functional inequalities, see e.g. [12], [8], [39], [34], [35], and of inequalities arising from the optimization of shapes and solutions of elliptic equations, see e.g. [18], [15]. The motivation of this interest probably resides in the recent work [62] where it has been proved a sharp quantitative version of the isoperimetric inequality.

Let us begin by explaining what we mean by quantitative stability of an inequality. The starting point is again an optimization problem of the form

$$
\min \{F(x): x \in X\} .
$$

Suppose that we have solved this problem, and found a class of minimizers of $F$ in X : $M \subset X$. Then we can translate our solution in terms of an inequality of the form

$$
\begin{equation*}
F(x)-F(y) \geq 0 \quad \forall x \in X, \quad y \in M . \tag{0.12}
\end{equation*}
$$

The quantitative stability question is then the following:
suppose that a point $z \in X$ is such that we almost get equality in (0.12). Can we claim then that $z$ is near the set of minimizers $M$ ?
The previous question, in the form it is stated, lacks at least formally into two points: first, we should formalize what we mean by almost getting equality; then we should specify in which sense we mean nearness. These are the crucial and starting points of any analysis about the quantitative version of an inequality, and can vary case by case.

To get into the concept of quantitative stability, we quote three examples of stability inequalities: the first one describing the state of the art for the isoperimetric inequality in quantitative version, the second one about the stability of some spectral problems, with particular regard to that related to the first eigenvalue of the Dirichlet Laplacian, and the third one regarding a class of functional inequalities, with particular attention to the Sobolev and Gagliardo-Nirenberg-Sobolev inequalities. After this overview, we pass to the description of the original results contained in the second part of the thesis: one stability problem about a class of isoperimetric inequalities, one related to the stability of the spectral problem of the Stekloff-Laplacian, and one which concerns the stability of the Gagliardo-Nirenberg-Sobolev inequalities.

## The isoperimetric inequality in quantitative form

We recall once again that the isoperimetric inequality states that if $E$ is a measurable subset of $\mathbb{R}^{d}$ of finite measure and if $B_{E}$ is a ball of the same measure of $E$, then $P(E) \geq P\left(B_{E}\right)$. Thanks to the scaling laws satisfied by the perimeter and by the $d$-dimensional Lebesgue measure, we can avoid to prescribe the measure of $E$

$$
|E|^{(1-d) / d} P(E) \geq d \omega_{d}^{1 / d}
$$

where $\omega_{d}$ is the measure of the unitary ball of $\mathbb{R}^{d}$. If we introduce the isoperimetric deficit of a set $E$ as

$$
\delta(E)=\frac{P(E)}{d \omega_{d}^{1 / d}|E|^{(d-1) / d}}-1
$$

then the isoperimetric inequality becomes

$$
\begin{equation*}
\delta(E) \geq 0 \tag{0.13}
\end{equation*}
$$

Then for a quantitative version of the isoperimetric inequality we mean an improvement of $(0.13)$ of the form

$$
\begin{equation*}
\delta(E) \geq \Phi\left(\operatorname{dist}\left(E, M_{E}\right)\right) \tag{0.14}
\end{equation*}
$$

where $M_{E}$ is the manifold composed of all the balls of $\mathbb{R}^{d}$ of measure $|E|$, $\operatorname{dist}(\cdot, M)$ is a suitable distance of $E$ from $M_{E}$ and $\Phi$ is a modulus of continuity. Often the distance adopted for the isoperimetric problem is the so-called Fraenkel asymmetry defined as follows:

$$
\mathcal{A}(E)=\inf \left\{\frac{|E \Delta B|}{|E|}: B \text { ball, }|B|=|E|\right\}
$$

where $E \Delta B=(E \backslash B) \cup(B \backslash E)$ denotes the symmetric difference between $E$ and $B$. Notice that the choice of the isoperimetric deficit is not free: it is just the (rescaled) excess offered by the inequality itself. This is not the case of the asymmetry, whose choice is a crucial starting point for the analysis of the stability. Several partial results in the direction of proving an inequality of the form (0.14) have been obtained. For instance in [60], where the distance taken in consideration is actually stronger than the Fraenkel asymmetry, but the result is valid only for sets whose boundary is a small $W^{1, \infty}$ perturbation of the boundary of the ball (the so-called nearly spherical sets). The first remarkable results valid for any set of finite perimeter has been obtained in [72] and [73], where a quantitative version of the stability inequality of the form

$$
\delta(E) \geq C(d) \mathcal{A}(E)^{4}
$$

is proved. The inconvenience in this inequality is the exponent 4. Indeed by computing the deficit and the asymmetry on a family of ellipsoids converging to the ball, the guess is that the expected optimal exponent should be 2 . This is actually the case, as proved recently in three works. The first one, in chronological order, takes its moves from the works [72] and [73] and exploits a clever symmetrization method. The second one [58] makes use of some results coming from the Mass Transport Theory and in particular of the proof of the isoperimetric inequality done by means of the Brenier map (see [110] for a comprehensive account on the subject). In this case the authors are able to prove the isoperimetric inequality also for anisotropic perimeters. The third paper [40], develops a strategy called Selection Principle which, combined with the results in [60] (the analysis of the nearly spherical sets) and classical regularity theory for sets of finite perimeter, brings again to the sharp exponent 2. It is worth mentioning that,
although all three proofs have their own interests, the Selection Principle developed in [40] has been successfully applied to prove the stability of several other inequalities, as in [15], [1], [53], [13], [14].

## Spectral inequalities in quantitative form

As mentioned, since the development of the techniques used for proofs of the sharp quantitative version of the isoperimetric inequality, there have been several authors who have tried to attack other inequalities. A class of inequalities which attracted some interest is that of the inequalities coming from the shape optimization theory, with particular regard to spectral problems. As a basic example we consider the first eigenvalue of the Dirichlet Laplacian, introduced in the previous subsection. We recall that for a given bounded set $\Omega$ with Lipschitz boundary, it is characterized as

$$
\lambda_{1}(\Omega)=\min \left\{\frac{\int_{\Omega}|\nabla u|^{2} d x}{\int_{\Omega} u^{2} d x}: 0 \neq u \in H_{0}^{1}(\Omega)\right\}
$$

The problem we introduce is then

$$
\min \left\{\lambda_{1}(\Omega):|\Omega|=c\right\}
$$

The solution of this problem can be stated by means of the celebrated Faber-Krahn inequality

$$
\lambda_{1}(\Omega) \geq \lambda_{1}(B)
$$

which holds for any ball $B$ such that $|\Omega|=|B|$. As in the case of the perimeter, the first eigenvalue of the Dirichlet-Laplacian satisfies a scaling law: $\lambda_{1}(t \cdot)=t^{-2} \lambda_{1}(\cdot)$ so that the previous inequality becomes, without prescribing the volume of the sets,:

$$
\begin{equation*}
|\Omega|^{d / 2} \lambda_{1}(\Omega) \geq \omega_{d}^{d / 2} \lambda_{1}\left(B_{1}\right), \tag{0.15}
\end{equation*}
$$

where $B_{1}$ is any ball of radius 1 . It is worth mentioning the connection for this inequality with the isoperimetric one. Indeed (0.15) is usually proved via the Schwarzsymmetrization and in particular by means of the Pólya-Szëgo inequality: given a measurable function $u$, let $u^{*}$ be the non-negative function which level sets $\left\{u^{*}>t\right\}$ are balls centred at the origin of measure $|\{|u|>t\}|$, then the Pólya-Szëgo inequality states that

$$
\int|\nabla u|^{2} d x \geq \int\left|\nabla u^{*}\right|^{2} d x
$$

Moreover, in [23] is proved that the equality can holds only if $u$ has not flat regions. These results immediately imply that the ball is the unique minimizer (up to translations and dilations) for $\lambda_{1}$. For a quantitative version of the Faber-Krahn inequality (0.15) we mean then an inequality of the form

$$
|\Omega|^{d / 2} \lambda_{1}(\Omega)-\omega_{d}^{d / 2} \lambda_{1}\left(B_{1}\right) \geq \sigma \mathcal{A}(\Omega)^{\alpha}
$$

where $\sigma$ and $\alpha$ are fixed constants depending on $d$. Partial results in this direction has been obtained (in any dimension) in [64] where, by exploiting the sharp quantitative form of the isoperimetric inequality, the authors get the non-sharp exponent 4 ( 3 in dimension $d=2$ ). For a sharp version of this inequality, that is with $\alpha=2$, we refer to the recent work [15] which suitably exploits a selection principle argument, in the spirit of [40].

We will not enter into further details in this introduction about quantitative versions for spectral inequalities, but we mention that many other problems have been afforded about higher eigenvalues of the Dirichlet-Laplacian or about spectral problems related to the Laplacian with different kind of boundary conditions. See for instance [18] and the references therein.

## Stability of functional inequalities

The last example we shall deal with and we briefly examine is that of the quantitative stability of functional inequalities. The principal examples in this area are given by the Sobolev and the Gagliardo-Nirenberg-Sobolev (briefly: GNS) inequalities. These inequalities, exhaustively described in Chapter 6, take the general form

$$
\begin{equation*}
\|u\|_{L^{d p /(d-p)}\left(\mathbb{R}^{d}\right)} \leq C(d, p)\|\nabla u\|_{L^{p}\left(\mathbb{R}^{d}\right)} \tag{0.16}
\end{equation*}
$$

for the Sobolev case, and

$$
\begin{equation*}
\|u\|_{L^{q}\left(\mathbb{R}^{d}\right)} \leq C(d, p, q, s)\|\nabla u\|_{L^{p}\left(\mathbb{R}^{d}\right)}^{\theta}\|u\|_{L^{s}\left(\mathbb{R}^{d}\right)}^{1-\theta} \tag{0.17}
\end{equation*}
$$

in the GNS case, for suitable parameters $d, p, q, s$ and $\theta$ (for precise definitions, we refer again to Chapter 6). The solutions and optimal constants for the Sobolev inequality are known after the works of Aubin and Talenti [6], [109]. In the case of the GNS inequalities, instead, we know the solutions and the optimal constants only for a given one-parameter family of exponents [48]. See also [43] for an approach to the proofs of such inequalities using the Mass Transport theory. An important result about the quantitative form for the Sobolev inequality has been given in [12], in which only the case of $p=2$ is considered and it is crucially exploited the fact the in this case the EulerLagrange equation associated to the Sobolev inequality is linear in the gradient part, and also that the set of solutions have an Hilbertian structure. The ideas developed in [12] have then been used to prove a stability version for a class of GNS inequality (again prescribing $p=2$ ) in the recent work [35].

After the proof of the sharp isoperimetric inequality in quantitative form, there have been several attempts to prove a sharp quantitative version for Sobolev and GNS inequalities. Regarding the Sobolev case successful attempts have been done in [63] for $W^{1,1}$ functions (actually also for functions of bounded variation). See also the recent work [54] where the case $p=2$ is considered. On the contrary, for the GNS case there have been proved quantitative versions only in the one-parameter class whose solutions are known or for the simplified structure obtained by prescribing $p=2$, see [34], [35].

It is important then to remark about the work [39], where a non-sharp version of the Sobolev inequality is proved by means of the sharp version of the isoperimetric inequality. Indeed it is well known that the Sobolev inequality is implied by the isoperimetric one, and the proof of this fact basically relies on the co-area formula and a rearrangement argument (see [2]). The proof in [39] makes use of this implication and then exploits the stability version of the isoperimetric inequality. One benefit of this idea is the geometric flavour of the proof, which turns out to be adaptable also to inequalities whose optimal cases are not explicit, but whose qualitative geometric aspects (as being radial functions) are known. This is the case, as we will see in Chapter 6, of the GNS inequalities.

## Results of Part II of the thesis

We describe hereafter the main results of the chapters contained in the second part of the thesis. Each one is somehow related to one of the three examples introduced above.

## Weighted isoperimetric inequalities

The fourth chapter is based on part of a joint work with Lorenzo Brasco and Guido De Philippis [16] and addresses the problem of the quantitative stability of weighted isoperimetric inequalities. Namely, given a Lipschitz set $E$ we consider its weighted perimeter

$$
P_{V}(E)=\int_{\partial E} V(x) d x
$$

where $V$ is a non-negative weight function, and we address the problem

$$
\begin{equation*}
\min \left\{P_{V}(E):|E|=c\right\} \tag{0.18}
\end{equation*}
$$

where $|\cdot|$ stands for the $d$-dimensional Lebesgue measure of a set. Such a problem has been solved for a class of radial functions $V$ in [9] where it is proved, thanks to a symmetrization technique, that under some suitable hypotheses on $V$ the ball centred at origin is the unique minimum. We first afford the same issue, offering a totally different proof of it (for the same class of functions $V$ ) by means of a sort of calibration technique, which allows us to get also optimal stability. Then we consider the analogous problem related to the exponential measures

$$
\min \left\{P_{w e^{V}}(E): \int_{E} e^{V} d x=c\right\}
$$

proving that under suitable hypotheses on $V$ and $w$ (having as main application the Gauss measure, i.e. $\left.V=-|x|^{2}\right)$ the right half spaces $\left\{\left(x_{1}, x^{\prime}\right) \in \mathbb{R} \times \mathbb{R}^{d-1}: x_{1}>t\right\}$ are solutions of the problem.

## Stability of the first Stekloff-Laplacian eigenvalue

The fifth chapter (as well as the fourth) is based on the work [16]. In it we consider the first non-trivial eigenvalue of the Stekloff-Laplacian which can be characterized, for a given Lipschitz set $\Omega$ as the number $\sigma_{2}(\Omega)$ such that there exists a (unique) solution of the PDE

$$
\begin{cases}-\Delta u=0 & \text { in } \Omega \\ \left\langle\nabla u, \nu_{\Omega}\right\rangle=\sigma_{2}(\Omega) u & \text { on } \partial \Omega\end{cases}
$$

where $\nu_{\Omega}$ is the outer normal of $\Omega$. Weinstock in dimension $2[111]^{3}$ and later Brock [20] in any dimension proved that among all the sets of prescribed measure $c>0$, the ball maximizes $\sigma_{2}(\cdot)$. This translates in the so-called Brock-Weinstock inequality

$$
\sigma_{2}(\Omega) \leq \sigma_{2}\left(B_{\Omega}\right) \quad B_{\Omega} \text { ball such that }\left|B_{\Omega}\right|=|\Omega|
$$

A crucial tool in the Brock's proof, is the following weighted isoperimetric inequality:

$$
\int_{\partial \Omega}|x|^{2} d \mathcal{H}^{d-1} \geq \int_{\partial B_{\Omega}}|x|^{2} d \mathcal{H}^{d-1}
$$

where $\Omega$ is any set and $B_{\Omega}$ is the ball centred at the origin such that $\left|B_{\Omega}\right|=|\Omega|$. In this chapter, as a consequence of the the stability of the weighted isoperimetric inequalities proved in Chapter 4, we prove the stability of $\sigma_{2}$ as well. Then we prove that the quantitative version of the inequality is sharp. This is done thanks to a long and delicate proof modelled on the corresponding one related to Neumann eigenvalues in [18].

## A reduction theorem for the stability of Gagliardo-Nirenberg-Sobolev inequalities

In the sixth chapter we deal with the Gagliardo-Nirenberg-Sobolev inequalities (0.17) and we prove a reduction theorem which entails that to prove a quantitative version for such inequalities, it is sufficient to prove it only for radial decreasing functions. Namely we address the problem of proving the following inequality

$$
\begin{equation*}
\delta_{G N S}(u):=\frac{\|\nabla u\|_{L^{p}\left(\mathbb{R}^{d}\right)}^{\theta}\|u\|_{L^{s}\left(\mathbb{R}^{d}\right)}^{1-\theta}}{G\|u\|_{L^{q}\left(\mathbb{R}^{d}\right)}}-1 \geq C \lambda_{G N S}(u)^{\alpha_{0}} \tag{0.19}
\end{equation*}
$$

where $G=G(d, p, q, s)$ is the optimal constant in (0.17), $C$ and $\alpha_{0}$ are positive constants (which do not depend on $u$ ), and

$$
\lambda(u)=\inf \left\{\frac{\|u-v\|_{L^{q}\left(\mathbb{R}^{d}\right)}}{\|u\|_{L^{q}\left(\mathbb{R}^{d}\right)}}:\|v\|_{L^{q}\left(\mathbb{R}^{d}\right)}=\|u\|_{L^{q}\left(\mathbb{R}^{d}\right)}, \quad v \text { is optimal for }(0.17)\right\}
$$

[^2]We prove that if (0.19) holds for radial decreasing functions, then, up to substitute $\alpha_{0}$ with another suitable positive constant $\alpha_{1}$, it holds true for any function. This allows us to simplify the problem in its general formulation, and in particular reduces a $d$-dimensional problem to a 1 -dimensional one. The results proved are valid also for the classes of parameters for which the optimal function of the GNS inequality (0.17) is not known. This is possible thanks to an adaptation of the symmetrization techniques developed in [39], where only the knowledge of the qualitative shape of the optimizer for the Sobolev functions (i.e. radial decreasing functions) is exploited.

## Other works

In the final chapter we report a recent note written in collaboration with Camillo De Lellis and Matteo Focardi [51], and not related to the rest of the thesis. In this, we consider the localized Mumford-Shah functional for a bounded open domain $\Omega$

$$
\operatorname{MS}(v)=\int_{\Omega}|\nabla v|^{2} d x+\mathcal{H}^{d-1}\left(S_{v} \cap \Omega\right), \quad \text { for } v \in S B V(\Omega) \text { and } A \subseteq \Omega \text { open. }
$$

where $S B V(\Omega)$ is the space of Special Bounded Variation functions and $S_{v}$ is the jump set of the function $v$ (see [2]). More precisely, we consider a local minimizer $u$ for MS, that is a function in $S B V$ such that $\operatorname{MS}(u) \leq \operatorname{MS}(v)$ for any $v$ such that $\{u \neq v\} \subset \subset \Omega$. Let us denote with $\Sigma_{u}$ the set of points of $S_{u}$ out of which $u$ is locally regular and let

$$
\Sigma_{u}^{\prime}=\left\{x \in \Sigma_{u}: \lim _{\rho \rightarrow 0} \frac{1}{\rho^{d-1}} \int_{B_{\rho}(x)}|\nabla u|^{2} d x=0\right\}
$$

A classical result due to L. Ambrosio, N. Fusco and J.E. Hutchinson [3, Theorem 5.6], states that the dimension of $\Sigma_{u}^{\prime}$ is less or equal than $d-2$. The proof is quite involved and makes use of the concept of Almgren minimizers (see [2]). In this chapter we offer a simplified proof of the same result exploiting a recent work of C. De Lellis and M. Focardi [50] where it is proved that if $u$ is a minimizer of MS, then the blow-up of its jump set in small gradient regime, that is $\Sigma_{u}^{\prime}$, is a Caccioppoli partition.

## Part I

# Optimization problems for shapes and for solutions of elliptic equations 

## Chapter 1

## Optimization problems related to Schrödinger operators

### 1.1 Introduction

This chapter is based on the joint work with Giuseppe Buttazzo, Augusto Gerolin and Bozhidar Velichkov [31]. The problem we consider takes the general form

$$
\begin{equation*}
\min \{F(V): V \in \mathcal{V}\}, \tag{1.1}
\end{equation*}
$$

where $F$ denotes a cost functional and $\mathcal{V}$ is an admissible class of functions $V: \mathbb{R}^{d} \rightarrow \mathbb{R}$. In this chapter, we will always suppose that $V$ is non-negative and we will refer to it as potential. The principal examples we shall take into account, and which are also the leading motivation for the study of problem (1.1) in our case, are problems related to the Schrödinger operator $\mathcal{S}_{V}=-\Delta+V(x)$. More precisely the two principal cases we study are the following.

Energy functionals Consider the PDE concerning the energy of $\mathcal{S}_{V}$, related to a suitable function $f$ :

$$
\begin{equation*}
-\Delta u+V(x) u=f, \quad u \in H_{0}^{1}(\Omega) . \tag{1.2}
\end{equation*}
$$

If such equation admits a solution $u_{V}$, which is the case under suitable assumption on $\Omega$ and $f$, for instance if $\Omega$ is bounded and $f \in L^{2}(\Omega)$, then one can consider $F(V)$ as the energy of $V$ :

$$
\begin{equation*}
\mathcal{E}_{f}(V)=\inf \left\{\int_{\Omega}\left(\frac{1}{2}|\nabla u|^{2}+\frac{1}{2} V(x) u^{2}-f(x) u\right) d x: u \in H_{0}^{1}(\Omega)\right\}=-\frac{1}{2} \int_{\Omega} u_{V} f . \tag{1.3}
\end{equation*}
$$

The last equality in the previous formula follows by multiplying by $u_{V}$ and integrating by part the Euler-Lagrange equation (1.2) of problem (1.3).

Spectral functionals The operator $\mathcal{S}_{V}=-\Delta+V(x)$ is positive definite (since $V \geq 0$ ) and symmetric. Moreover, under suitable assumption on $V$ (which we will always have
in force), its resolvent operator is a compact operator from $L^{2}(\Omega)$ to $L^{2}(\Omega)$, whence $\mathcal{S}_{V}$ admits a discrete, positive (and unbounded) spectrum

$$
\Lambda(V)=\left(\lambda_{1}(\Omega, V), \lambda_{2}(\Omega, V), \ldots\right)
$$

Since in all our analysis the set $\Omega$ will be fixed a priori, we will drop the dependence of $\Omega$ in the definition of $\lambda_{k}$ simply writing $\lambda_{k}(V)$. Then we can consider functionals of the form

$$
F(V)=\Phi(\Lambda(V))
$$

where $\Phi: \mathbb{R}^{\mathbb{N}} \rightarrow \mathbb{R}$ is an assigned function. A typical choice of such a function is given by $\Phi_{k}\left(\left(x_{1}, x_{2}, \ldots\right)\right)=x_{k}$, which, in terms of the functional $F$, corresponds to $F_{k}(V)=\lambda_{k}(V)$. In this case, we can express the functional also in a variational form:

$$
\begin{equation*}
\lambda_{k}(V)=\min \left\{\frac{\int_{\Omega}|\nabla u|^{2}+V(x) u d x}{\int_{\Omega} u^{2}}: u \in H_{0}^{1}(\Omega), \int_{\Omega} u u_{i} d x=0, i=1, \ldots, k-1\right\} \tag{1.4}
\end{equation*}
$$

where $u_{i}$ is an eigenfunction related to $\lambda_{i}$.
Although we do not intend to list the (many) known results (for which, a good guide is the book of Henrot [75, Chapter 8]), we mention that most of them are related to spectral problems, and in particular they are established for bounded domains $\Omega$. As an example we recall the following general result proved in [5].

Theorem 1.1. Let $\Omega$ be a $C^{1,1}$ bounded open set, $p>1$ and $\mathcal{V}$ be a closed bounded convex subset of $L^{p}(\Omega)$. Then, there exists a unique $\widetilde{V}$ which maximizes $\lambda_{1}(V)$ in the class $\mathcal{V}$.

Remark 1.2. The previous theorem addresses the maximization of the first eigenvalue $\lambda_{1}(V)$. The main reason is that the minimization of $\lambda_{1}$, as well as that of $\mathcal{E}_{f}$, often turns out to be a trivial or an ill-posed problem. For example, consider the constraint

$$
\mathcal{V}_{p}=\left\{V:\|V\|_{L^{p}(\Omega)} \leq 1\right\}
$$

Then the problem

$$
\min \left\{\lambda_{1}(V): V \in \mathcal{V}_{p}\right\}
$$

has the trivial solution $V=0$ since, by (1.4) we have $\lambda_{1}(V) \geq \lambda_{1}(0)=\lambda_{1}(\Omega)$. In the maximization case, however, the convexity of the constraint is not necessary. Indeed, thanks to the inequality $\lambda_{1}\left(V_{1}\right) \leq \lambda_{1}\left(V_{2}\right)$ if $V_{1} \leq V_{2}$ (due to (1.4)), maximizing $\lambda_{1}(V)$ on the constraint $\mathcal{V}_{p}$ is equivalent to maximize it under the constraint $\left\{V:\|V\|_{L^{p}(\Omega)}=1\right\}$. Also the minimization problem may turn out to be not interesting, even with a nonconvex constraint. For instance, consider a sequence of potentials $V_{n}$ which converges weakly in $L^{2}$ to 0 and let $u_{\Omega}$ be the Dirichlet eigenfunction of $\Omega$ (or equivalently, the Schrödinger eigenfunction related to $V=0$ ). Using $u_{\Omega}$ as a test function for the Rayleigh quotient in the definition of $\lambda_{1}\left(V_{n}\right)$ for each $V_{n}$, we get that

$$
\inf \left\{\lambda_{1}(V):\|V\|_{L^{p}(\Omega)}=1\right\}=\lambda(\Omega)
$$

But this implies that a minimizer $V$ should be null almost everywhere, which is impossible since such a function cannot have $L^{p}-$ norm equal to 1 . Nevertheless, we will see, in Section 1.4, a class of natural constraints whose spectral and energy minimization problems are solvable.

## Organization and main results of the chapter

The chapter is organized as follows. In Section 1.2 we introduce the notion of capacity, $\Gamma$-convergence, and $\gamma$-convergence. These tools, together with some related results, will be then exploited principally in the following Section 1.3.

In Section 1.3 we address the problem (1.1) in bounded domains where the constraint $\mathcal{V}$ is a weakly compact subset of $L^{p}\left(\mathbb{R}^{d}\right)$. As we have pointed out, many explicit minimization problems with such constraint are trivial. This is not the case if we consider maximization problems, see Remark 1.11. More precisely, in this section we prove two general results: in Theorem 1.10 we prove existence of solutions under some natural assumptions on the functional $F$ and on the constraint $\mathcal{V}$. These assumptions are satisfied by a large class of integral functionals and of spectral functionals. Then we shall see a version of Theorem 1.10 (Theorem 1.13) valid in the class of capacitary measures. These theorems may be seen as a generalization of Theorem 1.1. The main tools we exploit to prove these results are contained in the theory of the $\gamma$-convergence discussed in Section 1.2.

In Section 1.4 we consider problem (1.1), this time seen as an actual minimization problem, that is with $F$ non-negative, and, exploiting again some techniques borrowed by the $\gamma$-convergence theory, we prove an existence theorem (Theorem 1.16) for a wide class of non-convex constraints. Such theorem apply, for instance, to the case where the cost functional $F$ is a spectral functional or the Energy functional, see Subsection 1.4.1.

In Section 1.5, we consider the problem of maximizing $\mathcal{E}(V)$. Namely we consider, for $p \geq 1$, the following problem

$$
\max \left\{\mathcal{E}_{f}(V): \int_{\Omega} V^{p} d x \leq 1, \quad V \geq 0\right\}
$$

Following some ideas developed in [75, Theorem 8.2.3], where the same problem related to $\lambda_{1}(V)$ is considered, we prove existence (and uniqueness, if $p>1$ ) of optimal potentials (see Proposition 1.21 and Proposition 1.24). The class of constraints considered ( $p \geq 1$ ) is sharp, in the sense that for $p<1$ the problem has no solution, see Remark 1.26.

In Section 1.6 we consider the case $\Omega=\mathbb{R}^{d}$ and we prove existence of optimal potentials for the first eigenvalue $\lambda_{1}(V)$ and for the Energy function $\mathcal{E}_{f}(V)$ of the Schrödinger operator. For these two cases we are able to study qualitatively the shape of the optimal potentials, proving that they are radial functions and, in the minimization case, have compact support (see Proposition 1.31) while in the maximization case, they
have support $\mathbb{R}^{d}$. Eventually, thanks to the previous results concerning the supports of $\lambda_{1}(V)$, we are able to characterize, at least qualitatively, the shape of the minimizers of the second eigenvalue $\lambda_{2}(V)$ (Proposition 1.34).

### 1.2 Capacitary measures and $\gamma$-convergence

For a subset $E \subset \mathbb{R}^{d}$ its capacity is defined by
$\operatorname{cap}(E)=\inf \left\{\int_{\mathbb{R}^{d}}|\nabla u|^{2} d x+\int_{\mathbb{R}^{d}} u^{2} d x: u \in H^{1}\left(\mathbb{R}^{d}\right), u \geq 1\right.$ in a neighborhood of $\left.E\right\}$.
If a property $P(x)$ holds for all $x \in \Omega$, except for the elements of a set $E \subset \Omega$ of capacity zero, we say that $P(x)$ holds quasi-everywhere (shortly q.e.) in $\Omega$, whereas the expression almost everywhere (shortly a.e.) refers, as usual, to the Lebesgue measure, which we often denote by $|\cdot|$.

A subset $A$ of $\mathbb{R}^{d}$ is said to be quasi-open if for every $\varepsilon>0$ there exists an open subset $A_{\varepsilon}$ of $\mathbb{R}^{d}$, with $A \subset A_{\varepsilon}$, such that $\operatorname{cap}\left(A_{\varepsilon} \backslash A\right)<\varepsilon$. Similarly, a function $u: \mathbb{R}^{d} \rightarrow \mathbb{R}$ is said to be quasi-continuous (respectively quasi-lower semicontinuous) if there exists a decreasing sequence of open sets $\left(A_{n}\right)_{n}$ such that $\operatorname{cap}\left(A_{n}\right) \rightarrow 0$ and the restriction $u_{n}$ of $u$ to the set $A_{n}^{c}$ is continuous (respectively lower semicontinuous). It is well known (see for instance [57]) that every function $u \in H^{1}\left(\mathbb{R}^{d}\right)$ has a quasicontinuous representative $\widetilde{u}$, which is uniquely defined up to a set of capacity zero, and given by

$$
\widetilde{u}(x)=\lim _{\varepsilon \rightarrow 0} \frac{1}{\left|B_{\varepsilon}(x)\right|} \int_{B_{\varepsilon}(x)} u(y) d y,
$$

where $B_{\varepsilon}(x)$ denotes the ball of radius $\varepsilon$ centred at $x$. We identify the (a.e.) equivalence class $u \in H^{1}\left(\mathbb{R}^{d}\right)$ with the (q.e.) equivalence class of quasi-continuous representatives $\widetilde{u}$.

We denote by $\mathcal{M}^{+}\left(\mathbb{R}^{d}\right)$ the set of positive Borel measures on $\mathbb{R}^{d}$ (not necessarily finite or Radon) and by $\mathcal{M}_{\text {cap }}^{+}\left(\mathbb{R}^{d}\right) \subset \mathcal{M}^{+}\left(\mathbb{R}^{d}\right)$ the set of capacitary measures, i.e. the measures $\mu \in \mathcal{M}^{+}\left(\mathbb{R}^{d}\right)$ such that $\mu(E)=0$ for any set $E \subset \mathbb{R}^{d}$ of capacity zero. We note that when $\mu$ is a capacitary measure, the integral $\int_{\mathbb{R}^{d}}|u|^{2} d \mu$ is well-defined for each $u \in H^{1}\left(\mathbb{R}^{d}\right)$, i.e. if $\widetilde{u}_{1}$ and $\widetilde{u}_{2}$ are two quasi-continuous representatives of $u$, then $\int_{\mathbb{R}^{d}}\left|\widetilde{u}_{1}\right|^{2} d \mu=\int_{\mathbb{R}^{d}}\left|\widetilde{u}_{2}\right|^{2} d \mu$.

For a subset $\Omega \subset \mathbb{R}^{d}$, we define the Sobolev space $H_{0}^{1}(\Omega)$ as

$$
H_{0}^{1}(\Omega)=\left\{u \in H^{1}\left(\mathbb{R}^{d}\right): u=0 \text { q.e. on } \Omega^{c}\right\} .
$$

Alternatively, by using the capacitary measure $I_{\Omega}$ defined as

$$
I_{\Omega}(E)=\left\{\begin{array}{ll}
0 & \text { if } \operatorname{cap}(E \backslash \Omega)=0  \tag{1.5}\\
+\infty & \text { if } \operatorname{cap}(E \backslash \Omega)>0
\end{array} \quad \text { for every Borel set } E \subset \mathbb{R}^{d},\right.
$$

the Sobolev space $H_{0}^{1}(\Omega)$ can be defined as

$$
H_{0}^{1}(\Omega)=\left\{u \in H^{1}\left(\mathbb{R}^{d}\right): \int_{\mathbb{R}^{d}}|u|^{2} d I_{\Omega}<+\infty\right\} .
$$

More generally, for any capacitary measure $\mu \in \mathcal{M}_{\text {cap }}^{+}\left(\mathbb{R}^{d}\right)$, we define the space

$$
H_{\mu}^{1}=\left\{u \in H^{1}\left(\mathbb{R}^{d}\right): \int_{\mathbb{R}^{d}}|u|^{2} d \mu<+\infty\right\}
$$

which is a Hilbert space when endowed with the norm $\|u\|_{1, \mu}$, where

$$
\|u\|_{1, \mu}^{2}=\int_{\mathbb{R}^{d}}|\nabla u|^{2} d x+\int_{\mathbb{R}^{d}} u^{2} d x+\int_{\mathbb{R}^{d}} u^{2} d \mu
$$

If $u \notin H_{\mu}^{1}$, then we set $\|u\|_{1, \mu}=+\infty$.
For $\Omega \subset \mathbb{R}^{d}$, we define $\mathcal{M}_{\text {cap }}^{+}(\Omega)$ as the space of capacitary measures $\mu \in \mathcal{M}_{\text {cap }}^{+}\left(\mathbb{R}^{d}\right)$ such that $\mu(E)=+\infty$ for any set $E \subset \mathbb{R}^{d}$ such that $\operatorname{cap}(E \backslash \Omega)>0$. For $\mu \in \mathcal{M}_{\text {cap }}^{+}\left(\mathbb{R}^{d}\right)$, we denote with $H_{\mu}^{1}(\Omega)$ the space $H_{\mu \vee I_{\Omega}}^{1}=H_{\mu}^{1} \cap H_{0}^{1}(\Omega)$, where $a \vee b=\max \{a, b\}$.
Definition 1.3. Given a metric space $(X, d)$ and sequence of functionals $J_{n}: X \rightarrow$ $\mathbb{R} \cup\{+\infty\}$, we say that $J_{n} \Gamma$-converges, in the topology provided by $d$, to the functional $J: X \rightarrow \mathbb{R} \cup\{+\infty\}$, if the following two conditions are satisfied:

- for every sequence $x_{n}$ converging in to $x \in X$, we have

$$
J(x) \leq \underline{\lim }_{n \rightarrow \infty} J_{n}\left(x_{n}\right) ;
$$

- for every $x \in X$, there exists a sequence $x_{n}$ converging to $x$, such that

$$
J(x)=\lim _{n \rightarrow \infty} J_{n}\left(x_{n}\right) .
$$

For all details and properties of $\Gamma$-convergence we refer to [44]; here we simply recall that, whenever $J_{n} \Gamma$-converges to $J$,

$$
\min _{x \in X} J(x) \leq \underline{\lim }_{n \rightarrow \infty} \min _{x \in X} J_{n}(x) .
$$

Definition 1.4. We say that the sequence of capacitary measures $\mu_{n} \in \mathcal{M}_{\text {cap }}^{+}(\Omega)$, $\gamma$-converges to the capacitary measure $\mu \in \mathcal{M}_{\text {cap }}^{+}(\Omega)$ if the sequence of functionals $\|\cdot\|_{1, \mu_{n}} \Gamma$-converges to the functional $\|\cdot\|_{1, \mu}$ in $L^{2}(\Omega)$, i.e. if the following two conditions are satisfied:

- for every sequence $u_{n} \rightarrow u$ in $L^{2}(\Omega)$ we have

$$
\int_{\mathbb{R}^{d}}|\nabla u|^{2} d x+\int_{\mathbb{R}^{d}} u^{2} d \mu \leq \underline{\lim }_{n \rightarrow \infty}\left\{\int_{\mathbb{R}^{d}}\left|\nabla u_{n}\right|^{2} d x+\int_{\mathbb{R}^{d}} u_{n}^{2} d \mu_{n}\right\} ;
$$

- for every $u \in L^{2}(\Omega)$, there exists $u_{n} \rightarrow u$ in $L^{2}(\Omega)$ such that

$$
\int_{\mathbb{R}^{d}}|\nabla u|^{2} d x+\int_{\mathbb{R}^{d}} u^{2} d \mu=\lim _{n \rightarrow \infty}\left\{\int_{\mathbb{R}^{d}}\left|\nabla u_{n}\right|^{2} d x+\int_{\mathbb{R}^{d}} u_{n}^{2} d \mu_{n}\right\} .
$$

If $\mu \in \mathcal{M}_{\text {cap }}^{+}(\Omega)$ and $f \in L^{2}(\Omega)$ we define the functional $J_{\mu}(f, \cdot): L^{2}(\Omega) \rightarrow \mathbb{R} \cup\{+\infty\}$ by

$$
\begin{equation*}
J_{\mu}(f, u)=\frac{1}{2} \int_{\Omega}|\nabla u|^{2} d x+\frac{1}{2} \int_{\Omega} u^{2} d \mu-\int_{\Omega} f u d x \tag{1.6}
\end{equation*}
$$

If $\Omega \subset \mathbb{R}^{d}$ is a bounded open set, $\mu \in \mathcal{M}_{\text {cap }}^{+}(\Omega)$ and $f \in L^{2}(\Omega)$, then the functional $J_{\mu}(f, \cdot)$ has a unique minimizer $u \in H_{\mu}^{1}$ that verifies the PDE formally written as

$$
\begin{equation*}
-\Delta u+\mu u=f, \quad u \in H_{\mu}^{1}(\Omega) \tag{1.7}
\end{equation*}
$$

and whose precise meaning is given in the weak form

$$
\left\{\begin{array}{l}
\int_{\Omega}\langle\nabla u, \nabla \varphi\rangle d x+\int_{\Omega} u \varphi d \mu=\int_{\Omega} f \varphi d x, \quad \forall \varphi \in H_{\mu}^{1}(\Omega), \\
u \in H_{\mu}^{1}(\Omega) .
\end{array}\right.
$$

The resolvent operator of $-\Delta+\mu$, that is the map $\mathcal{R}_{\mu}$ that associates to every $f \in L^{2}(\Omega)$ the solution $u \in H_{\mu}^{1}(\Omega) \subset L^{2}(\Omega)$, is a compact linear operator in $L^{2}(\Omega)$ and so, it has a discrete spectrum

$$
0<\cdots \leq \Lambda_{k} \leq \cdots \leq \Lambda_{2} \leq \Lambda_{1} .
$$

Their inverses $1 / \Lambda_{k}$ are denoted by $\lambda_{k}(\mu)$ and are the eigenvalues of the operator $-\Delta+\mu$.

In the case $f=1$ the solution will be denoted by $w_{\mu}$ and when $\mu=I_{\Omega}$ we will use the notation $w_{\Omega}$ instead of $w_{I_{\Omega}}$. We also recall (see [25]) that if $\Omega$ is bounded, then the strong $L^{2}$-convergence of the minimizers $w_{\mu_{n}}$ to $w_{\mu}$ is equivalent to the $\gamma$-convergence of Definition 1.4.

Remark 1.5. An important well known characterization of the $\gamma$-convergence is the following: a sequence $\mu_{n} \gamma$-converges to $\mu$, if and only if, the sequence of resolvent operators $\mathcal{R}_{\mu_{n}}$ associated to $-\Delta+\mu_{n}$, converges (in the strong convergence of linear operators on $L^{2}$ ) to the resolvent $\mathcal{R}_{\mu}$ of the operator $-\Delta+\mu$. A consequence of this fact is that the spectrum of the operator $-\Delta+\mu_{n}$ converges (pointwise) to the one of $-\Delta+\mu$.

Remark 1.6. The space $\mathcal{M}_{\text {cap }}^{+}(\Omega)$ endowed with the $\gamma$-convergence is metrizable. If $\Omega$ is bounded, one may take $d_{\gamma}(\mu, \nu)=\left\|w_{\mu}-w_{\nu}\right\|_{L^{2}}$. Moreover, in this case, in [45] it is proved that the space $\mathcal{M}_{\text {cap }}^{+}(\Omega)$ endowed with the metric $d_{\gamma}$ is compact.

### 1.3 Existence of optimal potentials in $L^{p}(\Omega)$

In this section we consider the optimization problem

$$
\begin{equation*}
\min \left\{F(V): V: \Omega \rightarrow[0,+\infty], \int_{\Omega} V^{p} d x \leq 1\right\} \tag{1.8}
\end{equation*}
$$

where $p>0$ and $F(V)$ is a cost functional depending on the solution of some partial differential equation on $\Omega$. Typically, $F(V)$ is the minimum of some functional $J_{V}: H_{0}^{1}(\Omega) \rightarrow \mathbb{R}$ depending on $V$. A natural assumption in this case is the lower semicontinuity of the functional $F$ with respect to the $\gamma$-convergence, that is

$$
\begin{equation*}
F(\mu) \leq \underline{\lim }_{n \rightarrow \infty} F\left(\mu_{n}\right), \quad \text { whenever } \mu_{n} \rightarrow_{\gamma} \mu \tag{1.9}
\end{equation*}
$$

Proposition 1.7. Let $\Omega \subset \mathbb{R}^{d}$ and let $V_{n} \in L^{1}(\Omega)$ be a sequence weakly converging in $L^{1}(\Omega)$ to a function $V$. Then the capacitary measures $V_{n} d x \gamma$-converge to $V d x$.

Proof. We have to prove that the solutions $u_{n}=R_{V_{n}}(1)$ of

$$
\left\{\begin{array}{l}
-\Delta u_{n}+V_{n}(x) u_{n}=1 \\
u \in H_{0}^{1}(\Omega)
\end{array}\right.
$$

weakly converge in $H_{0}^{1}(\Omega)$ to the solution $u=R_{V}(1)$ of

$$
\left\{\begin{array}{l}
-\Delta u+V(x) u=1 \\
u \in H_{0}^{1}(\Omega)
\end{array}\right.
$$

or equivalently that the functionals

$$
J_{n}(u)=\int_{\Omega}|\nabla u|^{2} d x+\int_{\Omega} V_{n}(x) u^{2} d x
$$

$\Gamma\left(L^{2}(\Omega)\right)$-converge to the functional

$$
J(u)=\int_{\Omega}|\nabla u|^{2} d x+\int_{\Omega} V(x) u^{2} d x
$$

The $\Gamma$-liminf inequality (Definition 1.3) is immediate since, if $u_{n} \rightarrow u$ in $L^{2}(\Omega)$, we have

$$
\int_{\Omega}|\nabla u|^{2} d x \leq \underline{\lim }_{n \rightarrow \infty} \int_{\Omega}\left|\nabla u_{n}\right|^{2} d x
$$

by the lower semicontinuity of the $H^{1}(\Omega)$ norm with respect to the $L^{2}(\Omega)$-convergence, and

$$
\int_{\Omega} V(x) u^{2} d x \leq \underline{\lim }_{n \rightarrow \infty} \int_{\Omega} V_{n}(x) u_{n}^{2} d x
$$

by the strong-weak lower semicontinuity theorem for integral functionals (see for instance [27]).

Let us now prove the $\Gamma$-limsup inequality (see Definition 1.3) which consists, given $u \in H_{0}^{1}(\Omega)$, in constructing a sequence $u_{n} \rightarrow u$ in $L^{2}(\Omega)$ such that

$$
\begin{equation*}
\varlimsup_{n \rightarrow \infty} \int_{\Omega}\left|\nabla u_{n}\right|^{2} d x+\int_{\Omega} V_{n}(x) u_{n}^{2} d x \leq \int_{\Omega}|\nabla u|^{2} d x+\int_{\Omega} V(x) u^{2} d x . \tag{1.10}
\end{equation*}
$$

For every $t>0$ let $u^{t}=(u \wedge t) \vee(-t)$; then, by the weak convergence of $V_{n}$, for $t$ fixed we have

$$
\lim _{n \rightarrow \infty} \int_{\Omega} V_{n}(x)\left|u^{t}\right|^{2} d x=\int_{\Omega} V(x)\left|u^{t}\right|^{2} d x,
$$

and

$$
\lim _{t \rightarrow+\infty} \int_{\Omega} V(x)\left|u^{t}\right|^{2} d x=\int_{\Omega} V(x)|u|^{2} d x .
$$

Then, by a diagonal argument, we can find a sequence $t_{n} \rightarrow+\infty$ such that

$$
\lim _{n \rightarrow \infty} \int_{\Omega} V_{n}(x)\left|u^{t_{n}}\right|^{2} d x=\int_{\Omega} V(x)|u|^{2} d x .
$$

Taking now $u_{n}=u^{t_{n}}$, and noticing that for every $t>0$

$$
\int_{\Omega}\left|\nabla u^{t}\right|^{2} d x \leq \int_{\Omega}|\nabla u|^{2} d x,
$$

we obtain (1.10) and so the proof is complete.
In the case of weak* convergence of measures the statement of Proposition 1.7 is no longer true, as the following proposition shows.

Proposition 1.8. Let $\Omega \subset \mathbb{R}^{d}(d \geq 2)$ be a bounded open set and let $V, W \in L_{+}^{1}(\Omega)$ be two functions such that $V \geq W$. Then, there is a sequence $V_{n} \in L_{+}^{1}(\Omega)$, uniformly bounded in $L^{1}(\Omega)$, such that the sequence of measures $V_{n}(x) d x$ converges weakly* to $V(x) d x$ and $\gamma-$ converges to $W(x) d x$.

Proof. Without loss of generality we can suppose $\int_{\Omega}(V-W) d x=1$. Let $\mu_{n}$ be a sequence of probability measures on $\Omega$ weakly* converging to $(V-W) d x$ and such that each $\mu_{n}$ is a finite sum of Dirac masses. For each $n \in \mathbb{N}$ consider a sequence of positive functions $V_{n, m} \in L^{1}(\Omega)$ such that $\int_{\Omega} V_{n, m} d x=1$ and $V_{n, m} d x$ converges weakly* to $\mu_{n}$ as $m \rightarrow \infty$. Moreover, we choose $V_{n, m}$ as a convex combination of functions of the form $\left|B_{1 / m}\right|^{-1} \mathbf{1}_{B_{1 / m}\left(x_{j}\right)}$. Here $B_{r}(x)$ is the ball of centre $x$ and radius $r$, while $\mathbf{1}_{A}$ indicates the characteristic function of the set $A$.

We now prove that for fixed $n \in \mathbb{N},\left(V_{n, m}+W\right) d x \gamma$-converges, as $m \rightarrow \infty$, to $W d x$ or, equivalently, that the sequence $w_{W+V_{n, m}}$ converges in $L^{2}$ to $w_{W}$, as $m \rightarrow \infty$. Indeed, by the weak maximum principle, we have

$$
w_{W+I_{\Omega_{m}, n}} \leq w_{W+V_{n, m}} \leq w_{W}
$$

where $\Omega_{m, n}=\Omega \backslash \cup_{j} B_{1 / m}\left(x_{j}\right)$ and $I_{\Omega_{m, n}}$ is as in (1.5).
Since a point has zero capacity in $\mathbb{R}^{d}(d \geq 2)$ there exists a sequence $\phi_{m} \rightarrow 0$ strongly in $H^{1}\left(\mathbb{R}^{d}\right)$ with $\phi_{m}=1$ on $B_{1 / m}(0)$ and $\phi_{m}=0$ outside $B_{1 / \sqrt{m}}(0)$. We have

$$
\begin{align*}
\int_{\Omega}\left|w_{W}-w_{W+I_{\Omega_{m, n}}}\right|^{2} d x \leq & 2\left\|w_{W}\right\|_{L^{\infty}} \int_{\Omega}\left(w_{W}-w_{W+I_{\Omega_{m, n}}}\right) d x \\
= & 4\left\|w_{W}\right\|_{L^{\infty}}\left(E\left(W+I_{\Omega_{m, n}}\right)-E(W)\right)  \tag{1.11}\\
\leq & 4\left\|w_{W}\right\|_{L^{\infty}}\left(\int_{\Omega} \frac{1}{2}\left|\nabla w_{m}\right|^{2}+\frac{1}{2} W w_{m}^{2}-w_{m} d x\right. \\
& \left.\quad-\int_{\Omega} \frac{1}{2}\left|\nabla w_{W}\right|^{2}+\frac{1}{2} W w_{W}^{2}-w_{W} d x\right)
\end{align*}
$$

where $w_{m}$ is any function in $\in H_{0}^{1}\left(\Omega_{m, n}\right)$. Taking

$$
w_{m}(x)=w_{W}(x) \prod_{j}\left(1-\phi_{m}\left(x-x_{j}\right)\right)
$$

since $\phi_{m} \rightarrow 0$ strongly in $H^{1}\left(\mathbb{R}^{d}\right)$, it is easy to see that $w_{m} \rightarrow w_{W}$ strongly in $H^{1}(\Omega)$
 probability measures and the $\gamma$-convergence are both induced by metrics, a diagonal sequence argument brings to the conclusion.

We note that the hypotheses $V \geq W$ in Proposition 1.8 is necessary. Indeed, we have the following proposition, whose proof is contained in [33, Theorem 3.1] and we report it here for the sake of completeness.

Proposition 1.9. Let $\mu_{n} \in \mathcal{M}_{\text {cap }}^{+}(\Omega)$ be a sequence of capacitary Radon measures weakly* converging to the measure $\nu$ and $\gamma$-converging to the capacitary measure $\mu \in$ $\mathcal{M}_{\text {cap }}^{+}(\Omega)$. Then $\mu \leq \nu$ in $\Omega$.

Proof. We note that it is enough to show that $\mu(K) \leq \nu(K)$ whenever $K \subset \subset \Omega$ is a compact set. Let $u$ be a non-negative smooth function with compact support in $\Omega$ such that $u \leq 1$ in $\Omega$ and $u=1$ on $K$; we have

$$
\mu(K) \leq \int_{\Omega} u^{2} d \mu \leq \underline{\lim }_{n \rightarrow \infty} \int_{\Omega} u^{2} d \mu_{n}=\int_{\Omega} u^{2} d \nu \leq \nu(\{u>0\})
$$

Since $u$ is arbitrary, we have the conclusion by the Borel regularity of $\nu$.
Theorem 1.10. Let $F: L_{+}^{1}(\Omega) \rightarrow \mathbb{R}$ be a functional, lower semicontinuous with respect to the $\gamma$-convergence, and let $\mathcal{V}$ be a weakly $L^{1}(\Omega)$ compact set. Then the problem

$$
\begin{equation*}
\min \{F(V): V \in \mathcal{V}\} \tag{1.12}
\end{equation*}
$$

admits a solution.
$\operatorname{Proof}$. Let $\left(V_{n}\right)$ be a minimizing sequence in $\mathcal{V}$. By the compactness assumption on $\mathcal{V}$, we may assume that $V_{n}$ converges weakly $L^{1}(\Omega)$ to some $V \in \mathcal{V}$. By Proposition 1.7, we have that $V_{n} \gamma$-converges to $V$ and so, by the semicontinuity of $F$,

$$
F(V) \leq \underline{\lim }_{n \rightarrow \infty} F\left(V_{n}\right)
$$

which gives the conclusion.
Remark 1.11. Theorem 1.10 applies for instance to the maximization of integral functionals and to the spectral functionals considered in the introduction, as for instance the Dirichlet eigenvalues and the Energy Function; it is not difficult to show that they are continuous with respect to the $\gamma$-convergence.
Remark 1.12. In some special cases the solution of (1.8) can be written explicitly in terms of the solution of some partial differential equation on $\Omega$. This is the case of the Dirichlet Energy, that we shall discuss in Section 1.5, and of the first eigenvalue of the Dirichlet Laplacian $\lambda_{1}$ (see [75, Chapter 8]).

The compactness assumption on the admissible class $\mathcal{V}$ for the weak $L^{1}(\Omega)$ convergence in Theorem 1.10 is for instance satisfied if $\Omega$ has finite measure and $\mathcal{V}$ is a convex closed and bounded subset of $L^{p}(\Omega)$, with $p \geq 1$. In the case of measures an analogous result holds.

Theorem 1.13. Let $\Omega \subset \mathbb{R}^{d}$ be a bounded open set and let $F: \mathcal{M}_{\text {cap }}^{+}(\Omega) \rightarrow \mathbb{R}$ be a functional lower semicontinuous with respect to the $\gamma-$ convergence. Then the problem

$$
\begin{equation*}
\min \left\{F(\mu): \mu \in \mathcal{M}_{\text {cap }}^{+}(\Omega), \mu(\Omega) \leq 1\right\} \tag{1.13}
\end{equation*}
$$

admits a solution.
Proof. Let ( $\mu_{n}$ ) be a minimizing sequence. Then, up to a subsequence $\mu_{n}$ converges weakly* to some measure $\nu$ and $\gamma$-converges to some measure $\mu \in \mathcal{M}_{\text {cap }}^{+}(\Omega)$. By Proposition 1.9, we have that $\mu(\Omega) \leq \nu(\Omega) \leq 1$ and so $\mu$ is a solution of (1.13).

The following example shows that the optimal solution of problem (1.13) is not, in general, a function $V(x)$, even when the optimization criterion is the energy $\mathcal{E}_{f}$ introduced in (1.3). On the other hand, an explicit form for the optimal potential $V(x)$ will be provided in Proposition 1.21 and Proposition 1.24 assuming that $f$ is in $L^{2}(\Omega)$.
Example 1.14. Let $\Omega=(-1,1)$ and consider the functional

$$
F(\mu)=-\min \left\{\frac{1}{2} \int_{\Omega}\left|u^{\prime}\right|^{2} d x+\frac{1}{2} \int_{\Omega} u^{2} d \mu-u(0): u \in H_{0}^{1}(\Omega)\right\} .
$$

Then, for any $\mu$ such that $\mu(\Omega) \leq 1$, we have

$$
\begin{equation*}
F(\mu) \geq-\min \left\{\frac{1}{2} \int_{\Omega}\left|u^{\prime}\right|^{2} d x+\frac{1}{2}\left(\sup _{\Omega} u\right)^{2}-u(0): u \in H_{0}^{1}(\Omega), u \geq 0\right\} \tag{1.14}
\end{equation*}
$$

By a symmetrization argument, the minimizer $u$ of the right-hand side of (1.14) is radially decreasing; moreover, $u$ is linear on the set $u<M$, where $M=\sup u$, and so it is of the form

$$
u(x)= \begin{cases}\frac{M}{1-\alpha} x+\frac{M}{1-\alpha}, & x \in[-1,-\alpha],  \tag{1.15}\\ M, & x \in[-\alpha, \alpha], \\ -\frac{M}{1-\alpha} x+\frac{M}{1-\alpha}, & x \in[\alpha, 1],\end{cases}
$$

for some $\alpha \in[0,1]$. A straightforward computation gives $\alpha=0$ and $M=1 / 3$. Thus, $u$ is also the minimizer of

$$
F\left(\delta_{0}\right)=-\min \left\{\frac{1}{2} \int_{\Omega}\left|u^{\prime}\right|^{2} d x+\frac{1}{2} u(0)^{2}-u(0): u \in H_{0}^{1}(\Omega)\right\}
$$

and so $\delta_{0}$ is the solution of

$$
\min \{F(\mu): \mu(\Omega) \leq 1\}
$$

### 1.4 Existence of optimal potentials for unbounded constraints

In this section we consider the optimization problem

$$
\begin{equation*}
\min \{F(V): V \in \mathcal{V}\}, \tag{1.16}
\end{equation*}
$$

where $\mathcal{V}$ is an admissible class of non-negative Borel functions on the bounded open set $\Omega \subset \mathbb{R}^{d}$ and $F$ is a cost functional on the family of capacitary measures $\mathcal{M}_{\text {cap }}^{+}(\Omega)$. The admissible classes we study depend on a function $\Psi:[0,+\infty] \rightarrow[0,+\infty]$

$$
\mathcal{V}=\left\{V: \Omega \rightarrow[0,+\infty]: V \text { Lebesgue measurable, } \int_{\Omega} \Psi(V) d x \leq 1\right\}
$$

Before we state the main existence result of this section, we need the following lemma which is a particular case of [27, Theorem 2.3.1].

Lemma 1.15. Let $1<p, q<\infty$ and let $u_{n} \in L^{p}(\Omega)$ and $v_{n} \in L^{q}(\Omega)$ be two sequences of positive functions on the open set $\Omega \subset \mathbb{R}^{d}$ such that $u_{n}$ converges strongly in $L^{p}$ to $u \in L^{p}(\Omega)$ and $v_{n}$ converges weakly in $L^{q}$ to $v \in L^{q}(\Omega)$. Suppose that $H:[0,+\infty] \rightarrow$ $[0,+\infty]$ is a convex function. Then we have

$$
\int_{\Omega} u H(v) d x \leq \varliminf_{n \rightarrow \infty} \int_{\Omega} u_{n} H\left(v_{n}\right) d x .
$$

Theorem 1.16. Let $\Omega \subset \mathbb{R}^{d}$ be a bounded open set and $\Psi:[0,+\infty] \rightarrow[0,+\infty]$ a strictly decreasing function such that there exists $\varepsilon>0$ for which the function $x \rightarrow \Psi^{-1}\left(x^{1+\varepsilon}\right)$, defined on $[0,+\infty]$, is convex. Then, for any increasing functional $F: \mathcal{M}_{\text {cap }}^{+}(\Omega) \rightarrow \mathbb{R}$ which is lower semicontinuous with respect to the $\gamma$-convergence, the problem (1.16) has a solution.

Proof. Let $V_{n} \in \mathcal{V}$ be a minimizing sequence for problem (1.16). Then, $v_{n}:=\Psi\left(V_{n}\right)^{\frac{1}{1+\varepsilon}}$ is a bounded sequence in $L^{1+\varepsilon}(\Omega)$ and so, up to a subsequence, we have that $v_{n}$ converges weakly in $L^{1+\varepsilon}$ to some $v \in L^{1+\varepsilon}(\Omega)$. We will prove that $V:=\Psi^{-1}\left(v^{1+\varepsilon}\right)$ is a solution of (1.16). Clearly $V \in \mathcal{V}$ and so it remains to prove that $F(V) \leq \underline{\lim } F\left(V_{n}\right)$. In view of Remark 1.6, we can suppose that, up to a subsequence, $V_{n} \gamma$-converges to a capacitary measure $\mu \in \mathcal{M}_{\text {cap }}^{+}(\Omega)$. We claim that the following inequalities hold true:

$$
\begin{equation*}
F(V) \leq F(\nu) \leq \underline{\lim }_{n \rightarrow \infty} F\left(V_{n}\right) . \tag{1.17}
\end{equation*}
$$

In fact, the second inequality in (1.17) is the lower semicontinuity of $F$ with respect to the $\gamma$-convergence, while the first needs a more careful examination. By the definition of $\gamma$-convergence, we have that for any $u \in H_{0}^{1}(\Omega)$, there is a sequence $u_{n} \in H_{0}^{1}(\Omega)$ which converges to $u$ strongly in $L^{2}(\Omega)$ and is such that

$$
\begin{align*}
\int_{\Omega}|\nabla u|^{2} d x+\int_{\Omega} u^{2} d \mu & =\lim _{n \rightarrow \infty} \int_{\Omega}\left|\nabla u_{n}\right|^{2} d x+\int_{\Omega} u_{n}^{2} V_{n} d x \\
& =\lim _{n \rightarrow \infty} \int_{\Omega}\left|\nabla u_{n}\right|^{2} d x+\int_{\Omega} u_{n}^{2} \Psi^{-1}\left(v_{n}^{1+\varepsilon}\right) d x  \tag{1.18}\\
& \geq \int_{\Omega}|\nabla u|^{2} d x+\int_{\Omega} u^{2} \Psi^{-1}\left(v^{1+\varepsilon}\right) d x \\
& =\int_{\Omega}|\nabla u|^{2} d x+\int_{\Omega} u^{2} V d x,
\end{align*}
$$

where the inequality in (1.18) is due to Lemma 1.15. Thus, for any $u \in H_{0}^{1}(\Omega)$, we have that

$$
\int_{\Omega} u^{2} d \mu \geq \int_{\Omega} u^{2} V d x
$$

and so, $V \leq \mu$. Since $F$ is increasing, we obtain the first inequality in (1.17) which concludes the proof.

Remark 1.17. The condition on the function $\Psi$ in Theorem 1.16 is satisfied for instance by the following functions:

1. $\Psi(x)=x^{-p}$, for any $p>0$;
2. $\Psi(x)=e^{-\alpha x}$, for any $\alpha>0$.

### 1.4.1 Optimal potentials for the Dirichlet Energy and the first eigenvalue of the Dirichlet Laplacian

In the particular cases $F=\lambda_{1}$ and when $F=\mathcal{E}_{f}$, with $f \in L^{2}(\Omega)$ the solution of problem (1.16) can be expressed as a double minimization of a functional, where the
first minimization is on $H_{0}^{1}(\Omega)$ and the second on $\mathcal{V}$. We note that, by the variational formulation

$$
\begin{equation*}
\lambda_{1}(V)=\min \left\{\int_{\Omega}|\nabla u|^{2} d x+\int_{\Omega} u^{2} V d x: u \in H_{0}^{1}(\Omega), \int_{\Omega} u^{2} d x=1\right\} \tag{1.19}
\end{equation*}
$$

we can rewrite problem (1.16) as

$$
\begin{align*}
& \min \left\{\min _{\|u\|_{2}=1}\left\{\int_{\Omega}|\nabla u|^{2} d x+\int_{\Omega} u^{2} V d x\right\}: V \geq 0, \int_{\Omega} \Psi(V) d x \leq 1\right\} \\
&=\min _{\|u\|_{2}=1}\left\{\min \left\{\int_{\Omega}|\nabla u|^{2} d x+\int_{\Omega} u^{2} V d x: V \geq 0, \int_{\Omega} \Psi(V) d x \leq 1\right\}\right\} \tag{1.20}
\end{align*}
$$

By computing the Euler-Lagrange equation for $V$ with the constraint $\mathcal{V}$ ), and if $\Psi$ is differentiable with $\Psi^{\prime}$ invertible, we get that the second minimum in (1.20), for a generic $u \in H_{0}^{1}(\Omega)$ is achieved for

$$
\begin{equation*}
V=\left(\Psi^{\prime}\right)^{-1}\left(\Lambda_{u} u^{2}\right) \tag{1.21}
\end{equation*}
$$

where $\Lambda_{u}$ is a constant such that $\int_{\Omega} \Psi\left(\left(\Psi^{\prime}\right)^{-1}\left(\Lambda_{u} u^{2}\right)\right) d x=1$. Thus, the solution of the problem on the right hand side of (1.20) is given through the solution of

$$
\begin{equation*}
\min \left\{\int_{\Omega}|\nabla u|^{2} d x+\int_{\Omega} u^{2}\left(\Psi^{\prime}\right)^{-1}\left(\Lambda_{u} u^{2}\right) d x: u \in H_{0}^{1}(\Omega), \int_{\Omega} u^{2} d x=1\right\} \tag{1.22}
\end{equation*}
$$

Analogously, we obtain that the optimal potential for the Dirichlet Energy $\mathcal{E}_{f}$ is given by (1.21), where this time $u$ is a solution of

$$
\begin{equation*}
\min \left\{\int_{\Omega} \frac{1}{2}|\nabla u|^{2} d x+\int_{\Omega} \frac{1}{2} u^{2}\left(\Psi^{\prime}\right)^{-1}\left(\Lambda_{u} u^{2}\right) d x-\int_{\Omega} f u d x: u \in H_{0}^{1}(\Omega)\right\} \tag{1.23}
\end{equation*}
$$

Thus we obtain the following result.
Corollary 1.18. Under the assumptions of Theorem 1.16, for the functionals $F=\lambda_{1}$ and $F=\mathcal{E}_{f}$ there exists a solution of (1.16) given by $V=\left(\Psi^{\prime}\right)^{-1}\left(\Lambda_{u} u^{2}\right)$, where $u \in$ $H_{0}^{1}(\Omega)$ is a minimizer of (1.22), in the case $F=\lambda_{1}$, and of (1.23), in the case $F=\mathcal{E}_{f}$.
Example 1.19. If $\Psi(x)=x^{-p}$ with $p>0$, the optimal potentials for $\lambda_{1}$ and $\mathcal{E}_{f}$ are given by

$$
\begin{equation*}
V=\left(\int_{\Omega}|u|^{2 p /(p+1)} d x\right)^{1 / p} u^{-2 /(p+1)} \tag{1.24}
\end{equation*}
$$

where $u$ is the minimizer of (1.22) and (1.23), respectively. We also note that, in this case

$$
\int_{\Omega} u^{2}\left(\Psi^{\prime}\right)^{-1}\left(\Lambda_{u} u^{2}\right) d x=\left(\int_{\Omega}|u|^{2 p /(p+1)} d x\right)^{(1+p) / p}
$$

Example 1.20. If $\Psi(x)=e^{-\alpha x}$ with $\alpha>0$, the optimal potentials for $\lambda_{1}$ and $\mathcal{E}_{f}$ are given by

$$
\begin{equation*}
V=\frac{1}{\alpha}\left(\log \left(\int_{\Omega} u^{2} d x\right)-\log \left(u^{2}\right)\right), \tag{1.25}
\end{equation*}
$$

where $u$ is the minimizer of (1.22) and (1.23), respectively. We also note that, in this case

$$
\int_{\Omega} u^{2}\left(\Psi^{\prime}\right)^{-1}\left(\Lambda_{u} u^{2}\right) d x=\frac{1}{\alpha}\left(\int_{\Omega} u^{2} d x \int_{\Omega} \log \left(u^{2}\right) d x-\int_{\Omega} u^{2} \log \left(u^{2}\right) d x\right) .
$$

### 1.5 Maximization problems in $L^{p}$ concerning the Dirichlet Energy functional

In this section we study the a particular case of Theorem 1.10 where the functional $F$ is the energy function (with a minus sign) $-\mathcal{E}_{f}$ and $\Omega \subset \mathbb{R}^{d}$ is an open bounded set and $f \in L^{2}(\Omega)$. The natural constraint used for similar problems is, in analogy with the problem related to the classical Dirichlet energy of a set, is volume constraint. More precisely we are going to study the class of problem, depending on the parameter $p$,

$$
\begin{equation*}
\max \left\{\mathcal{E}_{f}(V): V \geq 0, \int_{\Omega} V^{p} \leq 1\right\} . \tag{1.26}
\end{equation*}
$$

To link this problem to Theorem 1.10 we only have to observe that the $\mathcal{E}_{f}$ is continuous with respect to the $\gamma$-convergence, and that we can always formulate a maximization problem on $\mathbb{R}$ as a minimization problem, simply changing the sign of the functional we wish to optimize. We recall that replacing max by min, makes problem (1.26) trivial, with the only solution $V \equiv 0$. Analogous results for $F(V)=-\lambda_{1}(V)$ has been treated in a series of papers. We cite for instance [74] and [5] in addition to the survey [75, Chapter 8]. We start by considering the case $p>1$.

Proposition 1.21. Let $\Omega \subset \mathbb{R}^{d}$ be a bounded open set, $1<p<\infty$ and $f \in L^{2}(\Omega)$. Then the problem (1.26) has the unique solution

$$
V_{p}=\left(\int_{\Omega}\left|u_{p}\right|^{2 p /(p-1)} d x\right)^{-1 / p}\left|u_{p}\right|^{-1+(p+1) /(p-1)},
$$

where $u_{p} \in H_{0}^{1}(\Omega) \cap L^{2 p /(p-1)}(\Omega)$ is the minimizer of the functional

$$
\begin{equation*}
J_{p}(u):=\frac{1}{2} \int_{\Omega}|\nabla u|^{2} d x+\frac{1}{2}\left(\int_{\Omega}|u|^{2 p /(p-1)} d x\right)^{(p-1) / p}-\int_{\Omega} u f d x . \tag{1.27}
\end{equation*}
$$

Moreover, we have $\mathcal{E}_{f}\left(V_{p}\right)=J_{p}\left(u_{p}\right)$.

Proof. To fix the notation let us define

$$
\mathcal{V}_{p}=\left\{V \geq 0: \int_{\Omega} V^{p} \leq 1\right\}
$$

We start noticing the trivial inequality

$$
\begin{equation*}
\max _{V \in \mathcal{V}_{p}} \min _{u \in H_{0}^{1}(\Omega)} \int_{\Omega}\left(\frac{1}{2}|\nabla u|^{2}+u^{2} V-u f\right) d x \leq \min _{u \in H_{0}^{1}(\Omega)} \max _{V \in \mathcal{V}_{p}} \int_{\Omega}\left(\frac{1}{2}|\nabla u|^{2}+u^{2} V-u f\right) d x \tag{1.28}
\end{equation*}
$$

Notice that by the Hölder inequality of parameter $\left(p, p^{\prime}\right)$ applied on the integral of $u^{2} V$, the maximum on the right hand side of (1.28) is finite. Moreover it is achieved by a function $V$ such that $\Lambda p V^{p-1}=u^{2}$, where $\Lambda$ is a suitable Lagrange multiplier. By the condition $\int_{\Omega} V^{p} d x=1$ we precisely gets

$$
V=\left(\int_{\Omega}|u|^{\frac{2 p}{p-1}} d x\right)^{1 / p}|u|^{\frac{2}{p-1}} .
$$

Substituting in (1.28), we obtain

$$
\begin{equation*}
\max \left\{\mathcal{E}_{f}(V): V \in \mathcal{V}\right\} \leq \min \left\{J_{p}(u): u \in H_{0}^{1}(\Omega)\right\} \tag{1.29}
\end{equation*}
$$

Let $u_{n}$ be a minimizing sequence for $J_{p}$. Since $\inf J_{p} \leq 0$, we can assume $J_{p}\left(u_{n}\right) \leq 0$ for each $n \in \mathbb{N}$. Thus, we have

$$
\begin{equation*}
\frac{1}{2} \int_{\Omega}\left|\nabla u_{n}\right|^{2} d x+\frac{1}{2}\left(\int_{\Omega}\left|u_{n}\right|^{2 p /(p-1)} d x\right)^{(p-1) / p} \leq \int_{\Omega} u_{n} f d x \leq C\|f\|_{L^{2}(\Omega)}\left\|\nabla u_{n}\right\|_{L^{2}}, \tag{1.30}
\end{equation*}
$$

where $C$ is a constant depending on $\Omega$. Thus we obtain

$$
\begin{equation*}
\int_{\Omega}\left|\nabla u_{n}\right|^{2} d x+\left(\int_{\Omega}\left|u_{n}\right|^{2 p /(p-1)} d x\right)^{(p-1) / p} \leq 4 C^{2}\|f\|_{L^{2}(\Omega)}^{2} \tag{1.31}
\end{equation*}
$$

and so, up to subsequence $u_{n}$ converges weakly in $H_{0}^{1}(\Omega)$ and $L^{2 p /(p-1)}(\Omega)$ to some $u_{p} \in H_{0}^{1}(\Omega) \cap L^{2 p /(p-1)}(\Omega)$. By the semicontinuity of the $L^{2}-$ norm of the gradient and the $L^{\frac{2 p}{p-1}}-$ norm and the fact that $\int_{\Omega} f u_{n} d x \rightarrow \int_{\Omega} f u_{p} d x$, as $n \rightarrow \infty$, we have that $u_{p}$ is a minimizer of $J_{p}$. By the strict convexity of $J_{p}$, we have that $u_{p}$ is unique. Moreover, by (1.30) and (1.31), $J_{p}\left(u_{p}\right)>-\infty$. Writing down the Euler-Lagrange equation for $u_{p}$, we obtain

$$
-\Delta u_{p}+\left(\int_{\Omega}\left|u_{p}\right|^{2 p /(p-1)} d x\right)^{-1 / p}\left|u_{p}\right|^{2 /(p-1)} u_{p}=f .
$$

Setting

$$
V_{p}=\left(\int_{\Omega}\left|u_{p}\right|^{2 p /(p-1)} d x\right)^{-1 / p}\left|u_{p}\right|^{2 /(p-1)},
$$

we have that $\int_{\Omega} V_{p}^{p} d x=1$ and $u_{p}$ is the solution of

$$
\begin{equation*}
-\Delta u_{p}+V_{p} u_{p}=f \tag{1.32}
\end{equation*}
$$

In particular, we have $J_{p}\left(u_{p}\right)=\mathcal{E}_{p}\left(V_{p}\right)$ and so $V_{p}$ solves (1.26). The uniqueness of $V_{p}$ follows by the uniqueness of $u_{p}$ and the characterization of the equality cases of the Hölder inequality

$$
\int_{\Omega} u^{2} V d x \leq\left(\int_{\Omega} V^{p} d x\right)^{1 / p}\left(\int_{\Omega}|u|^{2 p /(p-1)} d x\right)^{(p-1) / p} \leq\left(\int_{\Omega}|u|^{2 p /(p-1)} d x\right)^{(p-1) / p}
$$

When the functional $F$ is the energy $\mathcal{E}_{f}$, the existence result holds also in the case $p=1$, but its proof is a bit more complicated and is given in Proposition 1.24. It worth noticing the peculiar shape of the minimizer in this case given by (1.40). Before proving Proposition 1.24, we need some preliminary results. We also recall that analogous results has been obtained in the case $F=\lambda_{1}$ (see [75, Theorem 8.2.4]) and in the case $F=\mathcal{E}_{f}$, with the further request that $f \geq 0$. See [33].

Remark 1.22. Let $u_{p}$ be the minimizer of $J_{p}$, defined in (1.27). By (1.31), we have the estimate

$$
\begin{equation*}
\left\|\nabla u_{p}\right\|_{L^{2}(\Omega)}+\left\|u_{p}\right\|_{L^{2 p /(p-1)}(\Omega)} \leq 4 C^{2}\|f\|_{L^{2}(\Omega)} \tag{1.33}
\end{equation*}
$$

where $C$ is the constant from (1.30). Moreover, we have $u_{p} \in H_{l o c}^{2}(\Omega)$ and for each open set $\Omega^{\prime} \subset \subset \Omega$, there is a constant $C$ not depending on $p$ such that

$$
\left\|u_{p}\right\|_{H^{2}\left(\Omega^{\prime}\right)} \leq C\left(f, \Omega^{\prime}\right)
$$

Indeed, $u_{p}$ satisfies the PDE

$$
\begin{equation*}
-\Delta u+c|u|^{\alpha} u=f \tag{1.34}
\end{equation*}
$$

with $c>0$ and $\alpha=2 /(p-1)$, and standard elliptic regularity arguments (see [56, Section 6.3]) give that $u \in H_{l o c}^{2}(\Omega)$. To show that $\left\|u_{p}\right\|_{H^{2}\left(\Omega^{\prime}\right)}$ is bounded independently of $p$ we apply the Nirenberg operator $\partial_{k}^{h} u=\frac{u\left(x+h e_{k}\right)-u(x)}{h}$ on both sides of (1.34), and multiplying by $\phi^{2} \partial_{k}^{h} u$, where $\phi$ is an appropriate cut-off function which equals 1 on $\Omega^{\prime}$, we have

$$
\begin{align*}
\int_{\Omega} \phi^{2}\left|\nabla \partial_{k}^{h} u\right|^{2} d x+\int_{\Omega}\left\langle\nabla\left(\partial_{k}^{h} u\right), \nabla\left(\phi^{2}\right) \partial_{k}^{h} u\right\rangle d x+c(\alpha+1) & \int_{\Omega} \phi^{2}|u|^{\alpha}\left|\partial_{k}^{h} u\right|^{2} d x  \tag{1.35}\\
& =-\int f \partial_{k}^{h}\left(\phi^{2} \partial_{k}^{h} u\right) d x
\end{align*}
$$

for all $k=1, \ldots, d$. Some straightforward manipulations now give

$$
\begin{equation*}
\left\|\nabla^{2} u\right\|_{L^{2}\left(\Omega^{\prime}\right)}^{2} \leq \sum_{k=1}^{d} \int_{\Omega} \phi^{2}\left|\nabla \partial_{k} u\right|^{2} d x \leq C\left(\Omega^{\prime}\right)\left(\|f\|_{L^{2}\left(\left\{\phi^{2}>0\right\}\right)}+\|\nabla u\|_{L^{2}(\Omega)}\right) . \tag{1.36}
\end{equation*}
$$

### 1.5 Maximization problems in $L^{p}$ concerning the Dirichlet Energy functional

Lemma 1.23. Let $\Omega \subset \mathbb{R}^{d}$ be an open set and $f \in L^{2}(\Omega)$. Consider the functional $J_{1}: L^{2}(\Omega) \rightarrow \mathbb{R}$ defined by

$$
\begin{equation*}
J_{1}(u):=\frac{1}{2} \int_{\Omega}|\nabla u|^{2} d x+\frac{1}{2}\|u\|_{\infty}^{2}-\int_{\Omega} u f d x, \tag{1.37}
\end{equation*}
$$

Then, $J_{p} \Gamma$-converges in $L^{2}(\Omega)$ to $J_{1}$, as $p \rightarrow 1$, where $J_{p}$ is defined in (1.27).
Proof. Let $v_{n} \in L^{2}(\Omega)$ be a sequence of positive functions converging in $L^{2}$ to $v \in L^{2}(\Omega)$ and let $\alpha_{n} \rightarrow+\infty$. Then, we have that

$$
\begin{equation*}
\|v\|_{L^{\infty}(\Omega)} \leq \underline{\lim }_{n \rightarrow \infty}\left\|v_{n}\right\|_{L^{\alpha_{n}}(\Omega)} \tag{1.38}
\end{equation*}
$$

In fact, suppose first that $\|v\|_{L^{\infty}}=M<+\infty$ and let $\omega_{\varepsilon}=\{v>M-\varepsilon\}$, for any $\varepsilon>0$. Then, we have

$$
\underline{\lim }_{n \rightarrow \infty}\left\|v_{n}\right\|_{L^{\alpha_{n}}(\Omega)} \geq \lim _{n \rightarrow \infty}\left|\omega_{\varepsilon}\right|^{\left(1-\alpha_{n}\right) / \alpha_{n}} \int_{\omega_{\varepsilon}} v_{n} d x=\left|\omega_{\varepsilon}\right|^{-1} \int_{\omega_{\varepsilon}} v d x \geq M-\varepsilon
$$

and so, letting $\varepsilon \rightarrow 0$, we have $\underline{l i m}_{n \rightarrow \infty}\left\|v_{n}\right\|_{L^{\alpha_{n}}(\Omega)} \leq M$. If $\|v\|_{L^{\infty}(\Omega)}=+\infty$, then setting $\omega_{k}=\{v>k\}$, for any $k \geq 1$, and arguing as above, we obtain (1.38).
Let $u_{n} \rightarrow u$ in $L^{2}(\Omega)$. Then, by the semicontinuity of the $L^{2}$ norm of the gradient and (1.38) and the continuity of the term $\int_{\Omega} u f d x$, we have

$$
\begin{equation*}
J_{1}(u) \leq \varliminf_{n \rightarrow \infty} J_{p_{n}}\left(u_{n}\right), \tag{1.39}
\end{equation*}
$$

for any decreasing sequence $p_{n} \rightarrow 1$. On the other hand, for any $u \in L^{2}$, we have $J_{p_{n}}(u) \rightarrow J_{1}(u)$ as $n \rightarrow \infty$ and so we have the conclusion.

Proposition 1.24. Let $\Omega \subset \mathbb{R}^{d}$ be a bounded open set and $f \in L^{2}(\Omega)$. Then there is a unique solution of problem (1.26) with $p=1$, given by

$$
\begin{equation*}
V_{1}=\frac{1}{M}\left(\mathbf{1}_{\omega_{+}} f-\mathbf{1}_{\omega_{-}} f\right), \tag{1.40}
\end{equation*}
$$

where $M=\left\|u_{1}\right\|_{L^{\infty}(\Omega)}, \omega_{+}=\left\{u_{1}=M\right\}, \omega_{-}=\left\{u_{1}=-M\right\}$, being $u_{1} \in H_{0}^{1}(\Omega) \cap L^{\infty}(\Omega)$ the unique minimizer of the functional $J_{1}$, defined in (1.37). In particular, $\int_{\omega_{+}} f d x-$ $\int_{\omega_{-}} f d x=M, f \geq 0$ on $\omega_{+}$and $f \leq 0$ on $\omega_{-}$.
Proof. For any $u \in H_{0}^{1}(\Omega)$ and any $V \geq 0$ with $\int_{\Omega} V d x \leq 1$ we have

$$
\int_{\Omega} u^{2} V d x \leq\|u\|_{\infty}^{2} \int_{\Omega} V d x \leq\|u\|_{\infty}^{2},
$$

where for the sake of simplicity, we write $\|\cdot\|_{\infty}$ instead of $\|\cdot\|_{L^{\infty}(\Omega)}$. Arguing as in the proof of Proposition 1.21, we obtain the inequalities

$$
\begin{aligned}
& \frac{1}{2} \int_{\Omega}|\nabla u|^{2} d x+\frac{1}{2} \int_{\Omega} u^{2} V d x-\int_{\Omega} u f d x \leq J_{1}(u) \\
& \max \left\{\mathcal{E}_{f}(V): \int_{\Omega} V \leq 1\right\} \leq \min \left\{J_{1}(u): u \in H_{0}^{1}(\Omega)\right\}
\end{aligned}
$$

As in (1.30), we have that a minimizing sequence of $J_{1}$ is bounded in $H_{0}^{1}(\Omega) \cap L^{\infty}(\Omega)$ and thus by semicontinuity there is a minimizer $u_{1} \in H_{0}^{1}(\Omega) \cap L^{\infty}(\Omega)$ of $J_{1}$, which is also unique, by the strict convexity of $J_{1}$. Let $u_{p}$ denotes the minimizer of $J_{p}$ as in Proposition 1.21. Then, by Remark 1.22 , we have that the family $u_{p}$ is bounded in $H_{0}^{1}(\Omega)$ and in $H^{2}\left(\Omega^{\prime}\right)$ for each $\Omega^{\prime} \subset \subset \Omega$. Then, we have that each sequence $u_{p_{n}}$ has a subsequence converging weakly in $L^{2}(\Omega)$ to some $u \in H_{l o c}^{2}(\Omega) \cap H_{0}^{1}(\Omega)$. By Lemma 1.23, we have $u=u_{1}$ and so, $u_{1} \in H_{l o c}^{2}(\Omega) \cap H_{0}^{1}(\Omega)$. Thus $u_{p_{n}} \rightarrow u_{1}$ in $L^{2}(\Omega)$.

Let us define $M=\left\|u_{1}\right\|_{\infty}$ and $\omega=\omega_{+} \cup \omega_{-}$. We claim that $u_{1}$ satisfies, on $\Omega$ the PDE

$$
\begin{equation*}
-\Delta u+\mathbf{1}_{\omega} f=f \tag{1.41}
\end{equation*}
$$

Indeed, setting $\Omega_{t}=\Omega \cap\{|u|<t\}$ for $t>0$, we compute the variation of $J_{1}$ with respect to any function $\varphi \in H_{0}^{1}\left(\Omega_{M-\varepsilon}\right)$. Namely we consider functions of the form $\varphi=\psi w_{\varepsilon}$ where $w_{\varepsilon}$ is the solution of $-\Delta w_{\varepsilon}=1$ on $\Omega_{M-\varepsilon}$, and $w_{\varepsilon}=0$ on $\partial \Omega_{M-\varepsilon}$. Thus we obtain that $-\Delta u_{1}=f$ on $\Omega_{M-\varepsilon}$ and letting $\varepsilon \rightarrow 0$ we conclude, thanks to the Monotone Convergence Theorem, that

$$
-\Delta u_{1}=f \quad \text { on } \Omega_{M}=\Omega \backslash \omega
$$

Moreover, since $u_{1} \in H_{l o c}^{2}(\Omega)$, we have that $\Delta u_{1}=0$ on $\omega$ and so, we obtain (1.41). Since $u_{1}$ is the minimizer of $J_{1}$, we have that for each $\varepsilon \in \mathbb{R}, J_{1}\left((1+\varepsilon) u_{1}\right)-J_{1}\left(u_{1}\right) \geq 0$. Taking the derivative of this difference at $\varepsilon=0$, we obtain

$$
\begin{equation*}
\int_{\Omega}\left|\nabla u_{1}\right|^{2} d x+M^{2}=\int_{\Omega} f u_{1} d x \tag{1.42}
\end{equation*}
$$

By (1.41), we have $\int_{\Omega}\left|\nabla u_{1}\right|^{2} d x=\int_{\Omega \backslash \omega} f u_{1} d x$ and so

$$
\begin{equation*}
M=\int_{\omega_{+}} f d x-\int_{\omega_{-}} f d x \tag{1.43}
\end{equation*}
$$

Setting $V_{1}:=\frac{1}{M}\left(\mathbf{1}_{\omega_{+}} f-\mathbf{1}_{\omega_{-}} f\right)$, we have that $\int_{\Omega} V_{1} d x=1,-\Delta u_{1}+V_{1} u_{1}=f$ in $H^{-1}(\Omega)$ and

$$
J_{1}\left(u_{1}\right)=\frac{1}{2} \int_{\Omega}\left|\nabla u_{1}\right|^{2} d x+\frac{1}{2} \int_{\Omega} u_{1}^{2} V_{1} d x-\int_{\Omega} u_{1} f d x
$$

We are left to prove that $V_{1}$ is admissible, i.e. $V_{1} \geq 0$. To do this, consider $w_{\varepsilon}$ the Energy function of the quasi-open set $\{u<M-\varepsilon\}$ and let $\varphi=w_{\varepsilon} \psi$ where $\psi \in C_{c}^{\infty}\left(\mathbb{R}^{d}\right)$, $\psi \geq 0$. Since $\varphi \geq 0$, we get that

$$
0 \leq \lim _{t \rightarrow 0^{+}} \frac{J_{1}\left(u_{1}+t \varphi\right)-J_{1}\left(u_{1}\right)}{t}=\int_{\Omega}\left\langle\nabla u_{1}, \nabla \varphi\right\rangle d x-\int_{\Omega} f \varphi d x
$$

This inequality holds for any $\psi$ so that, integrating by parts, we obtain

$$
-\Delta u_{1}-f \geq 0
$$

almost everywhere on $\left\{u_{1}<M-\varepsilon\right\}$. In particular, since $\Delta u_{1}=0$ almost everywhere on $\omega_{-}=\{u=-M\}$, we obtain that $f \leq 0$ on $\omega_{-}$. Arguing in the same way, and considering test functions supported on $\left\{u_{1} \geq-M+\varepsilon\right\}$, we can prove that $f \geq 0$ on $\omega_{+}$. This implies $V_{1} \geq 0$ as required.

Remark 1.25. Under some additional assumptions on $\Omega$ and $f$ one can obtain some more precise regularity results for $u_{1}$. In fact, in [55, Theorem A1] it was proved that if $\partial \Omega \in C^{2}$ and if $f \in L^{\infty}(\Omega)$ is positive, then $u_{1} \in C^{1,1}(\bar{\Omega})$.

Remark 1.26. In the case $p<1$ problem (1.26) does not admit, in general, a solution, even for regular $f$ and $\Omega$. We give a counterexample in dimension one, which can be easily adapted to higher dimensions.

Let $\Omega=(0,1), f=1$, and let $x_{n, k}=k / n$ for any $n \in \mathbb{N}$ and $k=1, \ldots, n-1$. We define the (capacitary) measures

$$
\mu_{n}=\sum_{k=1}^{n-1}+\infty \delta_{x_{n, k}},
$$

where $\delta_{x}$ is the Dirac measure at the point $x$. Let $w_{n}$ be the minimizer of the functional $J_{\mu_{n}}(1, \cdot)$, defined in (1.6). Then $w_{n}$ vanishes at $x_{n, k}$, for $k=1, \ldots, n-1$, and so we have

$$
\mathcal{E}\left(\mu_{n}\right)=n \min \left\{\frac{1}{2} \int_{0}^{1 / n}\left|u^{\prime}\right|^{2} d x-\int_{0}^{1 / n} u d x: u \in H_{0}^{1}(0,1 / n)\right\}=-\frac{C}{n^{2}},
$$

where $C>0$ is a constant.
For any fixed $n$ and $j$, let $V_{j}^{n}$ be the sequence of positive functions such that $\int_{0}^{1}\left|V_{j}^{n}\right|^{p} d x=1$, defined by

$$
\begin{equation*}
V_{j}^{n}=C_{n} \sum_{k=1}^{n-1} j^{1 / p} \mathbf{1}_{\left[\frac{k}{n}-\frac{1}{j}, \frac{k}{n}+\frac{1}{j}\right]}<\sum_{k=1}^{n-1} I_{\left[\frac{k}{n}-\frac{1}{j}, \frac{k}{n}+\frac{1}{j}\right]}, \tag{1.44}
\end{equation*}
$$

where $C_{n}$ is a constant depending on $n$ and $I$ is as in (1.5). By the compactness of the $\gamma$-convergence, we have that, up to a subsequence, $V_{j}^{n} d x \gamma$-converges to some capacitary measure $\mu$ as $j \rightarrow \infty$. On the other hand it is easy to check that $\sum_{k=1}^{n-1} I_{\left[\frac{k}{n}-\frac{1}{j}, \frac{k}{n}+\frac{1}{j}\right]}(x) \gamma$-converges to $\mu_{n}$ as $j \rightarrow \infty$. By (1.44), we have that $\mu \leq \mu_{n}$. In order to show that $\mu=\mu_{n}$ it is enough to check that each non-negative function $u \in H_{0}^{1}((0,1))$, for which $\int u^{2} d \mu<+\infty$, vanishes at $x_{n, k}$ for $k=1, \ldots, n-1$. Suppose that $u(k / n)>0$. By the definition of the $\gamma$-convergence, there is a sequence $u_{j} \in H_{0}^{1}(\Omega)=H_{V_{j}^{n}}^{1}(\Omega)$ such that $u_{j} \rightarrow u$ weakly in $H_{0}^{1}(\Omega)$ and $\int u_{j}^{2} V_{j}^{n} d x \leq C$, for some constant $C$ not depending on $j \in \mathbb{N}$. Since $u_{j}$ are uniformly $1 / 2-$ Hölder continuous, we can suppose that $u_{j} \geq \varepsilon>0$ on some interval $I$ containing $k / n$. But then for
$j$ large enough $I$ contains $[k / n-1 / j, k / n+1 / j]$ so that

$$
C \geq \int_{0}^{1} u_{j}^{2} V_{j}^{n} d x \geq \int_{k / n-1 / j}^{k / n+1 / j} u_{j}^{2} V_{j}^{n} d x \geq 2 C_{n} \varepsilon^{2} j^{1 / p-1}
$$

which is a contradiction for $p<1$. Thus, we have that $\mu=\mu_{n}$ and so $V_{j}^{n} \gamma$-converges to $\mu_{n}$ as $j \rightarrow \infty$. In particular, $\mathcal{E}\left(\mu_{n}\right)=\lim _{j \rightarrow \infty} \mathcal{E}_{1}\left(V_{j}^{n}\right)$ and since the left-hand side converges to zero as $n \rightarrow \infty$, we can choose a diagonal sequence $V_{j_{n}}^{n}$ such that $\mathcal{E}\left(V_{j_{n}}^{n}\right) \rightarrow 0$ as $n \rightarrow \infty$. Since there is no admissible functional $V$ such that $\mathcal{E}_{1}(V)=0$, we have the conclusion.

### 1.6 Optimization problems in unbounded domains

In this section we consider optimization problems for which the domain region is the entire Euclidean space $\mathbb{R}^{d}$. General existence results, in the case when the design region $\Omega$ is unbounded, are hard to achieve since most of the cost functionals are not semicontinuous with respect to the $\gamma$-convergence in these domains. For example, it is not hard to check that if $\mu$ is a capacitary measure, infinite outside the unit ball $B_{1}$, then, for every sequence $\left(x_{n}\right)_{n}$ such that $\left|x_{n}\right| \rightarrow \infty$, the sequence of translated measures $\mu_{n}=\mu\left(\cdot+x_{n}\right) \gamma$-converges to the capacitary measure

$$
I_{\emptyset}(E)= \begin{cases}0, & \text { if } \operatorname{cap}(E)=0 \\ +\infty, & \text { if } \operatorname{cap}(E)>0\end{cases}
$$

Thus increasing and translation invariant functionals are never lower semicontinuous with respect to the $\gamma$-convergence. In some special cases, as the Dirichlet Energy or the first eigenvalue of the Dirichlet Laplacian, one can obtain existence results by more direct methods, as those in Proposition 1.21.

For a potential $V \geq 0$ and a function $f \in L^{q}\left(\mathbb{R}^{d}\right)$, we define the Dirichlet Energy as

$$
\begin{equation*}
\mathcal{E}_{f}(V)=\inf \left\{\int_{\mathbb{R}^{d}}\left(\frac{1}{2}|\nabla u|^{2}+\frac{1}{2} V(x) u^{2}-f(x) u\right) d x: u \in C_{c}^{\infty}\left(\mathbb{R}^{d}\right)\right\} . \tag{1.45}
\end{equation*}
$$

In some cases it is convenient to work with the space $\dot{H}^{1}\left(\mathbb{R}^{d}\right)$, obtained as the closure of $C_{c}^{\infty}\left(\mathbb{R}^{d}\right)$ with respect to the $L^{2}$ norm of the gradient, instead of the classical Sobolev space $H^{1}\left(\mathbb{R}^{d}\right)$. From now on, with the scope of lightening the notation, and since there is no risk of confusion, we will write $\|\cdot\|_{p}$ in place of $\|\cdot\|_{L^{p}\left(\mathbb{R}^{d}\right)}$ to indicate the $L^{p}$-norm of a function on $\mathbb{R}^{d}$. We recall that if $d \geq 3$, the Sobolev inequality

$$
\begin{equation*}
\|u\|_{2 d /(d-2)} \leq C_{d}\|\nabla u\|_{2}, \quad \forall u \in \dot{H}^{1}\left(\mathbb{R}^{d}\right), \tag{1.46}
\end{equation*}
$$

holds, while in the cases $d \leq 2$, we have respectively (see for instance [86])

$$
\begin{align*}
& \|u\|_{\infty} \leq\left(\frac{r+2}{2}\right)^{2 /(r+2)}\|u\|_{r}^{r /(r+2)}\left\|u^{\prime}\right\|_{2}^{2 /(r+2)}, \quad \forall r \geq 1, \forall u \in \dot{H}^{1}(\mathbb{R})  \tag{1.47}\\
& \|u\|_{r+2} \leq\left(\frac{r+2}{2}\right)^{2 /(r+2)}\|u\|_{r}^{r /(r+2)}\|\nabla u\|_{2}^{2 /(r+2)}, \quad \forall r \geq 1, \forall u \in \dot{H}^{1}\left(\mathbb{R}^{2}\right) \tag{1.48}
\end{align*}
$$

### 1.6.1 Optimal potentials in $L^{p}\left(\mathbb{R}^{d}\right)$

In this section we consider optimization problems for the Dirichlet Energy $\mathcal{E}_{f}$ among potentials $V \geq 0$ satisfying a constraint of the form $\|V\|_{p} \leq 1$. We note that the results contained in this section hold in a generic regular unbounded domain $\Omega$. Nevertheless, for the sake of simplicity, we restrict our attention to the case $\Omega=\mathbb{R}^{d}$.

Proposition 1.27. Let $p>1$ and let $q$ be in the interval with end-points $a=2 p /(p+1)$ and $b=\max \{1,2 d /(d+2)\}$ (with a included for every $d \geq 1$, and $b$ included for every $d \neq 2$ ). Then, for every $f \in L^{q}\left(\mathbb{R}^{d}\right)$, there is a unique solution of the problem

$$
\begin{equation*}
\max \left\{\mathcal{E}_{f}(V): V \geq 0, \int_{\mathbb{R}^{d}} V^{p} d x \leq 1\right\} \tag{1.49}
\end{equation*}
$$

Proof. Arguing as in Proposition 1.21, we have that for $p>1$ the optimal potential $V_{p}$ is given by

$$
\begin{equation*}
V_{p}=\left(\int_{\mathbb{R}^{d}}\left|u_{p}\right|^{2 p /(p-1)} d x\right)^{-1 / p}\left|u_{p}\right|^{2 /(p-1)} \tag{1.50}
\end{equation*}
$$

where $u_{p}$ is the solution of the problem

$$
\begin{array}{r}
\min \left\{\frac{1}{2} \int_{\mathbb{R}^{d}}|\nabla u|^{2} d x+\frac{1}{2}\left(\int_{\mathbb{R}^{d}}|u|^{2 p /(p-1)} d x\right)^{(p-1) / p}-\int_{\mathbb{R}^{d}} u f d x:\right.  \tag{1.51}\\
\left.u \in \dot{H}^{1}\left(\mathbb{R}^{d}\right) \cap L^{2 p /(p-1)}\left(\mathbb{R}^{d}\right)\right\}
\end{array}
$$

Thus, it is enough to prove that there exists a solution of (1.51). For a minimizing sequence $u_{n}$ we have

$$
\begin{equation*}
\frac{1}{2} \int_{\mathbb{R}^{d}}\left|\nabla u_{n}\right|^{2} d x+\frac{1}{2}\left(\int_{\mathbb{R}^{d}}\left|u_{n}\right|^{2 p /(p-1)} d x\right)^{(p-1) / p} \leq \int_{\mathbb{R}^{d}} u_{n} f d x \leq C\|f\|_{q}\left\|u_{n}\right\|_{q^{\prime}} \tag{1.52}
\end{equation*}
$$

Suppose that $d \geq 3$. Interpolating $q^{\prime}$ between $2 p /(p-1)$ and $2 d /(d-2)$ and using the Sobolev inequality (1.46), we obtain that there is a constant $C$, depending only on $p, d$ and $f$, such that

$$
\frac{1}{2} \int_{\mathbb{R}^{d}}\left|\nabla u_{n}\right|^{2} d x+\frac{1}{2}\left(\int_{\mathbb{R}^{d}}\left|u_{n}\right|^{2 p /(p-1)} d x\right)^{(p-1) / p} \leq C
$$

Thus we can suppose that $u_{n}$ converges weakly in $\dot{H}^{1}\left(\mathbb{R}^{d}\right)$ and in $L^{2 p /(p-1)}\left(\mathbb{R}^{d}\right)$ and so, the problem (1.51) has a solution. In the case $d \leq 2$, the claim follows since, by using (1.47), (1.48) and interpolating, we can still estimate $\left\|u_{n}\right\|_{q^{\prime}}$ by means of $\left\|\nabla u_{n}\right\|_{2}$ and $\left\|u_{n}\right\|_{2 p /(p-1)}$.

Repeating the argument of Section 1.5, one obtains an existence result for (1.49) in the case $p=1$, too.
Proposition 1.28. Let $f \in L^{q}\left(\mathbb{R}^{d}\right)$, where $q \in\left[1, \frac{2 d}{d+2}\right]$, if $d \geq 3$, and $q=1$, if $d=1,2$. Then there is a unique solution $V_{1}$ of problem (1.49) with $p=1$, which is given by

$$
V_{1}=\frac{f}{M}\left(\mathbf{1}_{\omega_{+}}-\mathbf{1}_{\omega_{-}}\right),
$$

where $M=\left\|u_{1}\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)}, \omega_{+}=\left\{u_{1}=M\right\}, \omega_{-}=\left\{u_{1}=-M\right\}$, and $u_{1}$ is the unique minimizer of

$$
\begin{equation*}
\min \left\{\frac{1}{2} \int_{\mathbb{R}^{d}}|\nabla u|^{2} d x+\frac{1}{2}\|u\|_{L^{\infty}}^{2}-\int_{\mathbb{R}^{d}} u f d x: u \in \dot{H}^{1}\left(\mathbb{R}^{d}\right) \cap L^{\infty}\left(\mathbb{R}^{d}\right)\right\} . \tag{1.53}
\end{equation*}
$$

In particular, $\int_{\omega_{+}} f d x-\int_{\omega_{-}} f d x=M, f \geq 0$ on $\omega_{+}$and $f \leq 0$ on $\omega_{-}$.
We note that, when $p=1$, the support of the optimal potential $V_{1}$ is contained in the support of the function $f$. This is not the case if $p>1$, as the following example shows.

Example 1.29. Let $f=\mathbf{1}_{B_{1}(0)}$ and $p>1$. By our previous analysis we know that there exist a solution $u_{p}$ for problem (1.51) and a solution $V_{p}$ for problem (1.49) given by (1.50). We note that $u_{p}$ is positive, radially decreasing and satisfies the equation

$$
-u^{\prime \prime}(r)-\frac{d-1}{r} u^{\prime}(r)+C u^{\alpha}=0, \quad r \in(1,+\infty),
$$

where $\alpha=2 p /(p-1)>2$ and $C$ is a positive constant. Thus, we have that

$$
u_{p}(r)=k r^{2 /(1-\alpha)},
$$

where $k$ is an explicit constant depending on $C, d$ and $\alpha$. In particular $u_{p}$ is not compactly supported on $\mathbb{R}^{d}$.

### 1.6.2 Optimal potentials in $\mathbb{R}^{d}$ with unbounded constraint

In this subsection we consider the problems

$$
\begin{align*}
& \min \left\{\mathcal{E}_{f}(V): V \geq 0, \int_{\mathbb{R}^{d}} V^{-p} d x \leq 1\right\},  \tag{1.54}\\
& \min \left\{\lambda_{1}(V): V \geq 0, \int_{\mathbb{R}^{d}} V^{-p} d x \leq 1\right\}, \tag{1.55}
\end{align*}
$$

for $p>0$ and $f \in L^{q}\left(\mathbb{R}^{d}\right)$. We will see in Proposition 1.30 that in order to have existence for (1.54) the parameter $q$ must satisfy some constraints, depending on the value of $p$ and on the dimension $d$. Namely, we need $q$ to satisfy one of the following conditions

$$
\begin{array}{r}
q \in\left[\frac{2 d}{d+2}, \frac{2 p}{p-1}\right], \text { if } d \geq 3 \text { and } p>1 \\
q \in\left[\frac{2 d}{d+2},+\infty\right], \text { if } d \geq 3 \text { and } p \leq 1 \\
q \in\left(1, \frac{2 p}{p-1}\right], \text { if } d=2 \text { and } p>1  \tag{1.56}\\
q \in(1,+\infty], \text { if } d=2 \text { and } p \leq 1 \\
q \in\left[1, \frac{2 p}{p-1}\right], \text { if } d=1 \text { and } p>1 \\
q \in[1,+\infty], \text { if } d=1 \text { and } p \leq 1
\end{array}
$$

We say that $q=q(p, d) \in[1,+\infty]$ is admissible if it satisfies (1.56). Note that $q=2$ is admissible for any $d \geq 1$ and any $p>0$.

Proposition 1.30. Let $p>0$ and $f \in L^{q}\left(\mathbb{R}^{d}\right)$, where $q$ is admissible in the sense of (1.56). Then the minimization problem (1.54) has a solution $V_{p}$ given by

$$
\begin{equation*}
V_{p}=\left(\int_{\mathbb{R}^{d}}\left|u_{p}\right|^{2 p /(p+1)} d x\right)^{1 / p}\left|u_{p}\right|^{-2 /(1+p)} \tag{1.57}
\end{equation*}
$$

where $u_{p}$ is a minimizer of

$$
\begin{array}{r}
\min \left\{\frac{1}{2} \int_{\mathbb{R}^{d}}|\nabla u|^{2} d x+\frac{1}{2}\left(\int_{\mathbb{R}^{d}}|u|^{2 p /(p+1)} d x\right)^{(p+1) / p}-\int_{\mathbb{R}^{d}} u f d x:\right.  \tag{1.58}\\
\left.u \in \dot{H}^{1}\left(\mathbb{R}^{d}\right),|u|^{2 p /(p+1)} \in L^{1}\left(\mathbb{R}^{d}\right)\right\}
\end{array}
$$

Moreover, if $p \geq 1$, then the functional in (1.58) is convex, its minimizer is unique and so is the solution of (1.54).

Proof. By means of (1.46), (1.47) and (1.48), and thanks to the admissibility of $q$, we get the existence of a solution of (1.58) through an interpolation argument similar to the one used in the proof of Proposition 1.27. The existence of an optimal potential follows reasoning as in Subsection 1.4.1.

In Example 1.29, we showed that the optimal potentials for (1.49), may be supported on the whole $\mathbb{R}^{d}$. The analogous question for the problem (1.54) is whether the optimal potentials given by (1.57) have a bounded set of finiteness $\left\{V_{p}<+\infty\right\}$. In order to
answer this question, it is sufficient to study the support of the solutions $u_{p}$ of (1.58), which solve the equation

$$
\begin{equation*}
-\Delta u+C_{p}|u|^{-2 /(p+1)} u=f \tag{1.59}
\end{equation*}
$$

where $C_{p}>0$ is a constant depending on $p$.
Proposition 1.31. Let $p>0$ and let $f \in L^{q}\left(\mathbb{R}^{d}\right)$, for $q>d / 2$, be a non-negative function with a compact support. Then every solution $u_{p}$ of problem (1.58) has compact support.

Proof. Without loss of generality we may assume that $f$ is supported in the unit ball of $\mathbb{R}^{d}$. We first prove the result when $f$ is radially decreasing. In this case $u_{p}$ is also radially decreasing and non-negative. Let $v$ be the function defined by $v(|x|)=u_{p}(x)$. Thus $v$ satisfies the equation

$$
\left\{\begin{array}{l}
-v^{\prime \prime}-\frac{d-1}{r} v^{\prime}+C_{p} v^{s}=0 \quad r \in(1,+\infty),  \tag{1.60}\\
v(1)=u_{p}(1),
\end{array}\right.
$$

where $s=(p-1) /(p+1)$ and $C_{p}>0$ is a constant depending on $p$. Since $v \geq 0$ and $v^{\prime} \leq 0$, we have that $v$ is convex. Moreover, since

$$
\int_{1}^{+\infty} v^{2} r^{d-1} d r<+\infty, \quad \int_{1}^{+\infty}\left|v^{\prime}\right|^{2} r^{d-1} d r<+\infty
$$

we have that $v, v^{\prime}$ and $v^{\prime \prime}$ vanish at infinity. Multiplying (1.60) by $v^{\prime}$ we obtain

$$
\left(\frac{v^{\prime}(r)^{2}}{2}-C_{p} \frac{v(r)^{s+1}}{s+1}\right)^{\prime}=-\frac{d-1}{r} v^{\prime}(r)^{2} \leq 0 .
$$

Thus the function $v^{\prime}(r)^{2} / 2-C_{p} v(r)^{s+1} /(s+1)$ is decreasing and vanishing at infinity and thus non-negative. Thus we have

$$
\begin{equation*}
-v^{\prime}(r) \geq C v(r)^{(s+1) / 2}, r \in(1,+\infty) \tag{1.61}
\end{equation*}
$$

where $C=\left(2 C_{p} /(s+1)\right)^{1 / 2}$. Arguing by contradiction, suppose that $v$ is strictly positive on $(1,+\infty)$. Dividing both sides of (1.61) and integrating, we have

$$
-v(r)^{(1-s) / 2} \geq A r+B
$$

where $A=2 C /(1-s)$ and $B$ is determined by the initial datum $v(1)$. This cannot occur, since the left hand side is negative, while the right hand side goes to $+\infty$, as $r \rightarrow+\infty$.

We now prove the result for a generic compactly supported and non-negative $f \in$ $L^{q}\left(\mathbb{R}^{d}\right)$. Since the solution $u_{p}$ of (1.58) is non-negative and is a weak solution of (1.59), we have that on each ball $B_{R} \subset \mathbb{R}^{d}, u_{p} \leq u$, where $u \in H^{1}\left(B_{R}\right)$ is the solution of

$$
-\Delta u=f \text { in } B_{R}, \quad u=u_{p} \text { on } \partial B_{R} .
$$

Since $f \in L^{d / 2}\left(\mathbb{R}^{d}\right)$, by $[67$, Theorem 9.11] and a standard bootstrap argument on the integrability of $u$, we have that $u$ is continuous on $B_{R / 2}$. As a consequence, $u_{p}$ is locally bounded in $\mathbb{R}^{d}$. In particular, it is bounded since $u_{p} \wedge M$, where $M=\left\|u_{p}\right\|_{L^{\infty}\left(B_{1}\right)}$, is a better competitor than $u_{p}$ in (1.58). Let $w$ be a radially decreasing minimizer of (1.58) with $f=\mathbf{1}_{B_{1}}$. Thus $w$ is a solution of the PDE

$$
-\Delta w+C_{p} w^{s}=\mathbf{1}_{B_{1}},
$$

in $\mathbb{R}^{d}$, where $C_{p}$ is as in (1.60). Then, the function $w_{t}(x)=t^{2 /(1-s)} w(x / t)$ is a solution of the equation

$$
-\Delta w_{t}+C_{p} w_{t}^{s}=t^{2 s /(1-s)} \mathbf{1}_{B_{t}} .
$$

Since $u_{p}$ is bounded, there exists some $t \geq 1$ large enough such that $w_{t} \geq u_{p}$ on the ball $B_{t}$. Moreover, $w_{t}$ minimizes (1.58) with $f=t^{2 s /(1-s)} \mathbf{1}_{B_{t}}$ and so $w_{t} \geq u_{p}$ on $\mathbb{R}^{d}$ (otherwise $w_{t} \wedge u_{p}$ would be a better competitor in (1.58) than $w_{p}$ ). The conclusion follows since, by the first step of the proof, $w_{t}$ has compact support.

The problems (1.55) and (1.54) are similar both in the questions of existence and the qualitative properties of the solutions.

Proposition 1.32. For every $p>0$ there is a solution of the problem (1.55) given by

$$
\begin{equation*}
V_{p}=\left(\int_{\mathbb{R}^{d}}\left|u_{p}\right|^{2 p /(p+1)} d x\right)^{1 / p}\left|u_{p}\right|^{-2 /(1+p)}, \tag{1.62}
\end{equation*}
$$

where $u_{p}$ is a radially decreasing minimizer of

$$
\begin{equation*}
\min \left\{\int_{\mathbb{R}^{d}}|\nabla u|^{2} d x+\left(\int_{\mathbb{R}^{d}}|u|^{2 p /(p+1)} d x\right)^{(p+1) / p}: u \in H^{1}\left(\mathbb{R}^{d}\right), \int_{\mathbb{R}^{d}} u^{2} d x=1\right\} . \tag{1.63}
\end{equation*}
$$

Moreover, $u_{p}$ has a compact support, hence the set $\left\{V_{p}<+\infty\right\}$ is a ball of finite radius in $\mathbb{R}^{d}$.

Proof. Let us first show that the minimum in (1.63) is achieved. Let $u_{n} \in H^{1}\left(\mathbb{R}^{d}\right)$ be a minimizing sequence of positive functions normalized in $L^{2}$. Note that by the PólyaSzegö inequality we may assume that each of these functions is radially decreasing in $\mathbb{R}^{d}$ and so we will use the identification $u_{n}=u_{n}(r)$. In order to prove that the minimum is achieved it is enough to show that the sequence $u_{n}$ converges in $L^{2}\left(\mathbb{R}^{d}\right)$. Indeed, since
$u_{n}$ is a radially decreasing minimizing sequence, there exists $C>0$ such that for each $r>0$ we have

$$
u_{n}(r)^{2 p /(p+1)} \leq \frac{1}{\left|B_{r}\right|} \int_{B_{r}} u_{n}^{2 p /(p+1)} d x \leq \frac{C}{r^{d}} .
$$

Thus, for each $R>0$, we obtain

$$
\begin{equation*}
\int_{B_{R}^{c}} u_{n}^{2} d x \leq C_{1} \int_{R}^{+\infty} r^{-d(p+1) / p} r^{d-1} d r=C_{2} R^{-1 / p} \tag{1.64}
\end{equation*}
$$

where $C_{1}$ and $C_{2}$ do not depend on $n$ and $R$. Since the sequence $u_{n}$ is bounded in $H^{1}\left(\mathbb{R}^{d}\right)$, it converges locally in $L^{2}\left(\mathbb{R}^{d}\right)$ and, by (1.64), this convergence is also strong in $L^{2}\left(\mathbb{R}^{d}\right)$. Thus, we obtain the existence of a radially symmetric and decreasing solution $u_{p}$ of (1.63) and so, of an optimal potential $V_{p}$ given by (1.62).

We now prove that the support of $u_{p}$ is a ball of finite radius. By the radial symmetry of $u_{p}$ we can write it in the form $u_{p}(x)=u_{p}(|x|)=u_{p}(r)$, where $r=|x|$. With this notation, $u_{p}$ satisfies the equation:

$$
-u_{p}^{\prime \prime}-\frac{d-1}{r} u_{p}^{\prime}+C_{p} u_{p}^{s}=\lambda u_{p},
$$

where $s=(p-1) /(p+1)<1$ and $C_{p}>0$ is a constant depending on $p$. Arguing as in Proposition 1.31, we obtain that, for $r$ large enough,

$$
-u_{p}^{\prime}(r) \geq\left(\frac{C_{p}}{s+1} u_{p}(r)^{s+1}-\frac{\lambda}{2} u_{p}(r)^{2}\right)^{1 / 2} \geq\left(\frac{C_{p}}{2(s+1)} u_{p}(r)^{s+1}\right)^{1 / 2}
$$

where, in the last inequality, we used the fact that $u_{p}(r) \rightarrow 0$, as $r \rightarrow \infty$, and $s+1<2$. Integrating both sides of the above inequality, we conclude that $u_{p}$ has a compact support.

Remark 1.33. We note that the solution $u_{p} \in H^{1}\left(\mathbb{R}^{d}\right)$ of (1.63) is the function for which the best constant $C$ in the interpolated Gagliardo-Nirenberg-Sobolev inequality

$$
\begin{equation*}
\|u\|_{2} \leq C\|\nabla u\|_{2}^{d /(d+2 p)}\|u\|_{2 p /(p+1)}^{2 p /(d+2 p)} \tag{1.65}
\end{equation*}
$$

is achieved. Indeed, for any $u \in H^{1}\left(\mathbb{R}^{d}\right)$ and any $t>0$, we define $u_{t}(x):=t^{d / 2} u(t x)$. Thus, we have that $\|u\|_{2}=\left\|u_{t}\right\|_{2}$, for any $t>0$. Moreover, up to a rescaling, we may assume that the function $g:(0,+\infty) \rightarrow \mathbb{R}$, defined by

$$
\begin{aligned}
g(t) & =\int_{\mathbb{R}^{d}}\left|\nabla u_{t}\right|^{2} d x+\left(\int_{\mathbb{R}^{d}}\left|u_{t}\right|^{2 p /(p+1)} d x\right)^{(p+1) / p} \\
& =t^{2} \int_{\mathbb{R}^{d}}|\nabla u|^{2} d x+t^{-d / p}\left(\int_{\mathbb{R}^{d}}|u|^{2 p /(p+1)} d x\right)^{(p+1) / p},
\end{aligned}
$$

achieves its minimum in the interval $(0,+\infty)$ and, moreover, we have

$$
\min _{t \in(0,+\infty)} g(t)=C\left(\int_{\mathbb{R}^{d}}|\nabla u|^{2} d x\right)^{d /(d+2 p)}\left(\int_{\mathbb{R}^{d}}|u|^{\frac{2 p}{p+1}} d x\right)^{2(p+1) /(d+2 p)}
$$

where $C$ is a constant depending on $p$ and $d$. In the case $u=u_{p}$, the minimum of $g$ is achieved for $t=1$ and so, we have that $u_{p}$ is a solution also of
$\min \left\{\left(\int_{\mathbb{R}^{d}}|\nabla u|^{2} d x\right)^{d /(d+2 p)}\left(\int_{\mathbb{R}^{d}}|u|^{2 p /(p+1)} d x\right)^{2(p+1) /(d+2 p)}: u \in H^{1}\left(\mathbb{R}^{d}\right), \int_{\mathbb{R}^{d}} u^{2} d x=1\right\}$,
which is just another form of (1.65).

### 1.6.3 Further remarks

We recall (see [26]) that the injection $H_{V}^{1}\left(\mathbb{R}^{d}\right) \hookrightarrow L^{2}\left(\mathbb{R}^{d}\right)$ is compact whenever the potential $V$ satisfies $\int_{\mathbb{R}^{d}} V^{-p} d x<+\infty$ for some $0<p \leq 1$. In this case the spectrum of the Schrödinger operator $-\Delta+V$ is discrete and we denote by $\lambda_{k}(V)$ its eigenvalues. The existence of an optimal potential for spectral optimization problems of the form

$$
\begin{equation*}
\min \left\{\lambda_{k}(V): V \geq 0, \int_{\mathbb{R}^{d}} V^{-p} d x \leq 1\right\} \tag{1.66}
\end{equation*}
$$

for general $k \in \mathbb{N}$, cannot be deduced by the direct methods used in Subsection 1.6.2. In this last section we make the following conjectures:

- For every $k \geq 1$, there is a solution $V_{k}$ of the problem (1.66).
- The set of finiteness $\left\{V_{k}<+\infty\right\}$, of the optimal potential $V_{k}$, is bounded.

In what follows, we prove an existence result in the case $k=2$. We first recall that, by Proposition 1.32, there exists an optimal potential $V_{p}$, for $\lambda_{1}$, such that the set of finiteness $\left\{V_{p}<+\infty\right\}$ is a ball. Thus, we have a situation analogous to the Faber-Krahn inequality, which states that the minimum

$$
\begin{equation*}
\min \left\{\lambda_{1}(\Omega): \Omega \subset \mathbb{R}^{d},|\Omega|=c\right\} \tag{1.67}
\end{equation*}
$$

is achieved for the ball of measure $c$. We recall that, starting from (1.67), one may deduce, by a simple argument (see for instance [75]), the Krahn-Szegö inequality, which states that the minimum

$$
\begin{equation*}
\min \left\{\lambda_{2}(\Omega): \Omega \subset \mathbb{R}^{d},|\Omega|=c\right\} \tag{1.68}
\end{equation*}
$$

is achieved for a disjoint union of equal balls. In the case of potentials one can find two optimal potentials for $\lambda_{1}$ with disjoint sets of finiteness and then apply the argument from the proof of the Krahn-Szegö inequality. In fact, we have the following result.

Proposition 1.34. There exists an optimal potential, solution of (1.66) with $k=2$. Moreover, any optimal potential is of the form $\min \left\{V_{1}, V_{2}\right\}$, where $V_{1}$ and $V_{2}$ are optimal potentials for $\lambda_{1}$ which have disjoint sets of finiteness $\left\{V_{1}<+\infty\right\} \cap\left\{V_{2}<+\infty\right\}=\emptyset$ and are such that $\int_{\mathbb{R}^{d}} V_{1}^{-p} d x=\int_{\mathbb{R}^{d}} V_{2}^{-p} d x=1 / 2$.
Proof. Given $V_{1}$ and $V_{2}$ as above, we prove that for every $V: \mathbb{R}^{d} \rightarrow[0,+\infty]$ with $\int_{\mathbb{R}^{d}} V^{-p} d x=1$, we have

$$
\lambda_{2}\left(\min \left\{V_{1}, V_{2}\right\}\right) \leq \lambda_{2}(V) .
$$

Indeed, let $u_{2}$ be the second eigenfunction of $-\Delta+V$. We first suppose that $u_{2}$ changes sign on $\mathbb{R}^{d}$ and consider the functions $V_{+}=\sup \left\{V, \infty_{\left\{u_{2} \leq 0\right\}}\right\}$ and $V_{-}=$ $\sup \left\{V, \infty_{\left\{u_{2} \geq 0\right\}}\right\}$ where, for any measurable $A \subset \mathbb{R}^{d}$, we set

$$
\infty_{A}(x)=\left\{\begin{array}{lc}
+\infty, & x \in A, \\
0, & x \notin A .
\end{array}\right.
$$

We note that

$$
1=\int_{\mathbb{R}^{d}} V^{-p} d x=\int_{\mathbb{R}^{d}} V_{+}^{-p} d x+\int_{\mathbb{R}^{d}} V_{-}^{-p} d x
$$

Moreover, on the sets $\left\{u_{2}>0\right\}$ and $\left\{u_{2}<0\right\}$, the following equations are satisfied:

$$
-\Delta u_{2}^{+}+V_{+} u_{2}^{+}=\lambda_{2}(V) u_{2}^{+}, \quad-\Delta u_{2}^{-}+V_{-} u_{2}^{-}=\lambda_{2}(V) u_{2}^{-},
$$

and so, multiplying respectively by $u_{2}^{+}$and $u_{2}^{-}$, we obtain that

$$
\begin{equation*}
\lambda_{2}(V) \geq \lambda_{1}\left(V_{+}\right), \quad \lambda_{2}(V) \geq \lambda_{1}\left(V_{-}\right), \tag{1.69}
\end{equation*}
$$

where we have equalities if, and only if, $u_{2}^{+}$and $u_{2}^{-}$are the first eigenfunctions corresponding to $\lambda_{1}\left(V_{+}\right)$and $\lambda_{1}\left(V_{-}\right)$. Let now $\widetilde{V}_{+}$and $\tilde{V}_{-}$be optimal potentials for $\lambda_{1}$ corresponding to the constraints

$$
\int_{\mathbb{R}^{d}} \widetilde{V}_{+}^{-p} d x=\int_{\mathbb{R}^{d}} V_{+}^{-p} d x, \quad \int_{\mathbb{R}^{d}} \tilde{V}_{-}^{-p} d x=\int_{\mathbb{R}^{d}} V_{-}^{-p} d x
$$

By Proposition 1.32, the sets of finiteness of $\widetilde{V}_{+}$and $\widetilde{V}_{-}$are compact, hence we may assume (up to translations) that they are also disjoint. By the monotonicity of $\lambda_{1}$, we have

$$
\max \left\{\lambda_{1}\left(V_{1}\right), \lambda_{1}\left(V_{2}\right)\right\} \leq \max \left\{\lambda_{1}\left(\widetilde{V}_{+}\right), \lambda_{1}\left(\widetilde{V}_{-}\right)\right\},
$$

and so we obtain

$$
\lambda_{2}\left(\min \left\{V_{1}, V_{2}\right\}\right) \leq \max \left\{\lambda_{1}\left(\widetilde{V}_{+}\right), \lambda_{1}\left(\widetilde{V}_{-}\right)\right\} \leq \max \left\{\lambda_{1}\left(V_{+}\right), \lambda_{1}\left(V_{-}\right)\right\} \leq \lambda_{2}(V),
$$

as required. If $u_{2}$ does not change sign, then we consider $V_{+}=\sup \left\{V, \infty_{\left\{u_{2}=0\right\}}\right\}$ and $V_{-}=\sup \left\{V, \infty_{\left\{u_{1}=0\right\}}\right\}$, where $u_{1}$ is the first eigenfunction of $-\Delta+V$. Then the claim follows by the same argument as above.

## Chapter 2

## Optimization problems for metric graphs

### 2.1 Introduction

This chapter is based on the paper [32], written in collaboration with Giuseppe Buttazzo and Bozhidar Velichkov. The issue we take into account, is an adaptation of the problem of minimizing the first eigenvalue and the Energy function of the Dirichlet Laplacian to the class of graphs. To give a precise meaning to the previous sentence we must define what is a graph, and introduce a suitable concept of differential operator on it. Initially, we define a graph $C$ in $\mathbb{R}^{d}$ to be simply a closed connected subset of $\mathbb{R}^{d}$ with finite 1 -dimensional Hausdorff measure $\mathcal{H}^{1}(C)$. Since such sets are rectifiable (see for instance [4]) the standard theory on rectifiable sets allows us to define all the variational tools that are usually defined in the Euclidean setting:

- Dirichlet integral $\int_{C} \frac{1}{2}\left|u^{\prime}\right|^{2} d \mathcal{H}^{1}$;
- Sobolev spaces

$$
\begin{aligned}
& H^{1}(C)=\left\{u \in L^{2}(C): \int_{C}\left|u^{\prime}\right|^{2} d \mathcal{H}^{1}<+\infty\right\}, \\
& H_{0}^{1}(C ; \mathcal{D})=\left\{u \in H^{1}(C): u=0 \text { on } \mathcal{D}\right\} ;
\end{aligned}
$$

- (Dirichlet) Energy

$$
\mathcal{E}(C ; \mathcal{D})=\inf \left\{\int_{C}\left(\frac{1}{2}\left|u^{\prime}\right|^{2}-u\right) d \mathcal{H}^{1}: u \in H_{0}^{1}(C, \mathcal{D})\right\}
$$

In particular, for a fixed set $\mathcal{D}$ consisting of $d$ points, $\mathcal{D}=\left\{D_{1}, \ldots, D_{d}\right\}$, we consider the shape optimization problem

$$
\begin{equation*}
\min \left\{\mathcal{E}(C ; \mathcal{D}): \mathcal{H}^{1}(C)=l, \mathcal{D} \subset C\right\} \tag{2.1}
\end{equation*}
$$

where the total length $l$ is fixed. Notice that in the problem above the unknown is the graph $C$ and no a priori constraints on its topology are imposed.

In spite of the fact that the optimization problem (2.1) looks very natural, we show that in general an optimal graph may not exist (see Example 2.23); this leads us to consider another, larger, admissible class consisting of the so-called metric graphs, for which the embedding into $\mathbb{R}^{d}$ is not required. The precise definition of a metric graph is given in Section 2.3; roughly speaking they are metric spaces induced by combinatorial graphs with weighted edges.

Our main result is an existence theorem for optimal metric graphs, where the cost functional is the extension of the Energy functional defined above. In Section 2.4 we show some explicit examples of optimal metric graphs. The last section contains some remarks on possible extensions of our main result to other similar problems and on some open questions.

### 2.2 Sobolev space and Dirichlet Energy of a rectifiable set

Let $C \subset \mathbb{R}^{d}$ be a closed connected set of finite length, i.e. $\mathcal{H}^{1}(C)<\infty$, where $\mathcal{H}^{1}$ denotes the one-dimensional Hausdorff measure. On the set $C$ we consider the metric

$$
d(x, y)=\inf \left\{\int_{0}^{1}|\dot{\gamma}(t)| d t: \gamma:[0,1] \rightarrow \mathbb{R}^{d} \text { Lipschitz, } \gamma([0,1]) \subset C, \gamma(0)=x, \gamma(1)=y\right\}
$$

which is finite since, by the First Rectifiability Theorem (see [4, Theorem 4.4.1]), there is at least one rectifiable curve in $C$ connecting $x$ to $y$. For any function $u: C \rightarrow \mathbb{R}$, Lipschitz with respect to the distance $d$ (we also use the term $d$-Lipschitz), we define the norm

$$
\|u\|_{H^{1}(C)}^{2}=\int_{C}|u(x)|^{2} d \mathcal{H}^{1}(x)+\int_{C}\left|u^{\prime}\right|(x)^{2} d \mathcal{H}^{1}(x)
$$

where

$$
\left|u^{\prime}\right|(x)=\varlimsup_{y \rightarrow x} \frac{|u(y)-u(x)|}{d(x, y)}
$$

The Sobolev space $H^{1}(C)$ is the closure of the $d$-Lipschitz functions on $C$ with respect to the norm $\|\cdot\|_{H^{1}(C)}$.

Remark 2.1. The inclusion $H^{1}(C) \subset C_{d}(C)$ is compact, where $C_{d}(C)$ indicates the space of real-valued continuous functions on $C$, with respect to the metric $d$. In fact, for each $x, y \in C$, there is a rectifiable curve $\gamma:[0, d(x, y)] \rightarrow C$ connecting $x$ to $y$,
which we may assume arc-length parametrized. Thus, for any $u \in H^{1}(C)$, we have that

$$
\begin{aligned}
|u(x)-u(y)| & \leq \int_{0}^{d(x, y)}\left|\frac{d}{d t} u(\gamma(t))\right| d t \\
& \leq d(x, y)^{1 / 2}\left(\int_{0}^{d(x, y)}\left|\frac{d}{d t} u(\gamma(t))\right|^{2} d t\right)^{1 / 2} \\
& \leq d(x, y)^{1 / 2}\left\|u^{\prime}\right\|_{L^{2}(C)}
\end{aligned}
$$

and so, $u$ is $1 / 2-$ Hölder continuous. On the other hand, for any $x \in C$, we have that

$$
\int_{C} u(y) d \mathcal{H}^{1}(y) \geq \int_{C}\left(u(x)-d(x, y)^{1 / 2}\left\|u^{\prime}\right\|_{L^{2}(C)}\right) d \mathcal{H}^{1}(y) \geq l u(x)-l^{3 / 2}\left\|u^{\prime}\right\|_{L^{2}(C)}
$$

where $l=\mathcal{H}^{1}(C)$. Thus, we obtain the $L^{\infty}$ bound

$$
\|u\|_{L^{\infty}} \leq l^{-1 / 2}\|u\|_{L^{2}(C)}+l^{1 / 2}\left\|u^{\prime}\right\|_{L^{2}(C)} \leq\left(l^{-1 / 2}+l^{1 / 2}\right)\|u\|_{H^{1}(C)} .
$$

and so, by the Ascoli-Arzelá Theorem, we have that the inclusion is compact.
Remark 2.2. By the same argument as in Remark 2.1 above, we have that for any $u \in H^{1}(C)$, the $(1,2)$-Poincaré inequality holds, i.e.

$$
\begin{equation*}
\int_{C}\left|u(x)-\frac{1}{l} \int_{C} u d \mathcal{H}^{1}\right| d \mathcal{H}^{1}(x) \leq l^{3 / 2}\left(\int_{C}\left|u^{\prime}\right|^{2} d \mathcal{H}^{1}\right)^{1 / 2} \tag{2.2}
\end{equation*}
$$

Moreover, if $u \in H^{1}(C)$ is such that $u(x)=0$ for some point $x \in C$, then we have the Poincaré inequality:

$$
\begin{equation*}
\|u\|_{L^{2}(C)} \leq l^{1 / 2}\|u\|_{L^{\infty}(C)} \leq l\left\|u^{\prime}\right\|_{L^{2}(C)} \tag{2.3}
\end{equation*}
$$

Since $C$ is supposed connected, by the Second Rectifiability Theorem (see [4, Theorem 4.4.8]) there exists a countable family of injective arc-length parametrized Lipschitz curves $\gamma_{i}:\left[0, l_{i}\right] \rightarrow C, i \in \mathbb{N}$ and an $\mathcal{H}^{1}-$ negligible set $N \subset C$ such that

$$
C=N \cup\left(\bigcup_{i} \operatorname{Im}\left(\gamma_{i}\right)\right),
$$

where $\operatorname{Im}\left(\gamma_{i}\right)=\gamma_{i}\left(\left[0, l_{i}\right]\right)$. By the chain rule (see Lemma 2.3 below) we have

$$
\left|\frac{d}{d t} u\left(\gamma_{i}(t)\right)\right|=\left|u^{\prime}\right|\left(\gamma_{i}(t)\right), \quad \forall i \in \mathbb{N}
$$

and so, we obtain for the norm of $u \in H^{1}(C)$ :

$$
\begin{equation*}
\|u\|_{H^{1}(C)}^{2}=\int_{C}|u(x)|^{2} d \mathcal{H}^{1}(x)+\sum_{i} \int_{0}^{l_{i}}\left|\frac{d}{d t} u\left(\gamma_{i}(t)\right)\right|^{2} d t \tag{2.4}
\end{equation*}
$$

Moreover, we have the inclusion

$$
\begin{equation*}
H^{1}(C) \subset \oplus_{i \in} H^{1}\left(\left[0, l_{i}\right]\right), \tag{2.5}
\end{equation*}
$$

which gives the reflexivity of $H^{1}(C)$ and the lower semicontinuity of the $H^{1}(C)$ norm, with respect to the strong convergence in $L^{2}(C)$.
Lemma 2.3. Let $\gamma:[0, l] \rightarrow \mathbb{R}^{d}$ be an injective arc-length parametrized Lipschitz curve with $\gamma([0, l]) \subset C$. Then we have

$$
\begin{equation*}
\left|\frac{d}{d t} u(\gamma(t))\right|=\left|u^{\prime}\right|(\gamma(t)), \quad \text { for } \mathcal{H}^{1}-\text { a.e. } t \in[0, l] \text {. } \tag{2.6}
\end{equation*}
$$

Proof. Let $u: C \rightarrow \mathbb{R}$ be a Lipschitz map with Lipschitz constant Lip $(u)$ with respect to the distance $d$. We prove that the chain rule (2.6) holds in all the points $t \in[0, l]$ which are Lebesgue points for $\left|\frac{d}{d t} u(\gamma(t))\right|$ and such that the point $\gamma(t)$ has density one, i.e.

$$
\begin{equation*}
\lim _{r \rightarrow 0} \frac{\mathcal{H}^{1}\left(C \cap B_{r}(\gamma(t))\right)}{2 r}=1, \tag{2.7}
\end{equation*}
$$

(thus almost every points, see for istance [92]) where $B_{r}(x)$ indicates the ball of radius $r$ in $\mathbb{R}^{d}$. Since, $\mathcal{H}^{1}$-almost all points $x \in C$ have this property, we obtain the conclusion. Without loss of generality, we consider $t=0$. Let us first prove that $\left|u^{\prime}\right|(\gamma(0)) \geq$ $\left|\frac{d}{d t} u(\gamma(0))\right|$. We have that

$$
\left|u^{\prime}\right|(\gamma(0)) \geq \overline{\lim }_{t \rightarrow 0} \frac{|u(\gamma(t))-u(\gamma(0))|}{d(\gamma(t), \gamma(0))}=\left|\frac{d}{d t} u(\gamma(0))\right|,
$$

since $\gamma$ is arc-length parametrized. On the other hand, we have

$$
\begin{align*}
&\left|u^{\prime}\right|(x)=\varlimsup_{\lim }^{y \rightarrow x} \\
&=\lim _{n \rightarrow \infty} \frac{|u(y)-u(x)|}{d(y, x)} \\
&=\lim _{n \rightarrow \infty} \frac{\left|u\left(\gamma_{n}\left(r_{n}\right)\right)-u(x)\right|}{r_{n}} \\
& \leq \lim _{n \rightarrow \infty} \frac{1}{r_{n}} \int_{0}^{r_{n}}\left|\frac{d}{d t} u\left(\gamma_{n}(0)\right)\right|  \tag{2.8}\\
&\hline(t)) \mid d t
\end{align*}
$$

where $y_{n} \in C$ is a sequence of points which realizes the limsup and $\gamma_{n}:\left[0, r_{n}\right] \rightarrow \mathbb{R}^{d}$ is a geodesic in $C$ connecting $x$ to $y_{n}$. Let $S_{n}=\left\{t: \gamma_{n}(t)=\gamma(t)\right\} \subset\left[0, r_{n}\right]$, then, we have

$$
\begin{align*}
\int_{0}^{r_{n}}\left|\frac{d}{d t} u\left(\gamma_{n}(t)\right)\right|^{2} d t & \leq \int_{S_{n}}\left|\frac{d}{d t} u(\gamma(t))\right|^{2} d t+\operatorname{Lip}(u)\left(r_{n}-\left|S_{n}\right|\right) \\
& \leq \int_{0}^{r_{n}}\left|\frac{d}{d t} u(\gamma(t))\right|^{2} d t+\operatorname{Lip}(u)\left(\mathcal{H}^{1}\left(B_{r_{n}}(\gamma(0)) \cap C\right)-2 r_{n}\right), \tag{2.9}
\end{align*}
$$

and so, since $\gamma(0)$ is of density 1 , we conclude applying this estimate to (2.8).

Given a set of points $\mathcal{D}=\left\{D_{1}, \ldots, D_{k}\right\} \subset \mathbb{R}^{d}$ we define the admissible class $\mathcal{A}(\mathcal{D} ; l)$ as the family of all closed connected sets $C$ containing $\mathcal{D}$ and of length $\mathcal{H}^{1}(C)=l$. For any $C \in \mathcal{A}(\mathcal{D} ; l)$ we consider the space of Sobolev functions which satisfy a Dirichlet condition at the points $D_{i}$ :

$$
H_{0}^{1}(C ; \mathcal{D})=\left\{u \in H^{1}(C): u\left(D_{j}\right)=0, j=1 \ldots, k\right\}
$$

which is well-defined by Remark 2.1. For the points $D_{i}$ we use the term Dirichlet points. The Dirichlet Energy of the set $C$ with respect to $D_{1}, \ldots, D_{k}$ is defined as

$$
\begin{equation*}
\mathcal{E}(C ; \mathcal{D})=\min \left\{J(u): u \in H_{0}^{1}(C ; \mathcal{D})\right\}, \tag{2.10}
\end{equation*}
$$

where

$$
\begin{equation*}
J(u)=\frac{1}{2} \int_{C}\left|u^{\prime}\right|(x)^{2} d \mathcal{H}^{1}(x)-\int_{C} u(x) d \mathcal{H}^{1}(x) . \tag{2.11}
\end{equation*}
$$

Remark 2.4. For any $C \in \mathcal{A}(\mathcal{D} ; l)$ there exists a unique minimizer of the functional $J: H_{0}^{1}(C ; \mathcal{D}) \rightarrow \mathbb{R}$. In fact, by Remark 2.1 we have that a minimizing sequence is bounded in $H^{1}$ and compact in $L^{2}$. The conclusion follows by the semicontinuity of the $L^{2}$ norm of the gradient, with respect to the strong $L^{2}$ convergence, which is an easy consequence of equation (2.4). The uniqueness follows by the strict convexity of the $L^{2}$ norm and the sub-additivity of the gradient $\left|u^{\prime}\right|$. We call the minimizer of $J$ the Energy function of $C$ with Dirichlet conditions in $D_{1}, \ldots, D_{k}$.

Remark 2.5. Let $u \in H^{1}(C)$ and $v: C \rightarrow \mathbb{R}$ be a positive Borel function. Applying the chain rule, as in (2.4), and the one dimensional co-area formula (see for instance [2]), we obtain a co-area formula for the functions $u \in H^{1}(C)$ :

$$
\begin{align*}
\int_{C} v(x)\left|u^{\prime}\right|(x) d \mathcal{H}^{1}(x) & =\sum_{i} \int_{0}^{l_{i}}\left|\frac{d}{d t} u\left(\gamma_{i}(t)\right)\right| v\left(\gamma_{i}(t)\right) d t \\
& =\sum_{i} \int_{0}^{+\infty}\left(\sum_{u \circ \gamma_{i}(t)=\tau} v \circ \gamma_{i}(t)\right) d \tau  \tag{2.12}\\
& =\int_{0}^{+\infty}\left(\sum_{u(x)=\tau} v(x)\right) d \tau .
\end{align*}
$$

### 2.2.1 Optimization problem for the Dirichlet Energy on the class of connected sets

We study the following shape optimization problem:

$$
\begin{equation*}
\min \{\mathcal{E}(C ; \mathcal{D}): C \in \mathcal{A}(\mathcal{D} ; l)\} \tag{2.13}
\end{equation*}
$$

where $\mathcal{D}=\left\{D_{1}, \ldots, D_{k}\right\}$ is a given set of points in $\mathbb{R}^{d}$ and $l$ is a prescribed length.

Remark 2.6. When $k=1$ problem (2.13) reads as

$$
\begin{equation*}
\mathcal{E}=\min \left\{\mathcal{E}(C ; D): \mathcal{H}^{1}(C)=l, D \in C\right\}, \tag{2.14}
\end{equation*}
$$

where $D \in \mathbb{R}^{d}$ and $l>0$. In this case the solution is a line of length $l$ starting from $D$ (see Figure 2.1). A proof of this fact, in a slightly different context, can be found in [59] and we report it here for the sake of completeness.


Figure 2.1: The optimal graph with only one Dirichlet point.
Let $C \in \mathcal{A}(D ; l)$ be a generic connected set and let $w \in H_{0}^{1}(C ; D)$ be its Energy function, i.e. the minimizer of $J$ on $C$. Let $v:[0, l] \rightarrow \mathbb{R}$ be such that $\mu_{w}(\tau)=\mu_{v}(\tau)$, where $\mu_{w}$ and $\mu_{v}$ are the distribution function of $w$ and $v$ respectively, defined by

$$
\mu_{w}(\tau)=\mathcal{H}^{1}(w \leq \tau)=\sum_{i} \mathcal{H}^{1}\left(w_{i} \leq \tau\right), \quad \mu_{v}(\tau)=\mathcal{H}^{1}(v \leq \tau) .
$$

It is easy to see that, by the Cavalieri Formula, $\|v\|_{L^{p}([0, l])}=\|w\|_{L^{p}(C)}$, for each $p \geq 1$. By the co-area formula (2.12)

$$
\begin{equation*}
\int_{C}\left|w^{\prime}\right|^{2} d \mathcal{H}^{1}=\int_{0}^{+\infty}\left(\sum_{w=\tau}\left|w^{\prime}\right|\right) d \tau \geq \int_{0}^{+\infty}\left(\sum_{w=\tau} \frac{1}{\left|w^{\prime}\right|}\right)^{-1} d \tau=\int_{0}^{+\infty} \frac{d \tau}{\mu_{w}^{\prime}(\tau)} \tag{2.15}
\end{equation*}
$$

where we used the Cauchy-Schwartz inequality and the identity

$$
\mu_{w}(t)=\mathcal{H}^{1}(\{w \leq t\})=\int_{w \leq t} \frac{\left|w^{\prime}\right|}{\left|w^{\prime}\right|} d s=\int_{0}^{t}\left(\sum_{w=s} \frac{1}{\left|w^{\prime}\right|}\right) d s
$$

which implies that $\mu_{w}^{\prime}(t)=\sum_{w=t} \frac{1}{\left|w^{\prime}\right|}$. The same argument applied to $v$ gives:

$$
\begin{equation*}
\int_{0}^{l}\left|v^{\prime}\right|^{2} d x=\int_{0}^{+\infty}\left(\sum_{v=\tau}\left|v^{\prime}\right|\right) d \tau=\int_{0}^{+\infty} \frac{d \tau}{\mu_{v}^{\prime}(\tau)} \tag{2.16}
\end{equation*}
$$

Since $\mu_{w}=\mu_{v}$, the conclusion follows.
The following theorem shows that it is enough to study the problem (2.13) on the class of finite graphs embedded in $\mathbb{R}^{d}$. Consider the subset $\mathcal{A}_{N}(\mathcal{D} ; l) \subset \mathcal{A}(\mathcal{D} ; l)$ of those sets $C$ for which there exists a finite family $\gamma_{i}:\left[0, l_{i}\right] \rightarrow \mathbb{R}, i=1, \ldots, n$ with $n \leq N$, of injective rectifiable curves such that $\cup_{i} \gamma_{i}\left(\left[0, l_{i}\right]\right)=C$ and $\gamma_{i}\left(\left(0, l_{i}\right)\right) \cap \gamma_{j}\left(\left(0, l_{j}\right)\right)=\emptyset$, for each $i \neq j$.

Theorem 2.7. Consider the set of distinct points $\mathcal{D}=\left\{D_{1}, \ldots, D_{k}\right\} \subset \mathbb{R}^{d}$ and $l>0$. We have that

$$
\begin{equation*}
\inf \{\mathcal{E}(C ; \mathcal{D}): C \in \mathcal{A}(\mathcal{D} ; l)\}=\inf \left\{\mathcal{E}(C ; \mathcal{D}): C \in \mathcal{A}_{N}(\mathcal{D} ; l)\right\} \tag{2.17}
\end{equation*}
$$

where $N=2 k-1$. Moreover, if $C$ is a solution of the problem (2.13), then there is also a solution $\widetilde{C}$ of the same problem such that $\widetilde{C} \in \mathcal{A}_{N}(\mathcal{D} ; l)$.
Proof. Consider a connected set $C \in \mathcal{A}(\mathcal{D} ; l)$. We show that there is a set $\widetilde{C} \in \mathcal{A}_{N}(\mathcal{D} ; l)$ such that $\mathcal{E}(\widetilde{C} ; \mathcal{D}) \leq \mathcal{E}(C ; \mathcal{D})$. Let $\eta_{1}:\left[0, a_{1}\right] \rightarrow C$ be a geodesic in $C$ connecting $D_{1}$ to $D_{2}$ and let $\eta_{2}:[0, a] \rightarrow C$ be a geodesic connecting $D_{3}$ to $D_{1}$. Let $a_{2}$ be the smallest real number such that $\eta_{2}\left(a_{2}\right) \in \eta_{1}\left(\left[0, a_{1}\right]\right)$. Then, consider the geodesic $\eta_{3}$ connecting $D_{4}$ to $D_{1}$ and the smallest real number $a_{3}$ such that $\eta_{3}\left(a_{3}\right) \in \eta_{1}\left(\left[0, a_{1}\right]\right) \cup \eta_{2}\left(\left[0, a_{2}\right]\right)$. Repeating this operation, we obtain a family of geodesics $\eta_{i}, i=1, \ldots, k-1$ which intersect each other in a finite number of points. Each of these geodesics can be decomposed in several parts according to the intersection points with the other geodesics (see Figure 2.2).


Figure 2.2: Construction of the set $C^{\prime}$.
So, we can consider a new family of geodesics (still denoted by $\eta_{i}$ ), $\eta_{i}:\left[0, l_{i}\right] \rightarrow C$, $i=1, \ldots, n$, which does not intersect each other in internal points. Note that, by an induction argument on $k \geq 2$, we have $n \leq 2 k-3$. Let $C^{\prime}=\cup_{i} \eta_{i}\left(\left[0, l_{i}\right]\right) \subset C$. By the Second Rectifiability Theorem (see [4, Theorem 4.4.8]), we have that

$$
C=C^{\prime} \cup E \cup \Gamma
$$

where $\mathcal{H}^{1}(E)=0$ and $\Gamma=\left(\bigcup_{j=1}^{+\infty} \gamma_{j}\right)$, where $\gamma_{j}:\left[0, l_{j}\right] \rightarrow C$ for $j \geq 1$ is a family of Lipschitz curves in $C$. Moreover, we can suppose that $\mathcal{H}^{1}\left(\Gamma \cap C^{\prime}\right)=0$. In fact, if $\mathcal{H}^{1}\left(\operatorname{Im}\left(\gamma_{j}\right) \cap C^{\prime}\right) \neq 0$ for some $j \in \mathbb{N}$, we consider the restriction of $\gamma_{j}$ to (the closure of) each connected component of $\gamma_{j}^{-1}\left(\mathbb{R}^{d} \backslash C^{\prime}\right)$.

Let $w \in H_{0}^{1}(C ; \mathcal{D})$ be the Energy function on $C$ and let $v:\left[0, \mathcal{H}^{1}(\Gamma)\right] \rightarrow \mathbb{R}$ be a monotone increasing function such that $|\{v \leq \tau\}|=\mathcal{H}^{1}(\{w \leq \tau\} \cap \Gamma)$. Reasoning as in Remark 2.6, we have that

$$
\begin{equation*}
\frac{1}{2} \int_{0}^{\mathcal{H}^{1}(\Gamma)}\left|v^{\prime}\right|^{2} d x-\int_{0}^{\mathcal{H}^{1}(\Gamma)} v d x \leq \frac{1}{2} \int_{\Gamma}\left|w^{\prime}\right|^{2} d \mathcal{H}^{1}-\int_{\Gamma} w d \mathcal{H}^{1} \tag{2.18}
\end{equation*}
$$

Let $\sigma:\left[0, \mathcal{H}^{1}(\Gamma)\right] \rightarrow \mathbb{R}^{d}$ be an injective arc-length parametrized curve such that $\operatorname{Im}(\sigma) \cap C^{\prime}=\sigma(0)=x^{\prime}$, where $x^{\prime} \in C^{\prime}$ is the point where $w_{\mid C^{\prime}}$ achieves its maximum. Let $\widetilde{C}=C^{\prime} \cup \operatorname{Im}(\sigma)$. Notice that $\widetilde{C}$ connects the points $D_{1}, \ldots, D_{k}$ and has length $\mathcal{H}^{1}(\widetilde{C})=\mathcal{H}^{1}\left(C^{\prime}\right)+\mathcal{H}^{1}(\operatorname{Im}(\sigma))=\mathcal{H}^{1}\left(C^{\prime}\right)+\mathcal{H}^{1}(\Gamma)=l$. Moreover, we have

$$
\begin{equation*}
\mathcal{E}(\widetilde{C} ; \mathcal{D}) \leq J(\widetilde{w}) \leq J(w)=\mathcal{E}(C ; \mathcal{D}) \tag{2.19}
\end{equation*}
$$

where $\widetilde{w}$ is defined by

$$
\widetilde{w}(x)= \begin{cases}w(x), & \text { if } x \in C^{\prime}  \tag{2.20}\\ v(t)+w\left(x^{\prime}\right)-v(0), & \text { if } x=\sigma(t)\end{cases}
$$

We have then (2.19), i.e. the energy decreases. We conclude by noticing that the point $x^{\prime}$ where we attach $\sigma$ to $C^{\prime}$ may be an internal point for $\eta_{i}$, i.e. a point such that $\eta_{i}^{-1}\left(x^{\prime}\right) \in\left(0, l_{i}\right)$. Thus, the set $\widetilde{C}$ is composed of at most $2 k-1$ injective arc-length parametrized curves which does not intersect in internal points, i.e. $\widetilde{C} \in \mathcal{A}_{2 k-1}(\mathcal{D} ; l)$.

Remark 2.8. Theorem 2.7 above provides a nice class of admissible sets, where to search a minimizer of the energy functional $\mathcal{E}$. Indeed, according to its proof, we may limit ourselves to consider only graphs $C$ such that:

1. $C$ is a tree, i.e. it does not contain any closed loop;
2. the Dirichlet points $D_{i}$ are vertices of degree one (endpoints) for $C$;
3. there are at most $k-1$ other vertices; if a vertex has degree three or more, we call it Kirchhoff point;
4. there is at most one vertex of degree one for $C$ which is not a Dirichlet point. In this vertex the energy function $w$ satisfies Neumann boundary condition $w^{\prime}=0$ and so we call it Neumann point.

The previous properties are also necessary conditions for the optimality of the graph $C$ (see Proposition 2.19 for more details).

As we show in Example 2.23, the problem (2.13) may not have a solution in the class of connected sets. It is worth noticing that the lack of existence only occurs for particular configurations of the Dirichlet points $D_{i}$ and not because of some degeneracy of the cost functional $\mathcal{E}$. In fact, we are able to produce other examples in which an optimal graph exists (see Section 2.4).

### 2.3 Sobolev space and Dirichlet Energy of a metric graph

Let $V=\left\{V_{1}, \ldots, V_{N}\right\}$ be a finite set and let $E \subset\left\{e_{i j}=\left\{V_{i}, V_{j}\right\}\right\}$ be a set of pairs of elements of $V$. We define combinatorial graph (or just graph) a pair $\Gamma=(V, E)$. We say the set $V=V(\Gamma)$ is the set of vertices of $\Gamma$ and the set $E=E(\Gamma)$ is the set of edges. We denote with $|E|$ and $|V|$ the cardinalities of $E$ and $V$ and with $\operatorname{deg}\left(V_{i}\right)$ the degree of the vertex $V_{i}$, i.e. the number of edges incident to $V_{i}$.

A path in the graph $\Gamma$ is a sequence $V_{\alpha_{0}}, \ldots, V_{\alpha_{n}} \in V$ such that for each $k=$ $0, \ldots, n-1$, we have that $\left\{V_{\alpha_{k}}, V_{\alpha_{k+1}}\right\} \in E$. With this notation, we say that the path connects $V_{i_{0}}$ to $V_{i_{\alpha}}$. The path is said to be simple if there are no repeated vertices in $V_{\alpha_{0}}, \ldots, V_{\alpha_{n}}$. We say that the graph $\Gamma=(V, E)$ is connected, if for each pair of vertices $V_{i}, V_{j} \in V$ there is a path connecting them. We say that the connected graph $\Gamma$ is a tree, if after removing any edge, the graph becomes not connected.

If we associate a non-negative length (or weight) to each edge, i.e. a map $l: E(\Gamma) \rightarrow$ $[0,+\infty)$, then we say that the couple $(\Gamma, l)$ determines a metric graph of length

$$
l(\Gamma):=\sum_{i<j} l\left(e_{i j}\right) .
$$

A function $u: \Gamma \rightarrow \mathbb{R}^{d}$ on the metric graph $\Gamma$ is a collection of functions $u_{i j}$ : $\left[0, l_{i j}\right] \rightarrow \mathbb{R}$, for $1 \leq i \neq j \leq d$, such that:

1. $u_{j i}(x)=u_{i j}\left(l_{i j}-x\right)$, for each $1 \leq i \neq j \leq d$,
2. $u_{i j}(0)=u_{i k}(0)$, for all $\{i, j, k\} \subset\{1, \ldots, d\}$,
where we used the notation $l_{i j}=l\left(e_{i j}\right)$. A function $u: \Gamma \rightarrow \mathbb{R}$ is said continuous $(u \in C(\Gamma))$, if $u_{i j} \in C\left(\left[0, l_{i j}\right]\right)$, for all $i, j \in\{1, \ldots, n\}$. We call $L^{p}(\Gamma)$ the space of $p$-summable functions $(p \in[1,+\infty))$, i.e. the functions $u=\left(u_{i j}\right)_{i j}$ such that

$$
\|u\|_{L^{p}(\Gamma)}^{p}:=\frac{1}{2} \sum_{i, j}\left\|u_{i j}\right\|_{L^{p}\left(0, l_{i j}\right)}^{p}<+\infty,
$$

where $\|\cdot\|_{L^{p}(a, b)}$ denotes the usual $L^{p}$ norm on the interval $[a, b]$. As usual, the space $L^{2}(\Gamma)$ has a Hilbert structure endowed by the scalar product:

$$
\langle u, v\rangle_{L^{2}(\Gamma)}:=\frac{1}{2} \sum_{i, j}\left\langle u_{i j}, v_{i j}\right\rangle_{L^{2}\left(0, l_{i j}\right)} .
$$

function defined on $\mathbb{R}^{d}$ to the set $\cup_{i, j} \gamma_{i j}\left(\left[0, l_{i j}\right]\right)$ : this is mainly due to the fact that we allow intersections.

We define the Sobolev space $H^{1}(\Gamma)$ as:

$$
\begin{equation*}
H^{1}(\Gamma)=\left\{u \in C(\Gamma): u_{i j} \in H^{1}\left(\left[0, l_{i j}\right]\right), \forall i, j \in\{1, \ldots, n\}\right\}, \tag{2.21}
\end{equation*}
$$

which is a Hilbert space with the norm

$$
\begin{equation*}
\|u\|_{H^{1}(\Gamma)}^{2}=\frac{1}{2} \sum_{i, j}\left\|u_{i j}\right\|_{H^{1}\left(\left[0, l_{i j}\right]\right)}^{2}=\frac{1}{2} \sum_{i, j}\left(\int_{0}^{l_{i j}}\left|u_{i j}\right|^{2} d x+\int_{0}^{l_{i j}}\left|u_{i j}^{\prime}\right|^{2} d x\right) \tag{2.22}
\end{equation*}
$$

Remark 2.9. Note that for $u \in H^{1}(\Gamma)$ the family of derivatives $\left(u_{i j}^{\prime}\right)_{1 \leq i \neq j \leq d}$ is not a function on $\Gamma$, since $u_{i j}^{\prime}(x)=\frac{\partial}{\partial x} u_{j i}\left(l_{i j}-x\right)=-u_{j i}^{\prime}\left(l_{i j}-x\right)$. Thus, we work with the function $\left|u^{\prime}\right|=\left(\left|u_{i j}^{\prime}\right|\right)_{1 \leq i \neq j \leq d} \in L^{2}(\Gamma)$.
Remark 2.10. The inclusions $H^{1}(\Gamma) \subset C(\Gamma)$ and $H^{1}(\Gamma) \subset L^{2}(\Gamma)$ are compact, since the corresponding inclusions, for each of the intervals $\left[0, l_{i j}\right]$, are compact. By the same argument, the $H^{1}$ norm is lower semicontinuous with respect to the strong $L^{2}$ convergence of the functions in $H^{1}(\Gamma)$.

For any subset $W=\left\{W_{1}, \ldots, W_{k}\right\}$ of the set of vertices $V(\Gamma)=\left\{V_{1}, \ldots, V_{N}\right\}$, we introduce the Sobolev space with Dirichlet boundary conditions on $W$ :

$$
\begin{equation*}
H_{0}^{1}(\Gamma ; W)=\left\{u \in H^{1}(\Gamma): u\left(W_{1}\right)=\cdots=u\left(W_{k}\right)=0\right\} \tag{2.23}
\end{equation*}
$$

Remark 2.11. Arguing as in Remark 2.1 we have that for each $u \in H_{0}^{1}(\Gamma ; W)$ and, more generally, for each $u \in H^{1}(\Gamma)$ such that $u\left(V_{\alpha}\right)=0$ for some $\alpha=1, \ldots, d$, the Poincaré inequality

$$
\begin{equation*}
\|u\|_{L^{2}(\Gamma)} \leq l^{1 / 2}\|u\|_{L^{\infty}} \leq l\left\|u^{\prime}\right\|_{L^{2}(\Gamma)} \tag{2.24}
\end{equation*}
$$

holds, where

$$
\left\|u^{\prime}\right\|_{L^{2}(\Gamma)}^{2}:=\int_{\Gamma}\left|u^{\prime}\right|^{2} d x:=\sum_{i, j} \int_{0}^{l_{i j}}\left|u_{i j}^{\prime}\right|^{2} d x
$$

On the metric graph $\Gamma$, we consider the Dirichlet Energy with respect to $W$ :

$$
\begin{equation*}
\mathcal{E}(\Gamma ; W)=\inf \left\{J(u): u \in H_{0}^{1}(\Gamma ; W)\right\} \tag{2.25}
\end{equation*}
$$

where the functional $J: H_{0}^{1}(\Gamma ; W) \rightarrow \mathbb{R}$ is defined by

$$
\begin{equation*}
J(u)=\frac{1}{2} \int_{\Gamma}\left|u^{\prime}\right|^{2} d x-\int_{\Gamma} u d x \tag{2.26}
\end{equation*}
$$

Lemma 2.12. Given a metric graph $\Gamma$ of length $l$ and Dirichlet points $\left\{W_{1}, \ldots, W_{k}\right\} \subset$ $V(\Gamma)=\left\{V_{1}, \ldots, V_{N}\right\}$, there is a unique function $w=\left(w_{i j}\right)_{1 \leq i \neq j \leq d} \in H_{0}^{1}(\Gamma ; W)$ which minimizes the functional $J$. Moreover, we have

1. for each $1 \leq i \neq j \leq d$ and each $t \in\left(0, l_{i j}\right),-w_{i j}^{\prime \prime}=1$;
2. at every vertex $V_{i} \in V(\Gamma)$, which is not a Dirichlet point, $w$ satisfies the Kirchhoff's law:

$$
\sum_{j} w_{i j}^{\prime}(0)=0
$$

where the sum is over all $j$ for which the edge $e_{i j}$ exists;

Furthermore, the conditions (i) and (ii) uniquely determine $w$.
Proof. The existence is a consequence of Remark 2.10 and the uniqueness is due to the strict convexity of the $L^{2}$ norm. for each $n \in \mathbb{N} \gamma_{i j}$ to a function $w_{i j}$. Defining $w=\left(w_{i j}\right)_{i j} \in H_{0}^{1}\left(\Gamma ; D_{1}, \ldots, D_{k}\right)$, by the lower semicontinuity of the $L^{2}$ norm of the gradient on each edge, we obtain the existence. For any $\varphi \in H_{0}^{1}(\Gamma ; W)$, we have that 0 is a critical point for the function

$$
\varepsilon \mapsto \frac{1}{2} \int_{\Gamma}\left|(w+\varepsilon \varphi)^{\prime}\right|^{2} d x-\int_{\Gamma}(w+\varepsilon \varphi) d x .
$$

Since $\varphi$ is arbitrary, we obtain the first claim. The Kirchhoff's law at the vertex $V_{i}$ follows by choosing $\varphi$ supported in a "small neighborhood" of $V_{i}$. The last claim is due to the fact that if $u \in H_{0}^{1}(\Gamma ; W)$ satisfies $(i)$ and $(i i)$, then it is an extremal for the convex functional $J$ and so, $u=w$.

Remark 2.13. As in Remark 2.5 we have that the co-area formula holds for the functions $u \in H^{1}(\Gamma)$ and any positive Borel (on each edge) function $v: \Gamma \rightarrow \mathbb{R}$ :

$$
\begin{align*}
\int_{\Gamma} v(x)\left|u^{\prime}\right|(x) d x & =\sum_{1 \leq i<j \leq d} \int_{0}^{l_{i j}}\left|u_{i j}^{\prime}(x)\right| v(x) d x \\
& =\sum_{1 \leq i<j \leq d} \int_{0}^{+\infty}\left(\sum_{u_{i j}(x)=\tau} v(x)\right) d \tau  \tag{2.27}\\
& =\int_{0}^{+\infty}\left(\sum_{u(x)=\tau} v(x)\right) d \tau .
\end{align*}
$$

### 2.3.1 Optimization problem for the Dirichlet Energy on the class of metric graphs

We say that the continuous function $\gamma=\left(\gamma_{i j}\right)_{1 \leq i \neq j \leq d}: \Gamma \rightarrow \mathbb{R}^{d}$ is an immersion of the metric graph $\Gamma$ into $\mathbb{R}^{d}$, if for each $1 \leq i \neq j \leq d$ the function $\gamma_{i j}:\left[0, l_{i j}\right] \rightarrow \mathbb{R}^{d}$ is an injective arc-length parametrized curve. We say that $\gamma: \Gamma \rightarrow \mathbb{R}^{d}$ is an embedding, if it is an immersion which is also injective, i.e. for any $i \neq j$ and $i^{\prime} \neq j^{\prime}$, we have

1. $\gamma_{i j}\left(\left(0, l_{i j}\right)\right) \cap \gamma_{i^{\prime} j^{\prime}}\left(\left[0, l_{i^{\prime} j^{\prime}}\right]\right)=\emptyset$,
2. $\gamma_{i j}(0)=\gamma_{i^{\prime} j^{\prime}}(0)$, if and only if, $i=i^{\prime}$.

Remark 2.14. Suppose that $\Gamma$ is a metric graph of length $l$ and that $\gamma: \Gamma \rightarrow \mathbb{R}^{d}$ is an embedding. Then the set $C:=\gamma(\Gamma)$ is rectifiable of length $\mathcal{H}^{1}(\gamma(\Gamma))=l$ and the spaces $H^{1}(\Gamma)$ and $H^{1}(C)$ are isometric as Hilbert spaces, where the isomorphism is given by the composition with the function $\gamma$.

Indeed, the topology of the embedded is uniquely determined by the topology of its representation in $\mathbb{R}^{d}$. This is not the case of the larger class of immersed graphs, where the intersections are allowed and new vertices may appear.

Consider a finite set of distinct points $\mathcal{D}=\left\{D_{1}, \ldots, D_{k}\right\} \subset \mathbb{R}^{d}$ and let $l \geq \operatorname{St}(\mathcal{D})$, where $S t(\mathcal{D})$ is the length of the Steiner set, the minimal among the ones connecting all the points $D_{i}$ (see [4] for more details on the Steiner problem). Consider the shape optimization problem:

$$
\begin{equation*}
\min \left\{\mathcal{E}(\Gamma ; \mathcal{V}): \Gamma \in C M G, l(\Gamma)=l, \mathcal{V} \subset V(\Gamma), \exists \gamma: \Gamma \rightarrow \mathbb{R}^{d} \text { immersion, } \gamma(\mathcal{V})=\mathcal{D}\right\} \tag{2.28}
\end{equation*}
$$

where $C M G$ indicates the class of connected metric graphs. Note that since $l \geq S t(\mathcal{D})$, there is a metric graph and an embedding $\gamma: \Gamma \rightarrow \mathbb{R}^{d}$ such that $\mathcal{D} \subset \gamma(V(\Gamma))$ and so the admissible set in the problem (2.28) is non-empty, as well as the admissible set in the problem
$\min \left\{\mathcal{E}(\Gamma ; \mathcal{V}): \Gamma \in C M G, l(\Gamma)=l, \mathcal{V} \subset V(\Gamma), \exists \gamma: \Gamma \rightarrow \mathbb{R}^{d}\right.$ embedding, $\left.\gamma(\mathcal{V})=\mathcal{D}\right\}$.
We will see in Theorem 2.18 that problem (2.28) admits a solution, while Example 2.23 shows that in general an optimal embedded graph for problem (2.29) may not exist.

Remark 2.15. By Remark 2.14 and by the fact that the functionals we consider are invariant with respect to the isometries of the Sobolev space, we have that the problems (2.13) and (2.29) are equivalent, i.e. if $\Gamma \in C M G$ and $\gamma: \Gamma \rightarrow \mathbb{R}^{d}$ is an embedding such that the pair $(\Gamma, \gamma)$ is a solution of (2.29), then the set $\gamma(\Gamma)$ is a solution of the problem (2.13). On the other hand, if $C$ is a solution of the problem (2.13), by Theorem 2.7, we can suppose that $C=\bigcup_{i=1}^{d} \gamma_{i}\left(\left[0, l_{i}\right]\right)$, where $\gamma_{i}$ are injective arc-length parametrized curves, which does not intersect internally. Thus, we can construct a metric graph $\Gamma$ with vertices the set of points $\left\{\gamma_{i}(0), \gamma_{i}\left(l_{i}\right)\right\}_{i=1}^{d} \subset \mathbb{R}^{d}$, and $d$ edges of lengths $l_{i}$ such that two vertices are connected by an edge, if and only if they are the endpoints of the same curve $\gamma_{i}$. The function $\gamma=\left(\gamma_{i}\right)_{i=1, \ldots, d}: \Gamma \rightarrow \mathbb{R}^{d}$ is an embedding by construction and by Remark 2.14, we have $\mathcal{E}(C ; \mathcal{D})=\mathcal{E}(\Gamma ; \mathcal{D})$.

Theorem 2.16. Let $\mathcal{D}=\left\{D_{1}, \ldots, D_{k}\right\} \subset \mathbb{R}^{d}$ be a finite set of points and let $l \geq S t(\mathcal{D})$ be a positive real number. Suppose that $\Gamma$ is a connected metric graph of length $l$, $\mathcal{V} \subset V(\Gamma)$ is a set of vertices of $\Gamma$ and $\gamma: \Gamma \rightarrow \mathbb{R}^{d}$ is an immersion (embedding) such that $\mathcal{D}=\gamma(\mathcal{V})$. Then there exist a connected metric graph $\widetilde{\Gamma}$ of at most $2 k$ vertices and $2 k-1$ edges, a set $\widetilde{\mathcal{V}} \subset V(\widetilde{\Gamma})$ of vertices of $\widetilde{\Gamma}$ and an immersion (embedding) $\widetilde{\gamma}: \widetilde{\Gamma} \rightarrow \mathbb{R}^{d}$ such that $\mathcal{D}=\widetilde{\gamma}(\widetilde{\mathcal{V}})$ and

$$
\begin{equation*}
\mathcal{E}(\widetilde{\Gamma} ; \widetilde{\mathcal{V}}) \leq \mathcal{E}(\Gamma ; \mathcal{V}) \tag{2.30}
\end{equation*}
$$

Proof. We repeat the argument from Theorem 2.7. We first construct a connected metric graph $\Gamma^{\prime}$ such that $V\left(\Gamma^{\prime}\right) \subset V(\Gamma)$ and the edges of $\Gamma^{\prime}$ are appropriately chosen paths in $\Gamma$. The edges of $\Gamma$, which are not part of any of these paths, are symmetrized
in a single edge, which we attach to $\Gamma^{\prime}$ in a point, where the restriction of $w$ to $\Gamma^{\prime}$ achieves its maximum, where $w$ is the Energy function for $\Gamma$.

Suppose that $V_{1}, \ldots, V_{k} \in \mathcal{V} \subset V(\Gamma)$ are such that $\gamma\left(V_{i}\right)=D_{i}, i=1, \ldots, k$. We start constructing $\Gamma^{\prime}$ by taking $\widetilde{\mathcal{V}}:=\left\{V_{1}, \ldots, V_{k}\right\} \subset V\left(\Gamma^{\prime}\right)$. Let $\sigma_{1}=\left\{V_{i_{0}}, V_{i_{1}}, \ldots, V_{i_{s}}\right\}$ be a path of different vertices (i.e. simple path) connecting $V_{1}=V_{i_{s}}$ to $V_{2}=V_{i_{0}}$ and let $\tilde{\sigma}_{2}=\left\{V_{j_{0}}, V_{j_{1}}, \ldots, V_{j_{t}}\right\}$ be a simple path connecting $V_{1}=V_{j_{t}}$ to $V_{3}=V_{j_{0}}$. Let $t^{\prime} \in\{1, \ldots, t\}$ be the smallest integer such that $V_{j_{t^{\prime}}} \in \sigma_{1}$. Then we set $V_{j_{t^{\prime}}} \in V\left(\Gamma^{\prime}\right)$ and $\sigma_{2}=\left\{V_{j_{0}}, V_{j_{1}}, \ldots, V_{j_{t^{\prime}}}\right\}$. Consider a simple path $\tilde{\sigma}_{3}=\left\{V_{m_{0}}, V_{m_{1}}, \ldots, V_{m_{r}}\right\}$ connecting $V_{1}=V_{m_{r}}$ to $V_{3}=V_{m_{0}}$ and the smallest integer $r^{\prime}$ such that $V_{m_{r^{\prime}}} \in \sigma_{1} \cup \sigma_{2}$. We set $V_{m_{r^{\prime}}} \in V\left(\Gamma^{\prime}\right)$ and $\sigma_{3}=\left\{V_{m_{0}}, V_{m_{1}}, \ldots, V_{m_{r^{\prime}}}\right\}$. We continue the operation until each of the points $V_{1}, \ldots, V_{k}$ is in some path $\sigma_{j}$. Thus we obtain the set of vertices $V\left(\Gamma^{\prime}\right)$. We define the edges of $\Gamma^{\prime}$ by saying that $\left\{V_{i}, V_{i^{\prime}}\right\} \in E\left(\Gamma^{\prime}\right)$ if there is a simple path $\sigma$ connecting $V_{i}$ to $V_{i^{\prime}}$ and which is contained in some path $\sigma_{j}$ from the construction above; the length of the edge $\left\{V_{i}, V_{i^{\prime}}\right\}$ is the sum of the lengths of the edges of $\Gamma$ which are part of $\sigma$. We notice that $\Gamma^{\prime} \in C M G$ is a tree with at most $2 k-2$ vertices and $2 k-2$ edges. Moreover, even if $\Gamma^{\prime}$ is not a subgraph of $\Gamma\left(E\left(\Gamma^{\prime}\right)\right.$ may not be a subset of $E(\Gamma)$ ), we have the inclusion $H^{1}\left(\Gamma^{\prime}\right) \subset H^{1}(\Gamma)$.

Consider the set $E^{\prime \prime} \subset E(\Gamma)$ composed of the edges of $\Gamma$ which are not part of any of the paths $\sigma_{j}$ from the construction above. We denote with $l^{\prime \prime}$ the sum of the lengths of the edges in $E^{\prime \prime}$. For any $e_{i j} \in E^{\prime \prime}$ we consider the restriction $w_{i j}:\left[0, l_{i j}\right] \rightarrow \mathbb{R}$ of the Energy function $w$ on $e_{i j}$. Let $v:\left[0, l^{\prime \prime}\right] \rightarrow \mathbb{R}$ be the monotone function defined by the equality $|\{v \geq \tau\}|=\sum_{e_{i j} \in E^{\prime \prime}}\left|\left\{w_{i j} \geq \tau\right\}\right|$. Using the co-area formula (2.27) and repeating the argument from Remark 2.14, we have that

$$
\begin{equation*}
\frac{1}{2} \int_{0}^{l^{\prime \prime}}\left|v^{\prime}\right|^{2} d x-\int_{0}^{l^{\prime \prime}} v(x) d x \leq \sum_{e_{i j} \in E^{\prime \prime}}\left(\frac{1}{2} \int_{0}^{l_{i j}}\left|w_{i j}^{\prime}\right|^{2} d x-\int_{0}^{l_{i j}} w_{i j} d x\right) \tag{2.31}
\end{equation*}
$$

Let $\widetilde{\Gamma}$ be the graph obtained from $\Gamma$ by creating a new vertex $W_{1}$ in the point, where the restriction $w_{\mid \Gamma^{\prime}}$ achieves its maximum, and another vertex $W_{2}$, connected to $W_{1}$ by an edge of length $l^{\prime \prime}$. It is straightforward to check that $\widetilde{\Gamma}$ is a connected metric tree of length $l$ and that there exists an immersion $\widetilde{\gamma}: \widetilde{\Gamma} \rightarrow \mathbb{R}^{d}$ such that $\mathcal{D}=\widetilde{\gamma}(\widetilde{\mathcal{V}})$. The inequality (2.30) follows since, by $(2.31), J(\widetilde{w}) \leq J(w)$, where $\widetilde{w}$ is defined as $w$ on the edges $E\left(\Gamma^{\prime}\right) \subset E(\widetilde{\Gamma})$ and as $v$ on the edge $\left\{W_{1}, W_{2}\right\}$.

Before we prove our main existence result, we need a preliminary Lemma.
Lemma 2.17. Let $\Gamma$ be a connected metric tree and let $\mathcal{V} \subset V(\Gamma)$ be a set of Dirichlet vertices. Let $w \in H_{0}^{1}(\Gamma ; \mathcal{V})$ be the Energy function on $\Gamma$ with Dirichlet conditions in $\mathcal{V}$, i.e. the function that realizes the minimum in the definition of $\mathcal{E}(\Gamma ; \mathcal{V})$. Then, we have the bound $\left\|w^{\prime}\right\|_{L^{\infty}} \leq l(\Gamma)$.

Proof. Up to adding vertices in the points where $\left|w^{\prime}\right|=0$, we can suppose that on each edge $e_{i j}:=\left\{V_{i}, V_{j}\right\} \in E(\Gamma)$ the function $w_{i j}:\left[0, l_{i j}\right] \rightarrow \mathbb{R}^{+}$is monotone. Moreover, up
to relabel the vertices of $\Gamma$ we can suppose that if $e_{i j} \in V(\Gamma)$ and $i<j$, then $w\left(V_{i}\right) \leq$ $w\left(V_{j}\right)$. Fix $V_{i}, V_{i^{\prime}} \in V(\Gamma)$ such that $e_{i i^{\prime}} \in E(\Gamma)$. We will prove that $\left|w_{i i^{\prime}}^{\prime}(0)\right| \leq l(\Gamma)$. It is enough to consider the case $i<i^{\prime}$, i.e. $w_{i i^{\prime}}^{\prime}(0)>0$. We construct the graph $\widetilde{\Gamma}$ inductively, as follows (see Figure 2.3):

1. $V_{i} \in V(\widetilde{\Gamma})$;
2. if $V_{j} \in V(\widetilde{\Gamma})$ and $V_{k} \in V(\Gamma)$ are such that $e_{j k} \in E(\Gamma)$ and $j<k$, then $V_{k} \in V(\widetilde{\Gamma})$ and $e_{j k} \in E(\widetilde{\Gamma})$.


Figure 2.3: The graph $\widetilde{\Gamma}$; the letter $\mathcal{N}$ labels the vertices in which $w^{\prime}=0$.
The graph $\widetilde{\Gamma}$ constructed by the above procedure and the restriction $\widetilde{w} \in H^{1}(\widetilde{\Gamma})$ of $w$ to $\widetilde{\Gamma}$ have the following properties:

1. On each edge $e_{j k} \in E(\widetilde{\Gamma})$, the function $\widetilde{w}_{j k} \geq 0$ is positive, monotone and $\widetilde{w}_{j k}^{\prime \prime}=$ $-1 ;$
2. $\widetilde{w}\left(V_{j}\right)>\widetilde{w}\left(V_{k}\right)$ whenever $e_{j k} \in E(\widetilde{\Gamma})$ and $j>k$;
3. if $V_{j} \in V(\widetilde{\Gamma})$ and $j>i$, then there is exactly one $k<j$ such that $e_{k j} \in E(\widetilde{\Gamma})$;
4. for $j$ and $k$ as in the previous point, we have that

$$
0 \leq \widetilde{w}_{k j}^{\prime}\left(l_{k j}\right) \leq \sum_{s} \widetilde{w}_{j s}^{\prime}(0),
$$

where the sum on the right-hand side is over all $s>j$ such that $e_{s j} \in E(\widetilde{\Gamma})$. If there are not such $s$, we have that $\widetilde{w}_{k j}^{\prime}\left(l_{k j}\right)=0$.

We prove that for any graph $\widetilde{\Gamma}$ and any function $\widetilde{w} \in H^{1}(\widetilde{\Gamma})$, for which the conditions $(a),(b),(c)$ and $(d)$ are satisfied, we have that

$$
\sum_{j} \widetilde{w}_{i j}^{\prime}(0) \leq l(\widetilde{\Gamma}),
$$

where the sum is over all $j \geq i$ and $e_{i j} \in E(\widetilde{\Gamma})$. It is enough to observe that each of the operations ( $i$ ) and (ii) described below, produces a graph which still satisfies (a), (b), (c) and (d). Let $V_{j} \in V(\widetilde{\Gamma})$ be such that for each $s>j$ for which $e_{j s} \in E(\widetilde{\Gamma})$, we have that $\widetilde{w}_{j s}^{\prime}\left(l_{j s}\right)=0$ and let $k<j$ be such that $e_{j k} \in E(\widetilde{\Gamma})$.

1. If there is only one $s>j$ with $e_{j s} \in E(\widetilde{\Gamma})$, then we erase the vertex $V_{j}$ and the edges $e_{k j}$ and $e_{j s}$ and add the edge $e_{k s}$ of length $l_{k s}:=l_{k j}+l_{j s}$. On the new edge we define $\widetilde{w}_{k s}:\left[0, l_{s k}\right] \rightarrow \mathbb{R}^{+}$as

$$
\widetilde{w}_{k s}(x)=-\frac{x^{2}}{2}+l_{k s} x+\widetilde{w}_{k j}(0),
$$

which still satisfies the conditions above since $\widetilde{w}_{k j}^{\prime}-l_{k j} \leq l_{j s}$, by (d), and $\widetilde{w}_{k s}^{\prime}=$ $l_{k s} \geq \widetilde{w}_{k j}^{\prime}(0)$.
2. If there are at least two $s>j$ such that $e_{j s} \in E(\widetilde{\Gamma})$, we erase all the vertices $V_{s}$ and edges $e_{j s}$, substituting them with a vertex $V_{S}$ connected to $V_{j}$ by an edge $e_{j S}$ of length

$$
l_{j S}:=\sum_{s} l_{j s}
$$

where the sum is over all $s>j$ with $e_{j s} \in E(\widetilde{\Gamma})$. On the new edge, we consider the function $\widetilde{w}_{j S}$ defined by

$$
\widetilde{w}_{j S}(x)=-\frac{x^{2}}{2}+l_{j S} x+\widetilde{w}\left(V_{j}\right)
$$

which still satisfies the conditions above since

$$
\sum_{s: s>j} \widetilde{w}_{j s}^{\prime}(0)=\sum_{s: s>j} l_{j s}=l_{j S}=\widetilde{w}_{j S}^{\prime}(0)
$$

We apply (i) and (ii) until we obtain a graph with vertices $V_{i}, V_{j}$ and only one edge $e_{i j}$ of length $l(\widetilde{\Gamma})$. The function we obtain on this graph is $-\frac{x^{2}}{2}+l(\widetilde{\Gamma}) x$ with derivative in 0 equal to $l(\widetilde{\Gamma})$. Since, after applying $(i)$ and (ii), the sum $\sum_{j>i} \widetilde{w}_{i j}^{\prime}(0)$ does not decrease, we get that the claim is true.
Theorem 2.18. Consider a set of distinct points $\mathcal{D}=\left\{D_{1}, \ldots, D_{k}\right\} \subset \mathbb{R}^{d}$ and a positive real number $l \geq S t(\mathcal{D})$. Then there exists a connected metric graph $\Gamma$, a set of vertices $\mathcal{V} \subset V(\Gamma)$ and an immersion $\gamma: \Gamma \rightarrow \mathbb{R}^{d}$ which are solution of the problem (2.28). Moreover, $\Gamma$ can be chosen to be a tree of at most $2 k$ vertices and $2 k-1$ edges.

Proof. Consider a minimizing sequence ( $\Gamma_{n}, \gamma_{n}$ ) of connected metric graphs $\Gamma_{n}$ and immersions $\gamma_{n}: \Gamma_{n} \rightarrow \mathbb{R}^{d}$. By Theorem 2.16, we can suppose that each $\Gamma_{n}$ is a tree with at most $2 k$ vertices and $2 k-1$ edges. Up to extracting a subsequence, we may
assume that the metric graphs $\Gamma_{n}$ are the same graph $\Gamma$ but with different lengths $l_{i j}^{n}$ of the edges $e_{i j}$. We can suppose that for each $e_{i j} \in E(\Gamma) l_{i j}^{n} \rightarrow l_{i j}$ for some $l_{i j} \geq 0$ as $n \rightarrow \infty$. We construct the graph $\widetilde{\Gamma}$ from $\Gamma$ identifying the vertices $V_{i}, V_{j} \in V(\Gamma)$ such that $l_{i j}=0$. The graph $\widetilde{\Gamma}$ is a connected metric tree of length $l$ and there is an immersion $\widetilde{\gamma}: \widetilde{\Gamma} \rightarrow \mathbb{R}^{d}$ such that $\mathcal{D} \subset \widetilde{\gamma}(\widetilde{\Gamma})$. In fact if $\left\{V_{1}, \ldots V_{N}\right\}$ are the vertices of $\Gamma$, up to extracting a subsequence, we can suppose that for each $i=1, \ldots, d \gamma_{n}\left(V_{i}\right) \rightarrow X_{i} \in \mathbb{R}^{d}$. We define $\widetilde{\gamma}\left(V_{i}\right):=X_{i}$ and $\gamma_{i j}:\left[0, l_{i j}\right] \rightarrow \mathbb{R}^{d}$ as any injective arc-length parametrized curve connecting $X_{i}$ and $X_{j}$, which exists, since

$$
l_{i j}=\lim l_{i j}^{n} \geq \lim \left|\gamma_{n}\left(V_{i}\right)-\gamma_{n}\left(V_{j}\right)\right|=\left|X_{i}-X_{j}\right| .
$$

To prove the theorem, it is enough to check that

$$
\mathcal{E}(\widetilde{\Gamma} ; \mathcal{V})=\lim _{n \rightarrow \infty} \mathcal{E}\left(\Gamma_{n} ; \mathcal{V}\right)
$$

Let $w^{n}=\left(w_{i j}^{n}\right)_{i j}$ be the Energy function on $\Gamma_{n}$. Up to a subsequence, we may suppose that for each $i=1, \ldots, N, w^{n}\left(V_{i}\right) \rightarrow a_{i} \in \mathbb{R}$ as $n \rightarrow \infty$. Moreover, by Lemma 2.17, we have that if $l_{i j}=0$, then $a_{i}=a_{j}$. On each of the edges $e_{i j} \in E(\widetilde{\Gamma})$, where $l_{i j}>0$, we define the function $w_{i j}:\left[0, l_{i j}\right] \rightarrow \mathbb{R}$ as the parabola such that $w_{i j}(0)=a_{i}$, $w_{i j}\left(l_{i j}\right)=a_{j}$ and $w_{i j}^{\prime \prime}=-1$ on $\left(0, l_{i j}\right)$. Then, we have

$$
\frac{1}{2} \int_{0}^{l_{i j}^{n}}\left|\left(w_{i j}^{n}\right)^{\prime}\right|^{2} d x-\int_{0}^{l_{i j}^{n}} w_{i j}^{n} d x \underset{n \rightarrow \infty}{\longrightarrow} \frac{1}{2} \int_{0}^{l_{i j}}\left|\left(w_{i j}\right)^{\prime}\right|^{2} d x-\int_{0}^{l_{i j}} w_{i j} d x,
$$

and so, it is enough to prove that $\widetilde{w}=\left(w_{i j}\right)_{i j}$ is the Energy function on $\widetilde{\Gamma}$, i.e. (by Lemma 2.12) that the Kirchoff's law holds in each vertex of $\widetilde{\Gamma}$. This follows since for each $1 \leq i \neq j \leq N$ we have

1. $\left(w_{i j}^{n}\right)^{\prime}(0) \rightarrow w_{i j}^{\prime}(0)$, as $n \rightarrow \infty$, if $l_{i j} \neq 0$;
2. $\left|\left(w_{i j}^{n}\right)^{\prime}(0)-\left(w_{i j}^{n}\right)^{\prime}\left(l_{i j}^{n}\right)\right| \leq l_{i j}^{n} \rightarrow 0$, as $n \rightarrow \infty$, if $l_{i j}=0$.

The proof is then concluded.
The proofs of Theorem 2.16 and Theorem 2.18 suggest that a solution $(\Gamma, \mathcal{V}, \gamma)$ of the problem (2.28) must satisfy some optimality conditions. We summarize these additional informations in the following Proposition.

Proposition 2.19. Consider a connected metric graph $\Gamma$, a set of vertices $\mathcal{V} \subset V(\Gamma)$ and an immersion $\gamma: \Gamma \rightarrow \mathbb{R}^{d}$ such that $(\Gamma, \mathcal{V}, \gamma)$ is a solution of the problem (2.28). Moreover, suppose that all the vertices of degree two are in the set $\mathcal{V}$. Then we have that:

1. the graph $\Gamma$ is a tree;
2. the set $\mathcal{V}$ has exactly $k$ elements, where $k$ is the number of Dirichlet points $\left\{D_{1}, \ldots, D_{k}\right\} ;$
3. there is at most one vertex $V_{j} \in V(\Gamma) \backslash \mathcal{V}$ of degree one;
4. if there is no vertex of degree one in $V(\Gamma) \backslash \mathcal{V}$, then the graph $\Gamma$ has at most $2 k-2$ vertices and $2 k-3$ edges;
5. if there is exactly one vertex of degree one in $V(\Gamma) \backslash \mathcal{V}$, then the graph $\Gamma$ has at most $2 k$ vertices and $2 k-1$ edges.

Proof. We use the notation $V(\Gamma)=\left\{V_{1}, \ldots, V_{N}\right\}$ for the vertices of $\Gamma$ and $e_{i j}$ for the edges $\left\{V_{i}, V_{j}\right\} \in E(\Gamma)$, whose lengths are denoted by $l_{i j}$. Moreover, we can suppose that for $j=1, \ldots, k$, we have $\gamma\left(V_{j}\right)=D_{j}$, where $D_{1}, \ldots, D_{k}$ are the Dirichlet points from problem (2.28) and so, $\left\{V_{1}, \ldots, V_{k}\right\} \subset \mathcal{V}$. Let $w=\left(w_{i j}\right)_{i j}$ be the Energy function on $\Gamma$ with Dirichlet conditions in the points of $\mathcal{V}$.

1. Suppose that we can remove an edge $e_{i j} \in E(\Gamma)$, such that the graph $\Gamma^{\prime}=$ $\left(V(\Gamma), E(\Gamma) \backslash e_{i j}\right)$ is still connected. Since $w_{i j}^{\prime \prime}=-1$ on $\left[0, l_{i j}\right]$ we have that at least one of the derivatives $w_{i j}^{\prime}(0)$ and $w_{i j}^{\prime}\left(l_{i j}\right)$ is not zero. We can suppose that $w_{i j}^{\prime}\left(l_{i j}\right) \neq 0$. Consider the new graph $\widetilde{\Gamma}$ to which we add a new vertex: $V(\widetilde{\Gamma})=V(\Gamma) \cup V_{0}$, then erase the edge $e_{i j}$ and create a new one $e_{i 0}=\left\{V_{i}, V_{0}\right\}$, of the same length, connecting $V_{i}$ to $V_{0}: E(\widetilde{\Gamma})=\left(E(\Gamma) \backslash e_{i j}\right) \cup e_{i 0}$. Let $\widetilde{w}$ be the Energy function on $\tilde{\Gamma}$ with Dirichlet conditions in $\mathcal{V}$. When seen as a subspaces of $\oplus_{i j} H^{1}\left(\left[0, l_{i j}\right]\right)$, we have that $H_{0}^{1}(\Gamma ; \mathcal{V}) \subset H_{0}^{1}(\widetilde{\Gamma} ; \mathcal{V})$ and so $\mathcal{E}(\widetilde{\Gamma} ; \mathcal{V}) \leq \mathcal{E}(\Gamma ; \mathcal{V})$, where the equality occurs, if and only if the Energy functions $w$ and $\widetilde{w}$ have the same components in $\oplus_{i j} H^{1}\left(\left[0, l_{i j}\right]\right)$. In particular, we must have that $w_{i j}=\widetilde{w}_{i 0}$ on the interval $\left[0, l_{i j}\right]$, which is impossible since $w_{i j}^{\prime}\left(l_{i j}\right) \neq 0$ and $\widetilde{w}_{i 0}^{\prime}\left(l_{i j}\right)=0$.
2. Suppose that there is a vertex $V_{j} \in \mathcal{V}$ with $j>k$ and let $\widetilde{w}$ be the Energy function on $\Gamma$ with Dirichlet conditions in $\left\{V_{1}, \ldots, V_{k}\right\}$. We have the inclusion $H_{0}^{1}(\Gamma ; \mathcal{V}) \subset$ $H_{0}^{1}\left(\Gamma ;\left\{V_{1}, \ldots, V_{k}\right\}\right)$ and so, the inequality $J(\widetilde{w})=\mathcal{E}\left(\Gamma ;\left\{V_{1}, \ldots, V_{k}\right\}\right) \leq \mathcal{E}(\Gamma ; \mathcal{V})=$ $J(w)$, which becomes an equality if and only if $\widetilde{w}=w$, which is impossible. Indeed, if the equality holds, then in $V_{j}, w$ satisfies both the Dirichlet condition and the Kirchoff's law. Since $w$ is positive, for any edge $e_{j i}$ we must have $w_{j i}(0)=$ $0, w_{j i}^{\prime}(0)=0, w_{j i}^{\prime \prime}=-1$ ad $w_{j i} \geq 0$ on $\left[0, l_{j i}\right]$, which is impossible.
3. Suppose that there are two vertices $V_{i}$ and $V_{j}$ of degree one, which are not in $\mathcal{V}$, i.e. $i, j>k$. Since $\Gamma$ is connected, there are two edges, $e_{i i^{\prime}}$ and $e_{j j^{\prime}}$ starting from $V_{i}$ and $V_{j}$ respectively. Suppose that the Energy function $w \in H_{0}^{1}\left(\Gamma ;\left\{V_{1}, \ldots, V_{k}\right\}\right)$ is such that $w\left(V_{i}\right) \geq w\left(V_{j}\right)$. We define a new graph $\tilde{\Gamma}$ by erasing the edge $e_{j j^{\prime}}$ and creating the edge $e_{i j}$ of length $l_{j j^{\prime}}$. On the new edge $e_{i j}$ we consider the function $w_{i j}(x)=w_{j j^{\prime}}(x)+w\left(V_{i}\right)-w\left(V_{j}\right)$. The function $\widetilde{w}$ on $\widetilde{\Gamma}$ obtained by this construction is such that $J(\widetilde{w}) \leq J(w)$, this concludes the proof.

The points $(i v)$ and $(v)$ follow by the construction in Theorem 2.16 and the previous claims ( $i$ ), (ii) and (iii).

Remark 2.20. Suppose that $V_{j} \in V(\Gamma) \backslash \mathcal{V}$ is a vertex of degree one and let $V_{i}$ be the vertex such that $e_{i j} \in E(\Gamma)$. Then the Energy function $w$ with Dirichlet conditions in $\mathcal{V}$ satisfies $w_{j i}^{\prime}(0)=0$. In this case, we call $V_{j}$ a Neumann vertex. By Proposition 2.19, an optimal graph has at most one Neumann vertex.

### 2.4 Some examples of optimal metric graphs

In this section we show three examples. In the first one we deal with two Dirichlet points, the second concerns three aligned Dirichlet points and the third one deals with the case in which the Dirichlet points are vertices of an equilateral triangle. In the first and the third one we find the minimizer explicitly as an embedded graph, while in the second one we limit ourselves to prove that there is no embedded minimizer of the energy, i.e. that problem (2.29) does not admit a solution.

In the following example we use a symmetrization technique similar to the one from Remark 2.6.

Example 2.21. Let $D_{1}$ and $D_{2}$ be two distinct points in $\mathbb{R}^{d}$ and let $l \geq\left|D_{1}-D_{2}\right|$ be a real number. Then the problem

$$
\begin{align*}
& \min \left\{\mathcal{E}\left(\Gamma ;\left\{V_{1}, V_{2}\right\}\right): \Gamma \in C M G, l(\Gamma)=l, \quad V_{1}, V_{2} \in V(\Gamma),\right. \\
& \text { exists } \left.\gamma: \Gamma \rightarrow \mathbb{R} \text { immersion, } \gamma\left(V_{1}\right)=D_{1}, \gamma\left(V_{2}\right)=D_{2}\right\} . \tag{2.32}
\end{align*}
$$

has a solution $(\Gamma, \gamma)$, where $\Gamma$ is a metric graph with vertices $V(\Gamma)=\left\{V_{1}, V_{2}, V_{3}, V_{4}\right\}$ and edges $E(\Gamma)=\left\{e_{13}=\left\{V_{1}, V_{3}\right\}, e_{23}=\left\{V_{2}, V_{3}\right\}, e_{43}=\left\{V_{4}, V_{3}\right\}\right\}$ of lengths $l_{13}=l_{23}=$ $\frac{1}{2}\left|D_{1}-D_{2}\right|$ and $l_{34}=l-\left|D_{1}-D_{2}\right|$, respectively. The map $\gamma: \Gamma \rightarrow \mathbb{R}^{d}$ is an embedding such that $\gamma\left(V_{1}\right)=D_{1}, \gamma\left(V_{2}\right)=D_{2}$ and $\gamma\left(V_{3}\right)=\frac{D_{1}+D_{2}}{2}$ (see Figure 2.4).


Figure 2.4: The optimal graph with two Dirichlet points.
To fix the notations, we suppose that $\left|D_{1}-D_{2}\right|=l-\varepsilon$. Let $u=\left(u_{i j}\right)_{i j}$ be the Energy function of a generic metric graph $\Sigma$ and immersion $\sigma: \Sigma \rightarrow \mathbb{R}^{d}$ with $D_{1}, D_{2} \in \sigma(V(\Sigma))$. Let $M=\max \{u(x): x \in \Sigma\}>0$. We construct a candidate $v \in H_{0}^{1}\left(\Gamma ;\left\{V_{1}, V_{2}\right\}\right)$ such that $J(v) \leq J(u)$, which immediately gives the conclusion.

We define $v$ through the following three increasing functions

$$
v_{13}=v_{23} \in H^{1}([0,(l-\varepsilon) / 2]), \quad v_{34} \in H^{1}([0, \varepsilon])
$$

with boundary values

$$
v_{13}(0)=v_{23}(0)=0, \quad v_{13}((l-\varepsilon) / 2)=v_{23}((l-\varepsilon) / 2)=v_{34}(0)=m<M
$$

and level sets uniquely determined by the equality $\mu_{u}=\mu_{v}$, where $\mu_{u}$ and $\mu_{v}$ are the distribution functions of $u$ and $v$ respectively, defined by

$$
\begin{aligned}
& \mu_{u}(t)=\mathcal{H}^{1}(\{u \leq t\})=\sum_{e_{i j} \in E(\Sigma)} \mathcal{H}^{1}\left(\left\{u_{i j} \leq t\right\}\right) \\
& \mu_{v}(t)=\mathcal{H}^{1}(\{v \leq t\})=\sum_{j=1,2,4} \mathcal{H}^{1}\left(\left\{v_{j 3} \leq t\right\}\right)
\end{aligned}
$$

As in Remark 2.6 we have $\|v\|_{L^{1}(\Gamma)}=\|u\|_{L^{1}(C)}$ and

$$
\begin{equation*}
\int_{\Sigma}\left|u^{\prime}\right|^{2} d x=\int_{0}^{M}\left(\sum_{u=\tau}\left|u^{\prime}\right|\right) d \tau \geq \int_{0}^{M} n_{u}^{2}(\tau)\left(\sum_{u=\tau} \frac{1}{\left|u^{\prime}\right|(\tau)}\right)^{-1} d \tau=\int_{0}^{M} \frac{n_{u}^{2}(\tau)}{\mu_{u}^{\prime}(\tau)} d \tau \tag{2.33}
\end{equation*}
$$

where $n_{u}(\tau)=\mathcal{H}^{0}(\{u=\tau\})$. The same argument holds for $v$ on the graph $\Gamma$ but, this time, with the equality sign:

$$
\begin{equation*}
\int_{\Gamma}\left|v^{\prime}\right|^{2} d x=\int_{0}^{M}\left(\sum_{v=\tau}\left|v^{\prime}\right|\right) d \tau=\int_{0}^{M} \frac{n_{v}^{2}(\tau)}{\mu_{v}^{\prime}(\tau)} d \tau \tag{2.34}
\end{equation*}
$$

since $\left|v^{\prime}\right|$ is constant on $\{v=\tau\}$, for every $\tau$. Then, in view of (2.33) and (2.34), to conclude it is enough to prove that $n_{u}(\tau) \geq n_{v}(\tau)$ for almost every $\tau$. To this aim we first notice that, by construction $n_{v}(\tau)=1$ if $\tau \in[m, M]$ and $n_{v}(\tau)=2$ if $\tau \in[0, m)$. Since $n_{u}$ is decreasing and greater than 1 on $[0, M]$, we only need to prove that $n_{u} \geq 2$ on $[0, m]$. To see this, consider two vertices $W_{1}, W_{2} \in V(\Sigma)$ such that $\sigma\left(W_{1}\right)=D_{1}$ and $\sigma\left(W_{2}\right)=D_{2}$. Let $\eta$ be a simple path connecting $W_{1}$ to $W_{2}$ in $\Sigma$. Since $\sigma$ is an immersion we know that the length $l(\eta)$ of $\eta$ is at least $l-\varepsilon$. By the continuity of $u$, we know that $n_{u} \geq 2$ on the interval $\left[0, \max _{\eta} u\right)$. Since $n_{v}=1$ on $[m, M]$, we need to show that $\max _{\eta} u \geq m$. Otherwise, we would have

$$
l(\eta) \leq\left|\left\{u \leq \max _{\eta} u\right\}\right|<|\{u \leq m\}|=|\{v \leq m\}|=\left|D_{1}-D_{2}\right| \leq l(\eta)
$$

which is impossible.
Remark 2.22. In the previous example the optimal metric graph $\Gamma$ is such that for any (admissible) immersion $\gamma: \Gamma \rightarrow \mathbb{R}^{d}$, we have $\left|\gamma\left(V_{1}\right)-\gamma\left(V_{3}\right)\right|=l_{13}$ and $\mid \gamma\left(V_{2}\right)-$ $\gamma\left(V_{3}\right) \mid=l_{23}$, i.e. the point $\gamma\left(V_{3}\right)$ is necessary the midpoint $\frac{D_{1}+D_{2}}{2}$, so we have a sort
of rigidity of the graph $\Gamma$. More generally, we say that an edge $e_{i j}$ is rigid, if for any admissible immersion $\gamma: \Gamma \rightarrow \mathbb{R}^{d}$, i.e. an immersion such that $\mathcal{D}=\gamma(\mathcal{V})$, we have $\left|\gamma\left(V_{i}\right)-\gamma\left(V_{j}\right)\right|=l_{i j}$, in other words the realization of the edge $e_{i j}$ in $\mathbb{R}^{d}$ via any immersion $\gamma$ is a segment. One may expect that in the optimal graph all the edges, except the one containing the Neumann vertex, are rigid. Unfortunately, we are able to prove only the weaker result that:

1. if the Energy function $w$, of an optimal metric graph $\Gamma$, has a local maximum in the interior of an edge $e_{i j}$, then the edge is rigid; if the maximum is global, then $\Gamma$ has no Neumann vertices;
2. if $\Gamma$ contains a Neumann vertex $V_{j}$, then $w$ achieves its maximum at it.

To prove the second claim, we just observe that if it is not the case, then we can use an argument similar to the one from point (iii) of Proposition 2.19, erasing the edge $e_{i j}$ containing the Neumann vertex $V_{j}$ and creating an edge of the same length that connects $V_{j}$ to the point, where $w$ achieves its maximum, which we may assume a vertex of $\Gamma$ (possibly of degree two).

For the first claim, we apply a different construction which involves a symmetrization technique. In fact, if the edge $e_{i j}$ is not rigid, then we can create a new metric graph of smaller energy, for which there is still an immersion which satisfies the conditions in problem (2.28). In this there are points $0<a<b<l_{i j}$ such that $l_{i j}-(b-a) \geq\left|\gamma\left(V_{i}\right)-\gamma\left(V_{j}\right)\right|$ and $\min _{[a, b]} w_{i j}=w_{i j}(a)=w_{i j}(b)<\max _{[a, b]} w_{i j}$. Since the edge is not rigid, there is an immersion $\gamma$ such that $\left|\gamma_{i j}(a)-\gamma_{i j}(b)\right|>|b-a|$. The problem (2.32) with $D_{1}=\gamma_{i j}(a)$ and $D_{2}=\gamma_{i j}(b)$ has as a solution the $T$-like graph described in Example 2.21. This shows that the original graph could not be optimal, which is a contradiction.

Example 2.23. Consider the set of points $\mathcal{D}=\left\{D_{1}, D_{2}, D_{3}\right\} \subset \mathbb{R}^{2}$ with coordinates respectively $(-1,0),(1,0)$ and $(n, 0)$, where $n$ is a positive integer. Given $l=(n+2)$, we aim to show that for $n$ large enough there is no solution of the optimization problem

$$
\begin{equation*}
\min \{\mathcal{E}(\Gamma ; \mathcal{V}): \Gamma \in C M G, l(\Gamma)=l, \mathcal{V} \subset V(\Gamma), \exists \gamma: \Gamma \rightarrow \mathbb{R} \text { embedding, } \mathcal{D}=\gamma(\mathcal{V})\} \tag{2.35}
\end{equation*}
$$

In fact, we show that all the possible solutions of the problem

$$
\begin{equation*}
\min \{\mathcal{E}(\Gamma ; \mathcal{V}): \Gamma \in C M G, l(\Gamma)=l, \mathcal{V} \subset V(\Gamma), \exists \gamma: \Gamma \rightarrow \mathbb{R} \text { immersion, } \mathcal{D}=\gamma(\mathcal{V})\} \tag{2.36}
\end{equation*}
$$

are metric graphs $\Gamma$ for which there is no embedding $\gamma: \Gamma \rightarrow \mathbb{R}^{2}$ such that $\mathcal{D} \subset \gamma(V(\Gamma))$. Moreover, there is a sequence of embedded metric graphs which is a minimizing sequence for the problem (2.36).

More precisely, we show that the only possible solution of (2.36) is one of the following metric trees:

1. $\Gamma_{1}$ with vertices $V\left(\Gamma_{1}\right)=\left\{V_{1}, V_{2}, V_{3}, V_{4}\right\}$ and edges $E\left(\Gamma_{1}\right)=\left\{e_{14}=\left\{V_{1}, V_{4}\right\}, e_{24}=\right.$ $\left.\left\{V_{2}, V_{4}\right\}, e_{34}=\left\{V_{3}, V_{4}\right\}\right\}$ of lengths $l_{14}=l_{24}=1$ and $l_{34}=n$, respectively. The set of vertices in which the Dirichlet condition holds is $\mathcal{V}_{1}=\left\{V_{1}, V_{2}, V_{3}\right\}$.
2. $\Gamma_{2}$ with vertices $V\left(\Gamma_{2}\right)=\left\{W_{i}\right\}_{i=1}^{6}$, and edges $E\left(\Gamma_{2}\right)=\left\{e_{14}, e_{24}, e_{35}, e_{45}, e_{56}\right\}$ ,where $e_{i j}=\left\{W_{i}, W_{j}\right\}$ for $1 \leq i \neq j \leq 6$ of lengths $l_{14}=1+\alpha, l_{24}=1-\alpha, l_{35}=$ $n-\beta, l_{45}=\beta-\alpha, l_{56}=\alpha$, where $0<\alpha<1$ and $\alpha<\beta<n$. The set of vertices in which the Dirichlet condition holds is $\mathcal{V}_{1}=\left\{V_{1}, V_{2}, V_{3}\right\}$. A possible immersion $\gamma$ is described in Figure 2.5.


Figure 2.5: The two candidates for a solution of (2.36).
We start showing that if there is an optimal metric graph with no Neumann vertex, then it must be $\Gamma_{1}$. In fact, by Proposition 2.19, we know that the optimal metric graph is of the form $\Gamma_{1}$, but we have no information on the lengths of the edges, which we set as $l_{i}=l\left(e_{i 4}\right)$, for $i=1,2,3$ (see Figure 2.6). We can calculate explicitly the minimizer of the Energy functional and the energy itself in function of $l_{1}, l_{2}$ and $l_{3}$.


Figure 2.6: A metric tree with the same topology as $\Gamma_{1}$.
The minimizer of the energy $w: \Gamma \rightarrow \mathbb{R}$ is given by the functions $w_{i}:\left[0, l_{i}\right] \rightarrow \mathbb{R}$, where $i=1,2,3$ and

$$
\begin{equation*}
w_{i}(x)=-\frac{x^{2}}{2}+a_{i} x . \tag{2.37}
\end{equation*}
$$

where

$$
\begin{equation*}
a_{1}=\frac{l_{1}}{2}+\frac{l_{2} l_{3}\left(l_{1}+l_{2}+l_{3}\right)}{2\left(l_{1} l_{2}+l_{2} l_{3}+l_{3} l_{1}\right)}, \tag{2.38}
\end{equation*}
$$

and $a_{2}$ and $a_{3}$ are defined by a cyclic permutation of the indices. As a consequence, we obtain that the derivative along the edge $e_{14}$ in the vertex $V_{4}$ is given by

$$
\begin{equation*}
w_{1}^{\prime}\left(l_{1}\right)=-l_{1}+a_{1}=-\frac{l_{1}}{2}+\frac{l_{2} l_{3}\left(l_{1}+l_{2}+l_{3}\right)}{2\left(l_{1} l_{2}+l_{2} l_{3}+l_{3} l_{1}\right)}, \tag{2.39}
\end{equation*}
$$

and integrating the Energy function $w$ on $\Gamma$, we obtain

$$
\begin{equation*}
\mathcal{E}\left(\Gamma ;\left\{V_{1}, V_{2}, V_{3}\right\}\right)=-\frac{1}{12}\left(l_{1}^{3}+l_{2}^{3}+l_{3}^{3}\right)-\frac{\left(l_{1}+l_{2}+l_{3}\right)^{2} l_{1} l_{2} l_{3}}{4\left(l_{1} l_{2}+l_{2} l_{3}+l_{3} l_{1}\right)} . \tag{2.40}
\end{equation*}
$$

Studying this function using Lagrange multipliers is somehow complicated due to the complexity of its domain. Thus we use a more geometric approach applying the symmetrization technique described in Remark 2.6 in order to select the possible candidates. We prove that if the graph is optimal, then all the edges must be rigid (this would force the graph to coincide with $\Gamma_{1}$ ). Suppose that the optimal graph $\Gamma$ is not rigid, i.e. there is a non-rigid edge. Then, for $n>4$, we have that $l_{2}<l_{1}<l_{3}$ and so, by (2.39), we obtain $w_{3}^{\prime}\left(l_{3}\right)<w_{1}^{\prime}\left(l_{1}\right)<w_{2}^{\prime}\left(l_{2}\right)$. As a consequence of the Kirchoff's law we have $w_{3}^{\prime}\left(l_{3}\right)<0$ and $w_{2}^{\prime}\left(l_{2}\right)>0$ and so, $w$ has a local maximum on the edge $e_{34}$ and is increasing on $e_{14}$. By Remark 2.22, we obtain that the edge $e_{34}$ is rigid.

We first prove that $w_{1}^{\prime}\left(l_{1}\right)>0$. In fact, if this is not the case, i.e. $w_{1}^{\prime}\left(l_{1}\right)<0$, by Remark 2.22, we have that the edges $e_{14}$ is also rigid and so, $l_{1}+l_{3}=\left|D_{1}-D_{3}\right|=n+1$, i.e. $l_{2}=1$. Moreover, by (2.39), we have that $w_{1}^{\prime}\left(l_{1}\right)<0$, if and only if $l_{1}^{2}>l_{2} l_{3}=l_{3}$. The last inequality does not hold for $n>11$, since, by the triangle inequality, $l_{2}+l_{3} \geq$ $\left|D_{2}-D_{3}\right|=n-1$, we have $l_{1} \leq 3$. Thus, for $n$ large enough, we have that $w$ is increasing on the edge $e_{14}$.

We now prove that the edges $e_{14}$ and $e_{24}$ are rigid. In fact, suppose that $e_{24}$ is not rigid. Let $a \in\left(0, l_{1}\right)$ and $b \in\left(0, l_{2}\right)$ be two points close to $l_{1}$ and $l_{2}$ respectively and such that $w_{14}(a)=w_{24}(b)<w\left(V_{4}\right)$ since $w_{14}$ and $w_{24}$ are strictly increasing. Consider the metric graph $\widetilde{\Gamma}$ whose vertices and edges are

$$
\begin{gathered}
V(\widetilde{\Gamma})=\left\{V_{1}=\widetilde{V}_{1}, V_{2}=\widetilde{V}_{2}, V_{3}=\widetilde{V}_{3}, V_{4}=\widetilde{V}_{4}, \widetilde{V}_{5}, \widetilde{V}_{6}\right\}, \\
E(\widetilde{\Gamma})=\left\{e_{15}, e_{25}, e_{45}, e_{34}, e_{46}\right\},
\end{gathered}
$$

where $e_{i j}=\left\{\widetilde{V}_{i}, \widetilde{V}_{j}\right\}$ and the lengths of the edges are respectively (see Figure 2.7)

$$
\widetilde{l}_{15}=a, \widetilde{l}_{25}=b, \widetilde{l}_{45}=l_{2}-b, \widetilde{l}_{34}=l_{3}, \widetilde{l}_{46}=l_{1}-a .
$$



Figure 2.7: The graph $\Gamma$ (on the left) and the modified one $\widetilde{\Gamma}$ (on the right).
The new metric graph is still a competitor in the problem (2.36) and there is a function $w \in H_{0}^{1}\left(\widetilde{\Gamma} ;\left\{V_{1}, V_{2}, V_{3}\right\}\right)$ such that $\mathcal{E}\left(\widetilde{\Gamma} ;\left\{V_{1}, V_{2}, V_{3}\right\}\right)<J(\widetilde{w})=J(w)$, which is a contradiction with the optimality of $\Gamma$. In fact, it is enough to define $\widetilde{w}$ as

$$
\widetilde{w}_{15}=\left.w_{14}\right|_{[0, a]}, \widetilde{w}_{25}=\left.w_{24}\right|_{[0, b]}, \widetilde{w}_{54}=\left.w_{24}\right|_{\left[b, l_{2}\right]}, \widetilde{w}_{34}=w_{34}, \widetilde{w}_{64}=\left.w_{14}\right|_{\left[a, l_{1}\right]},
$$

and observe that $\widetilde{w}$ is not the Energy function on the graph $\widetilde{\Gamma}$ since it does not satisfy the Neumann condition in $\widetilde{V}_{6}$. In the same way, if we suppose that $w_{14}$ is not rigid, we obtain a contradiction, and so all the three edges must be rigid, i.e. $\Gamma=\Gamma_{1}$.

In a similar way we prove that a metric graph $\Gamma$ with a Neumann vertex can be a solution of (2.36) only if it is of the same form as $\Gamma_{2}$. We proceed in two steps: first, we show that, for $n$ large enough, the edge containing the Neumann vertex has a common vertex with the longest edge of the graph; then we can conclude reasoning analogously to the previous case. Let $\Gamma$ be a metric graph with vertices $V(\Gamma)=\left\{V_{i}\right\}_{i=1}^{6}$, and edges $E(\Gamma)=\left\{e_{15}, e_{24}, e_{34}, e_{45}, e_{56}\right\}$, where $e_{i j}=\left\{V_{i}, V_{j}\right\}$ for $1 \leq i \neq j \leq 6$.

We prove that $w\left(V_{6}\right) \leq \max _{e_{34}} w$, i.e. the graph $\Gamma$ is not optimal, since, by Remark 2.22, the maximum of $w$ must be achieved in the Neumann vertex $V_{6}$ (the case $E(\Gamma)=$ $\left\{e_{14}, e_{25}, e_{34}, e_{45}, e_{56}\right\}$ is analogous). Let $w_{15}:\left[0, l_{15}\right] \rightarrow \mathbb{R}, w_{65}:\left[0, l_{65}\right] \rightarrow \mathbb{R}$ and $w_{34}:\left[0, l_{34}\right] \rightarrow \mathbb{R}$ be the restrictions of the Energy function $w$ of $\Gamma$ to the edges $e_{15}, e_{65}$ and $e_{34}$ of lengths $l_{15}, l_{65}$ and $l_{34}$, respectively. Let $u:\left[0, l_{15}+l_{56}\right] \rightarrow \mathbb{R}$ be defined as

$$
u(x)=\left\{\begin{array}{l}
w_{15}(x), x \in\left[0, l_{15}\right]  \tag{2.41}\\
w_{56}\left(x-l_{15}\right), x \in\left[l_{15} \cdot l_{15}+l_{56}\right]
\end{array}\right.
$$

If the metric graph $\Gamma$ is optimal, then the Energy function on $w_{54}$ on the edge $e_{45}$ must be decreasing and so, by the Kirchhoff's law in the vertex $V_{5}$, we have that $w_{15}^{\prime}\left(l_{15}\right)+w_{65}^{\prime}\left(l_{65}\right) \leq 0$, i.e. the left derivative of $u$ at $l_{15}$ is less than the right one:

$$
\partial_{-} u\left(l_{15}\right)=w_{15}^{\prime}\left(l_{15}\right) \leq w_{56}^{\prime}(0)=\partial_{+} u\left(l_{15}\right)
$$

By the maximum principle, we have that

$$
u(x) \leq \widetilde{u}(x)=-\frac{x^{2}}{2}+\left(l_{15}+l_{56}\right) x \leq \frac{1}{2}\left(l_{15}+l_{56}\right)^{2}
$$

On the other hand, $w_{34}(x) \geq v(x)=-\frac{x^{2}}{2}+\frac{l_{34}}{2} x$, again by the maximum principle on the interval $\left[0, l_{34}\right]$. Thus we have that

$$
\max _{x \in\left[0, l_{34}\right]} w_{34}(x) \geq \max _{x \in\left[0, l_{34}\right]} v(x)=\frac{1}{8} l_{34}^{2}>\frac{1}{2}\left(l_{15}+l_{56}\right)^{2} \geq w\left(V_{6}\right)
$$

for $n$ large enough.
Repeating the same argument, one can show that the optimal metric graph $\Gamma$ is not of the form $V(\Gamma)=\left\{V_{1}, V_{2}, V_{3}, V_{4}, V_{5}\right\}, E(\Gamma)=\left\{\left\{V_{1}, V_{4}\right\},\left\{V_{2}, V_{4}\right\},\left\{V_{3}, V_{4}\right\},\left\{V_{4}, V_{5}\right\}\right\}$.

Thus, we obtained that the if the optimal graph has a Neumann vertex, then the corresponding edge must be attached to the longest edge. To prove that it is of the same form as $\Gamma_{2}$, there is one more case to exclude, namely: $\Gamma$ defined by $V(\Gamma)=\left\{V_{1}, V_{2}, V_{3}, V_{4}, V_{5}\right\}, E(\Gamma)=\left\{\left\{V_{1}, V_{2}\right\},\left\{V_{2}, V_{4}\right\},\left\{V_{3}, V_{4}\right\},\left\{V_{4}, V_{5}\right\}\right\}$ (see Figure 2.8). By Example 2.21, the only possible candidate of this form is the graph with lengths
$l\left(\left\{V_{1}, V_{2}\right\}\right)=\left|D_{1}-D_{2}\right|=2, l\left(\left\{V_{2}, V_{4}\right\}\right)=\frac{n-1}{2}, l\left(\left\{V_{3}, V_{4}\right\}\right)=\frac{n-1}{2}, l\left(\left\{V_{4}, V_{5}\right\}\right)=2$. In this case, we compare the energy of $\Gamma$ and $\Gamma_{1}$ by an explicit calculation:

$$
\begin{equation*}
\mathcal{E}\left(\Gamma ;\left\{V_{1}, V_{2}, V_{3}\right\}\right)=-\frac{n^{3}-3 n^{2}+6 n}{24}>-\frac{n^{2}(n+1)^{2}}{12(2 n+1)}=\mathcal{E}\left(\Gamma_{1} ;\left\{V_{1}, V_{2}, V_{3}\right\}\right) \tag{2.42}
\end{equation*}
$$

for $n$ large enough.


Figure 2.8: The graph $\Gamma_{1}$ (on the left) has lower energy than the graph $\Gamma$ (on the right).
Before we pass to our last example, we need the following Lemma.
Lemma 2.24. Let $w_{a}:[0,1] \rightarrow \mathbb{R}$ be given by $w_{a}(x)=-\frac{x^{2}}{2}+$ ax, for some positive real number $a$. If $w_{a}(1) \leq w_{A}(1) \leq \max _{x \in[0,1]} w_{a}(x)$, then $J\left(w_{A}\right) \leq J\left(w_{a}\right)$, where $J(w)=\frac{1}{2} \int_{0}^{1}\left|w^{\prime}\right|^{2} d x-\int_{0}^{1} w d x$.

Proof. It follows by performing the explicit calculations.
Example 2.25. Let $D_{1}, D_{2}$ and $D_{3}$ be the vertices of an equilateral triangle of side 1 in $\mathbb{R}^{2}$, i.e.

$$
D_{1}=\left(-\frac{\sqrt{3}}{3}, 0\right), D_{2}=\left(\frac{\sqrt{3}}{6},-\frac{1}{2}\right), D_{3}=\left(\frac{\sqrt{3}}{6}, \frac{1}{2}\right)
$$

We study the problem (2.28) with $\mathcal{D}=\left\{D_{1}, D_{2}, D_{3}\right\}$ and $l>\sqrt{3}$. We show that the solutions may have different qualitative properties for different $l$ and that there is always a symmetry breaking phenomenon, i.e. the solutions does not have the same symmetries as the initial configuration $\mathcal{D}$. We first reduce our study to the following three candidates (see Figure 2.9):

1. The metric tree $\Gamma_{1}$, defined by with vertices $V(\Gamma)=\left\{V_{1}, V_{2}, V_{3}, V_{4}\right\}$ and edges $E(\Gamma)=\left\{e_{14}, e_{24}, e_{34}\right\}$, where $e_{i j}=\left\{V_{i}, V_{j}\right\}$ and the lengths of the edges are respectively $l_{24}=l_{34}=x, l_{14}=\frac{\sqrt{3}}{2}-\sqrt{x^{2}-\frac{1}{4}}$, for some $x \in[1 / 2,1 / \sqrt{3}]$. Note that the length of $\Gamma_{1}$ is less than $1+\sqrt{3} / 2$, i.e. it is a possible solution only for $l \leq 1+\sqrt{3} / 2$. The new vertex $V_{4}$ is of Kirchhoff type and there are no Neumann vertices. A symmetrization argument around the point $P$ shows that there is a better candidate among the graphs of the type described in point (2) below.
2. The metric tree $\Gamma_{2}$ with vertices $V=\left(V_{1}, V_{2}, V_{3}, V_{4}, V_{5}\right)$ and $E(\Gamma)=\left\{e_{14}, e_{24}, e_{34}, e_{45}\right\}$, where $e_{i j}=\left\{V_{i}, V_{j}\right\}$ and the lengths of the edges $l_{14}=l_{24}=l_{34}=1 / \sqrt{3}$, $l_{45}=l-\sqrt{3}$, respectively. The new vertex $V_{4}$ is of Kirchhoff type and $V_{5}$ is a Neumann vertex.
3. The metric tree $\Gamma_{3}$ with vertices $V(\Gamma)=\left\{V_{1}, V_{2}, V_{3}, V_{4}, V_{5}, V_{6}\right\}$ and edges $E(\Gamma)=$ $\left\{e_{15}, e_{24}, e_{34}, e_{45}, e_{56}\right\}$, where $e_{i j}=\left\{V_{i}, V_{j}\right\}$ and the lengths of the edges are $l_{24}=$ $l_{34}=x, l_{15}=\frac{l x}{2(2 l-3 x)}+\frac{\sqrt{3}}{4}-\frac{1}{4} \sqrt{4 x^{2}-1}, l_{45}=\frac{\sqrt{3}}{4}-\frac{l x}{2(2 l-3 x)}-\frac{1}{4} \sqrt{4 x^{2}-1}$ and $l_{56}=l-2 x-\sqrt{3} / 2+\frac{1}{2} \sqrt{4 x^{2}-1}$. The new vertices $V_{4}$ and $V_{5}$ are of Kirchhoff type and $V_{6}$ is a Neumann vertex.


Figure 2.9: The three competing graphs.
Suppose that the metric graph $\Gamma$ is optimal and has the same vertices and edges as $\Gamma_{1}$. Without loss of generality, we can suppose that the maximum of the Energy function $w$ on $\Gamma$ is achieved on the edge $e_{14}$. If $l_{24} \neq l_{34}$, we consider the metric graph $\widetilde{\Gamma}$ with the same vertices and edges as $\Gamma$ and lengths $\widetilde{l}_{14}=l_{14}, \widetilde{l}_{24}=\widetilde{l}_{34}=\left(l_{24}+l_{34}\right) / 2$. An immersion $\widetilde{\gamma}: \widetilde{\Gamma} \rightarrow \mathbb{R}^{2}$, such that $\widetilde{\gamma}\left(V_{j}\right)=D_{j}$, for $j=1,2,3$ still exists and the energy decreases, i.e. $\mathcal{E}\left(\widetilde{\Gamma} ;\left\{V_{1}, V_{2}, V_{3}\right\}\right)<\mathcal{E}\left(\Gamma ;\left\{V_{1}, V_{2}, V_{3}\right\}\right)$. In fact, let $v=\widetilde{w}_{24}=\widetilde{w}_{34}$ : $\left[0, \frac{l_{24}+l_{34}}{2}\right] \rightarrow \mathbb{R}$ be an increasing function such that $2|\{v \geq \tau\}|=\left|\left\{w_{24} \geq \tau\right\}\right|+\mid\left\{w_{34} \geq\right.$ $\tau\} \mid$. By the classical Pólya-Szegö inequality and by the fact that $w_{24}$ and $w_{34}$ are never constant in an open region, we obtain that

$$
J\left(\widetilde{w}_{24}\right)+J\left(\widetilde{w}_{34}\right)<J\left(w_{24}\right)+J\left(w_{34}\right)
$$

and so it is enough to construct a function $\widetilde{w}_{14}:\left[0, l_{14}\right] \rightarrow \mathbb{R}$ such that $\widetilde{w}_{14}\left(l_{14}\right)=$ $\widetilde{w}_{24}=\widetilde{w}_{34}$ and $J\left(\widetilde{w}_{14}\right) \leq J\left(w_{14}\right)$. Consider a function such that $\widetilde{w}_{14}^{\prime \prime}=-1, \widetilde{w}_{14}(0)=0$ and $\widetilde{w}_{14}\left(l_{14}\right)=\widetilde{w}_{24}\left(l_{2} 4\right)=\widetilde{w}_{34}\left(l_{34}\right)$. Since we have the inequality $w_{14}\left(l_{14}\right) \leq \widetilde{w}_{14}\left(l_{14}\right) \leq$ $\max _{\left[0, l_{14}\right]} w_{14}=\max _{\Gamma} w$, we can apply Lemma 2.24 and so, $J\left(\widetilde{w}_{14}\right) \leq J\left(w_{14}\right)$. Thus, we obtain that $l_{24}=l_{34}$ and that both the functions $w_{24}$ and $w_{34}$ are increasing (in particular, $l_{14} \geq l_{24}=l_{34}$ ). If the maximum of $w$ is achieved in the interior of the edge $e_{14}$ then, by Remark 2.22, the edge $e_{14}$ must be rigid and so, all the edges must
be rigid. Thus, $\Gamma$ coincides with $\Gamma_{1}$ for some $x \in\left(\frac{1}{2}, \frac{1}{\sqrt{3}}\right]$. If the maximum of $w$ is achieved in the vertex $V_{4}$, then applying one more time the above argument, we obtain $l_{14}=l_{24}=l_{34}=\frac{1}{\sqrt{3}}$, i.e. $\Gamma$ is $\Gamma_{1}$ corresponding to $x=\frac{1}{\sqrt{3}}$.

Suppose that the metric graph $\Gamma$ is optimal and that has the same vertices as $\Gamma_{2}$. If $w=\left(w_{i j}\right)_{i j}$ is the Energy function on $\Gamma$ with Dirichlet conditions in $\left\{V_{1}, V_{2}, V_{3}\right\}$, we have that $w_{14}, w_{24}$ an $w_{34}$ are increasing on the edges $e_{14}, e_{24}$ and $e_{34}$. As in the previous situation $\Gamma=\Gamma_{1}$, by a symmetrization argument, we have that $l_{14}=l_{24}=l_{34}$. Since any level set $\{w=\tau\}$ contains exactly 3 points, if $\tau<w\left(V_{4}\right)$, and 1 point, if $\tau \geq w\left(V_{4}\right)$, we can apply the same technique as in Example 2.21 to obtain that $l_{14}=l_{24}=l_{34}=\frac{1}{\sqrt{3}}$.

Suppose that the metric graph $\Gamma$ is optimal and that has the same vertices and edges as $\Gamma_{3}$. Let $w$ be the Energy function on $\Gamma$ with Dirichlet conditions in $\left\{V_{1}, V_{2}, V_{3}\right\}$. Since we assume $\Gamma$ optimal, we have that $w_{45}$ is increasing on the edge $e_{45}$ and $w\left(V_{5}\right) \geq w_{i j}$, for any $\{i, j\} \neq\{5,6\}$. Applying the symmetrization argument from the case $\Gamma=\Gamma_{1}$ and Lemma 2.24, we obtain that $l_{24}=l_{34}=x$ and that the functions $w_{24}=w_{34}$ are increasing on $\left[0, l_{24}\right]$. Let $a \in\left[0, l_{15}\right]$ be such that $w_{15}(a)=w\left(V_{4}\right)$. By a symmetrization argument, we have that necessarily $l_{15}-a=l_{45}$ an that $w_{45}(x)=w_{15}(x-a)$. Moreover, the edges $e_{15}$ and $e_{45}$ are rigid. Indeed, for any admissible immersion $\gamma=\left(\gamma_{i j}\right)_{i j}: \Gamma \rightarrow$ $\mathbb{R}^{2}$, we have that the graph $\widetilde{\Gamma}$ with vertices $V(\widetilde{\Gamma})=\left\{\widetilde{V}_{1}, V_{4}, V_{5}, V_{6}\right\}$ and edges $E(\widetilde{\Gamma})=$ $\left\{\left\{\widetilde{V}_{1}, V_{5}\right\},\left\{V_{4}, V_{5}\right\},\left\{V_{5}, V_{6}\right\}\right\}$, is a solution of the problem (2.32) with $D_{1}:=\gamma_{15}(a)$ and $D_{2}:=\gamma\left(V_{4}\right)$. By Example 2.21 and Remark 2.22, we have $\left|\gamma_{15}(a)-\gamma\left(V_{4}\right)\right|=2 l_{45}$ and, since this holds for every admissible $\gamma$, we deduce the rigidity of $e_{15}$ and $e_{45}$. Using this information one can calculate explicitly all the lengths of the edges of $\Gamma$ using only the parameter $x$, obtaining the third class of possible minimizers.


Figure 2.10: The optimal graphs for $l<1+\sqrt{3} / 2, l=1+\sqrt{3} / 2, l>1+\sqrt{3} / 2$ and $l \gg 1+\sqrt{3} / 2$.

An explicit estimate of the energy shows that:

1. If $\sqrt{3} \leq l \leq 1+\sqrt{3} / 2$, we have that the solution of the problem (2.28) with $\mathcal{D}=\left\{D_{1}, D_{2}, D_{3}\right\}$ is of the form $\Gamma_{1}$ (see Figure 2.10).
2. If $l>1+\sqrt{3} / 2$, then the solution of the problem (2.28) with $\mathcal{D}=\left\{D_{1}, D_{2}, D_{3}\right\}$
is of the form $\Gamma_{3}$.
In both cases, the parameter $x$ is uniquely determined by the total length $l$ and so, we have uniqueness up to rotation on $\frac{2 \pi}{3}$. Moreover, in both cases the solutions are metric graphs, for which there is an embedding $\gamma$ with $\gamma\left(V_{i}\right)=D_{i}$, i.e. they are also solutions of the problem (2.29) with $\mathcal{D}=\left\{D_{1}, D_{2}, D_{3}\right\}$ and $l \geq \sqrt{3}$.

### 2.5 Complements and further results

In this Section we present two generalizations of Theorem 2.18. The first one deals with a more general class of constraints $D_{1}, \ldots, D_{k}$ while in the second one we consider a larger class of admissible sets.

Corollary 2.26. Let $D_{1}, \ldots, D_{k}$ be $k$ disjoint compact sets in $\mathbb{R}^{d}$ and let $l \geq S t\left(d_{1}, \ldots, d_{k}\right)$, i.e. such that there exists a closed connected set $C$ of length $\mathcal{H}^{1}(C)=l$, which intersects all the sets $D_{1}, \ldots, D_{k}$. Then the optimization problem

$$
\begin{equation*}
\min \left\{\mathcal{E}(\Gamma ; \mathcal{V}): \Gamma \in C M G, l(\Gamma)=l, \mathcal{V} \subset V(\Gamma), \Gamma \in \operatorname{Adm}\left(\mathcal{V} ; D_{1}, \ldots, D_{k}\right)\right\} \tag{2.43}
\end{equation*}
$$

admits a solution, where we say that $\Gamma \in \operatorname{Adm}\left(\mathcal{V} ; D_{1}, \ldots, D_{k}\right)$, if there exists an immersion $\gamma: \Gamma \rightarrow \mathbb{R}^{d}$ such that for each $j=1, \ldots, k$ there is $V_{j} \in \mathcal{V}$ such that $\gamma\left(V_{j}\right) \in D_{j}$.
Proof. As in Theorem 2.16, we can restrict our attention to the connected metric trees $\Gamma$ with the same vertices $V(\Gamma)=\left\{V_{1}, \ldots, V_{N}\right\}$ and edges $E(\Gamma)=\left\{e_{i j}\right\}_{i j}$. Moreover, we can suppose that $\mathcal{V}=\left\{V_{1}, \ldots, V_{k}\right\}$ is fixed. By the compactness of the sets $D_{j}$, we can take a minimizing sequence $\Gamma_{n}$ and immersions $\gamma_{n}$ such that for each $j=1, \ldots, k$, we have $\gamma_{n}\left(V_{j}\right) \rightarrow X_{j} \in D_{j}$, as $n \rightarrow \infty$. The claim follows by the same argument as in Theorem 2.18.

Theorem 2.18 can be restated in the more general framework of the metric spaces of finite Hausdorff measure, which is the natural extension of the class of the one dimensional subspaces of $\mathbb{R}^{d}$ of finite length. In fact, for any compact connected metric space (shortly CCMS) $(C, d)$, we consider the one dimensional Hausdorff measure $\mathcal{H}_{d}^{1}$ with respect to the metric $d$ and the Sobolev space $H^{1}(C)$ obtained by the closure of the Lipschitz functions on $C$, with respect to the norm $\|u\|_{H^{1}(C)}^{2}=\|u\|_{L^{2}\left(\mathcal{H}_{d}^{1}\right)}^{2}+\left\|u^{\prime}\right\|_{L^{2}\left(\mathcal{H}_{d}^{1}\right)}^{2}$, where $u^{\prime}$ is defined as in the case $C \subset \mathbb{R}^{d}$. The energy $\mathcal{E}(C ; \mathcal{V})$ with respect to the set $\mathcal{V} \subset C$ is defined as in (2.10). As in the case of metric graphs, we define an immersion $\gamma: C \rightarrow \mathbb{R}^{d}$ as a continuous map such that for any arc-length parametrized curve $\eta:(-\varepsilon, \varepsilon) \rightarrow C$, we have that $\left|(\gamma \circ \eta)^{\prime}(t)\right|=1$ for almost every $t \in(-\varepsilon, \varepsilon)$. As a consequence of Theorem 2.18, we have the following:
Corollary 2.27. Consider the set of points $\mathcal{D}=\left\{D_{1}, \ldots, D_{k}\right\} \subset \mathbb{R}^{d}$ and a positive real number $l \geq S t\left(D_{1}, \ldots, D_{k}\right)$. Then the following optimization problem has solution:

$$
\begin{equation*}
\min \left\{\mathcal{E}(C ; \mathcal{V}):(C, d) \in C C M S, \mathcal{H}_{d}^{1}(C) \leq l, C \in \operatorname{Adm}\left(\mathcal{V} ; D_{1}, \ldots, D_{k}\right)\right\} \tag{2.44}
\end{equation*}
$$

where the admissible set $\operatorname{Adm}\left(\mathcal{V} ;\left\{D_{1}, \ldots, D_{k}\right\}\right)$ is the set of connected metric spaces, for which there exists an immersion $\gamma: \Gamma \rightarrow \mathbb{R}^{d}$ such that $\gamma(\mathcal{V})=\left\{D_{1}, \ldots, D_{k}\right\}$. Moreover, the solution of the problem (2.44) is a connected metric graph, which is a tree of at most $2 k$ vertices and $2 k-1$ edges.
Proof. Repeating the construction from Theorem 2.7, we can restrict our attention to the class of metric graphs. The conclusion follows from Theorem 2.18.

The results from Theorem 2.7 and Theorem 2.18, hold also for other cost functionals as, for example, the first eigenvalue of the Dirichlet Laplacian:

$$
\begin{equation*}
\lambda_{1}(\Gamma ; \mathcal{V})=\min \left\{\int_{\Gamma}\left|u^{\prime}\right|^{2} d x: u \in H_{0}^{1}(\Gamma), \int_{\Gamma} u^{2} d x=1\right\} \tag{2.45}
\end{equation*}
$$

where $\Gamma$ is a metric graph and $\mathcal{V} \subset V(\Gamma)$ is a set of vertices, where a Dirichlet boundary conditions are imposed. Reasoning as in Remark 2.6, we have that among all connected metric graphs (shortly, CMG) of fixed length $l$ and with at least one Dirichlet vertex, the one with the lowest first eigenvalue is given by the segment $[0, l]$, with Dirichlet condition in 0 . Moreover, for any pair $D_{1}, D_{2} \in \mathbb{R}^{d}$ and any $l \geq\left|D_{1}-D_{2}\right|=: l-\epsilon$ the solution of

$$
\begin{equation*}
\min \left\{\lambda_{1}(\Gamma ; \mathcal{V}): \Gamma \in C M G, l(\Gamma)=l, \mathcal{V} \subset V(\Gamma), \exists \gamma: \Gamma \rightarrow \mathbb{R}^{d} \text { immersion, } \gamma(\mathcal{V})=\mathcal{D}\right\} \tag{2.46}
\end{equation*}
$$

is the graph described in Figure 2.4, i.e. the solution of (2.32) from Example 2.21. In the case when the set $\mathcal{D}$ is given by three points disposed in the vertices of an equilateral triangle, the solutions of (2.46) are quantitatively the same (see Figure 2.10 ) as the solutions of (2.36) from Example 2.25. In general, we have the following existence result

Theorem 2.28. Consider a set of distinct points $\mathcal{D}=\left\{D_{1}, \ldots, D_{k}\right\} \subset \mathbb{R}^{d}$ and a positive real number $l \geq S t(\mathcal{D})$. Then there exists a connected metric graph $\Gamma$, a set of vertices $\mathcal{V} \subset V(\Gamma)$ and an immersion $\gamma: \Gamma \rightarrow \mathbb{R}^{d}$ which are solution of the problem (2.46). Moreover, $\Gamma$ can be chosen to be a tree of at most $2 k$ vertices and $2 k-1$ edges.

Proof. The proof is identical to the one of Theorem 2.18.
Remark 2.29. The question of existence of an optimal graph is open for general cost functionals $J$ spectral type, i.e. $J=F\left(\lambda_{1}, \lambda_{2}, \ldots, \lambda_{k}, \ldots\right)$, where $F: \mathbb{R} \rightarrow \mathbb{R}$ is a real function and $\lambda_{k}$ is the $k$-th eigenvalue of the Dirichlet Laplacian:

$$
\begin{equation*}
\lambda_{k}(\Gamma ; \mathcal{V})=\min _{K \subset H_{0}^{1}(\Gamma)} \max \left\{\int_{\Gamma}\left|u^{\prime}\right|^{2} d x: u \in K, \int_{\Gamma} u^{2} d x=1\right\} \tag{2.47}
\end{equation*}
$$

where the minimum is over all $k$ dimensional subspaces $K$ of $H_{0}^{1}(\Gamma)$. In fact, the crucial point in the proof of Theorem 2.18 is the reduction to the class of connected metric trees with number of vertices bounded by some universal constant. This reduction becomes a rather involved question even for the simplest spectral functionals $\lambda_{k}$ for $k \geq 2$.

## Chapter 3

## A non-local isoperimetric problem

### 3.1 Introduction

This chapter is based on a work (in preparation), in collaboration with Michael Goldman and Matteo Novaga [70].

In a series of recent works, among which we cite [80], [81], [41] and [66], and references therein, are investigated functionals of the form

$$
\begin{equation*}
\mathcal{E}(E)=P(E)+\mathcal{N} \mathcal{L}(E) \tag{3.1}
\end{equation*}
$$

where $E \subset \mathbb{R}^{d}, P(\cdot)$ is the perimeter functional and $\mathcal{N} \mathcal{L}(\cdot)$ is a non-local term which usually is translation and rotation invariant, and repulsive. In [80] and [81] it is considered the case

$$
\begin{equation*}
\mathcal{N} \mathcal{L}_{Q}(E)=\int_{E \times E} \frac{Q^{2}}{|x-y|^{\alpha}} d x d y, \quad 0<\alpha<d, \tag{3.2}
\end{equation*}
$$

where $Q$ is a real, non-negative number. The study of functionals of this type under volume constraint is often motivated by physical reasons. For example, when the dimension $d$ is 3 (or greater) and $\alpha=d-2$, the functional $P+\mathcal{N} \mathcal{L}_{Q}$ is related to the question of which shapes can assume a conductive liquid drop once it is provided with a uniform charge $Q$.
Notice that if $Q=0$ we fall in the classical isoperimetric problem. Moreover an easy computation shows that if $Q_{n} \rightarrow 0$ as $n \rightarrow \infty$, then $\mathcal{N} \mathcal{L}_{Q_{n}} \Gamma$-converges, in the $L^{1}$ topology of the sets, to 0 . These observations suggest that if the charge $Q$ is small enough, the perimeter plays the leading role. In [80] and [81] it is indeed proved that for small charges the non-local term is actually irrelevant and thus the minimizers are balls. It is worth mentioning that the isoperimetric inequality in sharp quantitative
form (see [62], [58] or [40]),

$$
\frac{P(E)-P\left(B_{E}\right)}{P\left(B_{E}\right)} \geq c(d)\left(\frac{\left|E \Delta B_{E}\right|}{|E|}\right)^{2}
$$

where $B_{E}$ is a suitable ball of measure $|E|$, together with the regularity theory of the perimeter quasi-minimizers (see [92], [2]), plays a crucial role in the proof. Another remarkable (but not surprising) result is that if $Q$ is big enough, the roles of $P$ and $\mathcal{N} \mathcal{L}_{Q}$ are overturned. In particular, there is a $Q_{0}$ such that if $Q>Q_{0}$ a minimum cannot be connected and, since the non-local term is repulsive, it cannot exist (indeed two connected components would tend to have infinite mutual distance). Moreover, empirical laws suggest that given a body $E$, the charge $Q$ tends to distribute on the boundary $\partial E$. This brings to consider the case where the charge is a priori distributed on the boundary of the set:

$$
\begin{equation*}
\mathcal{N} \mathcal{L}_{Q}^{\prime}(E)=\int_{\partial E \times \partial E} \frac{Q^{2}}{|x-y|^{\alpha}} d x d y \tag{3.3}
\end{equation*}
$$

In [66] it is taken into account the aforementioned case and it is proved that if one prescribes the potential ${ }^{1}$

$$
v_{\alpha}(x)=\int_{\partial E} \frac{d y}{|x-y|^{\alpha}}
$$

to be constant on the boundary of $E$, then the ball is a local minimum, where the locality is settled among under small $C^{2, \beta}$ perturbations of the ball, for some $\beta>0$.
Motivated by these results, we consider the case where the total charge $Q$ is fixed, but it is free to move in $E$. More precisely we consider the problem of minimizing, under volume constraint, the quantity

$$
P(E)+\min _{\mu(E)=Q} \int_{\mathbb{R}^{d} \times \mathbb{R}^{d}} \frac{d \mu(x) d \mu(y)}{|x-y|^{\alpha}}=P(E)+\mathcal{Q}_{\alpha}(E)
$$

As we will see in Section 3.2, depending on the value of $\alpha$, the charge measure can distributes on the whole $E$ or just in its boundary $\partial E$.

A remarkable case among those where the charge distributes on the boundary of the set, is the harmonic one: $\alpha=d-2$.

## Organization and main results of the chapter

In Section 3.2 we set the main definitions and preliminary results needed in the rest of the chapter. In particular we analyse the Riesz potential related to a measure $\mu$ and a parameter $\alpha$ :

$$
\begin{equation*}
v_{\alpha}^{\mu}(x)=\int_{\mathbb{R}^{d}} \frac{d \mu(y)}{|x-y|^{\alpha}} \tag{3.4}
\end{equation*}
$$

[^3]and its associated potential energy
$$
\int_{\mathbb{R}^{d}} v_{\alpha}^{\mu}(x) d \mu(x)
$$

Many results of the section are well known in literature and are reported for the reader convenience.

Section 3.3 is devoted to the proof of Theorem 3.21 and Theorem 3.22, where the following problem

$$
\begin{equation*}
\min _{E \subset \Omega,}^{|E|=1},\left\{P(E)+\mathcal{Q}_{\alpha}(E)\right\} \tag{3.5}
\end{equation*}
$$

is considered. Precisely, in Theorem 3.21 we prove that problem (3.5) does not admit a minimizer when $\Omega=\mathbb{R}^{d}$. Then, in Theorem 3.22 , we deal with a bounded, regular domain $\Omega$. In this case we show that problem (3.5) can be decoupled into the isoperimetric problem and that of minimizing $\mathcal{Q}_{\alpha}$ on $\Omega$ (see equations (3.15) and (3.16)). This means that the problem is ill-posed and a minimum cannot exist in general, see Remark 3.23.

In Section 3.4 we prove that if we impose some regularity (in our case, the $\delta$-ball condition, see Definition 3.16, and the connectedness of the admissible sets) to the class of minimization, the minimum is achieved.

Eventually, in Section 3.5 we prove that the ball is a local minimum, i.e. it is the minimum among all sets whose boundary is a $C^{1,1}$ perturbation of the boundary of the ball itself.

### 3.2 The Riesz potential energy

In this section we study the main properties of the Riesz potential (3.4) and of the interaction energy (see Definition 3.1 below). Many of the results of this section are probably well known in literature. An interesting and useful guide to this subject may be the book [84].

In the following, given an open set $\Omega \subset \mathbb{R}^{d}$, we denote by $\mathcal{M}(\Omega)$ the set of all Borel measures with support in $\Omega$.

Definition 3.1. Let $d \geq 2$ and $\alpha>0$. For any Borel measures $\mu$ and $\nu$ with $\mu\left(\mathbb{R}^{d}\right)=$ $\nu\left(\mathbb{R}^{d}\right)=1$, we define the interaction energy between $\mu$ and $\nu$ by

$$
\mathcal{Q}_{\alpha}(\mu, \nu):=\int_{\mathbb{R}^{d} \times \mathbb{R}^{d}} \frac{d \mu(x) d \mu(y)}{|x-y|^{\alpha}}
$$

When $\mu=\nu$, we simply write $\mathcal{Q}_{\alpha}(\mu):=\mathcal{Q}_{\alpha}(\mu, \mu)$. When $\mu=f \mathcal{H}^{d}\left\llcorner E\right.$ and $\nu=g \mathcal{H}^{d}\llcorner E$, we denote $\mathcal{Q}_{\alpha}^{E}(\mu, \nu)=\mathcal{Q}_{\alpha}^{E}(f, g)$ (and when $f=g$ we denote it by $\mathcal{Q}_{\alpha}^{E}(f)$ ). When $\mu=f \mathcal{H}^{d-1}\left\llcorner\partial E\right.$ and $\nu=g \mathcal{H}^{d-1}\left\llcorner\partial E\right.$ we simply write $\mathcal{Q}_{\alpha}(\mu, \nu)=\mathcal{Q}_{\alpha}^{\partial E}(f, g)$ (and when $f=g$ we denote it by $\left.\mathcal{Q}_{\alpha}^{\partial E}(f)\right)$.

Definition 3.2. Let $d \geq 2$ and $\alpha>0$. Then for every Borel set $A$ we define the Riesz potential energy of $A$ by

$$
\begin{equation*}
\mathcal{Q}_{\alpha}(A):=\inf \left\{\mathcal{Q}_{\alpha}(\mu): \mu \in \mathcal{M}\left(\mathbb{R}^{d}\right), \mu(A)=1\right\} . \tag{3.6}
\end{equation*}
$$

Remark 3.3. Notice that by changing $\mu$ in $Q \mu$, for every charge $Q>0$ and every Borel set $A$, there holds

$$
Q^{2} \mathcal{Q}_{\alpha}(A)=\inf \left\{\mathcal{Q}_{\alpha}(\mu): \mu \in \mathcal{M}\left(\mathbb{R}^{d}\right), \mu(A)=Q\right\}
$$

Lemma 3.4. If $E$ is compact then the infimum in (3.6) is achieved.
Proof. The kernel $1 /|x|^{\alpha}$ is positive and lower semicontinuous. Thus the energy

$$
\int_{\bar{E} \times \bar{E}} \frac{d \mu(x) d \mu(y)}{|x-y|^{\alpha}}
$$

is semicontinuous for the weak* topology. Moreover the constraint in the definition of $\mathcal{Q}_{\alpha}$ is weakly* compact (since $E$ is bounded), whence it follows that a minimizer $\mu$ exists.

Remark 3.5. When the set is unbounded, there does not always exist an optimal measure $\mu$. It is for example possible to construct a set $E$ of finite volume with $\mathcal{Q}_{\alpha}(E)=$ 0 . Indeed, for $\alpha \in(0, d-1)$ and $\gamma \in\left(\frac{1}{d-1},+\infty\right)$, consider the set $E=\left\{\left(x, x^{\prime}\right) \in\right.$ $\mathbb{R} \times \mathbb{R}^{d-1}:\left|x^{\prime}\right| \leq 1$ and $\left.\left|x^{\prime}\right| \leq \frac{1}{|x|^{\gamma}}\right\}$ then $E$ has finite volume and taking $N$ balls of radius $r=N^{-\beta}$, at mutual distance $\ell=N^{\frac{\beta}{\gamma}-1}$ inside $E$ with charge $1 / N$ distributed uniformly on each ball, we get

$$
\mathcal{Q}_{\alpha}(E) \leq c\left(N \frac{1}{N^{2}} r^{-\alpha}+N^{2} \frac{1}{N^{2}} \ell^{-\alpha}\right) \leq C\left(N^{\alpha \beta-1}+N^{\left(1-\frac{\beta}{\gamma}\right) \alpha}\right)
$$

for suitable constants $c, C>0$, so that if $\frac{1}{d-1}<\gamma<\beta<\frac{1}{\alpha}$, we obtain $\mathcal{Q}_{\alpha}(E)=0$.
Definition 3.6. Given a non-negative Borel measure $\mu$ on $\mathbb{R}^{d}$ and $\alpha \in(0, d-1)$, we define the function

$$
v_{\alpha}^{\mu}(x)=\int_{\mathbb{R}^{d}} \frac{d \mu(y)}{|x-y|^{\alpha}}=\mu * k_{\alpha}(x)
$$

where $k_{\alpha}(x)=|x|^{-\alpha}$. We will sometime drop the dependence of $\mu$ and/or $\alpha$ in the definition of $v_{\alpha}^{\mu}$ and we will refer to it as potential.

Lemma 3.7. Let $E$ be a bounded Borel set and let $\mu$ be a minimizer of $\mathcal{Q}_{\alpha}(E)$. Then $v^{\mu}=\mathcal{Q}_{\alpha}(E)$ q.e. on $\operatorname{spt}(\mu)$.

Proof. It follows by the minimality of $\mu$, by computing the first variation of $\mathcal{Q}_{\alpha}(E)$.

The following lemma is proven in [84, Chapter $I]$.
Lemma 3.8. For every $\alpha>0$ and every measure $\mu$, the function $v_{\alpha}^{\mu}$ is lower semicontinuous

We recall another important result which will be exploited in Section 3.4 (for its proof we refer to [84, Theorem 1.15] or [86, Corollary 5.10]).

Theorem 3.9. For any signed measure $\mu$ and for any $\alpha \in(0, d)$, there holds,

$$
\mathcal{Q}_{\alpha}(\mu)=\int_{\mathbb{R}^{d}}\left(v_{\alpha / 2}^{\mu}(x)\right)^{2} d x
$$

and therefore,

$$
\mathcal{Q}_{\alpha}(\mu) \geq 0 .
$$

Moreover equality holds if and only if $\mu=0$.
Remark 3.10. A useful consequence of Theorem 3.9, is that the functional $\mathcal{Q}_{\alpha}(\cdot, \cdot)$ is a positive, bilinear operator on the product space of Borel measures on $\mathbb{R}^{d}, \mathcal{M}\left(\mathbb{R}^{d}\right) \times$ $\mathcal{M}\left(\mathbb{R}^{d}\right)$. In particular it satisfies the Cauchy-Schwarz inequality

$$
\begin{equation*}
\mathcal{Q}_{\alpha}(\mu, \nu) \leq \mathcal{Q}_{\alpha}(\mu, \mu)^{1 / 2} \mathcal{Q}(\nu, \nu)^{1 / 2} \tag{3.7}
\end{equation*}
$$

Some properties of the potential energy $\mathcal{Q}_{\alpha}$ are in order.
Lemma 3.11. For every bounded Borel set $E$ the measure minimizing $\mathcal{Q}_{\alpha}(E)$ is unique.
Proof. Let $\mu$ and $\nu$ be two minimizing measures. Then since

$$
\mathcal{Q}(\mu-\nu, \nu)=\int_{\bar{E}} d(\mu-\nu)(x) \int_{\bar{E}} \frac{d \nu(y)}{|x-y|^{\alpha}}=\int_{\bar{E}} v_{\alpha}^{\nu}(x) d(\mu-\nu)(x)=\mathcal{Q}_{\alpha}(E)(\mu(E)-\nu(E))=0,
$$

we find that

$$
\begin{aligned}
\mathcal{Q}_{\alpha}(\mu, \mu)=\mathcal{Q}_{\alpha}((\mu-\nu)+\nu,(\mu-\nu)+\nu) & =\mathcal{Q}_{\alpha}(\mu-\nu, \mu-\nu)+\mathcal{Q}_{\alpha}(\nu, \nu)-2 \mathcal{Q}_{\alpha}(\mu-\nu, \nu) \\
& =\mathcal{Q}_{\alpha}(\mu-\nu, \mu-\nu)+\mathcal{Q}_{\alpha}(\mu, \mu)
\end{aligned}
$$

from which $\mathcal{Q}_{\alpha}(\mu-\nu)=0$. Hence, by Theorem 3.9, $\mu=\nu$.
We will also make use of the following density result which is an adaptation of [84, Theorem 1.11, Lemma 1.2].

Proposition 3.12. Let $E$ be a smooth connected closed set of $\mathbb{R}^{d}$. Then for every $\alpha \in(0, d-2)$,

$$
\mathcal{Q}_{\alpha}(E)=\inf \left\{\mathcal{Q}_{\alpha}^{E}(f): \mu=f d x, f \in L^{\infty}(E), \int_{E} f d x=1\right\}
$$

Proof. Let $\mu$ be such that $\mu(\bar{E})=1, \operatorname{spt}(\mu) \subset \bar{E}$ and $\mathcal{Q}_{\alpha}(\mu)<+\infty$. Then for $\varepsilon>0$ consider the measure $\mu_{\varepsilon} d x$ where

$$
\mu_{\varepsilon}(x):=\frac{1}{\left|B_{\varepsilon}(x) \cap E\right|} \int_{B_{\varepsilon}(x)} d \mu(y)
$$

Since $\left\|\mu_{\varepsilon}\right\|_{L^{\infty}(E)} \leq \frac{1}{\min _{x \in E}\left|B_{\varepsilon}(x) \cap E\right|} \leq C \frac{1}{\varepsilon^{d}}$, we are left to prove that $\mathcal{Q}_{\alpha}^{E}\left(\mu_{\varepsilon}\right) \rightarrow \mathcal{Q}_{\alpha}(\mu)$. By Theorem 3.9, $\mathcal{Q}_{\alpha}^{E}\left(\mu_{\varepsilon}\right)=\int_{\mathbb{R}^{d}}\left(v_{\alpha / 2}^{\mu_{\varepsilon}}\right)^{2}(x) d x$. Let us show that for all $x \in \mathbb{R}^{d}$,

$$
v_{\alpha / 2}^{\mu_{\varepsilon}}(x) \leq C v_{\alpha / 2}^{\mu}(x) \quad \text { and } \quad \lim _{\varepsilon \rightarrow 0} v_{\alpha / 2}^{\mu_{\varepsilon}}(x)=v_{\alpha / 2}^{\mu}(x)
$$

from which we can conclude using the Dominated Convergence Theorem. Let us start by noticing that for $x \in \mathbb{R}^{d}$,

$$
\begin{aligned}
v_{\alpha / 2}^{\mu_{\varepsilon}}(x) & =\int_{E} \int_{E} \frac{1}{\left|B_{\varepsilon}(y) \cap E\right|} 1_{B_{\varepsilon}}(y-z) \frac{d \mu(z)}{|x-y|^{\alpha / 2}} d y \\
& =\int_{E}\left(\int_{B_{\varepsilon}(z) \cap E} \frac{1}{\left|B_{\varepsilon}(y) \cap E\right|} \frac{|x-z|^{\alpha / 2}}{|x-y|^{\alpha / 2}} d y\right) \frac{d \mu(z)}{|x-z|^{\alpha / 2}} \\
& \leq \int_{E}\left(\frac{C}{\varepsilon^{d}} \int_{B_{\varepsilon}(z)} \frac{|x-z|^{\alpha / 2}}{|x-y|^{\alpha / 2}} d y\right) \frac{d \mu(z)}{|x-z|^{\alpha / 2}}
\end{aligned}
$$

As in the proof of [84, Theorem 1.11], we observe that the function $(x, z) \rightarrow \frac{1}{\varepsilon^{d}} \int_{B_{\varepsilon}(z)} \frac{|x-z|^{\alpha / 2}}{|x-y|^{\alpha / 2}} d y$ is uniformly bounded in $(x, z, \varepsilon)$ so that $v_{\alpha / 2}^{\mu_{\varepsilon}}(x) \leq C v_{\alpha / 2}^{\mu}(x)$ for some universal constant $C>0$. Consider now a point $x \in \mathbb{R}^{d}$ such that $v_{\alpha / 2}^{\mu}<+\infty$ then for every $\delta>0$ we can find a ball $B_{\eta}(x)$ such that $v_{\alpha / 2}^{\mu^{\prime}}<\delta$ where $\mu^{\prime}=\mu\left\llcorner B_{\eta}(x)\right.$. By the previous computations, we then find $v_{\alpha / 2}^{\left(\mu^{\prime}\right)_{\varepsilon}}(x) \leq C \delta$. Since $\lim _{\varepsilon \rightarrow 0} v_{\alpha / 2}^{\left(\mu-\mu^{\prime}\right)_{\varepsilon}}(x)=v_{\alpha / 2}^{\mu-\mu^{\prime}}(x)$, there holds

$$
\begin{aligned}
v_{\alpha / 2}^{\mu}(x) & =v_{\alpha / 2}^{\mu^{\prime}}(x)+v_{\alpha / 2}^{\mu-\mu^{\prime}}(x) \leq \delta+\lim _{\varepsilon \rightarrow 0} v_{\alpha / 2}^{\left(\mu-\mu^{\prime}\right)_{\varepsilon}}(x) \\
& \leq(1+C) \delta+\varliminf_{\varepsilon \rightarrow 0} v_{\alpha / 2}^{\mu_{\varepsilon}}(x) \leq(1+C) \delta+\varlimsup_{\varepsilon \rightarrow 0} v_{\alpha / 2}^{\mu_{\varepsilon}}(x) \\
& \leq(1+C) \delta+\varlimsup_{\varepsilon \rightarrow 0} v_{\alpha / 2}^{\left(\mu^{\prime}\right)_{\varepsilon}}(x)+\varlimsup_{\varepsilon \rightarrow 0} v_{\alpha / 2}^{\left(\mu-\mu^{\prime}\right)_{\varepsilon}}(x) \\
& \leq 2(1+C) \delta+v_{\alpha / 2}^{\mu}(x)
\end{aligned}
$$

so that letting $\delta \rightarrow 0$ we find that $\lim _{\varepsilon \rightarrow 0} v_{\alpha / 2}^{\mu_{\varepsilon}}(x)=v_{\alpha / 2}^{\mu}(x)$ as claimed.

With the same ideas, the following proposition can be proven.

Proposition 3.13. Let $E$ be a smooth connected closed set of $\mathbb{R}^{d}$ then for every $\alpha \in$ $[d-2, d-1)$,

$$
\mathcal{Q}_{\alpha}(\partial E)=\inf \left\{\mathcal{Q}_{\alpha}^{\partial E}(f): \mu=f d \mathcal{H}^{d-1}, f \in L^{\infty}(\partial E), \int_{\partial E} f d \mathcal{H}^{d-1}=1\right\}
$$

Lemma 3.14. Let $\alpha \in(0, d-1)$. For every bounded open set $E$, the minimizer $\mu$ of $\mathcal{Q}_{\alpha}(E)$ satisfies:

- if $\alpha \leq d-2$ then $\operatorname{spt}(\mu) \subset \partial E$, and thus $\mathcal{Q}_{\alpha}(E)=\mathcal{Q}_{\alpha}(\partial E)$,
- if $\alpha>d-2$ then $\operatorname{spt}(\mu)=\bar{E}$.

Moreover, when $\alpha \geq d-2, v_{\alpha}^{\mu}=\mathcal{Q}_{\alpha}(E)$ on $\bar{E}$.
Proof. The case $\alpha \leq d-2$ can be found in [84, pages 132 and 162] but we give a proof for the reader's convenience. Let $k_{\alpha}=1 /|x|^{\alpha}$ then:

- $\Delta k_{\alpha}>0$ for $|x|>0$ and $\alpha>d-2$,
- $\Delta k_{d-2}=0$ for $|x|>0$,
- $\Delta k_{\alpha}<0$ for $\alpha<d-2$ in $\mathbb{R}^{d} \backslash \operatorname{spt}(\mu)$.

Thus for every connected set $E$ :

- $\Delta v_{\alpha}^{\mu}>0$ in $\mathbb{R}^{d} \backslash \operatorname{spt}(\mu)$ for $\alpha>d-2$,
. $-\Delta v_{d-2}^{\mu}=\mu$,
- $\Delta v_{\alpha}^{\mu}<0$ in $\mathbb{R}^{d} \backslash \operatorname{spt}(\mu)$ for $\alpha<d-2$.

Moreover, $\Delta v_{\mu}^{\alpha}<0$ in $\mathbb{R}^{d}$ for $\alpha<d-2$, that is $v_{\alpha}^{\mu}$ is strictly superharmonic. Since $v_{\alpha}^{\mu}=\mathcal{Q}_{\alpha}(E)$ in $\operatorname{spt}(\mu) \subset \bar{E}$ and $v_{\alpha}^{\mu} \geq \mathcal{Q}_{\alpha}(E)$ in $\bar{E}$,

- for $\alpha>d-2, v_{\alpha}^{\mu}=\mathcal{Q}_{\alpha}(E)$ in $\bar{E}$ and $\operatorname{spt}(\mu)=\bar{E}$ (since otherwise $v_{a}^{\mu}$ would have a maximum out of $\operatorname{spt}(\mu)$ which would contradict $\Delta v_{\alpha}^{\mu}>0$ in $\left.\mathbb{R}^{d} \backslash \operatorname{spt}(\mu)\right)$.
- for $\alpha=d-2, v_{d-2}^{\mu}=\mathcal{Q}_{\alpha}(E)$ in $\bar{E}$ since otherwise it would have a maximum in the open set $\bar{E} \backslash \operatorname{spt}(\mu)$ which would contradict the fact that $\Delta v_{d-2}^{\mu}=0$ in $\bar{E} \backslash \operatorname{spt}(\mu)$. Now since $0=-\Delta v_{d-2}^{\mu}=\mu$ in $\bar{E}$, it implies that $\operatorname{spt}(\mu) \subset \partial E$. Since $v_{d-2}^{\mu}=\mathcal{Q}_{\alpha}(E)$ in $\bar{E}, v_{d-2}^{\mu} \geq \mathcal{Q}_{\alpha}(E)$ in $\partial E$ and by lower semicontinuity of $v_{d-2}^{\mu}$, we find that $v_{d-2}^{\mu}=\mathcal{Q}_{\alpha}(E)$ on $\partial E$.
- for $\alpha<d-2, v_{\alpha}^{\mu}$ has necessarily a minimum in $\bar{E}$ on the boundary of $E$ and by superharmonicity, $v_{\alpha}^{\mu}>\mathcal{Q}_{\alpha}(E)$ in $E$ so that $\operatorname{spt}(\mu) \subset\left\{v_{\alpha}^{\mu}=\mathcal{Q}_{\alpha}(E)\right\} \subset \partial E$.

Lemma 3.15. Let $\alpha \in(0, d-2]$. Then the uniform measure on the sphere $\partial B$

$$
d \mathcal{U}_{B}=\frac{1}{P(B)} d \mathcal{H}^{d-1}\llcorner\partial B
$$

is the unique optimizer for $\mathcal{Q}_{\alpha}(\partial B)$. Similarly, for $d-1>\alpha>d-2$, the uniform measure on the ball $B$,

$$
d \tilde{\mathcal{U}}_{B}=\frac{1}{|B|} d \mathcal{H}^{d}\llcorner B
$$

is the unique optimizer for $\mathcal{Q}_{\alpha}(B)$.
Proof. Let $\mu_{B}$ be a minimum for the problem $\mathcal{Q}_{\alpha}(\partial B)$. By Lemma 3.14 we know that $v^{\mu_{B}}=\mathcal{Q}_{\alpha}(\partial B)$ on $\partial B$. Thus for every $\bar{x}, x \in \partial B$ we have

$$
\begin{aligned}
v_{B}(\bar{x}) & =\int_{\partial B} \frac{d \mu_{B}(y)}{|x-y|^{\alpha}}=\frac{1}{P(B)} \int_{\partial B} \int_{\partial B} \frac{d \mu_{B}(y) d \mathcal{H}^{d-1}(x)}{|x-y|^{\alpha}} \\
& =\int_{\partial B}\left(\frac{1}{P(B)} \int_{\partial B} \frac{d \mathcal{H}^{d-1}(x)}{|x-y|^{\alpha}}\right) d \mu_{B}(y) \\
& =\frac{1}{P(B)} \int_{\partial B} \frac{d \mathcal{H}^{d-1}(x)}{|x-y|^{\alpha}} .
\end{aligned}
$$

Hence

$$
\mathcal{Q}_{\alpha}(\partial B)=\int_{\partial B} v_{B}(\bar{x}) d \mu_{B}(\bar{x})=\frac{1}{P(B)} \int_{\partial B} \frac{d \mathcal{H}^{d-1}(x)}{|x-y|^{\alpha}}=\int_{\partial B \times \partial B} \frac{d \mathcal{U}_{B}(x) d \mathcal{U}_{B}(y)}{|x-y|^{\alpha}},
$$

which means that $\mathcal{U}_{B}$ is an optimizer for $\mathcal{Q}_{\alpha}(\partial B)$. The proof for the cases $\alpha \in[d-$ $2, d-1$ ) is analogous.

Definition 3.16. Given $\delta>0$, we say that $E$ satisfies the $\delta$-ball condition if for any $x \in \partial E$, there are two balls of radius $\delta$, one contained in $E$ and one in $E^{c}$, both tangent to $\partial E$ in $x$.

Remark 3.17. A set which satisfies the $\delta$-ball condition has $C^{1,1}$ boundary, see [49].
Lemma 3.18. Let $\delta>0$, then every set $E \in \mathcal{K}_{\delta}^{c o}$ with $|E|=m$ satisfies

$$
\operatorname{diam}(E) \leq \sqrt{d} 2^{d+2} \frac{m}{|B|} \delta^{1-d}
$$

Proof. Consider the tiling of $\mathbb{R}^{d}$ given by $\left[0,2 \delta\left[^{d}+2 \delta \mathbb{Z}^{d}\right.\right.$ and for $k \in \mathbb{Z}^{d}$ let $C_{k}=2 \delta k+$ $\left[0,2 \delta\left[^{d}\right.\right.$. For every $k \in \mathbb{Z}^{d}$ such that $C_{k} \cap E \neq \emptyset$, let $B_{\delta}\left(x_{k}\right)$ be a ball of radius $\delta$ such
that $B_{\delta}\left(x_{k}\right) \subset E$ and $B_{\delta}\left(x_{k}\right) \cap C_{k} \neq \emptyset$. the existence of such a ball is guaranteed by the $\delta$-ball condition. Any such a ball can intersect at most $2^{d}$ cubes $C_{j}$ so that

$$
\sharp\left\{k \in \mathbb{Z}^{d}: E \cap C_{k} \neq \emptyset\right\}=\frac{1}{\left|B_{\delta}\right|} \sum_{k: C_{k} \cap E \neq \emptyset}\left|B_{\delta}\left(x_{k}\right)\right| \leq \frac{2^{d}}{\left|B_{\delta}\right|}|E|,
$$

where $\sharp A$ stands for the cardinality of the set $A$. Since $E$ is connected we can assume that $E \subset\left[0,4 \delta \frac{2^{d}}{\left|B_{\delta}\right|} m\right]^{d}$, and so

$$
\operatorname{diam}(E) \leq \sqrt{d} 2^{d+2} \frac{m}{|B|} \delta^{1-d}
$$

Proposition 3.19. Let $d \geq 3, \alpha=d-2, \delta>0$ and $E \subset \mathbb{R}^{d}$ be a bounded set which satisfies the $\delta$-ball condition. Then the optimal measure $\mu$ for $\mathcal{Q}_{\alpha}(E)=\mathcal{Q}_{\alpha}(\partial E)$ can be written as $\mu=f \mathcal{H}^{d-1}\left\llcorner\partial E\right.$ with $\|f\|_{L^{\infty}(\partial E)} \leq \mathcal{Q}_{\alpha}(E)(d-2) \delta^{-1}$.

Proof. By Lemma 3.14 we know that an optimizer $\mu$ is concentrated on $\partial E$. Denote by $v=v_{d-2}^{\mu}$ the potential related to $\mu$ on $E$. By Lemma 3.14 , we get that $v=\mathcal{Q}_{\alpha}(E)$ on $E$, and that $-\Delta v=\mu$. By classical elliptic regularity (see for instance [67, Corollary 8.36]), $v$ is regular in $\mathbb{R}^{d} \backslash E$ (and at least $C^{1, \beta}$ up to the boundary). Consider now a point $x \in \partial E$ and let $y \in E$ be such that the ball $B_{\delta}(y)$ is contained in $E$ and is tangent to $\partial E$ in $x$. The existence of such an $y$ is guaranteed by the $\delta$-ball condition satisfied by $E$. Let $u$ be a function satisfying

$$
\Delta u=0 \quad \text { in } B_{\delta}^{c}(y) ; \quad u=v(x)=\mathcal{Q}_{\alpha}(E) \quad \text { on } \quad \partial B_{\delta}(y)
$$

Notice that $u(z)=\frac{\mathcal{Q}_{\alpha}(E) \delta^{d-2}}{|z-y|^{d-2}}$ out of $B_{\delta}(y)$. By the maximum principle for harmonic functions, $u \leq \mathcal{Q}_{\alpha}(E)$ on $\partial E$. Thus, again by the maximum principle, applied to $u-v$, we get that $v \geq u$ on $\mathbb{R}^{d} \backslash E$. Since $u(x)=v(x)$, we get

$$
\begin{equation*}
|\nabla v(x)| \leq|\nabla u(x)|=\mathcal{Q}_{\alpha}(E)(d-2) \delta^{-1} \tag{3.8}
\end{equation*}
$$

Let us prove that $\mu=|\nabla v| \mathcal{H}^{d-1}\llcorner\partial E$. For this, let $x \in \partial E$ and $r>0$ and consider a test function $\varphi \in C_{c}^{\infty}\left(B_{r}(x)\right)$. Then

$$
\begin{aligned}
\int_{\partial E \cap B_{r}(x)} \varphi d \mu & =-\int_{B_{r}(x)} \varphi \Delta v=\int_{B_{r}(x)}\langle\nabla \varphi, \nabla v\rangle d y \\
& =\int_{B_{r}(x) \cap E^{c}}\langle\nabla \varphi, \nabla v\rangle d y=\int_{\partial E} \varphi\left\langle\nabla v, \nu^{E}\right\rangle d \mathcal{H}^{d-1}
\end{aligned}
$$

where $\nu^{E}$ is the external normal to $E$. Since $v$ is constant on $\partial E$, its tangential derivative is zero; moreover, since $v<\mathcal{Q}_{\alpha}(E)$ on $\mathbb{R}^{d} \backslash \bar{E}$ we have that $\left\langle\nabla v, \nu^{E}\right\rangle \geq 0$.

Therefore, $\left\langle\nabla v, \nu^{E}\right\rangle=|\nabla v|$ on $\partial E$. Putting all this together, we find that for every test function $\varphi$,

$$
\int_{\partial E \cap B_{r}(x)} \varphi d \mu=\int_{\partial E} \varphi|\nabla v| d \mathcal{H}^{d-1}
$$

which is equivalent to the claim $\mu=|\nabla v| \mathcal{H}^{d-1}\llcorner\partial E$.

### 3.3 Relaxation of the problem

Definition 3.20. Let $d \geq 2$ and $\alpha>0$. For every $Q>0$ and every Borel set $E \subset \mathbb{R}^{d}$ we define the fuctionals,

$$
\begin{equation*}
\mathcal{F}(E)=P(E)+Q^{2} \mathcal{Q}_{\alpha}(E) \tag{3.9}
\end{equation*}
$$

and

$$
\begin{equation*}
\mathcal{G}(E)=P(E)+Q^{2} \mathcal{Q}_{\alpha}(\partial E) \tag{3.10}
\end{equation*}
$$

Notice that for $\alpha \in(0, d-2$ ], by Lemma 3.14, the functionals $\mathcal{F}$ and $\mathcal{G}$ coincide.

In this section we consider a closed, connected, regular set $\Omega \subset \mathbb{R}^{d}$ (not necessarily bounded) of measure $|\Omega|>m$ and address the following problems:

$$
\begin{equation*}
\inf _{|E|=m, E \subset \Omega} \mathcal{F}(E) \tag{3.11}
\end{equation*}
$$

and

$$
\begin{equation*}
\inf _{|E|=m, E \subset \Omega} \mathcal{G}(E) \tag{3.12}
\end{equation*}
$$

where the (implicit) parameter $\alpha$ belongs to ( $0, d-1$ ). Let us start by investigating the case $\Omega=\mathbb{R}^{d}$ (despite for simplicity we will consider only this case, the proof of the following theorem can be easily adapted to any unbounded regular set).

Theorem 3.21. For every $\alpha \in(0, d-1)$,

$$
\inf _{|E|=m} \mathcal{F}(E)=\inf _{|E|=m} \mathcal{G}(E)=\min _{|E|=m} P(E)=m^{d-1} P(B)
$$

In other words, (3.11) and (3.12) do not admit a minimizer.
Proof. Let $N \in \mathbb{N}$ and $\beta>0$ (to be fixed later). Consider $N$ balls of radius $N^{-\beta}$ infinitely far away from each other and put on each of these balls a charge $\frac{1}{N}$. Let $V_{N}:=N r^{d}|B|$ be their total volume and consider the set $E$ given by the union of these balls with a (non-charged) ball of volume $m-V_{N}$. We want to choose $\beta$ in such a way that

$$
\begin{equation*}
\lim _{N \rightarrow+\infty} N r^{d-1}=0 \quad \text { and } \quad \lim _{N \rightarrow+\infty} \frac{1}{N^{2}} \frac{1}{r^{\alpha}}=0 \tag{3.13}
\end{equation*}
$$

Indeed, with (3.13) in force, we would get that $V_{N} \rightarrow 0$ and
$m^{d-1} P(B) \leq P(E)+Q^{2} \mathcal{Q}_{\alpha}(E) \leq\left(m-V_{N}\right)^{d-1} P(B)+N r^{d-1}+\frac{Q^{2}}{N^{2}} \frac{C}{r^{\alpha}} \rightarrow m^{d-1} P(B)$.
Conditions (3.13) are satisfied as soon as $\frac{1}{d-1}<\beta<\frac{2}{\alpha}$. Such choice of $\beta$ is allowed since $0<\alpha<d-1$.

We now investigate what happens when $\Omega$ is bounded.
Theorem 3.22. Let $E_{0}$ be a solution of the constrained isoperimetric problem

$$
\begin{equation*}
\min \{P(E): E \subset \Omega,|E|=m\} \tag{3.14}
\end{equation*}
$$

where $\Omega$ is a regular set. Then, for $\alpha \in(d-2, d-1)$, we have

$$
\begin{equation*}
\inf _{|E|=m, E \subset \Omega} \mathcal{F}(E)=P\left(E_{0}\right)+\mathcal{Q}_{\alpha}(\Omega) . \tag{3.15}
\end{equation*}
$$

Similarly, for $\alpha \in(0, d-2]$,

$$
\begin{equation*}
\inf _{|E|=m, E \subset \Omega} \mathcal{G}(E)=P\left(E_{0}\right)+\mathcal{Q}_{\alpha}(\partial \Omega) . \tag{3.16}
\end{equation*}
$$

Proof. We only prove (3.15) since (3.16) can be proved exactly in the same way. We divide the proof into three steps.

Step 1. For $\varepsilon>0$ and $f \in L^{\infty}(A)$, with $f \geq 0$ and $\int_{A} f=1$, we shall construct a measure $\tilde{\mu}_{\varepsilon}$ with $\operatorname{spt}\left(\tilde{\mu}_{\varepsilon}\right) \subset A, \tilde{\mu}_{\varepsilon}(A)=1$, satisfying

$$
\begin{equation*}
P\left(\operatorname{spt}\left(\tilde{\mu}_{\varepsilon}\right)\right) \leq \varepsilon \tag{3.17}
\end{equation*}
$$

and

$$
\begin{equation*}
\mathcal{Q}_{\alpha}\left(\tilde{\mu}_{\varepsilon}\right) \leq \mathcal{Q}_{\alpha}^{A}(f)+\varepsilon \tag{3.18}
\end{equation*}
$$

Let $\delta>\lambda>0$ be small parameters to be fixed later and consider the tiling of the space given by $[0, \lambda)^{d}+\lambda \mathbb{Z}^{d}$. For every $k \in \mathbb{Z}^{d}$ such that $\left(\lambda k+[0, \lambda)^{d}\right) \cap A \neq \emptyset$, we let $C_{k}:=\lambda k+[0, \lambda)^{d}$ and let $x_{k}$ be the centre of $C_{k}$. Notice that the number $N$ of such squares $C_{k}$ is bounded by $C(A) \lambda^{-d}$. Letting $f_{k}:=\int_{C_{k}} f d x$, it holds

$$
\begin{align*}
\sum_{\left|x_{k}-x_{j}\right| \geq 2 \delta} \frac{f_{k} f_{j}}{\left|x_{k}-x_{j}\right|^{\alpha}} & =\sum_{\left|x_{k}-x_{j}\right| \geq 2 \delta} \int_{C_{k} \times C_{j}} \frac{f(x) f(y)}{|x-y|^{\alpha}} \frac{|x-y|^{\alpha}}{\left|x_{k}-x_{j}\right|^{\alpha}} d x d y \\
& \leq \sum_{\left|x_{k}-x_{j}\right| \geq 2 \delta} \int_{C_{k} \times C_{j}} \frac{f(x) f(y)}{|x-y|^{\alpha}} \frac{\left(\left|x_{k}-x_{j}\right|+2 \lambda\right)^{\alpha}}{\left|x_{k}-x_{j}\right|^{\alpha}} d x d y  \tag{3.19}\\
& \leq \sum_{\left|x_{k}-x_{j}\right| \geq 2 \delta} \int_{C_{k} \times C_{j}} \frac{f(x) f(y)}{|x-y|^{\alpha}}\left(1+C(\alpha) \frac{\lambda}{\delta}\right) d x d y \\
& \leq \mathcal{Q}_{\alpha}^{A}(f)\left(1+C(\alpha) \frac{\lambda}{\delta}\right) .
\end{align*}
$$

where we used the fact

$$
\sum_{\left|x_{k}-x_{j}\right| \geq 2 \delta} \int_{C_{k} \times C_{j}} \frac{f(x) f(y)}{|x-y|^{\alpha}} \leq \int_{A \times A} \frac{f(x) f(y)}{|x-y|^{\alpha}}=\mathcal{Q}_{\alpha}^{A}(f)<\infty
$$

Let now $r=(\lambda / 2)^{\beta}$, with $\beta>1$. If $\operatorname{dist}\left(x_{k}, \mathbb{R}^{d} \backslash A\right) \leq r$, we replace the point $x_{k}$ with a point $\tilde{x}_{k} \in C_{j(k)}$, with $\left|\tilde{x}_{k}-x_{j(k)}\right| \geq \lambda / 4$, where $C_{j(k)} \subset A$ is a cube adjacent to $C_{k}$. For simplicity of notation, we still denote $\tilde{x}_{k}$ by $x_{k}$. We consider $N$ balls of radius $r$ centred at the points $x_{k}$, and we set

$$
\tilde{\mu}_{\varepsilon}:=\sum_{k} \frac{f_{k}}{\left|B_{r}\right|} \chi_{B_{r}\left(x_{k}\right)}
$$

Notice that, by construction, $\operatorname{spt}\left(\tilde{\mu}_{\varepsilon}\right) \subset A$ and $\tilde{\mu}_{\varepsilon}(A)=\int_{A} f=1$. Then we get

$$
\begin{aligned}
\mathcal{Q}_{\alpha}\left(\tilde{\mu}_{\varepsilon}\right) & =\sum_{j, k} \frac{f_{k} f_{j}}{\left|B_{r}\right|^{2}} \int_{B_{r}\left(x_{j}\right) \times B_{r}\left(x_{k}\right)} \frac{d x d y}{|x-y|^{\alpha}} \\
& =\sum_{k} \frac{f_{k}^{2}}{\left|B_{r}\right|^{2}} \int_{B_{r}\left(x_{k}\right) \times B_{r}\left(x_{k}\right)} \frac{d x d y}{|x-y|^{\alpha}} \\
& +\sum_{\left|x_{j}-x_{k}\right|<2 \delta, k \neq j} \frac{f_{k} f_{j}}{\left|B_{r}\right|^{2}} \int_{B_{r}\left(x_{j}\right) \times B_{r}\left(x_{k}\right)} \frac{d x d y}{|x-y|^{\alpha}} \\
& +\sum_{\left|x_{j}-x_{k}\right| \geq 2 \delta} \frac{f_{k} f_{j}}{\left|B_{r}\right|^{2}} \int_{B_{r}\left(x_{j}\right) \times B_{r}\left(x_{k}\right)} \frac{d x d y}{|x-y|^{\alpha}} \\
& =I_{1}+I_{2}+I_{3} .
\end{aligned}
$$

Moreover we have

$$
\begin{equation*}
I_{1} \leq C N\|f\|_{L^{\infty}(A)}^{2}\left|C_{k}\right|^{2} \frac{1}{r^{\alpha}} \leq C\|f\|_{L^{\infty}(A)}^{2} \lambda^{d-\alpha \beta} \tag{3.20}
\end{equation*}
$$

and

$$
\begin{equation*}
I_{2} \leq C \delta^{d} N^{2}\|f\|_{L^{\infty}(A)}^{2}\left|C_{k}\right|^{2} \frac{1}{\lambda^{\alpha}} \leq C\|f\|_{L^{\infty}(A)}^{2} \frac{\delta^{d}}{\lambda^{\alpha}} \tag{3.21}
\end{equation*}
$$

Eventually, from (3.19) it follows

$$
\begin{align*}
I_{3} & =\sum_{\left|x_{j}-x_{k}\right| \geq 2 \delta} \frac{f_{k} f_{j}}{\left|x_{k}-x_{j}\right|^{\alpha}} \frac{1}{\left|B_{r}\right|^{2}} \int_{B_{r}\left(x_{j}\right) \times B_{r}\left(x_{k}\right)} \frac{\left|x_{k}-x_{j}\right|^{\alpha}}{|x-y|^{\alpha}} d x d y \\
& \leq \sum_{\left|x_{k}-x_{j}\right| \geq 2 \delta} \frac{f_{k} f_{j}}{\left|x_{k}-x_{j}\right|^{\alpha}}\left(1+C(\alpha) \frac{r}{\delta}\right)  \tag{3.22}\\
& \leq \mathcal{Q}_{\alpha}^{A}(f)\left(1+C(\alpha) \frac{\lambda}{\delta}\right)\left(1+C(\alpha) \frac{r}{\delta}\right) \\
& \leq \mathcal{Q}_{\alpha}^{A}(f)+C(\alpha) \mathcal{Q}_{\alpha}^{A}(f) \frac{\lambda}{\delta}
\end{align*}
$$

Letting $\lambda=\delta^{\gamma}$, from (3.20), (3.21), (3.22), we then get

$$
\mathcal{Q}_{\alpha}\left(\tilde{\mu}_{\varepsilon}\right)=I_{1}+I_{2}+I_{3} \leq \mathcal{Q}_{\alpha}^{A}(f)+C(\alpha) \mathcal{Q}_{\alpha}^{A}(f) \delta^{\gamma-1}+C\|f\|_{L^{\infty}(A)}^{2}\left(\delta^{\gamma(d-\alpha \beta)}+\delta^{d-\alpha \gamma}\right)
$$

Choosing $1<\beta<d / \alpha$ and $1<\gamma<d / \alpha$, for $\delta$ small enough we obtain (3.18).
Let us show that (3.17) holds as well. To this aim, we notice that

$$
\begin{equation*}
P\left(\operatorname{spt}\left(\tilde{\mu}_{\varepsilon}\right)\right) \leq C N r^{d-1}=C N \lambda^{\beta(d-1)}=C \lambda^{\beta(d-1)-d} \tag{3.23}
\end{equation*}
$$

Hence, for $\lambda$ small enough, (3.17) follows from (3.23) by choosing $d / \alpha>\beta>d /(d-1)$, which is possible since $\alpha<d-1$.

Step 2. Let now $E_{0}$ be a solution of the constrained isoperimetric problem (3.14), and let

$$
E_{\varepsilon}:=\left(E_{0} \cup \bigcup_{k} B_{r}\left(x_{k}\right)\right) \backslash B_{\eta} \quad \mu_{\varepsilon}:=\frac{\tilde{\mu}_{\varepsilon}\left\llcorner E_{\varepsilon}\right.}{1-\tilde{\mu}_{\varepsilon}\left(B_{\eta}\right)}
$$

where $B_{\eta} \subset E_{0}$ is a ball such that $\left|E_{\varepsilon}\right|=m$. Notice that $\operatorname{spt}\left(\mu_{\varepsilon}\right) \subset E_{\varepsilon}$ and $\mu_{\varepsilon}\left(E_{\varepsilon}\right)=1$. By (3.23) we have

$$
\left|B_{\eta}\right|^{\frac{d-1}{d}} \leq\left|\bigcup_{k} B_{r}\left(x_{k}\right)\right|^{\frac{d-1}{d}} \leq C P\left(\bigcup_{k} B_{r}\left(x_{k}\right)\right) \leq C \lambda^{\beta(d-1)-d}
$$

so that $\eta \leq C \lambda^{\beta-1}$. In particular, recalling (3.18), for $\lambda$ sufficiently small the measure $\mu_{\varepsilon}$ satisfies

$$
\begin{equation*}
\mathcal{Q}_{\alpha}\left(\mu_{\varepsilon}\right) \leq \mathcal{Q}_{\alpha}\left(\tilde{\mu}_{\varepsilon}\right)+\varepsilon \leq \mathcal{Q}_{\alpha}^{A}(f)+2 \varepsilon \tag{3.24}
\end{equation*}
$$

From (3.24) we then get

$$
\begin{equation*}
\varlimsup_{\varepsilon \rightarrow 0} P\left(E_{\varepsilon}\right)+Q^{2} \mathcal{Q}_{\alpha}\left(\mu_{\varepsilon}\right)=P\left(E_{0}\right)+Q^{2} \mathcal{Q}_{\alpha}^{A}(f) \tag{3.25}
\end{equation*}
$$

Step 3. By Proposition 3.12, for any $\varepsilon>0$ we can find a function $f \in L^{\infty}(A)$ such that $\int_{A} f=1$ and $\mathcal{Q}_{\alpha}^{A}(f) \leq \mathcal{Q}_{\alpha}(A)+\varepsilon$. Thus (3.15) follows by (3.25) and a diagonal argument.

Remark 3.23. An interpretation of Theorem 3.22 is that the two problems in the definition of $\mathcal{F}$, the isoperimetric one and the charge-minimizing one, are actually decoupled. A consequence of this fact is that the minimization problem

$$
\min \left\{P(E)+Q^{2} \int_{E \times E} \frac{d \mu(x) d \mu(y)}{|x-y|^{\alpha}}:|E|=m, E \subset \Omega \mu(E)=1\right\}
$$

is ill-posed. This is due by the fact that the perimeter is invariant under perturbation by sets of null $L^{1}$ volume, while the Energy functional is sensible to perturbations which change only the capacity of the set, see [84, Chapter 2].

### 3.4 Existence and characterization of minimizers under a regularity condition

In the previous section we have seen that we cannot hope to get existence for problem (3.11) without some further assumption on the class of minimization. In this section we investigate the existence question in the class $K_{\delta}$.
We consider the four problems

$$
\begin{align*}
& \min \left\{\mathcal{F}(E):|E|=m, E \in \mathcal{K}_{\delta}^{c o}\right\},  \tag{3.26}\\
& \min \left\{\mathcal{G}(E):|E|=m, E \in \mathcal{K}_{\delta}^{c o}\right\},  \tag{3.27}\\
& \min \left\{\mathcal{F}(E):|E|=m, E \in \mathcal{K}_{\delta}\right\}, \tag{3.28}
\end{align*}
$$

and

$$
\begin{equation*}
\min \left\{\mathcal{G}(E):|E|=m, E \in \mathcal{K}_{\delta}\right\} . \tag{3.2}
\end{equation*}
$$

Up to rescaling, we can assume that $|E|=|B|$ where $B$ is the unit ball. Indeed, for (3.26) for example, for every $E$ of volume $m$ and every $f$ minimizing $\mathcal{Q}_{\alpha}(\partial E)$, considering $\tilde{E}=\left(\frac{|B|}{m}\right)^{1 / d} E$ and $\tilde{f}(x)=\left(\frac{|B|}{m}\right)^{\frac{1}{d(d-1)}} f\left(\left(\frac{m}{|B|}\right)^{1 / d} x\right)$ we have $|\tilde{E}|=|B|$, $\int_{\partial \tilde{E}} \tilde{f}=1$ and

$$
P(E)+Q^{2} \mathcal{Q}_{\alpha}^{\partial E}(f)=\left(\frac{m}{|B|}\right)^{d-1 / d}\left(P(\tilde{E})+Q^{2}\left(\frac{|B|}{m}\right)^{\frac{d-1}{d}-\alpha} \mathcal{Q}_{\alpha}^{\partial \tilde{E}}(\tilde{f})\right) .
$$

Given a set $E$ with $|E|=|B|$, we let $\delta P(E):=P(E)-P(B)$ be the (non-rescaled) isoperimetric deficit.

Theorem 3.24. Problem (3.26) has solution.
Proof. Let $E_{n} \in \mathcal{K}_{\delta}^{c o}$ be a minimizing sequence. And let $\mu_{n}$ be the corresponding optimal measures for $\mathcal{Q}_{\alpha}\left(E_{n}\right)$. We can then assume that

$$
\delta P\left(E_{n}\right) \leq Q^{2} \mathcal{Q}_{\alpha}(B),
$$

therefore $P\left(E_{n}\right)$ is uniformly bounded. By Lemma 3.18, the sets $E_{n}$ are also uniformly bounded so that by $B V$ compactness, there exists a subsequence converging in $L^{1}$ to some $E$ with $|E|=m$. Similarly, up to subsequence, $\mu_{n}$ is weakly* converging to some probability measure $\mu$. Let us prove that $E_{n}$ converges to $E$ also in the Kuratowski convergence, or equivalently, in the Hausdorff metric (see for instance [4]). Namely we have to check the following two conditions:
(i) $x_{n} \rightarrow x, \quad x_{n} \in E_{n} \Rightarrow x \in E$;
(ii) $x \in E \Rightarrow \exists x_{n} \in E_{n}$ such that $x_{n} \rightarrow x$.

### 3.4 Existence and characterization of minimizers under a regularity condition

The second condition is an easy consequence of the $L^{1}$-convergence. To prove the first one, we notice that by the $\delta$-ball condition, up to choose a radius $r$ small enough, there exists a constant $c=c(d, \delta)>0$ such that $\left|B_{r}\left(x_{n}\right) \cap E_{n}\right| \geq c r^{d}$, which implies, together with the $L^{1}$-convergence, statement ( $i$ ). Similarly one can also prove the Hausdorff convergence of $\partial E_{n}$ to $\partial E$. Since the set $\mathcal{K}_{\delta}^{c o}$ is stable under Hausdorff convergence, $E \in \mathcal{K}_{\delta}^{c o}$.
Since the perimeter $P$ is lower semicontinuous under $L^{1}$ convergence and since $\mathcal{Q}_{\alpha}(\mu)$ is lower semicontinuous under weak* convergence,

$$
\varliminf_{n \rightarrow+\infty} P\left(E_{n}\right)+Q^{2} \mathcal{Q}_{\alpha}\left(\mu_{n}\right) \geq P(E)+Q^{2} \mathcal{Q}_{\alpha}(\mu),
$$

which concludes the proof.
Similarly, one can show the following theorem.
Theorem 3.25. Problem (3.27) has solution.
Corollary 3.26. There exists $Q_{0}$ such that for every $Q \leq Q_{0}$, problems (3.28) and (3.29) have a solution.

Proof. We will only consider (3.28), since (3.29) can be treated in a similar way. As noticed before, for every minimizing sequence $E_{n} \in \mathcal{K}_{\delta}$, we can assume that there holds

$$
\delta P\left(E_{n}\right) \leq Q^{2} \mathcal{Q}_{\alpha}(B),
$$

so that by the quantitative isoperimetric inequality [62], up to a translation we can assume that

$$
\left|B \Delta E_{n}\right|^{2} \leq C(d) Q^{2} \mathcal{Q}_{\alpha}(B)
$$

so that $\left|E_{n} \cap B^{c}\right| \leq C Q$ but since every connected component of $E_{n}$ has volume at least $\left|B_{\delta}\right|$ by the $\delta$-ball condition, we see that for $Q$ small enough, $E_{n}$ must be connected. Thus the existence of minimizers follows as in Theorem 3.24.

It is natural to expect that for large enough charge $Q$, it is more favorable to have two connected components rather than one, which leads to non-existence of minimizers in $\mathcal{K}_{\delta}$. Let us prove that it is indeed the case (at least for certain $\alpha$ ). We start with the following lemma.

Lemma 3.27. Let $\alpha>0$ and let $E$ be a bounded Borel set. Then

$$
\mathcal{Q}_{\alpha}(E) \geq \frac{1}{\operatorname{diam}(E)^{\alpha}} .
$$

This implies

$$
\begin{equation*}
\min \left\{P(E)+Q^{2} \mathcal{Q}_{\alpha}(E):|E|=|B|, E \in \mathcal{K}_{\delta}^{c o}\right\} \geq P(B)+\left(\sqrt{d} 2^{d+2}\right)^{-\alpha} Q^{2} \delta^{(d-1) \alpha} \tag{3.30}
\end{equation*}
$$

and similarly

$$
\begin{equation*}
\min \left\{P(E)+Q^{2} \mathcal{Q}_{\alpha}(\partial E):|E|=|B|, E \in \mathcal{K}_{\delta}^{c o}\right\} \geq P(B)+\left(\sqrt{d} 2^{d+2}\right)^{-\alpha} Q^{2} \delta^{(d-1) \alpha} . \tag{3.31}
\end{equation*}
$$

Proof. Let $\mu$ be any positive measure with support in $E$ such that $\mu(E)=1$. Then

$$
\mathcal{Q}_{\alpha}(E) \geq \int_{E \times E} \frac{d \mu(x) d \mu(y)}{|x-y|^{\alpha}} \geq \int_{E \times E} \frac{d \mu(x) d \mu(y)}{\operatorname{diam}(E)^{\alpha}}=\frac{1}{\operatorname{diam}(E)^{\alpha}} .
$$

An application of Lemma 3.18 and the isoperimetric inequality lead to (3.30) and to (3.31).

Theorem 3.28. Let $\alpha<2$. Then there exists $\delta_{0}>0$ such that for every $\delta \leq \delta_{0}$, there exists $Q_{0}(\delta)$ such that for every $Q \geq Q_{0}(\delta)$ problems (3.28) and (3.29) do not have a solution.

Proof. If there exists a minimizer, then it must be connected. Thus, in order to prove non-existence, it is enough to construct a competitor $E \in \mathcal{K}_{\delta}$ with energy less than

$$
P(B)+\left(\sqrt{d} 2^{d+2}\right)^{-\alpha} Q^{2} \delta^{(d-1) \alpha}
$$

which bounds from below the energy of any set in $\mathcal{K}_{\delta}^{c o}$. To this aim, consider the set $E$ given by $N=\frac{|B|}{\delta^{d}}$ balls of radius $\delta$. Up to increase their mutual distances, we can suppose that the Riesz potential energy of $E$ is only made of the self interactions of each ball with itself. Considering (3.28), we have

$$
P(E)+Q^{2} \mathcal{Q}_{\alpha}(E)=N \delta^{d-1} P(B)+N^{2} Q^{2} \mathcal{Q}_{\alpha}\left(B_{\delta}\right)=|B| P(B) \delta^{-1}+\frac{\mathcal{Q}_{\alpha}(B)}{|B|^{2}} Q^{2} \delta^{2 d-\alpha}
$$

We see that if $2 d-\alpha>(d-1) \alpha$, i.e. $2>\alpha$, then for $\delta$ small enough there holds $\frac{\mathcal{Q}_{\alpha}(B)}{|B|^{2}} \delta^{2 d-\alpha}<\left(\sqrt{d} 2^{d+2}\right)^{-\alpha} \delta^{(d-1) \alpha}$. Thus, for $Q$ large enough,

$$
P(E)+Q^{2} \mathcal{Q}_{\alpha}(E)<P(B)+\left(\sqrt{d} 2^{d+2}\right)^{-\alpha} Q^{2} \delta^{(d-1) \alpha} .
$$

The non-existence for problem (3.29) follows similarly.

Remark 3.29. Notice that in dimension 3, the previous theorem contains the Coulomb interaction potential case, $\alpha=1, d=3$.

### 3.5 Stability of the ball

In this section we prove that in the harmonic case $\alpha=d-2$, the ball is an optimizer for problem (3.11) among sets in the family of the nearly spherical sets belonging to $\mathcal{K}_{\delta}^{c o}$ introduced in Definition 3.16, that is, the sets which are a small $W^{1, \infty}$ perturbation of the ball and that satisfy the $\delta$-ball condition.

Let $E$ be such that $|E|=|B|$ and $\partial E$ can be written as a graph over $\partial B$. In polar coordinates we have

$$
E=\{R(x) x: R(x)=1+\varphi(x), x \in \partial B\}
$$

The condition $|E|=|B|$ then is equivalent to

$$
\int_{\partial B}\left((1+\varphi(x))^{d}-1\right) d \mathcal{H}^{d-1}(x)=0
$$

which implies that if $\|\varphi\|_{L^{\infty}(\partial B)}$ is small enough, then

$$
\begin{equation*}
\int_{\partial B} \varphi d \mathcal{H}^{d-1}=O\left(\|\varphi\|_{L^{2}(\partial B)}^{2}\right) \tag{3.32}
\end{equation*}
$$

Letting

$$
\bar{\varphi}=\frac{1}{|\partial B|} \int_{\partial B} \varphi d \mathcal{H}^{d-1}
$$

the Poincaré Inequality gives

$$
\begin{align*}
\int_{\partial B}|\nabla \varphi|^{2} d \mathcal{H}^{d-1} & \geq C \int_{\partial B}|\varphi-\bar{\varphi}|^{2} d \mathcal{H}^{d-1}=C \int_{\partial B} \varphi^{2} \mathcal{H}^{d-1}-\frac{C}{|\partial B|}\left(\int_{\partial B} \varphi\right)^{2} d \mathcal{H}^{d-1} \\
& =C \int_{\partial B} \varphi^{2} d \mathcal{H}^{d-1}-\frac{C}{4|\partial B|}\left(\int_{\partial B} \varphi^{2} d \mathcal{H}^{d-1}\right)^{2} \\
& \geq \frac{3}{4} C \int_{\partial B} \varphi^{2} d \mathcal{H}^{d-1} \tag{3.33}
\end{align*}
$$

as soon as

$$
\begin{equation*}
\int_{\partial B} \varphi^{2} d \mathcal{H}^{d-1} \leq 1 \tag{3.34}
\end{equation*}
$$

for some constant $C$ depending only on the dimension $d$.
Up to translation, we can also assume that the barycentre of $E$ coincides with that of $B$, say 0 . This implies that

$$
\begin{equation*}
\left|\int_{\partial B} x \varphi(x) d \mathcal{H}^{d-1}(x)\right|=O\left(\|\varphi\|_{L^{2}(\partial B)}^{2}\right) \tag{3.35}
\end{equation*}
$$

Lemma 3.30. Suppose that $\varphi: \partial B \rightarrow \mathbb{R}^{d}$ parametrizes $\partial E$ and satisfies (3.32), (3.33), (3.35) and (3.34). Then we have

$$
\begin{equation*}
\delta P(E) \geq c_{0} \int_{\partial B}|\nabla \varphi|^{2} d \mathcal{H}^{d-1} \geq c_{1} \int_{\partial B}|\varphi|^{2} d \mathcal{H}^{d-1}=c_{2}\left|\int_{\partial B} \varphi d \mathcal{H}^{d-1}\right| . \tag{3.36}
\end{equation*}
$$

Proof. The first inequality is quite a well known fact, and we refer to [60] or [40] for its proof. The second one is exactly (3.33), while the third one follows from (3.32).

A consequence of Lemma 3.30 is the following corollary.
Corollary 3.31. Suppose that $\partial E$ is parametrized on $\partial B$ by a function $\varphi$ which satisfies the hypotheses of Lemma 3.30. Then there exists a positive constant $C=C(d)$ such that

$$
\begin{equation*}
\left|\mathcal{Q}_{\alpha}^{\partial B}(\varphi)\right| \leq C \delta P(E), \tag{3.37}
\end{equation*}
$$

and, for any positive constant $\lambda$,

$$
\begin{equation*}
\left|\mathcal{Q}_{\alpha}^{\partial B}(\lambda, \varphi)\right| \leq C \lambda \delta P(E) . \tag{3.38}
\end{equation*}
$$

Proof. Inequality (3.38) is an immediate consequence of (3.36). Concerning the first one we have, by the Hölder inequality and the Fubini Theorem,

$$
\begin{aligned}
& \mathcal{Q}_{\alpha}^{\partial B}(\varphi)=\int_{\partial B \times \partial B} \frac{\varphi(x) \varphi(y)}{|x-y|^{\alpha}} d \mathcal{H}^{d-1}(x) d \mathcal{H}^{d-1}(y) \\
& \leq\left(\int_{\partial B \times \partial B} \frac{\varphi(x)^{2}}{|x-y|^{\alpha}} d \mathcal{H}^{d-1}(x) d \mathcal{H}^{d-1}(y)\right)^{1 / 2}\left(\int_{\partial B \times \partial B} \frac{\varphi(y)^{2}}{|x-y|^{\alpha}} d \mathcal{H}^{d-1}(x) d \mathcal{H}^{d-1}(y)\right)^{1 / 2} \\
& \leq C(\alpha) \int_{\partial B} \varphi(x)^{2} d \mathcal{H}^{d-1}(x) .
\end{aligned}
$$

So (3.37) follows again from (3.36).
The following (technical) lemma will be exploited in Proposition 3.35.
Lemma 3.32. Let $E=\{R(x) x: R(x)=1+\varphi(x), x \in \partial B\}$ and let $g \in L^{\infty}(\partial B)$, then there exists $\varepsilon_{0}(\alpha, d)$ and a constant $C=C(\alpha, d)>0$ such that if $\|\varphi\|_{W^{1, \infty}(\partial B)} \leq \varepsilon_{0}$, then

$$
\begin{align*}
& \left|\int_{\partial B \times \partial B}\left(\frac{1}{|R(x)-R(y)|^{\alpha}}-\frac{\left(1-\frac{\alpha}{2} \varphi(x)\right)\left(1-\frac{\alpha}{2} \varphi(y)\right)}{|x-y|^{\alpha}}\right) g(x) g(y) d \mathcal{H}^{d-1}(x) d \mathcal{H}^{d-1}(y)\right| \\
& \quad \leq C(\alpha, d)\|g\|_{L^{\infty}(\partial B)}^{2} \delta P(E) . \tag{3.39}
\end{align*}
$$

Proof. First, notice that since $|x|=|y|=1$, we get

$$
\begin{equation*}
|R(x) x-R(y) y|^{2}=|x-y|^{2}(1+\varphi(x)+\varphi(y)+\varphi(x) \varphi(y)+\psi(x, y)) \tag{3.40}
\end{equation*}
$$

where $\psi(x, y)=\frac{(\varphi(x)-\varphi(y))^{2}}{|x-y|^{2}}$. Thus, for any $x, y \in \partial B$,

$$
\begin{equation*}
|R(x) x-R(y) y|^{-\alpha}=\frac{\left(1-\frac{\alpha}{2} \varphi(x)\right)\left(1-\frac{\alpha}{2} \varphi(y)\right)+\frac{\alpha(4-\alpha)}{4} \varphi(x) \varphi(y)-\frac{\alpha}{2}(\psi(x, y)+\eta(x, y))}{|x-y|^{\alpha}} \tag{3.41}
\end{equation*}
$$

where

$$
\eta(x, y) \leq C\left(\varphi^{2}(x)+\varphi^{2}(y)+\psi^{2}(x, y)\right) .
$$

By (3.39) and (3.41) we get

$$
\begin{align*}
& \int_{\partial B \times \partial B}\left(\frac{1}{|R(x)-R(y)|^{\alpha}}-\frac{\left(1-\frac{\alpha}{2} \varphi(x)\right)\left(1-\frac{\alpha}{2} \varphi(y)\right)}{|x-y|^{\alpha}}\right) g(x) g(y) d \mathcal{H}^{d-1}(x) d \mathcal{H}^{d-1}(y) \\
& =\frac{\alpha(4-\alpha)}{4} \int_{\partial B \times \partial B} \frac{\varphi(x) \varphi(y)}{|x-y|^{\alpha}} g(x) g(y) d \mathcal{H}^{d-1}(x) d \mathcal{H}^{d-1}(y) \\
& -\frac{\alpha}{2} \int_{\partial B \times \partial B} \frac{\psi(x, y)+\eta(x, y)}{|x-y|^{\alpha}} g(x) g(y) d \mathcal{H}^{d-1}(x) d \mathcal{H}^{d-1}(y) . \tag{3.42}
\end{align*}
$$

By (3.37),

$$
\int_{\partial B \times \partial B} \frac{\varphi(x) \varphi(y)}{|x-y|^{\alpha}} d \mathcal{H}^{d-1}(x) d \mathcal{H}^{d-1}(y)=\mathcal{Q}_{\alpha}^{\partial B}(\varphi) \leq C \delta P(E) .
$$

Furthermore, since

$$
\psi(x, y) \leq C\|\nabla \varphi\|_{L^{\infty}(\partial B)}^{2}
$$

and $\eta(x, y) \leq C \varphi^{2}(x)+\varphi^{2}(y)+\psi(x, y)$, for a suitable $C$, we only have to check that

$$
\int_{\partial B \times \partial B} \frac{\psi(x, y)}{|x-y|^{\alpha}} d \mathcal{H}^{d-1}(x) d \mathcal{H}^{d-1}(y) \leq C \delta P(E)
$$

For $x, y$ in $\partial B$ let us denote by $\Gamma_{x, y}$ the geodesic going from $x$ to $y$ and by $\ell(x, y)$ the geodesic distance between $x$ and $y$ (that is the length of $\Gamma_{x, y}$ ). Notice that on $\partial B$, the euclidean distance and $\ell$ are equivalent so that proving the previous inequality is equivalent to prove that it holds

$$
\int_{\partial B \times \partial B} \ell(x, y)^{-(\alpha+2)}(\varphi(x)-\varphi(y))^{2} \leq C \delta P(E) .
$$

We have

$$
\begin{aligned}
\int_{\partial B \times \partial B} & \ell(x, y)^{-(\alpha+2)}(\varphi(x)-\varphi(y))^{2} \\
& \leq c(d) \int_{\partial B \times \partial B} \ell(x, y)^{-(\alpha+1)} \int_{\Gamma_{x, y}}|\nabla \varphi|^{2}(z) d z d \mathcal{H}^{d-1}(x) d \mathcal{H}^{d-1}(y) \\
& =c(d) \int_{\partial B} \int_{0}^{2 \pi} t^{-(\alpha+1)} t^{d-1}\left(\int_{\{\ell(x, z) \leq t\}}|\nabla \varphi|^{2}(z) d \mathcal{H}^{d-1}(z)\right) d t d \mathcal{H}^{d-1}(x) \\
& =c(d) \int_{0}^{2 \pi} t^{(d-1)-(\alpha+1)}\left(\int_{\partial B} \int_{\{\ell(x, z) \leq t\}}|\nabla \varphi|^{2}(z) d \mathcal{H}^{d-1}(x) d \mathcal{H}^{d-1}(z)\right) d t \\
& =c(d) \mathcal{H}^{d-2}\left(\mathbb{S}^{d-2}\right) \int_{0}^{2 \pi} t^{(d-1)-\alpha}\left(\int_{\partial B}|\nabla \varphi|^{2}(z) d \mathcal{H}^{d-1}(z)\right) d t \\
& =c(d) \mathcal{H}^{d-2}\left(\mathbb{S}^{d-2}\right) \int_{0}^{2 \pi} t^{(d-1)-\alpha} d t\left(\int_{\partial B}|\nabla \varphi|^{2}(z) d \mathcal{H}^{d-1}(z)\right) \\
& \leq C(d) \delta P(E)
\end{aligned}
$$

where $\mathbb{S}^{d-2}$ is the $(d-2)$-dimensional sphere and where we used that $\alpha<d-1$.
Before we pass to our main stability estimates, let us recall the following simple interpolation inequality
Lemma 3.33. For every $0 \leq p<q<r<+\infty$, there exists a constant $C(r, p, q)$ such that for every $\varphi \in H^{r}\left(\mathbb{R}^{d}\right)$, it holds

$$
\begin{equation*}
\|\varphi\|_{H^{q}} \leq C\left(\|\varphi\|_{H^{r}}\right)^{\frac{r-q}{r-p}}\left(\|\varphi\|_{H^{p}}\right)^{\frac{q-p}{r-p}} . \tag{3.43}
\end{equation*}
$$

Remark 3.34. In the previous lemma we adopted the notation $\|u\|_{H^{p}}:=\left\||\xi|^{p} \hat{u}\right\|_{L^{2}\left(\mathbb{R}^{d}\right)}$ and $H^{p}\left(\mathbb{R}^{d}\right):=\left\{u \in L^{2}\left(\mathbb{R}^{d}\right):\|u\|_{H^{p}}<\infty\right\}$. Such notation will be used only in the remaining part of this chapter.
Proof. Using Fourier transform, for every $\varphi \in H^{r}$ and every $\lambda>0$, there holds

$$
\begin{aligned}
\|\varphi\|_{H^{q}}^{2} & =\int_{\mathbb{R}^{d}}|\hat{\varphi}|^{2}|\xi|^{2 q} d \xi=\int_{|\xi| \leq \lambda}|\hat{\varphi}|^{2}|\xi|^{2 p}|\xi|^{2(q-p)} d \xi+\int_{|\xi| \geq \lambda}|\hat{\varphi}|^{2}|\xi|^{2 r}|\xi|^{2(q-r)} d \xi \\
& \leq \lambda^{2(q-p)}\|\varphi\|_{H^{p}}^{2}+\lambda^{-2(r-q)}\|\varphi\|_{H^{r}}^{2} .
\end{aligned}
$$

Optimizing in $\lambda$ yields (3.43).
Proposition 3.35. Let $\alpha \in(0, d-1)$ and $r>\alpha / 2$. For $\partial E=\{R(x) x: R(x)=1+$ $\varphi(x), x \in \partial B\}$ and $f \in L^{\infty}(\partial E)$, there exists $\varepsilon_{0}(\alpha, d)$ and a constant $C=C(\alpha, d)>0$ such that if $\|\varphi\|_{W^{1, \infty}(\partial B)} \leq \varepsilon_{0}$ and $\|\varphi\|_{H^{r}} \leq \varepsilon_{0}$ then

$$
\mathcal{Q}_{\alpha}^{\partial E}(f)-\mathcal{Q}_{\alpha}^{\partial B}(\bar{f}) \geq-C\|f\|_{L^{\infty}(\partial E)}^{2} \delta P(E)
$$

where $\bar{f}=\frac{1}{P(E)} \int_{\partial E} f d \mathcal{H}^{d-1}$.

Proof. We have

$$
\begin{align*}
& \int_{\partial E \times \partial E} \frac{f(x) f(y)}{|x-y|^{\alpha}} d \mathcal{H}^{d-1}(x) d \mathcal{H}^{d-1}(y) \\
& =\int_{\partial B \times \partial B} \frac{f(R(x) x) f(R(y) y) R(x)^{d-2} R(y)^{d-2} \sqrt{R(x)^{2}+|\nabla R(x)|^{2}} \sqrt{R(y)^{2}+|\nabla R(y)|^{2}}}{|R(x) x-R(y) y|^{\alpha}} \tag{3.44}
\end{align*}
$$

Let $g(x):=f(R(x) x) R(x)^{d-2} \sqrt{R(x)^{2}+|\nabla R(x)|^{2}}$, then (3.44) can be rewritten as

$$
\begin{equation*}
\mathcal{Q}_{\alpha}^{\partial E}(f)=\int_{\partial B \times \partial B} \frac{g(x) g(y)}{|R(x)-R(y)|^{\alpha}} d \mathcal{H}^{d-1}(x) d \mathcal{H}^{d-1}(y) . \tag{3.45}
\end{equation*}
$$

We shall assume that $E$ is close enough to $B$ so that

$$
\begin{equation*}
\|g\|_{L^{\infty}(\partial B)} \leq 2\|f\|_{L^{\infty}(\partial E)} . \tag{3.46}
\end{equation*}
$$

A simple computation shows that there is a positive constant $C=C(\alpha, d)$ such that letting $\bar{g}:=\frac{1}{P(B)} \int_{\partial B} g d \mathcal{H}^{d-1}$, there holds

$$
\begin{equation*}
\mathcal{Q}_{\alpha}^{\partial B}(\bar{g})-\mathcal{Q}_{\alpha}^{\partial E}(\bar{f}) \leq C\|f\|_{L^{\infty}(\partial E)^{2}}^{2} \delta(E), \tag{3.47}
\end{equation*}
$$

so, up to sum and subtract $\mathcal{Q}_{\alpha}^{\partial E}(\bar{f})$, thanks to (3.46) and (3.47), the proposition is proven if we can show that

$$
\mathcal{Q}_{\alpha}^{\partial E}(f) \geq \mathcal{Q}_{\alpha}^{\partial B}(\bar{g})-\|g\|_{L^{\infty}(\partial E)}^{2} \delta P(E) .
$$

By Lemma 3.32, we obtain

$$
\mathcal{Q}_{\alpha}^{\partial E}(f)=\mathcal{Q}_{\alpha}^{\partial B}\left(g\left(1-\frac{\alpha}{2} \varphi\right)\right)+\mathcal{R}(f, \varphi)
$$

with

$$
|\mathcal{R}(f, \varphi)| \leq c\|f\|_{L^{\infty}(\partial E)}^{2} \delta P(E)
$$

Hence we get

$$
\begin{equation*}
\mathcal{Q}_{\alpha}^{\partial E}(f) \geq \mathcal{Q}_{\alpha}^{\partial B}\left(g\left(1-\frac{\alpha}{2} \varphi\right)\right)-c\|f\|_{L^{\infty}(\partial E)}^{2} \delta P(E) \tag{3.48}
\end{equation*}
$$

We now estimate

$$
\begin{align*}
\mathcal{Q}_{\alpha}^{\partial B}\left(g\left(1-\frac{\alpha}{2} \varphi\right)\right)= & \mathcal{Q}_{\alpha}^{\partial B}\left(g\left(1-\frac{\alpha}{2} \varphi\right), g\left(1-\frac{\alpha}{2} \varphi\right)\right) \\
= & \mathcal{Q}_{\alpha}^{\partial B}(g, g)-\alpha \mathcal{Q}_{\alpha}^{\partial B}(g, g \varphi)+\frac{\alpha^{2}}{4} \mathcal{Q}_{\alpha}^{\partial B}(g \varphi, g \varphi) \\
= & \mathcal{Q}_{\alpha}^{\partial B}(\bar{g}, \bar{g})+\mathcal{Q}_{\alpha}^{\partial B}(g-\bar{g}, g-\bar{g})-\alpha \mathcal{Q}_{\alpha}^{\partial B}(g-\bar{g}, g \varphi)-\alpha \mathcal{Q}_{\alpha}^{\partial B}(\bar{g}, g \varphi) \\
& +\frac{\alpha^{2}}{4} \mathcal{Q}_{\alpha}^{\partial B}(\bar{g} \varphi, \bar{g} \varphi)+\frac{\alpha^{2}}{2} \mathcal{Q}_{\alpha}^{\partial B}(\bar{g} \varphi,(g-\bar{g}) \varphi)+\frac{\alpha^{2}}{4} \mathcal{Q}_{\alpha}^{\partial B}((g-\bar{g}) \varphi,(g-\bar{g}) \varphi) \\
= & \mathcal{Q}_{\alpha}^{\partial B}(\bar{g})+\mathcal{Q}_{\alpha}^{\partial B}(g-\bar{g})+\frac{\alpha^{2}}{4} \mathcal{Q}_{\alpha}^{\partial B}((g-\bar{g}) \varphi)-\alpha \mathcal{Q}_{\alpha}^{\partial B}(g-\bar{g},(g-\bar{g}) \varphi) \\
& -\alpha \mathcal{Q}_{\alpha}^{\partial B}(\bar{g},(g-\bar{g}) \varphi)-\alpha \mathcal{Q}_{\alpha}^{\partial B}(g-\bar{g}, \bar{g} \varphi)+\frac{\alpha^{2}}{2} \mathcal{Q}_{\alpha}^{\partial B}(\bar{g} \varphi,(g-\bar{g}) \varphi) \\
& -\alpha \mathcal{Q}_{\alpha}^{\partial B}(\bar{g}, \bar{g} \varphi)+\frac{\alpha^{2}}{4} \mathcal{Q}_{\alpha}^{\partial B}(\bar{g} \varphi) . \tag{3.49}
\end{align*}
$$

The last two terms in the right hand side of (3.49), in view of Lemma 3.30 and the bilinearity of $\mathcal{Q}_{\alpha}^{\partial B}$, satisfy:

$$
\begin{equation*}
-\mathcal{Q}_{\alpha}^{\partial B}(\bar{g}, \bar{g} \varphi)+\frac{\alpha}{2} \mathcal{Q}_{\alpha}^{\partial B}(\bar{g} \varphi) \geq-c \bar{g}^{2} \delta P(E) \tag{3.50}
\end{equation*}
$$

By inequality (3.7) and Young's inequality, we get that for any functions $h_{1}$ and $h_{2}$ we have

$$
\begin{equation*}
\mathcal{Q}_{\alpha}^{\partial B}\left(h_{1}, h_{2}\right) \leq \mathcal{Q}_{\alpha}^{\partial B}\left(h_{1}, h_{1}\right)^{\frac{1}{2}} \mathcal{Q}_{\alpha}^{\partial B}\left(h_{2}, h_{2}\right)^{\frac{1}{2}} \leq \varepsilon \mathcal{Q}_{\alpha}^{\partial B}\left(h_{1}, h_{1}\right)+\frac{1}{4 \varepsilon} \mathcal{Q}_{\alpha}^{\partial B}\left(h_{2}, h_{2}\right) \tag{3.51}
\end{equation*}
$$

for any $\varepsilon>0$. In particular, applying such inequality to the functions $h_{1}=g-\bar{g}$ and $h_{2}=(g-\bar{g}) \varphi$ on the fourth term in the right hand side of (3.49), then on the sixth one, with $h_{1}=g-\bar{g}$ and $h_{2}=\bar{g} \varphi$ and exploiting (3.50), we get that there exists a positive constant $C$ such that

$$
\begin{align*}
& \mathcal{Q}_{\alpha}^{\partial B}\left(g\left(1-\frac{\alpha}{2} \varphi\right)\right)-\mathcal{Q}_{\alpha}^{\partial B}(\bar{g}) \\
& \geq C\left(\frac{1}{2} \mathcal{Q}_{\alpha}^{\partial B}(g-\bar{g})-\mathcal{Q}_{\alpha}^{\partial B}(\bar{g},(g-\bar{g}) \varphi)-\mathcal{Q}_{\alpha}^{\partial B}((g-\bar{g}) \varphi)-\bar{g}^{2} \delta P(E)\right) \tag{3.52}
\end{align*}
$$

Again by Lemma 3.30, we have that

$$
-\mathcal{Q}_{\alpha}^{\partial B}((g-\bar{g}) \varphi) \geq-\|g\|_{L^{\infty}(\partial B)}^{2} \mathcal{Q}_{\alpha}^{\partial B}(\varphi) \geq-C\|g\|_{L^{\infty}(\partial B)}^{2} \delta P(E)
$$

Let us show that the term $\mathcal{Q}_{\alpha}^{\partial B}(\bar{g},(g-\bar{g}) \varphi)$ can be estimated by the term $\mathcal{Q}_{\alpha}^{\partial B}(g-\bar{g})$. Let $\widetilde{\varphi}: \mathbb{R}^{d} \rightarrow \mathbb{R}$ be a regular extension of $\varphi$, and let $g=(g-\bar{g}) d \mathcal{H}^{d-1}\llcorner\partial B$. By a

Fourier transform we get

$$
\begin{aligned}
\mathcal{Q}_{\alpha}^{\partial B}(\bar{g},(g-\bar{g}) \varphi) & =\bar{g} \int_{\partial B}(g-\bar{g}) \varphi=\bar{g} \int_{\mathbb{R}^{d}} \widehat{\widetilde{\varphi}} \widehat{g} \\
& \leq \bar{g}\left(\int_{\mathbb{R}^{d}} \widehat{\widetilde{\varphi}}^{2}|\xi|^{d-\alpha}\right)^{\frac{1}{2}}\left(\int_{\mathbb{R}^{d}} \frac{\widehat{g}^{2}}{|\xi|^{d-\alpha}}\right)^{\frac{1}{2}} \\
& =\bar{g}\|\widetilde{\varphi}\|_{H^{\frac{d-\alpha}{2}}} \mathcal{Q}_{\alpha}^{\partial B}(g-\bar{g}, g-\bar{g})^{\frac{1}{2}} \\
& \leq C(\alpha, d) \bar{g}\|\varphi\|_{H^{\frac{d-\alpha}{2}}} \mathcal{Q}_{\alpha}^{\partial B}(g-\bar{g})^{\frac{1}{2}} .
\end{aligned}
$$

We remark now that, if

$$
\begin{equation*}
\mathcal{Q}_{\alpha}^{\partial B}(\bar{g},(g-\bar{g}) \varphi) \leq \frac{1}{2} \mathcal{Q}_{\alpha}^{\partial B}(g-\bar{g}), \tag{3.53}
\end{equation*}
$$

then we would get

$$
\mathcal{Q}_{\alpha}^{\partial B}\left(g\left(1-\frac{\alpha}{2} \varphi\right)\right)-\mathcal{Q}_{\alpha}^{\partial B}(\bar{g}) \geq-C\|\bar{g}\|_{L^{\infty}(\partial B)}^{2} \delta P(E)
$$

which would conclude the proof. On the other hand if (3.53) does not hold, then

$$
\mathcal{Q}_{\alpha}^{\partial B}(g-\bar{g})<C(\alpha, d) \bar{g}\|\varphi\|_{H^{\frac{d-\alpha}{2}}} \mathcal{Q}_{\alpha}^{\partial B}(g-\bar{g})^{\frac{1}{2}}
$$

which implies

$$
\mathcal{Q}_{\alpha}^{\partial B}(g-\bar{g})^{\frac{1}{2}}<C \bar{g}\|\varphi\|_{H^{\frac{d-\alpha}{2}}},
$$

so that

$$
\mathcal{Q}_{\alpha}^{\partial B}(\bar{g},(g-\bar{g}) \varphi) \leq C \bar{g}\|\widetilde{\varphi}\|_{H^{\frac{d-\alpha}{2}}} \mathcal{Q}_{\alpha}^{\partial B}(g-\bar{g})^{\frac{1}{2}} \leq C \bar{g}^{2}\|\varphi\|_{H^{\frac{d-\alpha}{2}}}^{2}
$$

If $\frac{d-\alpha}{2} \leq 1$ then using (3.43) with $p=0, q=\frac{d-\alpha}{2}$ and $r=1$, up to extend again $\varphi$ to a regular function on $\mathbb{R}^{d}$, we obtain

$$
\|\varphi\|_{H^{\frac{d-\alpha}{2}}}^{2} \leq c\left(\|\varphi\|_{H^{1}}^{2}\right)^{1-\alpha / 2}\left(\|\varphi\|_{L^{2}\left(\mathbb{R}^{d}\right)}^{2}\right)^{\frac{d-\alpha}{2}} \leq c\|\varphi\|_{H^{1}}^{2}+\|\varphi\|_{L^{2}\left(\mathbb{R}^{d}\right)}^{2} \leq C \delta P(E) .
$$

Otherwise, we use (3.43) with $p=1, q=\frac{d-\alpha}{2}$ and $r$ to obtain

$$
\|\varphi\|_{H^{\frac{d-\alpha}{2}}}^{2} \leq c\left(\|\varphi\|_{H^{1}}^{2}\right)^{\frac{\alpha / 2-1}{r-1}}\left(\|\varphi\|_{H^{r}}^{2}\right)^{\frac{r-\alpha / 2}{r-1}} \leq C \delta P(E)
$$

which concludes the proof.
Corollary 3.36. Let $d \geq 3$. Then for any $\delta>0$ there exists a $Q_{\delta}>0$ such that if $Q<Q_{\delta}$ then the ball is the unique minimizer of problem (3.10) with $\alpha=d-2$ among the sets in $\mathcal{K}_{\delta}$ with charge $Q$.

Proof. Let $Q>0$ and let $E_{Q}$ be a minimizer in $\mathcal{K}_{\delta}$. Since $\left|E_{Q} \Delta B\right|^{2} \leq C \delta P\left(E_{Q}\right) \leq$ $C Q^{2} \mathcal{Q}_{\alpha}(B), E_{Q}$ converges in $L^{1}$ to $B$ as $Q \rightarrow 0$. As before, thanks to the $\delta$-ball condition, there is also convergence in the sense of Hausdorff of $E_{Q}$ and $\partial E_{Q}$. By the $\delta$-ball condition again and the Hausdorff convergence of the boundaries, for $Q$ small enough, $\partial E_{Q}$ is a graph over $\partial B$ of some $C^{1,1}$ function. Moreover, by the $\delta$-ball condition we have that $\left\|\kappa_{E_{Q}}\right\|_{L^{\infty}\left(\partial E_{Q}\right)} \leq \frac{1}{\delta}$, where $\kappa_{A}(x)$ is the mean curvature of the boundary of the set $A$ in $x$, so that by classical Elliptic Regularity Theory (see for instance [67]), $\partial E_{Q}$ is $C^{1, \beta}$ with uniform $C^{1, \beta}$ bound depending only on $\delta$. From this we see that if $\partial E_{Q}=\left\{\left(1+\varphi_{Q}(x)\right) x: x \in \partial B\right\}$ then $\left\|\varphi_{Q}\right\|_{C^{1, \beta}(\partial B)}$ is converging to 0 . We can thus assume that $\varphi_{Q}$ satisfies the conditions of Proposition 3.35.
For simplicity of notations we drop the index $Q$ in the rest of the proof. Let $\mu=$ $f d \mathcal{H}^{d-1}\left\llcorner\partial E\right.$ be the minimizer of $\mathcal{Q}_{\alpha}(E)$. Since $\mathcal{Q}_{\alpha}(E) \leq P(B)+Q^{2} \mathcal{Q}_{\alpha}(B)$, by Proposition 3.19, $\|f\|_{L^{\infty}(\partial E)} \leq(d-2) \delta^{-1}\left(P(B)+Q^{2} \mathcal{Q}_{\alpha}(B)\right)$. On the other hand, since $\int_{\partial E} f=1$, we have

$$
\bar{f}=\frac{1}{P(E)}
$$

Recall that by Lemma 3.15, the optimal measure for $\mathcal{Q}_{\alpha}(B)$ is given by $\frac{\mathcal{H}^{d-1} L \partial B}{P(B)}$. By the minimality of $E$ we then have

$$
\begin{aligned}
\delta P(E) & =P(E)-P(B) \leq Q^{2}\left(\mathcal{Q}_{\alpha}(B)-\mathcal{Q}_{\alpha}(E)\right) \\
& =Q^{2}\left(\mathcal{Q}_{\alpha}^{\partial B}(\bar{f})-\mathcal{Q}_{\alpha}^{\partial E}(f)+\mathcal{Q}_{\alpha}^{\partial B}(1 / P(B))-\mathcal{Q}_{\alpha}^{\partial B}(1 / P(E))\right)
\end{aligned}
$$

A simple computation shows that

$$
\mathcal{Q}_{\alpha}^{\partial B}(1 / P(B))-\mathcal{Q}_{\alpha}^{\partial B}(1 / P(E)) \leq C^{2} \delta P(E)
$$

for a suitable positive constant $C=C(\alpha, d)$. So, by Proposition 3.35 we have that

$$
\delta P(E) \leq C Q^{2} \delta P(E)\left(1+\|f\|_{L^{\infty}(\partial E)}^{2}\right) \leq C Q^{2} \delta P(E)
$$

which implies that $E=B$ for $Q$ small enough.

## Part II

Quantitative stability problems

## Chapter 4

## Weighted isoperimetric inequalities in quantitative form

### 4.1 Introduction

This chapter is devoted to the study of two classes of weighted perimeters, and their related isoperimetric problems. Section 4.2 is based on the joint work with Lorenzo Brasco and Guido De Philippis [16].

We recall that the perimeter, at least for Lipschitz sets is simply the $d-1$ dimensional measure of the topological boundary:

$$
P(E)=\int_{\partial E} d \mathcal{H}^{d-1}(x)
$$

To introduce the issue of this chapter, we recall (once again) the isoperimetric inequality

$$
P(E) \geq P(B) \quad \text { for any ball } B \text { of measure }|E|
$$

Exploiting the rescaling law of the perimeter $P(t E)=t^{d-1} P(E)$ and of the $d$-dimensional Lebesgue measure $|t E|=t^{d}|E|$, the isoperimetric inequality can be stated avoiding to prescribe the measure of the sets $E$ and $B$ :

$$
|E|^{\frac{1-d}{d}} P(E) \geq d \omega_{d}^{1 / d}
$$

where $\omega_{d}$ is the Lebesgue measure of the unit ball in $\mathbb{R}^{d}$. We remark that, also if the Lebesgue measure is a priori fixed, the family of minimizer is not unique. This is clearly due to the invariance under translation of the perimeter. We introduce now a notion of perimeter which in general does not share with the classical one such property: the weighted perimeter.

Definition 4.1 (Weighted perimeter). Let $V: \mathbb{R}^{d} \rightarrow[0, \infty)$ be a non-negative Borel function. Then for every open bounded Lipschitz set $E \subset \mathbb{R}^{d}$, its weighted perimeter is
given by

$$
P_{V}(E)=\int_{\partial E} V(x) d \mathcal{H}^{d-1}(x)
$$

With this definition in force, we can write down the main problem we shall investigate in this chapter: given a positive constant $c$, find (if there) a solution to the problem

$$
\begin{equation*}
\min \left\{P_{V}(E):|E|=c\right\} \tag{4.1}
\end{equation*}
$$

We stress that the study of (4.1) is slightly related to a family of problems about weighted isoperimetric inequalities, as for instance those considered in [65]. However the issue we study presents a fundamental difference: the measure constraint and the weight assigned to the perimeter are not homogeneous! In order to study problem (4.1), we impose the weight $V$ to satisfy some conditions. A natural hypothesis for $V$ is to be a radial function, that is $V(x)=v(|x|)$ for some $v:[0, \infty) \rightarrow[0, \infty]$. Clearly with this sole requirement we cannot even expect to get existence for problem (4.1). Under further assumptions on $V$, in [9] it has been proved that the ball centred at the origin is the unique minimizer, or equivalently, that

$$
\begin{equation*}
P_{V}(E) \geq P_{V}\left(B_{r_{E}}(0)\right) \tag{4.2}
\end{equation*}
$$

where $r_{E}$ is such that $|E|=\left|B_{r_{E}}\right|$. In Section 4.2, Theorem 4.2, we obtain an improvement of the same result: a quantitative stability version of the inequality. More precisely we prove that, under suitable assumptions on the weight $V$, the inequality

$$
\begin{equation*}
P_{V}(E) \geq P_{V}\left(B_{r_{E}}(0)\right)+c_{d, V,|E|}\left(\frac{\left|E \Delta B_{r_{E}}(0)\right|}{|E|}\right)^{2} \tag{4.3}
\end{equation*}
$$

holds true, where $c_{d, V,|E|}$ is a constant. Clearly such result entails the weighted isoperimetric inequality (4.2). Moreover, in Subsection 4.2 .1 we shall prove that such inequality is sharp, in the sense that the exponent 2 cannot be substituted by any lower exponent.

Our proof of (4.3) (and thus of (4.2)), based on a sort of calibration method, is completely different from the one in [9], which is based on symmetrization techniques. Even if the hypotheses imposed on $V$ in both proofs (our and that in [9]) seem to be needed just for technical reasons, such conditions turn out to be (basically) the same; this may suggest that they could be optimal (see Remark 4.3). Despite this, we are not able to prove this fact by means of a counterexample.

In Section 4.3 we study a family of weighted measures linked to exponential measures. Namely we deal with the problem

$$
\begin{equation*}
\min \left\{P_{w e^{V}}(E): \int_{E} e^{V} d x=c\right\} \tag{4.4}
\end{equation*}
$$

In particular we prove in Theorem 4.5 (respectively Theorem 4.7) that, under suitable assumptions on $w$ and $V$, the half-spaces of the form $\left\{\left(x_{1}, x^{\prime}\right) \in \mathbb{R} \times \mathbb{R}^{d-1}: x_{1}>t\right\}$
(respectively $\left\{\left(x_{1}, x^{\prime}\right) \in \mathbb{R} \times \mathbb{R}^{d-1}: x_{1}<t\right\}$ ) are the only minimizer of problem (4.4). The proof is based on the calibration technique developed in the case of the Lebesgue measure, and exploits the particular structure of exponential measures. For this latter problem we will not address the stability issue.

### 4.2 Weighted isoperimetric inequality for the Lebesgue measure

The main result of this section is the following theorem. Later, in Subsection 4.2.1, we will prove that the exponent 2 appearing at the right-hand side of (4.6) is sharp.

Theorem 4.2. Let $d \geq 2$ and $V:[0, \infty) \rightarrow[0, \infty)$ be a weight function such that $V \in C^{2}((0, \infty))$ and satisfying the following properties:
$V(0)=0 \quad$ and $\quad W(t):=V^{\prime}(t)+(d-1) \frac{V(t)}{t} \quad$ is such that $\quad W^{\prime}(t)>0, \quad t>0$.
Then for every open bounded set $E \subset \mathbb{R}^{d}$ with Lipschitz boundary, we have

$$
\begin{equation*}
P_{V}(E) \geq d \omega_{d}^{1 / d}|E|^{1-\frac{1}{d}}\left[V\left(\left(\frac{|E|}{\omega_{d}}\right)^{\frac{1}{d}}\right)+c_{d, V,|E|}\left(\frac{|E \Delta B|}{|E|}\right)^{2}\right], \tag{4.6}
\end{equation*}
$$

where $B$ is the ball centred at the origin and such that $|B|=|E|$. Here $c_{d, V,|E|}$ is a constant depending on $d$, the weight $V$ and the measure of $E$, defined by

$$
c_{d, V,|E|}=\frac{1}{4}\left(\min _{t \in\left[r_{E}, r_{E} \sqrt[d]{2}\right]} W^{\prime}(t)\right) \frac{\sqrt[d]{2}-1}{d}\left(\frac{|E|}{\omega_{d}}\right)^{\frac{2}{d}}
$$

where for simplicity we set

$$
\begin{equation*}
r_{E}:=\left(\frac{|E|}{\omega_{d}}\right)^{\frac{1}{d}} \tag{4.7}
\end{equation*}
$$

Proof. Let $B$ be the ball centred at the origin and having radius $r_{E}$, so that $|B|=|E|$. The key idea of the proof is to use a sort of calibration technique, adapted to the case of weighted perimeters; related ideas can be found in the recent paper [82]. Namely, we consider the following vector field

$$
x \mapsto V(|x|) \frac{x}{|x|}, \quad x \in \mathbb{R}^{d} \backslash\{0\},
$$

whose divergence is given by

$$
\operatorname{div}\left(V(|x|) \frac{x}{|x|}\right)=V^{\prime}(|x|)+(d-1) \frac{V(|x|)}{|x|}=W(|x|), \quad x \in \mathbb{R}^{d} \backslash\{0\}
$$

Recall that, by hypotheses, $W$ is an increasing function. Integrating $W$ on $E$ and then applying the Divergence Theorem, we obtain

$$
\int_{E} W(|x|) d x=\int_{\partial E} V(|x|)\left\langle\frac{x}{|x|}, \nu_{E}(x)\right\rangle d \mathcal{H}^{d-1} \leq P_{V}(E),
$$

and in the same way, integrating on $B$ we get

$$
\int_{B} W(|x|) d x=\int_{\partial B} V(|x|) d \mathcal{H}^{d-1}=P_{V}(B) .
$$

Subtracting $P_{V}(B)$ from the previous inequality, we then obtain

$$
\int_{E} W(|x|) d x-\int_{B} W(|x|) d x \leq P_{V}(E)-P_{V}(B) .
$$

We now observe that thanks to the fact that $|B|=|E|$, we have $|E \backslash B|=|B \backslash E|$ and then

$$
\begin{aligned}
\int_{E} W(|x|) d x-\int_{B} W(|x|) d x & =\int_{E \backslash B} W(|x|) d x-\int_{B \backslash E} W(|x|) d x \\
& =\int_{E \backslash B}\left[W(|x|)-W\left(r_{E}\right)\right] d x-\int_{B \backslash E}\left[W(|x|)-W\left(r_{E}\right)\right] d x \\
& =\int_{E \Delta B}\left|W(|x|)-W\left(r_{E}\right)\right| d x=: \mathcal{R}(E),
\end{aligned}
$$

where in the last equality we used the monotone behaviour of $W$. Resuming, we have obtained the following

$$
\begin{equation*}
P_{V}(E)-P_{V}(B) \geq \mathcal{R}(E), \tag{4.8}
\end{equation*}
$$

and the right-hand side is just the integral of a given function over the region $E \Delta B$, so very likely this gives the desired estimate (4.6). Note that since the functional $\mathcal{R}$ is non-negative, we already proved the isoperimetric inequality. In order to make this precise, let us introduce the radius

$$
r_{2}=\left(r_{E}^{d}+\frac{|E \backslash B|}{\omega_{d}}\right)^{\frac{1}{d}}
$$

and the annular region

$$
T=\left\{x \in \mathbb{R}^{d}: r_{E}<|x|<r_{2}\right\},
$$

which by construction satisfies $|T|=|E \backslash B|=|B \backslash E|$. Notice that

$$
r_{2} \leq r_{E} \sqrt[d]{2}
$$

Using the monotonicity of the function $t \mapsto W(t)$, we get

$$
\begin{aligned}
\mathcal{R}(E) & =\int_{\left\{x \in E:|x|>r_{E}\right\}}\left[W(|x|)-W\left(r_{E}\right)\right] d x+\int_{\left\{x \notin E:|x|<r_{E}\right\}}\left[W\left(r_{E}\right)-W(|x|)\right] d x \\
& \geq \int_{T}\left[W(|x|)-W\left(r_{E}\right)\right] d x
\end{aligned}
$$

so that

$$
\begin{equation*}
\mathcal{R}(E) \geq d \omega_{d} \int_{r_{E}}^{r_{2}}\left[W(\varrho)-W\left(r_{E}\right)\right] \varrho^{d-1} d \varrho \tag{4.9}
\end{equation*}
$$

Thanks to the hypotheses $W^{\prime}(t)>0$ if $t>0$, if we set

$$
c_{1}=\min _{t \in\left[r_{E}, r_{E} \sqrt[d]{2}\right]} W^{\prime}(t)
$$

this is a strictly positive constant, depending on $d, V$ and $|E|$, then from (4.9) we can infer

$$
\mathcal{R}(E) \geq d \omega_{d} c_{1} \int_{r_{E}}^{r_{2}}\left(\varrho-r_{E}\right) \varrho^{d-1} d \varrho
$$

We now develope the computations for this integral: keeping into account that $|E|=$ $\omega_{d} r_{E}^{d}$, we have

$$
\begin{aligned}
\int_{r_{E}}^{r_{2}}\left(\varrho-r_{E}\right) \varrho^{d-1} d \varrho & =\frac{r_{2}^{d+1}-r_{E}^{d+1}}{d+1}-r_{E} \frac{r_{2}^{d}-r_{E}^{d}}{d} \\
& =r_{E}^{d+1}\left[\frac{1}{d+1}\left(\left(1+\frac{|E \backslash B|}{|E|}\right)^{\frac{d+1}{d}}-1\right)-\frac{1}{d} \frac{|E \backslash B|}{|E|}\right]
\end{aligned}
$$

Let us now focus on the function $\varphi(t)=(1+t)^{\alpha}-1$, for $t \in[0,1]$ and with $1<\alpha<2$ : we have the following elementary estimate

$$
(1+t)^{\alpha}-1 \geq \alpha t+c_{2} t^{2}, \quad t \in[0,1]
$$

with constant $c_{2}$ given by

$$
c_{2}=\frac{\alpha}{4}\left(2^{\alpha-1}-1\right)>0
$$

Applying this inequality with the choices $t=|E \backslash B| /|E|$ and $\alpha=1+1 / d$, we then obtain

$$
\int_{r_{E}}^{r_{2}}\left(\varrho-r_{E}\right) \varrho^{d-1} d \varrho \geq r_{E}^{d+1} \frac{\sqrt[d]{2}-1}{d}\left(\frac{|E \backslash B|}{|E|}\right)^{2}
$$

Thus, we arrive at the following estimate

$$
P_{V}(E)-P_{V}(B) \geq \mathcal{R}(E) \geq d \omega_{d} r_{E}^{d+1} \frac{C}{4}\left(\frac{|E \Delta B|}{|E|}\right)^{2}
$$

where we have set

$$
C=\left(\min _{t \in\left[r_{E}, r_{E} \sqrt[d]{2}\right]} W^{\prime}(t)\right) \frac{\sqrt[d]{2}-1}{d}
$$

This finally gives (4.6), keeping into account that

$$
P_{V}(B)=d \omega_{d}^{1 / d}|E|^{(d-1) / d} V\left(r_{E}\right)
$$

and recalling the definition of $r_{E}$.
Remark 4.3 (Assumptions on the weight $V$ ). In [9] it is proven under the following assumptions

$$
V \text { strictly increasing } \quad \text { and } \quad v(t):=V\left(t^{1 / d}\right) t^{1-1 / d}, \quad t \geq 0 \quad \text { convex, (4.10) }
$$

the following sharp lower bound for the weighted perimeter

$$
\begin{equation*}
P_{V}(E) \geq d \omega_{d}^{1 / d}|E|^{1-\frac{1}{d}} V\left(\left(\frac{|E|}{\omega_{d}}\right)^{\frac{1}{d}}\right) \tag{4.11}
\end{equation*}
$$

with equality if and only if $E$ is a ball centred at the origin. This precisely implies that the ball centred at the origin is the only minimizer of $P_{V}$, under volume constraint. It is not difficult to see that this is slightly more general than our (4.5), since (4.10) is equivalent to require that $W$ is non-decreasing. Anyway, our hypotheses could be somehow relaxed: first of all, from the estimate (4.9), we easily see that our proof still characterizes the ball as the unique isoperimetric set, simply requiring that $W$ is strictly increasing, in particular avoiding the requirement $W^{\prime}>0$ and the $C^{2}$ regularity of $V$. Secondly, a closer inspection of our proof reveals that it provides the stronger lower bound

$$
\begin{equation*}
P_{V}(E)-P_{V}(B) \geq \frac{1}{2} \int_{\partial E}\left|\nu_{E}(x)-\frac{x}{|x|}\right|^{2} V(|x|) \mathcal{H}^{d-1}+\mathcal{R}(E) \tag{4.12}
\end{equation*}
$$

Then a characterization of equality cases and a stability estimate seems still feasible, by simply requiring $W$ non-decreasing (as in [9]) and exploiting the first term in the right-hand side of (4.12). A stability estimate of this type - i.e. containing the $L^{2}$ distance of the normal versors - has been derived in [61] for the standard isoperimetric inequality. However, in our case some additional difficulties arise, due to the presence of the weight $V$.

In connection with our later application in Chapter 5, a significant instance of function $V$ satisfying (4.5) is given by any strictly convex power function, i.e. $V(|x|)=$ $|x|^{p}$ with $p>1$. In this case, we use the distinguished notation

$$
P_{p}(E)=\int_{\partial E}|x|^{p} \mathcal{H}^{d-1}
$$

and occasionally we will call $P_{p}(E)$ the $p$-perimeter of $E$. We have $P_{p}(\lambda E)=\lambda^{p+d-1} P_{p}(E)$, for every $\lambda>0$, which implies in particular that the shape functional

$$
E \mapsto|E|^{(1-d-p) / d} P_{p}(E),
$$

is dilation invariant. Then inequality (4.11) can be equivalently written in scaling invariant form as

$$
\begin{equation*}
|E|^{\frac{1-p-d}{d}} P_{p}(E) \geq d \omega_{d}^{\frac{1-p}{d}} \tag{4.13}
\end{equation*}
$$

As a corollary of the previous theorem, we have the following quantitative improvement of (4.13).

Corollary 4.4. Let $p>1$, then for every open bounded set $E \subset \mathbb{R}^{d}$ with Lipschitz boundary, we have

$$
\begin{equation*}
|E|^{\frac{1-p-d}{d}} P_{p}(E) \geq d \omega_{d}^{\frac{1-p}{d}}\left[1+c_{d, p}\left(\frac{|E \Delta B|}{|E|}\right)^{2}\right] \tag{4.14}
\end{equation*}
$$

where $B$ is the ball centred at the origin such that $|E|=|B|$ and $c_{d, p}$ is a constant depending only on $d$ and $p$, given by

$$
c_{d, p}=\frac{(d+p-1)(p-1)}{4} \frac{\sqrt[d]{2}-1}{d}\left(\min _{t \in[1, \sqrt[d]{2}]} t^{p-2}\right)
$$

Proof. We start observing that if $V(t)=t^{p}$, then

$$
W(t)=(d+p-1) t^{p-1} \quad \text { and } \quad W^{\prime}(t)=(d+p-1)(p-1) t^{p-2} .
$$

In particular, using the homogeneity of $W^{\prime}$ we get that
$\min _{t \in\left[r_{E}, r_{E} \sqrt[d]{2}\right]} W^{\prime}(t)=r_{E}^{p-2} \min _{t \in[1, \sqrt[d]{2}]} W^{\prime}(t)=\left(\frac{|E|}{\omega_{d}}\right)^{\frac{p-2}{d}}(d+p-1)(p-1)\left(\min _{t \in[1, \sqrt[d]{2}]} t^{p-2}\right)$.
Then in order to obtain (4.14), it is sufficient to insert the previous into (4.6), to use that

$$
V\left(\left(\frac{|E|}{\omega_{d}}\right)^{\frac{1}{d}}\right)=\left(\frac{|E|}{\omega_{d}}\right)^{\frac{p}{d}}
$$

and to divide both members of (4.6) by $|E|^{(p+d-1) / d}$.

### 4.2.1 Nearly spherical ellipsoids

We now show that the exponent 2 for the term $|E \Delta B|$ in inequality (4.6) is optimal. To this aim, we simply exhibit for every radius $R$ a sequence of sets $E_{\varepsilon}^{R}$, such that $\left|E_{\varepsilon}^{R}\right|=\omega_{d} R^{d}$ and that

$$
\begin{equation*}
\varlimsup_{\varepsilon \rightarrow 0} \frac{P_{V}\left(E_{\varepsilon}^{R}\right)-P_{V}\left(B_{R}\right)}{\left|B_{R} \Delta E_{\varepsilon}^{R}\right|^{2}} \leq C, \tag{4.15}
\end{equation*}
$$

where $B_{R}$ is the ball of radius $R$ centred in the origin. For the sake of simplicity, we confine ourselves to consider the case $d=2$ : the very same arguments still work for every $d \geq 3$.

First of all, we aim to prove (4.15) for $R=1$, then we will show how to obtain it for a general $R>0$. Let us consider the following family of ellipses

$$
E_{\varepsilon}=\left\{\left(x \sqrt{1+\varepsilon}, \frac{y}{\sqrt{1+\varepsilon}}\right): x^{2}+y^{2}<1\right\}
$$

whose boundary can be parametrized by

$$
\gamma_{\varepsilon}(\vartheta)=\left(\sqrt{1+\varepsilon} \cos \vartheta, \frac{1}{\sqrt{1+\varepsilon}} \sin \vartheta\right), \quad \vartheta \in[0,2 \pi] .
$$

Observe that by construction we have $\left|E_{\varepsilon}\right|=\left|B_{1}\right|=\pi$, since

$$
E_{\varepsilon}=\mathcal{M}_{\varepsilon}\left(B_{1}\right)
$$

with $\mathcal{M}_{\varepsilon}: \mathbb{R}^{2} \rightarrow \mathbb{R}^{2}$ linear application, having (with a slight abuse of notation) $\operatorname{det} \mathcal{M}_{\varepsilon}=1$. Now, we need to expand the term

$$
P_{V}\left(E_{\varepsilon}\right)=\int_{0}^{2 \pi} V\left(\left|\gamma_{\varepsilon}(\vartheta)\right|\right)\left|\gamma_{\varepsilon}^{\prime}(\vartheta)\right| d \vartheta
$$

At this aim we use the following second-order Taylor expansions for $\left|\gamma_{\varepsilon}\right|,\left|\gamma_{\varepsilon}^{\prime}\right|$ and $V\left(\left|\gamma_{\varepsilon}\right|\right)$ :

$$
\begin{aligned}
\left|\gamma_{\varepsilon}(\vartheta)\right| & =(1+\varepsilon)^{-1 / 2} \sqrt{1+2 \varepsilon \cos ^{2} \vartheta+\varepsilon^{2} \cos ^{2} \vartheta} \\
& \simeq 1+\varepsilon\left(\cos ^{2} \vartheta-\frac{1}{2}\right)+\frac{\varepsilon^{2}}{2}\left(\frac{3}{4}-\cos ^{4} \vartheta\right)
\end{aligned}
$$

and similarly

$$
\left|\gamma_{\varepsilon}^{\prime}(\vartheta)\right| \simeq 1+\varepsilon\left(\sin ^{2} \vartheta-\frac{1}{2}\right)+\frac{\varepsilon^{2}}{2}\left(\frac{3}{4}-\sin ^{4} \vartheta\right),
$$

while
$V\left(\left|\gamma_{\varepsilon}(\vartheta)\right|\right) \simeq V(1)+\varepsilon V^{\prime}(1)\left[\cos ^{2} \vartheta-\frac{1}{2}\right]+\frac{\varepsilon^{2}}{2}\left[V^{\prime}(1)\left(\frac{3}{4}-\cos ^{4} \vartheta\right)+V^{\prime \prime}(1)\left(\frac{1}{2}-\cos ^{2} \vartheta\right)^{2}\right]$.


Figure 4.1: The family of ellipses $E_{\varepsilon}$.
Thus we have the following second-order expansion for the integrand defining $P_{V}\left(E_{\varepsilon}\right)$ :

$$
\begin{aligned}
V\left(\left|\gamma_{\varepsilon}(\vartheta)\right|\right)\left|\gamma_{\varepsilon}^{\prime}(\vartheta)\right| & \simeq V(1)+\varepsilon\left[V^{\prime}(1)\left(\cos ^{2} \vartheta-\frac{1}{2}\right)+V(1)\left(\sin ^{2} \vartheta-\frac{1}{2}\right)\right] \\
& +\varepsilon^{2}\left[V^{\prime}(1)\left(\cos ^{2} \vartheta-\frac{1}{2}\right)\left(\sin ^{2} \vartheta-\frac{1}{2}\right)+\frac{V(1)}{2}\left(\frac{3}{4}-\sin ^{4} \vartheta\right)\right. \\
& \left.+\frac{V^{\prime \prime}(1)}{2}\left(\frac{1}{2}-\cos ^{2} \vartheta\right)^{2}+\frac{V^{\prime}(1)}{2}\left(\frac{3}{4}-\cos ^{4} \vartheta\right)\right] .
\end{aligned}
$$

Finally, we observe that

$$
\int_{0}^{2 \pi}\left(\cos ^{2} \vartheta-\frac{1}{2}\right) d \vartheta=\int_{0}^{2 \pi}\left(\sin ^{2} \vartheta-\frac{1}{2}\right) d \vartheta=0
$$

and

$$
\int_{0}^{2 \pi}\left(\cos ^{2} \vartheta-\frac{1}{2}\right)^{2} d \vartheta=-\int_{0}^{2 \pi}\left(\cos ^{2} \vartheta-\frac{1}{2}\right)\left(\sin ^{2} \vartheta-\frac{1}{2}\right) d \vartheta=\frac{\pi}{4}
$$

while

$$
\int_{0}^{2 \pi}\left(\frac{3}{4}-\cos ^{4} \vartheta\right) d \vartheta=\int_{0}^{2 \pi}\left(\frac{3}{4}-\sin ^{4} \vartheta\right) d \vartheta=\frac{3}{4} \pi
$$

Summarizing, we have obtained

$$
\begin{equation*}
P_{V}\left(E_{\varepsilon}\right)-P_{V}\left(B_{1}\right) \simeq \frac{\pi}{8} \varepsilon^{2}\left[3 V(1)+V^{\prime}(1)+V^{\prime \prime}(1)\right] . \tag{4.16}
\end{equation*}
$$

On the other hand it is easily seen that $\left|E_{\varepsilon} \Delta B_{1}\right|=O(\varepsilon)$, thus we get (4.15) for $R=1$.
To obtain this result for a generic $R>0$, we notice that for every set $E$,

$$
P_{V}(R E)=R P_{V_{R}}(E)
$$

where $V_{R}(t)=V(R t), t \geq 0$. Hence, if we set $\widetilde{E}_{\varepsilon}:=R E_{\varepsilon}$ we have

$$
\begin{aligned}
P_{V}\left(\widetilde{E}_{\varepsilon}\right)-P_{V}\left(B_{R}\right) & =R\left[P_{V_{R}}\left(E_{\varepsilon}\right)-P_{V_{R}}\left(B_{1}\right)\right] \\
& \simeq \varepsilon^{2} \frac{\pi R}{8}\left[3 V(R)+R V^{\prime}(R)+R^{2} V^{\prime \prime}(R)\right],
\end{aligned}
$$

thanks to (4.16), thus giving (4.15) also in the general case. Observe that (4.5) implies that

$$
R^{2} V^{\prime \prime}(R)+R V^{\prime}(R)>V(R)
$$

and thus in particular

$$
3 V(R)+R V^{\prime}(R)+R^{2} V^{\prime \prime}(R)>4 V(R)>0
$$

### 4.3 Weighted isoperimetric inequalities for exponential measures

In this section we adapt the calibration method developed in last section to another class of isoperimetric problems, related this time to a Gauss-type family of measures, instead of the Lebesgue one. In order to introduce our main results, stated in Theorem 4.5 and Theorem 4.7, we introduce some notations. Let $\mu$ be a finite positive Radon measure on $\mathbb{R}^{d}$ and let $A$ be a Borel subset of $\mathbb{R}^{d}$, the right rearrengement of the set $A$, denoted by $R_{A}^{\mu}$, is the open right half-space $\left\{\left(x_{1}, x^{\prime}\right) \in \mathbb{R} \times \mathbb{R}^{d-1}: x_{1}>t_{A}\right\}$, having the same measure of $A: \mu\left(R_{A}\right)=\mu(A)$. Notice that, if $d \mu=f d x$ for some integrable function $f$, then the value $t_{A}$ is uniquely determined.
Given a measurable function $V: \mathbb{R}^{d} \rightarrow \mathbb{R}$ we denote by $\gamma$ the measure whose density equals $e^{V}$, and for any measurable set $E \subset \mathbb{R}^{d}$ we define the $V$-volume of $E$ as

$$
\begin{equation*}
\gamma(E)=\int_{E} e^{V(x)} d x \tag{4.17}
\end{equation*}
$$

In what follows we will often adopt the notation $x=\left(x_{1}, x^{\prime}\right) \in \mathbb{R} \times \mathbb{R}^{d-1}$. Consider now a Borel weight function $w: \mathbb{R} \rightarrow[0,+\infty]$ and define, for any open set $A$ with Lipschitz boundary, the weighted $V$-perimeter as follows:

$$
P_{w e^{V}}(A)=\int_{\partial A} w\left(x_{1}\right) e^{V(x)} d \mathcal{H}^{d-1}(x) .
$$

In the following theorem we show that, under further conditions on both $w$ and $V$, right half-spaces are minimizers of such a perimeter among the sets of fixed measure with respect to the measure $\gamma$.

Theorem 4.5. Let $A \subset \mathbb{R}^{d}$ be a Lipschitz set, and $w: \mathbb{R} \rightarrow \mathbb{R}^{+}, V: \mathbb{R}^{d} \rightarrow \mathbb{R}$ be $C^{1}$ functions satisfying the following assumption:
(i) $\int_{A} e^{V} d x$ and $\int_{A} w\left(x_{1}\right) e^{V} d x$ are finite,
(ii) $g(x):=-w^{\prime}\left(x_{1}\right)-w\left(x_{1}\right) \partial_{1} V(x)$ depends only on $x_{1}$, and it is a decreasing function on the real line.
(iii) $\int_{x_{1} \geq t} e^{V} d x<+\infty, \int_{x_{1} \geq t} w e^{V} d x<+\infty$ for any $t \in \mathbb{R}$.

Then

$$
P_{w e^{V}}(A) \geq P_{w e^{V}}\left(R_{A}^{\mu}\right)
$$

where $R_{A}$ is the right half-space $R_{A}=\left\{\left(x_{1}, x^{\prime}\right)=: x_{1}>t_{A}\right\}$ which satisfies

$$
\gamma\left(R_{A}\right)=\gamma(A)
$$

Proof. Let $e_{1}=(1,0, \ldots, 0) \in \mathbb{R}^{d}$ and consider the vector field $-e_{1} w\left(x_{1}\right) e^{V}(x)$. Its divergence is given by

$$
\operatorname{div}\left(-e_{1} w\left(x_{1}\right) e^{V}(x)\right)=\left(-w^{\prime}\left(x_{1}\right)-w\left(x_{1}\right) \partial_{1} V(x)\right) e^{V(x)}=g(x) e^{V(x)}
$$

By an application of the Divergence Theorem we have

$$
\begin{align*}
\int_{A} g(x) d \gamma(x) & =\int_{A} \operatorname{div}\left(-e_{1} w\left(x_{1}\right) e^{V(x)}\right) d x \\
& =\int_{\partial A} w\left(x_{1}\right) e^{V(x)}\left\langle\nu_{A}(x),-e_{1}\right\rangle d \mathcal{H}^{d-1}(x)  \tag{4.18}\\
& \leq \int_{\partial A} w\left(x_{1}\right) e^{V(x)} d \mathcal{H}^{d-1}(x)=P_{w e^{V}}(A)
\end{align*}
$$

Let $t_{A}$ be a real number such that the right half-space $R_{A}=\left\{\left(x_{1}, x^{\prime}\right): x_{1} \geq t_{A}\right\}$ satisfies $\gamma\left(R_{A}\right)=\gamma(A)$. Then, since the outer normal of $R_{A}$ is the constant vector field $-e_{1}$, we have that the inequality in (4.18) turns into an equality if tested on $R_{A}$ instead of $A$. Hence we get

$$
P_{w e^{V}}(A)-P_{w e^{V}}\left(R_{A}\right) \geq \int_{A} g(x) d \gamma(x)-\int_{R_{A}} g(x) d \gamma(x)
$$

Since $\gamma(A)=\gamma\left(R_{A}\right)$, then it holds $\gamma\left(A \backslash R_{A}\right)=\gamma\left(R_{A} \backslash A\right)$ as well. Thus

$$
\begin{array}{r}
\int_{A} g(x) d \gamma(x)-\int_{R_{A}} g(x) d \gamma(x)=\int_{A \backslash R_{A}} g(x) d \gamma(x)-\int_{R_{A} \backslash A} g(x) d \gamma(x) \\
=\int_{A \backslash R_{A}}\left(g(x)-g\left(t_{A} e_{1}\right)\right) d \gamma(x)-\int_{R_{A} \backslash A}\left(g(x)-g\left(t_{A} e_{1}\right)\right) d \gamma(x)
\end{array}
$$

Since every $x \in A \backslash R_{A}$ (respectively $x \in R_{A} \backslash A$ ) satisfies $\left\langle x, e_{1}\right\rangle<t_{A}$ (respectively $>t_{A}$ ), by condition (ii) we deduce

$$
\begin{aligned}
P_{w e^{V}}(A)-P_{w e^{V}}\left(R_{A}\right) & \geq \int_{A \backslash R_{A}}\left|g(x)-g\left(t_{A} e_{1}\right)\right| d \gamma(x)+\int_{R_{A} \backslash A}\left|g(x)-g\left(t_{A} e_{1}\right)\right| d \gamma(x) \\
& =\int_{A \Delta R_{A}}\left|g(x)-g\left(t_{A} e_{1}\right)\right| d \gamma(x) \geq 0
\end{aligned}
$$

where $A \Delta R_{A}=\left(A \backslash R_{A}\right) \cup\left(R_{A} \backslash A\right)$ stands for the symmetric difference between $A$ and $R_{A}$.

Condition (ii) of the previous theorem is satisfied by functions of the form

$$
V\left(x_{1}, x^{\prime}\right)=V_{1}\left(x_{1}\right)+V_{2}\left(x^{\prime}\right),
$$

where the real function $V_{1} \in C^{2}(\mathbb{R})$ and, together with the weight $w \geq 0$ satisfies the ordinary differential inequality

$$
\begin{equation*}
w^{\prime \prime}+V_{1}^{\prime \prime} w+V_{1}^{\prime} w^{\prime} \geq 0 . \tag{4.19}
\end{equation*}
$$

Although Theorem 4.5 is stated in the whole space $\mathbb{R}^{d}$, inequality (4.19) is seldom (if ever) satisfied by functions which fulfil also the integrability property (i). This is not the case if we restrict our attention to the half space

$$
\Omega=\left\{\left(x_{1}, x^{\prime}\right) \in \mathbb{R} \times \mathbb{R}^{d-1}: x_{1}>0\right\} .
$$

In this case we can find a big class of functions $w$ and $V$ satisfying (4.19).
Corollary 4.6. Let $w \in C^{2}([0, \infty],[0, \infty])$ and $V=V_{1}\left(x_{1}\right)+V_{2}\left(x^{\prime}\right): \mathbb{R}^{d} \rightarrow \mathbb{R}$ be such that $V_{1} \in C^{1}(\mathbb{R})$ and

$$
w^{\prime \prime}+w^{\prime} V_{1}^{\prime}+w V_{1}^{\prime \prime} \geq 0
$$

Suppose moreover that the functions $w$ and $V$ satisfy the integrability condition (i) in Theorem 4.5. Then the solution of problem

$$
\min \left\{P_{w e}{ }^{V}(A): 1_{A} \in L^{1}\left(\Omega, w e^{V}\right), \int_{A} e^{V}=c\right\}
$$

is given by the right space $R_{c}=\left\{x_{1} \geq t_{c}\right\}$ such that $\int_{R_{c}} w e^{V}=c$.
A class of particular interest is that where $V(x)=-|x|^{2}$, that is where $e^{V} d x$ corresponds to the (non-rescaled) Gauss measure. In this case there is a large class of real functions $w$ which adapt to the hypotheses of Corollary 4.6. This is the case, for instance of $w(t)=e^{-t}, t^{-\alpha}$ with $\alpha \geq 1$ and $w(t)=-\log (t) \mathbf{1}_{(0,1]}$. Similar results have been considered in [21] where the authors deal with the case $w(t)=t^{k}, k>0$ and the substantial difference with our problem is that the space $\Omega$ is weighted as the perimeter.

If we suitably overturn the hypotheses of Theorem 4.5, we get an analogous reversed result, related to left half spaces.

Theorem 4.7. Let $A \subset \mathbb{R}^{d}$ be a Lipschitz set, and $w: \mathbb{R} \rightarrow \mathbb{R}^{+}, V: \mathbb{R}^{d} \rightarrow \mathbb{R}$ be $C^{1}$ functions satisfying the assumptions (i) and (ii) of Theorem 4.5 and
$(\text { iii) })^{\prime} \int_{x_{1} \leq t} e^{V} d x<+\infty, \int_{x_{1} \leq t} w e^{V} d x<+\infty$ for any $t \in \mathbb{R}$.
Then

$$
P_{w e^{V}}(A) \geq P_{w e^{V}}\left(L_{A}\right),
$$

where $L_{A}$ is the left half space $L_{A}=\left\{\left(x_{1}, x^{\prime}\right)=: x_{1}<t_{A}\right\}$ which satisfies

$$
\gamma\left(R_{A}\right)=\gamma(A)
$$

Proof. It is analogous to the proof of Theorem 4.5, considering this time the vector field $e_{1} w e^{V}$.

The same considerations made for Theorem 4.5 hold also in this case. Again some interesting examples can be made considering sets lying in the right half space $\Omega$. Observe that this request is not completely analogous to the previous one, where $\Omega$ was itself an optimal subspace. Indeed if we restrict as before our attention to the right half space $\Omega=\left\{x_{1}>0\right\}$, the minimizer are vertical strips of the form $\left\{\left(x_{1}, x^{\prime}\right) \in \mathbb{R} \times \mathbb{R}^{d-1}\right.$ : $\left.0<x_{1}<b\right\}$.

Corollary 4.8. Let $w \in C^{2}([0, \infty],[0, \infty])$ and $V=V_{1}\left(x_{1}\right)+V_{2}\left(x^{\prime}\right): \mathbb{R}^{d} \rightarrow \mathbb{R}$ be such that $V_{1} \in C^{1}(\mathbb{R})$ and

$$
w^{\prime \prime}+w^{\prime} V_{1}^{\prime}+w V_{1}^{\prime \prime} \geq 0
$$

Suppose moreover that the functions $w$ and $V$ satisfy the integrability condition (iii)' of Theorem 4.7. Then the solution of problem

$$
\min \left\{P_{w e^{V}}(A): \mathbf{1}_{A} \in L^{1}\left(\Omega, w \gamma_{d}\right), \gamma_{d}(A)=c\right\}
$$

is given by the strip $S_{c}=\left\{0<x_{1}<t_{c}\right\}$ such that $\int_{S_{c}} w e^{V}=c$.
In this case instances of weights $V$ are those of the form $V\left(x_{1}, x^{\prime}\right)=x_{1}^{2}+V_{2}\left(x^{\prime}\right)$, for a suitable integrable function $V_{2}$. These examples are slightly related to the recent work [22], where the weight considered is $V(x)=|x|^{2}$, the space (and not only the perimeter) is endowed with the measure $\varphi=x_{1}^{k} e^{|x|^{2}}, k>0$ and the minimizer of the isoperimetric problem

$$
\min \left\{\int_{A} \varphi(x) d x: A \subset \Omega, \int_{A} \varphi(x) d x=c\right\}
$$

is the centred half ball $B_{c}=\left\{x=\left(x_{1}, x^{\prime}\right) \in \mathbb{R} \times \mathbb{R}^{d-1}: x_{1}>0,|x|<r_{c}\right\}$ such that $\int_{B_{c}} \varphi(x) d x=c$. In our case some examples of weights $w$ which fulfil the hypotheses of Corollary 4.8 are $w(t)=e^{-t^{\alpha}}$, with $\alpha>1$ and $w(t)=t^{k}, k \in(0,1)$.

## Chapter 5

## Stability for the first Stekloff-Laplacian eigenvalue

### 5.1 Introduction

The scope of this chapter, based on the joint work with Lorenzo Brasco and Guido De Philippis [16], is to study the stability of the spectral problem related to the first nontrivial Stekloff eigenvalue. The first non-trivial Stekloff eigenvalue can be formulated as follows

$$
\sigma_{2}(E)=\inf _{u \in W^{1,2}(E) \backslash\{0\}}\left\{\frac{\int_{E}|\nabla u(x)|^{2} d x}{\int_{\partial E} u(x)^{2} d \mathcal{H}^{d-1}}: \int_{\partial E} u(x) d \mathcal{H}^{d-1}=0\right\},
$$

i.e. $1 / \sigma_{2}(E)$ is the best constant in the Poincaré-Wirtinger trace inequality

$$
\begin{equation*}
\int_{\partial E}\left|u(x)-\left(\int_{\partial E} u(x)\right)\right|^{2} d \mathcal{H}^{d-1} \leq C_{E} \int_{E}|\nabla u(x)|^{2} d x, \quad u \in W^{1,2}(E) . \tag{5.1}
\end{equation*}
$$

The Brock-Weinstock inequality asserts that in the class of sets with given volume, $\sigma_{2}$ is maximized by a ball, i.e.

$$
\begin{equation*}
|E|^{1 / d} \sigma_{2}(E) \leq|B|^{1 / d} \sigma_{2}(B) \tag{5.2}
\end{equation*}
$$

where $B$ is any ball and equality holds if and only if $E$ itself is a ball. Notice that the quantity $|E|^{1 / d} \sigma_{2}(E)$ is scaling invariant. The main result of this chapter is a sharp quantitative version of (5.2):

Theorem 5.1 (Quantitative Brock-Weinstock inequality). For every open bounded Lipschitz set $E \subset \mathbb{R}^{d}$, there holds

$$
\begin{equation*}
|E|^{1 / d} \sigma_{2}(E) \leq|B|^{1 / d} \sigma_{2}(B)\left[1-\alpha_{d} \mathcal{A}(E)^{2}\right] \tag{5.3}
\end{equation*}
$$

where $\alpha_{d}$ is an explicit dimensional constant and $\mathcal{A}(E)$ is the Fraenkel asymmetry

$$
\mathcal{A}(E)=\inf \left\{\frac{\left\|1_{E}-1_{B}\right\|_{L^{1}\left(\mathbb{R}^{d}\right)}}{|E|}:|E|=|B|\right\} .
$$

To be precise, we will prove, in Theorem 5.6, a slightly stronger result, which entails Theorem 5.1. Some words on the proof of this result are in order: the maximality of the ball for $\sigma_{2}$ is a consequence of a particular case of the weighted isoperimetric inequality analysed in Section 4.2 of the previous chapter. Namely, the crucial point is that the ball centred at the origin (uniquely) minimizes the shape functional

$$
E \mapsto \int_{\partial E}|x|^{2} d \mathcal{H}^{d-1}
$$

among sets with given measure, result that, as we have seen in Chapter 4, is proved in [9]. This further isoperimetric characterization of the ball is the key tool in Brock's proof of (5.2) in [20]: then in order to derive (5.3), we can take advantage of the stability result proved in the previous chapter for such a weighted perimeter. Namely, we will make use of the following result (stated in this chapter as a lemma, but which is a direct consequence of Theorem 4.2 and, in particular, of Corollary 4.4). .

Lemma 5.2. For every open bounded Lipschitz set $E \subset \mathbb{R}^{d}$, we have

$$
\int_{\partial E}|x|^{2} d \mathcal{H}^{d-1} \geq \int_{\partial B}|x|^{2} d \mathcal{H}^{d-1}\left[1+\beta_{d}\left(\frac{|E \Delta B|}{|E|}\right)^{2}\right]
$$

where $\beta_{d}$ is an explicit dimensional constant, $B$ is the ball centred at the origin such that $|B|=|E|$ and $E \Delta B$ denotes the symmetric difference.

Concerning the sharpness of the exponent 2 for the Fraenkel asymmetry in (5.3), we stress that its proof (Theorem 5.9) is much longer than the same result for weighted perimeters proved in Subsection 4.2.1 of the previous chapter. A possible explanation is the following: since, differently from the perimeter, the Stekloff eigenvalues do not have a straightforward geometrical meaning, it is much more complicated to understand how they are affected by deformations of an optimal shape. If the eigenvalue is differentiable in the sense of the shape derivative (see [77]) - like in the case of the first Dirichlet eigenvalue $\lambda_{1}$ - one can use the following argument. Any perturbation of the type $E_{\varepsilon}:=(\operatorname{Id}+\varepsilon X)(B)$, for some smooth vector field $X$, should provide a Taylor expansion of the form

$$
\begin{equation*}
|E|^{2 / d} \lambda_{1}\left(E_{\varepsilon}\right) \simeq|B|^{2 / d} \lambda_{1}(B)+O\left(\varepsilon^{2}\right), \quad \varepsilon \ll 1, \tag{5.4}
\end{equation*}
$$

since the first derivative of $|\cdot|^{2 / d} \lambda_{1}(\cdot)$ has to vanish at the minimum "point" $B$. Then one observes that for such a family of sets, the Fraenkel asymmetry satisfies $\mathcal{A}\left(E_{\varepsilon}\right)=$
$O(\varepsilon)$. This explains, for instance, why the following inequality was expected (and in fact proved in the recent work [15]) in the (sharp) form

$$
|E|^{2 / d} \lambda_{1}(E) \geq|B|^{2 / d} \lambda_{1}(B)\left[1+c_{d} \mathcal{A}(E)^{2}\right] .
$$

For the case of the first non-trivial Stekloff eigenvalue $\sigma_{2}$, proving the sharpness is more complicated: indeed, the most basic example - nearly spherical ellipsoids - leads to an expansion with a non-trivial first order term, i.e.

$$
\left|E_{\varepsilon}\right|^{1 / d} \sigma_{2}\left(E_{\varepsilon}\right) \simeq|B|^{1 / d} \sigma_{2}(B)+O(\varepsilon) .
$$

The same phenomenon has already been observed in [18, Section 5] for the Neumann case. A possible explanation for this fact is the following: at the maximum point, i.e. for a ball $B$, the eigenvalue $\sigma_{2}$ is multiple and thus not differentiable (see [75, Section 2]). Roughly speaking, this implies that along some "directions" (i.e. for some deformations of the ball) the functional $\sigma_{2}$ could have a non-trivial "super-differential". In order to show that the exponent 2 in (5.3) is indeed sharp, one has to exclude that this happens for every direction: namely, one has to exhibit a particular family of deformations $E_{\varepsilon}$ for which a correct expansion like (5.4) is guaranteed. We will achieve our aim by suitably refining a construction introduced in $[18$, Section 6$]$ to solve the same problem in the Neumann case.

The chapter is organized as follows: in Section 5.2 we recall some basic facts about Stekloff eigenvalues, as well as the spectral optimization problems we are concerned with. Thanks to our quantitative estimates for weighted perimeters, we will prove that optimal shapes for these spectral problems are stable (Section 5.3). The corresponding stability estimates happen to be sharp as well, as shown in the final Section 5.4.

### 5.2 Spectral optimization for Stekloff eigenvalues

We start recalling some basic definitions. We will refer to [75, Chapter 7] as the reference to this subject.

Let $E \subset \mathbb{R}^{d}$ be an open bounded set with Lipschitz boundary. Thanks to the compactness of the embedding of $W^{1,2}(E)$ into $L^{2}(\partial E)$, we have that the resolvant operator $\mathcal{R}: L^{2}(\partial E) \rightarrow L^{2}(\partial E)$ defined by

$$
\mathcal{R} g \in W^{1,2}(E) \quad \text { solves in weak sense } \quad\left\{\begin{array}{rll}
-\Delta u=0, & \text { in } E, \\
\left\langle\nabla u, \nu_{E}\right\rangle & =g & \text { on } \partial E,
\end{array}\right.
$$

is a compact, symmetric and positive linear operator. Hence $\mathcal{R}$ has a discrete spectrum, made only of real positive eigenvalues accumulating at 0 . As a consequence, we have that the following boundary value problem for harmonic functions

$$
\left\{\begin{array}{rlrl}
-\Delta u & =0, & & \text { in } E, \\
\left\langle\nabla u, \nu_{E}\right\rangle & =\sigma u, & \text { on } \partial E,
\end{array}\right.
$$

has non-trivial solutions only for a discrete set of values $\sigma_{1}(E) \leq \sigma_{2}(E) \leq \sigma_{3}(E) \ldots$ accumulating at $\infty$ : these are the so-called Stekloff eigenvalues of $E$. Here solutions are intended in the usual weak sense, i.e.

$$
\int_{E}\langle\nabla u(x), \nabla \varphi(x)\rangle d x=\sigma_{k}(E) \int_{\partial E} u(x) \varphi(x) d x, \quad \text { for every } \varphi \in W^{1,2}(E), k \in \mathbb{N} .
$$

The corresponding solutions $\left\{\xi_{k}\right\}_{k \geq 1}$ are called eigenfunctions of the Stekloff-Laplacian and they give an orthonormal basis of $L^{2}(\partial E)$, once renormalized by $\left\|\xi_{k}\right\|_{L^{2}(\partial E)}=1$, for every $k \geq 1$. Throughout the next sections we will use the classical convention of counting the eigenvalues with their multiplicities: this means that if for a certain $k \in \mathbb{N} \backslash\{0\}$, there exist $m$ linearly independent non-trivial solutions for $\sigma_{k}(E)$, then we will write $\sigma_{k}(E)=\sigma_{k+1}(E)=\cdots=\sigma_{k+m-1}(E)$.

Observe that if $E$ has $k$ connected components $E_{1}, \ldots, E_{k}$, then $\sigma_{1}(E)=\cdots=\sigma_{k}(E)=$ 0 and the corresponding renormalized eigenfunctions are given by

$$
\xi_{i}(x)=\frac{1_{E_{i}}(x)}{\sqrt{\mathcal{H}^{d-1}\left(\partial E_{i}\right)}}, \quad i=1, \ldots, k
$$

In particular the first Stekloff eigenvalue of a set is always trivial and corresponds to constant functions. For this reason, given $k \in \mathbb{N} \backslash\{0\}$, we always have that

$$
0=\inf \left\{\sigma_{k}(E):|E|=c\right\}
$$

and this infimum is attained for every open set having $k$ connected components.
Remark 5.3. For what follows, it is important to remark that the functions $\left\{\xi_{k}\right\}_{k \geq 2}$ also give an orthogonal basis for the following closed subspace of $W^{1,2}(E)$
$\operatorname{Har}(E)=\left\{u \in W^{1,2}(E): \int_{\partial E} u=0\right.$ and $\int_{E}\langle\nabla u, \nabla \varphi\rangle=0$ for every $\left.\varphi \in W_{0}^{1,2}(E)\right\}$,
on which $u \mapsto\|\nabla u\|_{L^{2}}$ and $u \mapsto\|u\|_{W^{1,2}}$ are equivalent norms, thanks to the PoincaréWirtinger inequality (5.1) and to inequality

$$
\|u\|_{L^{2}(E)} \leq C_{E}\left(\|\nabla u\|_{L^{2}(E)}+\|u\|_{L^{2}(\partial E)}\right), \quad u \in W^{1,2}(E)
$$

which can be proved by means of a standard compactness argument. Notice that for every $u \in \operatorname{Har}(E)$, its Dirichlet integral can be written as

$$
\begin{equation*}
\int_{E}|\nabla u(x)|^{2} d x=\sum_{k \geq 2} \alpha_{k}^{2} \sigma_{k}(E), \quad \text { where } \quad \alpha_{k}=\int_{\partial E} \xi_{k}(x) u(x) d \mathcal{H}^{d-1} \tag{5.6}
\end{equation*}
$$

For any ball $B$ of radius $R$, its first non-trivial Stekloff eigenvalue is given by

$$
\sigma_{2}(B)=\frac{1}{R}
$$

which corresponds to the eigenfunctions $\xi_{i}(x)=x_{i-1}$, with $i=2, \ldots, d+1$, i.e. the eigenvalue $\sigma_{2}(B)$ has multiplicity $d$. Also, we notice that the shape functional $E \mapsto$ $|E|^{1 / d} \sigma_{2}(E)$ is scaling invariant, thus in particular

$$
|B|^{1 / d} \sigma_{2}(B)=\omega_{d}^{1 / d}
$$

for any ball $B$. About the first non-trivial Stekloff eigenvalue of a set $E$, we have the following sharp estimate, first derived in [111] for dimension $d=2$ and then generalized to any dimension in [20].
Brock-Weinstock inequality. For every $E \subset \mathbb{R}^{d}$ open bounded set with Lipschitz boundary, we have

$$
\begin{equation*}
|E|^{1 / d} \sigma_{2}(E) \leq \omega_{d}^{1 / d} \tag{5.7}
\end{equation*}
$$

and equality holds if and only if $E$ is a ball. In other words, for every $c>0$ the unique solution of the following spectral optimization problem

$$
\max \left\{\sigma_{2}(E):|E| \geq c\right\}
$$

is given by a ball of measure $c$.
Remark 5.4. As already remarked in the Introduction, $1 / \sigma_{2}(E)$ can be characterized as the sharp constant in the Poincaré-Wirtinger trace inequality (5.1). We then notice that the Brock-Weinstock inequality can be extended to any set supporting such an inequality and for which the trace of a $W^{1,2}$ function is well-defined: in these cases, it is still meaningful speaking of $\sigma_{2}(E)$, though the embedding $W^{1,2}(E) \hookrightarrow L^{2}(\partial E)$ could not be compact and hence its Stekloff-Laplacian could have a continuous spectrum.

Actually, the Brock-Weinstock inequality is a straightforward consequence of a stronger estimate proved by Brock in [20], involving the first $d$ non-trivial Stekloff eigenvalues: namely, for every $E \subset \mathbb{R}^{d}$ bounded open set with Lipschitz boundary, we have

$$
\begin{equation*}
\frac{1}{|E|^{1 / d}} \sum_{i=2}^{d+1} \frac{1}{\sigma_{i}(E)} \geq \frac{d}{\omega_{d}^{1 / d}}, \tag{5.8}
\end{equation*}
$$

i.e. any ball minimizes the sum of the reciprocal of the first $d$ non-trivial Stekloff eigenvalues, among sets of given measure.

Remark 5.5. In the case of convex sets, an even stronger estimate is possible [78]: the ball maximizes the product of the first $d$ non-trivial Stekloff eigenvalues, under measure constraint

$$
\begin{equation*}
|E| \prod_{i=2}^{d+1} \sigma_{i}(E) \leq \omega_{d} \tag{5.9}
\end{equation*}
$$

A simple application of the arithmetic-geometric mean inequality shows that the previous implies (5.8): it should be noticed that in dimension $d=2$, the convexity assumption can be dropped (see [76]), while for higher dimensions it is still an open problem to know whether (5.9) holds for all sets or not.

### 5.3 The stability issue

The main goal of this section is to show how (5.8) and (5.7) can be improved by means of a quantitative stability estimate. At this aim, for every $E \subset \mathbb{R}^{d}$ open set with finite measure, we recall the definition of Fraenkel asymmetry

$$
\mathcal{A}(E):=\inf \left\{\frac{\left\|1_{E}-1_{B}\right\|_{L^{1}\left(\mathbb{R}^{d}\right)}}{|E|}: B \text { ball with }|B|=|E|\right\}
$$

i.e. this is the distance in the $L^{1}$ sense of a generic set $E$ from the "manifold" of balls, renormalized in order to make it scaling invariant: observe that $0 \leq \mathcal{A}(E)<2$. Then the main result of this section is the following quantitative improvement of (5.8).

Theorem 5.6. For every $E \subset \mathbb{R}^{d}$ open bounded set with Lipschitz boundary, we have

$$
\begin{equation*}
\frac{1}{|E|^{1 / d}} \sum_{i=2}^{d+1} \frac{1}{\sigma_{i}(E)} \geq \frac{d}{\omega_{d}^{1 / d}}\left[1+c_{d, 2} \mathcal{A}(E)^{2}\right] \tag{5.10}
\end{equation*}
$$

where

$$
c_{d, 2}=\frac{d+1}{d} \frac{\sqrt[d]{2}-1}{4} .
$$

Proof. We start reviewing the proof of Brock in [20]: the first step is to have a variational characterization for the sum of inverses of eigenvalues. In the case of Stekloff eigenvalues, the following formula holds (see [79, Theorem 1], for example):

$$
\sum_{i=2}^{d+1} \frac{1}{\sigma_{i}(E)}=\max _{\left(v_{2}, \ldots, v_{d+1}\right) \in \mathcal{I}} \sum_{i=2}^{d+1} \int_{\partial E} v_{i}(x)^{2} d \mathcal{H}^{d-1}
$$

where the set of admissible functions is given by

$$
\mathcal{I}=\left\{\left(v_{2}, \ldots, v_{d+1}\right) \in\left(W^{1,2}(E)\right)^{d}: \int_{\partial E} v_{i}(x) d \mathcal{H}^{d-1}=0, \int_{E}\left\langle\nabla v_{i}(x), \nabla v_{j}(x)\right\rangle d x=\delta_{i j}\right\} .
$$

Observe that the quantities $\sigma_{i}(E)$ are invariant under translations, so without loss of generality we can suppose that the barycentre of $\partial E$ is in the origin, i.e.

$$
\int_{\partial E} x_{i} d \mathcal{H}^{d-1}=0, \quad i=1, \ldots, d .
$$

This implies that the eigenfunctions $\xi_{i}$ relative to $\sigma_{2}(B)=\cdots=\sigma_{d+1}(B)$ are admissible in the previous maximization problem, thus as admissible functions we take

$$
v_{i}(x)=\frac{x_{i-1}}{\sqrt{|E|}}, \quad i=2, \ldots, d+1
$$

In this way, we obtain

$$
\frac{1}{|E|^{1 / d}} \sum_{i=2}^{d+1} \frac{1}{\sigma_{i}(E)} \geq \frac{1}{|E|^{1+1 / d}} \int_{\partial E}|x|^{2} d \mathcal{H}^{d-1}=|E|^{-\frac{d+1}{d}} P_{2}(E)
$$

which implies

$$
\frac{1}{|E|^{1 / d}} \sum_{i=2}^{d+1} \frac{1}{\sigma_{i}(E)}-\frac{d}{\omega_{d}^{1 / d}} \geq|E|^{-\frac{d+1}{d}} P_{2}(E)-\frac{d}{\omega_{d}^{1 / d}}
$$

This means that the deficit of this spectral inequality is controlling from above the deficit of the 2 -perimeter. Thus it is sufficient to use the quantitative estimate (5.2) for the $2-$ perimeter, so to obtain

$$
\frac{1}{|E|^{1 / d}} \sum_{i=2}^{d+1} \frac{1}{\sigma_{i}(E)}-\frac{d}{\omega_{d}^{1 / d}} \geq \frac{d}{\omega_{d}^{1 / d}} c_{d, 2}\left(\frac{|E \Delta B|}{|E|}\right)^{2}
$$

where $B$ is the ball centred at the origin and such that $|E|=|B|$. Using the definition of $\mathcal{A}(E)$, we can conclude the proof.

A straightforward consequence of the previous result is the following quantitative version of the Brock-Weinstock inequality.
Corollary 5.7. For every $E \subset \mathbb{R}^{d}$ open bounded set with Lipschitz boundary, we have

$$
\begin{equation*}
|E|^{1 / d} \sigma_{2}(E) \leq \omega_{d}^{1 / d}\left[1-\delta_{d} \mathcal{A}(E)^{2}\right] \tag{5.11}
\end{equation*}
$$

where $\delta_{d}$ is a constant depending only on the dimension, given by

$$
\delta_{d}=\frac{1}{8} \min \left\{1, \frac{d+1}{d}(\sqrt[d]{2}-1)\right\}
$$

Proof. First of all, we can suppose that

$$
\begin{equation*}
|E|^{1 / d} \sigma_{2}(E) \geq \frac{1}{2} \omega_{d}^{1 / d} \tag{5.12}
\end{equation*}
$$

otherwise estimate (5.11) is trivially true with constant $\delta_{d}=1 / 8$, just by using the fact that $\mathcal{A}(E)<2$. So, let us suppose that (5.12) holds true: since $\sigma_{2}(E) \leq \sigma_{i}(E)$ for every $i \geq 3$, from (5.10) we can infer

$$
\frac{d}{|E|^{1 / d} \sigma_{2}(E)} \geq \frac{d}{\omega_{d}^{1 / d}}\left[1+c_{d, 2} \mathcal{A}(E)^{2}\right]
$$

which can be rewritten as

$$
|E|^{1 / d} \sigma_{2}(E)\left[1+c_{d, 2} \mathcal{A}(E)^{2}\right] \leq \omega_{d}^{1 / d} .
$$

The previous formula easily implies (5.11), thanks to (5.12).
Remark 5.8. In the next section we will prove that both the estimates derived in Theorem 5.6 and Corollary 5.7 are sharp. We point out that defining the two deficit functionals

$$
\begin{equation*}
\operatorname{Inv}(E):=\frac{|B|^{1 / d}}{d|E|^{1 / d}} \sum_{i=2}^{d+1} \frac{\sigma_{2}(B)}{\sigma_{i}(E)}-1 \quad \text { and } \quad B W(E):=\frac{|B|^{1 / d} \sigma_{2}(B)}{|E|^{1 / d} \sigma_{2}(E)}-1, \tag{5.13}
\end{equation*}
$$

we have that

$$
c_{d, 2} \mathcal{A}(E)^{2} \leq \operatorname{Inv}(E) \leq B W(E),
$$

where in the first inequality we used Theorem 5.6. Then if one can prove that the exponent 2 for $\mathcal{A}(E)$ is sharp in the quantitative Brock-Weinstock inequality, this will automatically prove the optimality of the power 2 for inequality (5.10).

### 5.4 Sharpness of the quantitative Brock-Weinstock inequality

In this section, we will show the sharpness of the quantitative Brock-Weinstock inequality (5.11): as remarked, this in turn will give the sharpness of (5.10) as well. Namely, we are going to prove the following result.

Theorem 5.9. There exists a family $\left\{E_{\varepsilon}\right\}_{\varepsilon>0}$ of smooth sets approaching the ball $B$ of unit radius in such a way that

$$
\begin{equation*}
\mathcal{A}\left(E_{\varepsilon}\right) \simeq \frac{\left|E_{\varepsilon} \Delta B\right|}{\left|E_{\varepsilon}\right|} \simeq \varepsilon \quad \text { and } \quad B W\left(E_{\varepsilon}\right) \simeq \varepsilon^{2}, \quad \varepsilon \ll 1, \tag{5.14}
\end{equation*}
$$

where $B W(E)$ is defined by (5.13).
The rest of this section is devoted to construct such a family of deformations $E_{\varepsilon}$. Since the whole construction is quite complicate, for the sake of readability we will divide it into 4 main steps.

### 5.4.1 Step 1: setting of the construction and basic properties

In what follows, $B \subset \mathbb{R}^{d}$ stands for the open unit ball, centred at the origin. We consider a general nearly circular domain, given by

$$
E_{\varepsilon}=\left\{x \in \mathbb{R}^{d}: x=0 \text { or }|x|<1+\varepsilon \psi(x /|x|)\right\},
$$

where $\psi \in C^{\infty}(\partial B)$ satisfies the following assumptions.

Key assumptions. For every $a \in \mathbb{R}^{d}$, there holds

$$
\begin{equation*}
\int_{\partial B} \psi(x) d \mathcal{H}^{d-1}=0, \quad \int_{\partial B}\langle a, x\rangle \psi(x) d \mathcal{H}^{d-1}=0 \tag{5.15}
\end{equation*}
$$

and

$$
\begin{equation*}
\int_{\partial B}\langle a, x\rangle^{2} \psi(x) d \mathcal{H}^{d-1}=0 . \tag{5.16}
\end{equation*}
$$

We start with a basic result of geometric type.
Lemma 5.10. Let $\psi \in C^{\infty}(\partial B)$ satisfying (5.15). Then

$$
\begin{equation*}
\left|E_{\varepsilon}\right|-|B| \simeq \varepsilon^{2} \quad \text { and } \quad \mathcal{A}\left(E_{\varepsilon}\right) \simeq \varepsilon \simeq \frac{\left|E_{\varepsilon} \Delta B\right|}{\left|E_{\varepsilon}\right|} . \tag{5.17}
\end{equation*}
$$

Proof. Using polar coordinates, the measure $\left|E_{\varepsilon}\right|$ can be expressed as follows

$$
\begin{aligned}
\left|E_{\varepsilon}\right| & =\frac{1}{d} \int_{\partial B}(1+\varepsilon \psi(x))^{d} d \mathcal{H}^{d-1} \\
& \simeq|B|+\varepsilon \int_{\partial B} \psi(x) d \mathcal{H}^{d-1}+\varepsilon^{2} \frac{d-1}{2} \int_{\partial B} \psi(x)^{2} d \mathcal{H}^{d-1},
\end{aligned}
$$

which gives the first relation in (5.17), thanks to the fact that $\psi$ has zero-mean on $\partial B$.
For the second one, we start observing that $\left|E_{\varepsilon} \Delta B\right| \simeq \varepsilon$ : still using polar coordinates, we get

$$
\begin{aligned}
\left|E_{\varepsilon} \Delta B\right| & =\frac{1}{d} \int_{\{x \in \partial B: \psi(x)>0\}}\left[(1+\varepsilon \psi(x))^{d}-1\right] d \mathcal{H}^{d-1} \\
& +\frac{1}{d} \int_{\{x \in \partial B: \psi(x)<0\}}\left[1-(1+\varepsilon \psi(x))^{d}\right] d \mathcal{H}^{d-1} \simeq \varepsilon \int_{\partial B}|\psi(x)| d \mathcal{H}^{d-1} .
\end{aligned}
$$

Now, let $B\left(x_{0}, r_{\varepsilon}\right)$ be a ball realizing the asymmetry, i.e. such that $\mathcal{A}\left(E_{\varepsilon}\right)\left|E_{\varepsilon}\right|=$ $\left|E_{\varepsilon} \Delta B\left(x_{0}, r_{\varepsilon}\right)\right|$ : in particular, we have $\left|B\left(x_{0}, r_{\varepsilon}\right)\right|=\left|E_{\varepsilon}\right|$. It is easily seen that

$$
\begin{equation*}
\mathcal{A}\left(E_{\varepsilon}\right) \leq c \frac{\left|E_{\varepsilon} \Delta B\right|}{\left|E_{\varepsilon}\right|} \tag{5.18}
\end{equation*}
$$

for some constant $c$ independent of $\varepsilon$ : indeed, by definition of $\mathcal{A}\left(E_{\varepsilon}\right)$ and triangular inequality, we get

$$
\mathcal{A}\left(E_{\varepsilon}\right) \leq \frac{\left|E_{\varepsilon} \Delta B\left(0, r_{\varepsilon}\right)\right|}{\left|E_{\varepsilon}\right|} \leq \frac{\left|E_{\varepsilon} \Delta B\right|}{\left|E_{\varepsilon}\right|}+\frac{\left|B \Delta B\left(0, r_{\varepsilon}\right)\right|}{\left|E_{\varepsilon}\right|} \leq c \frac{\left|E_{\varepsilon} \Delta B\right|}{\left|E_{\varepsilon}\right|},
$$

since $\left|B \Delta B\left(0, r_{\varepsilon}\right)\right|=\left|\left|B\left(x_{0}, r_{\varepsilon}\right)\right|-|B|\right| \simeq \varepsilon^{2}$, while $\left|E_{\varepsilon} \Delta B\right| \simeq \varepsilon$.

Using the symmetries of $B$ and (5.15), for every $a \in \mathbb{R}^{d}$ we can infer

$$
\int_{E_{\varepsilon}}\langle a, y\rangle d y=\frac{1}{d+1} \int_{\partial B}(1+\varepsilon \psi(x))^{d+1}\langle a, x\rangle d \mathcal{H}^{d-1} \simeq \varepsilon^{2} \frac{d}{2} \int_{\partial B} \psi(x)^{2}\langle a, x\rangle d \mathcal{H}^{d-1} .
$$

Choosing $a=\mathbf{e}_{i}$, i.e. the coordinate directions, the previous formula implies that the barycentre of $E_{\varepsilon}$ coincides with the origin, up to an error of order $\varepsilon^{2}$. Since the barycentre of $B\left(x_{0}, r_{\varepsilon}\right)$ is given by its centre $x_{0}$, we then get

$$
\begin{aligned}
\left|B\left(x_{0}, r_{\varepsilon}\right)\right|\left|x_{0}\right|=\left|\int_{B\left(x_{0}, r_{\varepsilon}\right)} y d y\right| & \leq\left|\int_{B\left(x_{0}, r_{\varepsilon}\right)} y d y-\int_{E_{\varepsilon}} y d y\right|+\left|\int_{E_{\varepsilon}} y d y\right| \\
& \leq \int_{B\left(x_{0}, r_{\varepsilon}\right) \Delta E_{\varepsilon}}|y| d y+\left|\int_{E_{\varepsilon}} y d y\right| \leq C\left|B\left(x_{0}, r_{\varepsilon}\right) \Delta E_{\varepsilon}\right|+C \varepsilon^{2}
\end{aligned}
$$

for some constant $C$ independent of $\varepsilon$. In other words, we get $\left|x_{0}\right| \leq C \mathcal{A}\left(E_{\varepsilon}\right)+C \varepsilon^{2}-$ possibly with a different constant $C$, but still independent of $\varepsilon$ - then we can estimate

$$
\begin{equation*}
\frac{\left|B \Delta E_{\varepsilon}\right|}{\left|E_{\varepsilon}\right|} \leq \frac{\left|E_{\varepsilon} \Delta B\left(x_{0}, r_{\varepsilon}\right)\right|}{\left|E_{\varepsilon}\right|}+\frac{\left|B\left(x_{0}, r_{\varepsilon}\right) \Delta B\right|}{\left|E_{\varepsilon}\right|} \leq \mathcal{A}\left(E_{\varepsilon}\right)+C^{\prime}\left|x_{0}\right|+C^{\prime \prime} \varepsilon^{2}, \tag{5.19}
\end{equation*}
$$

for some $C^{\prime}, C^{\prime \prime}$ not depending on $\varepsilon$ : here we used that $\left|B\left(x_{0}, r_{\varepsilon}\right) \Delta B\right|$ is comparable to the distance of their centres - that is, comparable to $\left|x_{0}\right|$ - up to an error of order $\varepsilon^{2}$, due to the difference of the measures. It is only left to use the estimate on $\left|x_{0}\right|$ in (5.19), in conjunction with (5.18) and the fact that $\left|E_{\varepsilon} \Delta B\right| \simeq \varepsilon$ : we then get

$$
\frac{1}{c^{\prime}} \varepsilon \leq \mathcal{A}\left(E_{\varepsilon}\right)+C \varepsilon^{2} \leq c^{\prime} \varepsilon
$$

for some $c^{\prime}>1$, which finally gives $\mathcal{A}\left(E_{\varepsilon}\right) \simeq \varepsilon$, as desired.
Remark 5.11 (Meaning of the key assumptions). We point out that conditions (5.15) and (5.16) are equivalent to require that $\psi$ is orthogonal in the $L^{2}(\partial B)$ sense to the first three eigenspace of the Laplace-Beltrami operator on $\partial B$, i.e. to spherical harmonics of order 0,1 and 2 respectively (see [99] for a comprehensive account on spherical harmonics). Each of these conditions will play a precise role in our construction: thanks to the previous result, the first one implies that $E_{\varepsilon}$ has the same measure as $B$, up to an error of order $\varepsilon^{2}$. The second condition in (5.15) implies that $E_{\varepsilon}$ has the same barycentre as $B$, still up to an error of order $\varepsilon^{2}$ : then this order coincides with the magnitude of $\mathcal{A}\left(E_{\varepsilon}\right)^{2}$. Eventually, recalling that every Stekloff eigenfunction $\xi$ relative to $\sigma_{2}(B)$ has the form $\xi(x)=\langle a, x\rangle$, condition (5.16) implies

$$
\begin{equation*}
\int_{\partial B} \psi(x)|\xi(x)|^{2} d \mathcal{H}^{d-1}=0 \quad \text { and } \quad \int_{\partial B} \psi(x)\left|\nabla_{\tau} \xi(x)\right|^{2} d \mathcal{H}^{d-1}=0 \tag{5.20}
\end{equation*}
$$

where $\nabla_{\tau}$ is the tangential gradient. Relations (5.20) will be crucially exploited in order to prove that $\sigma_{2}(B)-\sigma_{2}\left(E_{\varepsilon}\right) \simeq \varepsilon^{2}$.

Let us fix now an eigenfunction $u_{\varepsilon}$ for $\sigma_{2}\left(E_{\varepsilon}\right)$, normalized in such a way that

$$
\begin{equation*}
\int_{\partial E_{\varepsilon}} u_{\varepsilon}(x)^{2} d \mathcal{H}^{d-1}=1 \quad \text { and } \quad \int_{E_{\varepsilon}}\left|\nabla u_{\varepsilon}(x)\right|^{2} d x=\sigma_{2}\left(E_{\varepsilon}\right) . \tag{5.21}
\end{equation*}
$$

Remark 5.12. Thanks to the fact that $\partial E_{\varepsilon}$ is of class $C^{\infty}$, we obtain that $u_{\varepsilon} \in$ $C^{\infty}\left(\overline{E_{\varepsilon}}\right)$. Moreover, the domains $E_{\varepsilon}$ are uniformly of class $C^{k}$, for every $k \geq 0$, hence we can assume the functions $u_{\varepsilon}$ to satisfy uniform $C^{k}$ estimates, i.e.

$$
\begin{equation*}
\left\|u_{\varepsilon}\right\|_{C^{k}\left(\overline{E_{\varepsilon}}\right)} \leq H_{k} \tag{5.22}
\end{equation*}
$$

for some constants $H_{k}>0$ depending only on $k \in \mathbb{N}$.
We now give the basic estimate of $\sigma_{2}(B)$ from above in terms of $\sigma_{2}\left(E_{\varepsilon}\right)$ : this is the cornerstone of the whole construction.

Lemma 5.13. Let $\varepsilon_{0} \ll 1$, there exist two functions $d, Q:\left[0, \varepsilon_{0}\right] \rightarrow \mathbb{R}$ with

$$
\lim _{\varepsilon \rightarrow 0}(|d(\varepsilon)|+|Q(\varepsilon)|)=0
$$

and a constant $K>0$ such that for every $\varepsilon$, we have

$$
\begin{equation*}
\sigma_{2}(B) \leq \frac{\sigma_{2}\left(E_{\varepsilon}\right)+d(\varepsilon)}{1+Q(\varepsilon)-K \varepsilon^{2}} \tag{5.23}
\end{equation*}
$$

Proof. Since we want to compare $\sigma_{2}\left(E_{\varepsilon}\right)$ with $\sigma_{2}(B)$, we have to suitably adapt the eigenfuction $u_{\varepsilon}$, in order to let it be admissible for the Rayleigh quotient defining $\sigma_{2}(B)$. To do so, we start considering a $C^{k}$ extension $\widetilde{u}_{\varepsilon}$ of $u_{\varepsilon}$ with $k=[d / 2]+3$ to the larger set ${ }^{1}$

$$
D_{\varepsilon}=\left\{x:|x| \leq 1+\varepsilon\|\psi\|_{L^{\infty}(\partial B)}\right\} \supset \overline{B \cup E_{\varepsilon}},
$$

and we can make such an extension in such a way that

$$
\begin{equation*}
\left\|\widetilde{u}_{\varepsilon}\right\|_{C^{k}\left(D_{\varepsilon}\right)} \leq K\left\|u_{\varepsilon}\right\|_{C^{k}\left(E_{\varepsilon}\right)} . \tag{5.24}
\end{equation*}
$$

Then, we estimate the mean value of this extension on the boundary $\partial B$ : we set

$$
\delta:=\int_{\partial B} \widetilde{u}_{\varepsilon}(x) d \mathcal{H}^{d-1},
$$

and we define the application $\phi_{\varepsilon}: \partial B \rightarrow \partial E_{\varepsilon}$, given by

$$
\begin{equation*}
\phi_{\varepsilon}(x)=x+\varepsilon \psi(x) x, \quad x \in \partial B \tag{5.25}
\end{equation*}
$$

Observe that we have

$$
\widetilde{u}_{\varepsilon}\left(\phi_{\varepsilon}(x)\right)=u_{\varepsilon}\left(\phi_{\varepsilon}(x)\right), \quad x \in \partial B,
$$

[^4]so that our uniform estimates (5.22) and (5.24) yield
\[

$$
\begin{equation*}
\widetilde{u}_{\varepsilon}(x)=u_{\varepsilon}\left(\phi_{\varepsilon}(x)\right)+O(\varepsilon), \quad x \in \partial B . \tag{5.26}
\end{equation*}
$$

\]

Using this information in the definition of $\delta$, we get

$$
\delta=\int_{\partial B} u_{\varepsilon}\left(\phi_{\varepsilon}(x)\right) d \mathcal{H}^{d-1}+O(\varepsilon)=\int_{\partial B} u_{\varepsilon}\left(\phi_{\varepsilon}(x)\right) J_{\varepsilon}(x) d \mathcal{H}^{d-1}+O(\varepsilon)
$$

where in the last equality we have set

$$
J_{\varepsilon}(x)=(1+\varepsilon \psi(x))^{d-2} \sqrt{(1+\varepsilon \psi(x))^{2}+\varepsilon^{2}\left|\nabla_{\tau} \psi(x)\right|^{2}}, \quad x \in \partial B,
$$

and we used the following straightforward estimate

$$
\begin{equation*}
\left\|J_{\varepsilon}(y)-1\right\|_{L^{\infty}(\partial B)}=O(\varepsilon), \tag{5.27}
\end{equation*}
$$

the quantity $\nabla_{\tau} \psi$ being the tangential gradient of $\psi$ on $\partial B$. With the change of variable $y=\phi_{\varepsilon}(x)$, we then arrive at

$$
\begin{equation*}
\delta=\frac{1}{\mathcal{H}^{d-1}(\partial B)} \int_{\partial E_{\varepsilon}} u_{\varepsilon}(y) d \mathcal{H}^{d-1}+O(\varepsilon)=O(\varepsilon), \tag{5.28}
\end{equation*}
$$

thanks to the fact that $\int_{\partial E_{\varepsilon}} u_{\varepsilon}=0$. We are now ready to define an admissible function for $\sigma_{2}(B)$ : we set

$$
\begin{equation*}
v_{\varepsilon}:=\widetilde{u}_{\varepsilon} \cdot 1_{\bar{B}}-\delta, \tag{5.29}
\end{equation*}
$$

and we immediately notice that

$$
\begin{equation*}
\left\|v_{\varepsilon}\right\|_{C^{k}(\bar{B})} \leq K(d), \tag{5.30}
\end{equation*}
$$

thanks to (5.22), (5.24) and (5.28) (recall that $k$ depends only on $d$ ). In words, $v_{\varepsilon}$ is the original eigenfunction $u_{\varepsilon}$ extended to the whole $D_{\varepsilon}$, then restricted to the ball $B$ and finally vertically translated in order to satisfy the zero-mean condition on $\partial B$. By its very definition and using (5.28), we immediately observe that

$$
\begin{equation*}
\left|\int_{\partial B} v_{\varepsilon}^{2}-\int_{\partial B} \widetilde{u}_{\varepsilon}^{2}\right|=\left|-2 \delta \int_{\partial B} \widetilde{u}_{\varepsilon}+\delta^{2} \mathcal{H}^{d-1}(\partial B)\right|=\delta^{2} \mathcal{H}^{d-1}(\partial B) \leq K \varepsilon^{2} . \tag{5.31}
\end{equation*}
$$

Now we set

$$
d(\varepsilon):=\int_{B \backslash E_{\varepsilon}}\left|\nabla v_{\varepsilon}\right|^{2}-\int_{E_{\varepsilon} \backslash B}\left|\nabla u_{\varepsilon}\right|^{2},
$$

so that we can write

$$
\begin{equation*}
\int_{B}\left|\nabla v_{\varepsilon}(x)\right|^{2}=\int_{E_{\varepsilon}}\left|\nabla u_{\varepsilon}\right|^{2}+d(\varepsilon)=\sigma_{2}\left(E_{\varepsilon}\right)+d(\varepsilon), \tag{5.32}
\end{equation*}
$$

where we used that $\nabla v_{\varepsilon}=\nabla u_{\varepsilon}$ on $B \cap E_{\varepsilon}$. Moreover, using (5.26) and (5.31), we have

$$
\begin{equation*}
\int_{\partial B} v_{\varepsilon}(x)^{2} \geq \int_{\partial B} \widetilde{u}_{\varepsilon}(x)^{2}-K \varepsilon^{2}=\int_{\partial E_{\varepsilon}} u_{\varepsilon}(x)^{2}+Q(\varepsilon)-K \varepsilon^{2}=1+Q(\varepsilon)-K \varepsilon^{2} \tag{5.33}
\end{equation*}
$$

having defined

$$
Q(\varepsilon):=\int_{\partial B} \widetilde{u}_{\varepsilon}(x)^{2}-\int_{\partial E_{\varepsilon}} u_{\varepsilon}(x)^{2}
$$

We are now able to estimate $\sigma_{2}(B)$ : since

$$
\sigma_{2}(B) \leq \frac{\int_{B}\left|\nabla v_{\varepsilon}(x)\right|^{2} d x}{\int_{\partial B} v_{\varepsilon}(x)^{2} d \mathcal{H}^{d-1}}
$$

using (5.32) and (5.33), we finally obtain (5.23).
Remark 5.14. Thanks to the uniform estimates (5.22) with $k=0,1$ and to (5.27), it is immediate to infer

$$
\begin{equation*}
|d(\varepsilon)| \leq K \varepsilon, \quad|Q(\varepsilon)| \leq K \varepsilon \tag{5.34}
\end{equation*}
$$

which inserted in (5.23) gives the easy estimate

$$
\sigma_{2}(B) \leq \sigma_{2}\left(E_{\varepsilon}\right)+K \varepsilon
$$

possibly with a different constant $K>0$.
The previous observation shows that in order to exhibit the sharp decay rate of the deficit along the sequence $E_{\varepsilon}$, we need a precise control of the decay rate of the error terms $d$ and $Q$. Indeed, each estimate on them automatically translates into an estimate of the same order for $\sigma_{2}(B)-\sigma_{2}\left(E_{\varepsilon}\right)$. Let us state precisely this observation, whose proof is immediate from (5.23).

Lemma 5.15. There exist two constants $C_{1}$ and $C_{2}$ such that

$$
\left|\sigma_{2}(B)-\sigma_{2}\left(E_{\varepsilon}\right)\right| \leq C_{1}(|d(\varepsilon)|+|Q(\varepsilon)|)+C_{2} \varepsilon^{2}, \quad \text { for every } \varepsilon \ll 1
$$

Keeping in mind Corollary 5.7 and (5.17), we know that

$$
\begin{equation*}
C_{3} \varepsilon^{2} \leq B W\left(E_{\varepsilon}\right) \leq C_{4}\left|\sigma_{2}(B)-\sigma_{2}\left(E_{\varepsilon}\right)\right|+C_{5} \varepsilon^{2} \tag{5.35}
\end{equation*}
$$

hence to conclude the optimality of the exponent 2 in (5.11) one would like to enforce (5.34), proving that

$$
|d(\varepsilon)|+|Q(\varepsilon)| \leq K \varepsilon^{2}
$$

### 5.4.2 Step 2: improving the decay rate

In order to gain this improvement, the following Lemma will be of crucial importance. This guarantees that if the distance in $C^{1}$ between $v_{\varepsilon}$ and the eigenspace corresponding to $\sigma_{2}(B)$ has a certain rate of decaying at 0 , then the decays of $d(\varepsilon)$ and $Q(\varepsilon)$ are improved of the same order. It is precisely here, in the proof of this result, that the Key Assumption (5.16) on $\psi$ will heavily come into play.

Lemma 5.16. Let $\omega:[0,1] \rightarrow \mathbb{R}^{+}$be a continuous function such that $t^{2} / K \leq \omega(t) \leq$ $K \sqrt{t}$. Suppose that for every $\varepsilon \ll 1$, there exists an eigenfunction $\xi_{\varepsilon}$ for $\sigma_{2}(B)$ such that

$$
\begin{equation*}
\left\|v_{\varepsilon}-\xi_{\varepsilon}\right\|_{C^{1}(\bar{B})} \leq C \omega(\varepsilon) \tag{5.36}
\end{equation*}
$$

for some constant $C$ independent of $\varepsilon$. Then there exists a constant $C_{6}$, still independent of $\varepsilon$, such that

$$
|d(\varepsilon)|+|Q(\varepsilon)| \leq C_{6} \omega(\varepsilon) \varepsilon \quad \text { for every } \varepsilon \ll 1
$$

Proof. We start estimating the term $|d(\varepsilon)|$ : the computations are similar to that in [18], but we have to pay attention to some extra terms, which come from the fact that we are facing a Stekloff problem.

Using the uniform estimates (5.22) and recalling the definition (5.25) of $\phi_{\varepsilon}$, we have

$$
\left|\nabla u_{\varepsilon}(x)\right|^{2}=\left|\nabla u_{\varepsilon}\left(\phi_{\varepsilon}\left(\frac{x}{|x|}\right)\right)\right|^{2}+O(\varepsilon), \quad \text { for every } x \in E_{\varepsilon} \backslash B
$$

and observe that, alignedting the gradient in its radial and tangential components, the right-hand side can be written as

$$
\begin{aligned}
\left|\nabla u_{\varepsilon}\left(\phi_{\varepsilon}(x /|x|)\right)\right|^{2} & =\left|\partial_{\varrho} u_{\varepsilon}\left(\phi_{\varepsilon}(x /|x|)\right)\right|^{2}+\frac{1}{(1+\varepsilon \psi(x /|x|))^{2}}\left|\nabla_{\tau} u_{\varepsilon}\left(\phi_{\varepsilon}(x /|x|)\right)\right|^{2} \\
& =\left|\partial_{\varrho} u_{\varepsilon}\left(\phi_{\varepsilon}(x /|x|)\right)\right|^{2}+\left|\nabla_{\tau} u_{\varepsilon}\left(\phi_{\varepsilon}(x /|x|)\right)\right|^{2}+O(\varepsilon)
\end{aligned}
$$

Using once again (5.22), the latter in turn can be estimated as follows

$$
\left|\partial_{\varrho} u_{\varepsilon}\left(\phi_{\varepsilon}(x /|x|)\right)\right|^{2}+\left|\nabla_{\tau} u_{\varepsilon}\left(\phi_{\varepsilon}(x /|x|)\right)\right|^{2}=\sigma_{2}\left(E_{\varepsilon}\right)^{2}\left|u_{\varepsilon}(x /|x|)\right|^{2}+\left|\nabla_{\tau} u_{\varepsilon}(x /|x|)\right|^{2}+O(\varepsilon) .
$$

Notice that we also used that $u_{\varepsilon}$ satisfies the boundary condition

$$
\left\langle\nabla u_{\varepsilon}(x), \nu_{E_{\varepsilon}}(x)\right\rangle=\sigma_{2}\left(E_{\varepsilon}\right) u_{\varepsilon}(x), \quad x \in \partial E_{\varepsilon}
$$

and that the normal vector on $\partial E_{\varepsilon}$ is radial up to an error of order $\varepsilon$, since we have

$$
\nu_{E_{\varepsilon}}(x)=\frac{(1+\varepsilon \psi(x /|x|)) x /|x|-\varepsilon \nabla_{\tau} \psi(x /|x|)}{\sqrt{(1+\varepsilon \psi(x /|x|))^{2}+\left|\nabla_{\tau} \psi(x /|x|)\right|^{2}}}=\frac{x}{|x|}+O(\varepsilon), \quad x \in \partial E_{\varepsilon}
$$

Therefore, recalling also that $\left|E_{\varepsilon} \backslash B\right| \simeq \varepsilon$, one obtains

$$
\begin{align*}
\int_{E_{\varepsilon} \backslash B}\left|\nabla u_{\varepsilon}(x)\right|^{2} d x & =\varepsilon \int_{\partial B \cap\{\psi>0\}} \psi(x)\left[\sigma_{2}\left(E_{\varepsilon}\right)^{2} u_{\varepsilon}(x)^{2}+\left|\nabla_{\tau} u_{\varepsilon}(x)\right|^{2}\right] d \mathcal{H}^{d-1}+O\left(\varepsilon^{2}\right) \\
& =\varepsilon \int_{\partial B \cap\{\psi>0\}} \psi(x)\left[\sigma_{2}(B)^{2} v_{\varepsilon}(x)^{2}+\left|\nabla_{\tau} v_{\varepsilon}(x)\right|^{2}\right] d \mathcal{H}^{d-1}+O\left(\varepsilon^{2}\right), \tag{5.37}
\end{align*}
$$

where the last equality comes from the fact that $v_{\varepsilon}=u_{\varepsilon}$ on $E_{\varepsilon} \cap B$ up to the additive constant $\delta$, which is of order $\varepsilon$ thanks to (5.28), and from the fact that $\left|\sigma_{2}(B)-\sigma_{2}\left(E_{\varepsilon}\right)\right| \leq$ $C \varepsilon$. In the very same way, recalling that by definition of $v_{\varepsilon}$ one has

$$
\nabla v_{\varepsilon}\left(\phi_{\varepsilon}(x)\right)=\nabla u_{\varepsilon}\left(\phi_{\varepsilon}(x)\right), \quad \text { for every } x \in \partial B \backslash E_{\varepsilon}
$$

and that the uniform estimates holds also for $v_{\varepsilon}$ by (5.30), one gets

$$
\begin{equation*}
\int_{B \backslash E_{\varepsilon}}\left|\nabla v_{\varepsilon}(x)\right|^{2} d x=-\varepsilon \int_{\partial B \cap\{\psi>0\}} \psi(x)\left[\sigma_{2}(B)^{2} v_{\varepsilon}(x)^{2}+\left|\nabla_{\tau} v_{\varepsilon}(x)\right|^{2}\right] d \mathcal{H}^{d-1}+O\left(\varepsilon^{2}\right) \tag{5.38}
\end{equation*}
$$

Finally, recalling the definition of $d(\varepsilon)$, from (5.36), (5.37) and (5.38) one obtains

$$
\begin{aligned}
|d(\varepsilon)| & \leq \varepsilon \sigma_{2}(B)^{2}\left|\int_{\partial B} \psi(x) v_{\varepsilon}(x)^{2} d \mathcal{H}^{d-1}\right| \\
& +\left.\varepsilon\left|\int_{\partial B} \psi(x)\right| \nabla_{\tau} v_{\varepsilon}(x)\right|^{2} d \mathcal{H}^{d-1} \mid+O\left(\varepsilon^{2}\right) \\
& =\varepsilon \sigma_{2}(B)^{2}\left|\int_{\partial B} \psi(x) \xi_{\varepsilon}(x)^{2} d \mathcal{H}^{d-1}\right| \\
& +\left.\varepsilon\left|\int_{\partial B} \psi(x)\right| \nabla_{\tau} \xi_{\varepsilon}(x)\right|^{2} d \mathcal{H}^{d-1} \mid+C^{\prime} \varepsilon \omega(\varepsilon)+O\left(\varepsilon^{2}\right) \leq \widetilde{C} \varepsilon \omega(\varepsilon)
\end{aligned}
$$

where in the last estimate we used property (5.16).
We now come to the estimate of $|Q(\varepsilon)|$ : remember that this is given by

$$
Q(\varepsilon)=\int_{\partial B}\left[\widetilde{u}_{\varepsilon}(x)^{2}-\widetilde{u}_{\varepsilon}(x+\varepsilon \psi(x) x)^{2} J_{\varepsilon}(x)\right] d \mathcal{H}^{d-1}
$$

i.e. this error term contains a boundary integral, then estimates are a bit different from the Neumann case treated in [18].

In order to handle this term $Q$, for ease of computations it could be more useful to rewrite it as follows

$$
Q(\varepsilon)=Q_{1}(\varepsilon)+Q_{2}(\varepsilon)
$$

where we set

$$
Q_{1}(\varepsilon):=\int_{\partial B}\left[\widetilde{u}_{\varepsilon}(x)^{2}-\widetilde{u}_{\varepsilon}\left(\phi_{\varepsilon}(x)\right)^{2}\right] d \mathcal{H}^{d-1}
$$

and

$$
Q_{2}(\varepsilon):=\int_{\partial B} \widetilde{u}_{\varepsilon}\left(\phi_{\varepsilon}(x /|x|)\right)^{2}\left[1-J_{\varepsilon}(x)\right] d \mathcal{H}^{d-1} .
$$

Let us start with $Q_{1}(\varepsilon)$ : by construction $\nabla \widetilde{u}_{\varepsilon}(x)=\nabla v_{\varepsilon}(x)$, then using the uniform estimates (5.22), (5.24) and the hypotheses (5.36), we have

$$
\begin{aligned}
\left|Q_{1}(\varepsilon)\right| & =\left|\int_{\partial B}\left[\widetilde{u}_{\varepsilon}(x)^{2}-\widetilde{u}_{\varepsilon}\left(\phi_{\varepsilon}(x)\right)^{2}\right] d \mathcal{H}^{d-1}\right| \\
& \leq 2 \varepsilon\left|\int_{\partial B} \widetilde{u}_{\varepsilon}(x) \partial_{\varrho} \widetilde{u}_{\varepsilon}(x) \psi(x) d \mathcal{H}^{d-1}\right|+O\left(\varepsilon^{2}\right) \\
& \leq 2 \varepsilon\left|\int_{\partial B} \xi_{\varepsilon}(x) \partial_{\varrho} \xi_{\varepsilon}(x) \psi(x) d \mathcal{H}^{d-1}\right|+C \omega(\varepsilon) \varepsilon \\
& =2 \varepsilon \sigma_{2}(B)\left|\int_{\partial B} \xi_{\varepsilon}(x)^{2} \psi(x) d \mathcal{H}^{d-1}\right|+C \omega(\varepsilon) \varepsilon,
\end{aligned}
$$

which yields the estimate $\left|Q_{1}(\varepsilon)\right| \leq C \omega(\varepsilon) \varepsilon$, again thanks to property (5.16). Observe that in the last equality we have exploited the fact that $\xi_{\varepsilon}$ satisfies the Stekloff boundary condition. Finally, it is left to estimate the term $Q_{2}(\varepsilon)$ : first of all, we have

$$
1-J_{\varepsilon}(x)=-(d-1) \varepsilon \psi(\vartheta)+O\left(\varepsilon^{2}\right)
$$

while using the definition of $v_{\varepsilon}$, the uniform estimates (5.22) and (5.24) and the fact that $\delta=O(\varepsilon)$, we get

$$
\widetilde{u}_{\varepsilon}\left(\phi_{\varepsilon}(x)\right)=\widetilde{u}_{\varepsilon}(x)+O(\varepsilon)=v_{\varepsilon}(x)+\delta+O(\varepsilon)=v_{\varepsilon}(x)+O(\varepsilon), \quad x \in \partial B .
$$

Inserting these into the definition of $Q_{2}(\varepsilon)$ and using (5.36), we finally obtain

$$
\begin{aligned}
\left|Q_{2}(\varepsilon)\right| & \leq(d-1) \varepsilon\left|\int_{\partial B} v_{\varepsilon}(x)^{2} \psi(x) d \mathcal{H}^{d-1}\right|+O\left(\varepsilon^{2}\right) \\
& \leq(d-1) \varepsilon\left|\int_{\partial B} \xi_{\varepsilon}(x)^{2} \psi(x) d \mathcal{H}^{d-1}\right|+C \omega(\varepsilon) \varepsilon,
\end{aligned}
$$

which concludes the proof, again thanks to property (5.16).
Remark 5.17. Observe that if on the contrary $\psi$ violates condition (5.16), we cannot assure that all the first-order term in the previous estimates cancel out: then we would not get any improvement on $d$ and $Q$. For example, for the case of the ellipsoids $E_{\varepsilon}$ considered in Section 4.2.1, their boundaries can be described as follows

$$
\partial E_{\varepsilon}=\left\{y=\varrho_{\varepsilon}(x) x \in \mathbb{R}^{2}: x \in \partial B \quad \text { and } \quad \varrho_{\varepsilon}(x)=\sqrt{(1+\varepsilon) x_{1}^{2}+\frac{x_{2}^{2}}{1+\varepsilon}}\right\}
$$

and observe that

$$
\varrho_{\varepsilon}(x) \simeq 1+\varepsilon\left(x_{1}^{2}-x_{2}^{2}\right), \quad x \in \partial B .
$$

It is not difficult to see that $\psi(x)=x_{1}^{2}-x_{2}^{2}$ does not satisfy (5.16): and in fact, in analogy with the Neumann case (see [18, Section 5]), one can show that

$$
\sigma_{2}(B)-\sigma_{2}\left(E_{\varepsilon}\right) \simeq \varepsilon
$$

i.e. ellipsoids do not exhibit the sharp decay rate for the Brock-Weinstock inequality.

### 5.4.3 Step 3: nearness estimates

Thanks to the previous step, we know that to improve (5.34) it is sufficient to estimate the $C^{1}$ distance of $v_{\varepsilon}$ from the eigenspace relative to $\sigma_{2}(B)$, in terms of $\varepsilon$ : the main point is that we can perform such an estimation, in terms of $|d(\varepsilon)|$ and $|Q(\varepsilon)|$ themselves. This is the content of the third step.

We start with an easy $W^{1,2}(B)$ estimate, whose proof is based on a Fourier decomposition on the basis $\left\{\xi_{k}\right\}_{k \geq 2}$ of Stekloff eigenfunctions for $B$ : the idea is quite the same as in [18], but an extra difficulty arises, since we cannot directly decompose $v_{\varepsilon}$ in $W^{1,2}$ on the basis $\left\{\xi_{k}\right\}_{k \geq 2}$. Rather, we have to project it on the space of harmonic functions and to control, in terms of $\varepsilon$, both the Dirichlet integral of this projection and the distance between $v_{\varepsilon}$ and the space of harmonic functions.
Lemma 5.18. For every $\varepsilon \ll 1$, there exists an eigenfunction $\xi_{\varepsilon}$ relative to $\sigma_{2}(B)$ such that

$$
\begin{equation*}
\left\|v_{\varepsilon}-\xi_{\varepsilon}\right\|_{W^{1,2}(B)} \leq C \sqrt{|d(\varepsilon)|+|Q(\varepsilon)|}+K \varepsilon, \quad \text { for every } \varepsilon \ll 1 \tag{5.39}
\end{equation*}
$$

for some constant $C$ independent of $\varepsilon$.
Proof. First of all, let us set $f_{\varepsilon}:=\Delta v_{\varepsilon}=\Delta \widetilde{u}_{\varepsilon}$. Thanks to the fact that $\widetilde{u}_{\varepsilon}$ is a $C^{k}$ extension of $u_{\varepsilon}$ and that the latter is harmonic on $E_{\varepsilon} \cap B$, we get that $f_{\varepsilon}$ is a $C^{k-2}$ function on $B$ such that

$$
f_{\varepsilon}(x)=0, \quad x \in E_{\varepsilon} \cap B .
$$

Moreover, on $B \backslash E_{\varepsilon}$ we have

$$
\begin{align*}
\left|f_{\varepsilon}(x)\right| & \leq\left|f_{\varepsilon}\left(\phi_{\varepsilon}\left(\frac{x}{|x|}\right)\right)\right|+\left\|\nabla f_{\varepsilon}\right\|_{L^{\infty}(B)}\left|\phi_{\varepsilon}\left(\frac{x}{|x|}\right)-x\right| \\
& =\left\|\nabla f_{\varepsilon}\right\|_{L^{\infty}(B)}\left|\phi_{\varepsilon}\left(\frac{x}{|x|}\right)-x\right| \leq C\left|\phi_{\varepsilon}\left(\frac{x}{|x|}\right)-\frac{x}{|x|}\right| \leq C \varepsilon\|\psi\|_{L^{\infty}}, \quad x \in B \backslash E_{\varepsilon}, \tag{5.40}
\end{align*}
$$

so that in conclusion $\left\|f_{\varepsilon}\right\|_{L^{\infty}(B)} \leq C \varepsilon$. We now introduce the harmonic projection $\varphi_{\varepsilon}$ of $v_{\varepsilon}$, i.e. $\varphi_{\varepsilon}$ solves

$$
\left\{\begin{aligned}
\Delta \varphi_{\varepsilon} & =0, \quad \text { in } B, \\
\varphi_{\varepsilon} & =v_{\varepsilon},
\end{aligned} \quad \text { on } \partial B,\right.
$$

and observe that we have

$$
\begin{equation*}
\left\|v_{\varepsilon}-\varphi_{\varepsilon}\right\|_{W^{1,2}(B)} \leq C\left\|f_{\varepsilon}\right\|_{L^{2}(B)} \leq C \varepsilon \tag{5.41}
\end{equation*}
$$

where we used the previous estimate on $f_{\varepsilon}$. Since $\varphi_{\varepsilon}$ is harmonic and $v_{\varepsilon}-\varphi_{\varepsilon} \in W_{0}^{1,2}(B)$, we obtain following estimate on the Dirichlet integrals:

$$
\left\|\nabla v_{\varepsilon}-\nabla \varphi_{\varepsilon}\right\|_{L^{2}(B)}^{2}=\int_{B}\left|\nabla v_{\varepsilon}(x)\right|^{2} d x-\int_{B}\left\langle\nabla v_{\varepsilon}, \nabla \varphi_{\varepsilon}\right\rangle d x
$$

Keeping into account (5.41), we finally obtain

$$
\begin{equation*}
\left.\left|\int_{B}\right| \nabla v_{\varepsilon}(x)\right|^{2} d x-\int_{B}\left|\nabla \varphi_{\varepsilon}(x)\right|^{2} d x \mid \leq C \varepsilon^{2} . \tag{5.42}
\end{equation*}
$$

Since $\varphi_{\varepsilon} \in \operatorname{Har}(B)$ - remember the definition (5.5) - we can use a spectral decomposition for it and write

$$
\varphi_{\varepsilon}=\sum_{k \geq 2} \alpha_{k}(\varepsilon) \xi_{k}, \quad \text { where } \quad \alpha_{k}(\varepsilon)=\int_{\partial B} \varphi_{\varepsilon}(x) \xi_{k}(x) d \mathcal{H}^{d-1}, k \geq 2
$$

then

$$
\left\|\varphi_{\varepsilon}\right\|_{L^{2}(\partial B)}^{2}=\sum_{k \geq 2} \alpha_{k}(\varepsilon)^{2} \quad \text { and } \quad\left\|\nabla \varphi_{\varepsilon}\right\|_{L^{2}(B)}^{2}=\sum_{k \geq 2} \sigma_{k}(B) \alpha_{k}(\varepsilon)^{2},
$$

where for the second decomposition we used (5.6). By (5.31) and the definition of $Q(\varepsilon)$, we have

$$
\begin{aligned}
\left|\int_{\partial B} v_{\varepsilon}(x)^{2}-1\right| & \leq\left|\int_{\partial B} \widetilde{u}_{\varepsilon}(x)^{2}-\int_{\partial E_{\varepsilon}} u_{\varepsilon}(x)^{2}\right|+\left|\int_{\partial B} v_{\varepsilon}(x)^{2}-\int_{\partial B} \widetilde{u}_{\varepsilon}(x)^{2}\right| \\
& \leq|Q(\varepsilon)|+K \varepsilon^{2},
\end{aligned}
$$

and since $\varphi_{\varepsilon}=v_{\varepsilon}$ on $\partial B$, the previous implies

$$
\left|\left\|\varphi_{\varepsilon}\right\|_{L^{2}(\partial B)}^{2}-1\right| \leq|Q(\varepsilon)|+K \varepsilon^{2} .
$$

In particular, we get

$$
\left|\sum_{k=2}^{d+1} \alpha_{k}(\varepsilon)^{2}-1\right| \leq \sum_{k \geq d+2} \alpha_{k}(\varepsilon)^{2}+|Q(\varepsilon)|+K \varepsilon^{2},
$$

and multiplying both members by $\sigma_{2}(B)$ we have

$$
\begin{equation*}
\sigma_{2}(B)\left|\sum_{k=2}^{d+1} \alpha_{k}(\varepsilon)^{2}-1\right| \leq \sigma_{2}(B) \sum_{k \geq d+2} \alpha_{k}(\varepsilon)^{2}+c_{1}|Q(\varepsilon)|+K \varepsilon^{2} . \tag{5.43}
\end{equation*}
$$

On the other hand, by (5.32) and (5.42) we have

$$
\begin{aligned}
\left|\left\|\nabla \varphi_{\varepsilon}\right\|_{L^{2}(B)}^{2}-\sigma_{2}(B)\right| & \leq\left|\left\|\nabla v_{\varepsilon}\right\|_{L^{2}(B)}^{2}-\sigma_{2}(B)\right|+\left|\left\|\nabla v_{\varepsilon}\right\|_{L^{2}(B)}^{2}-\left\|\nabla \varphi_{\varepsilon}\right\|_{L^{2}(B)}^{2}\right| \\
& \leq\left|\sigma_{2}\left(E_{\varepsilon}\right)-\sigma_{2}(B)\right|+|d(\varepsilon)|+C \varepsilon^{2} \\
& \leq C(|d(\varepsilon)|+|Q(\varepsilon)|)+K \varepsilon^{2},
\end{aligned}
$$

which can be rewritten as

$$
\left|\sigma_{2}(B)\left(\sum_{k=2}^{d+1} \alpha_{k}(\varepsilon)^{2}-1\right)+\sum_{k \geq d+2} \sigma_{k}(B) \alpha_{k}(\varepsilon)^{2}\right| \leq c_{2}(|d(\varepsilon)|+|Q(\varepsilon)|)+K \varepsilon^{2},
$$

and this implies

$$
\begin{equation*}
\sum_{k \geq d+2} \sigma_{k}(B) \alpha_{k}(\varepsilon)^{2} \leq c_{2}(|d(\varepsilon)|+|Q(\varepsilon)|)+K \varepsilon^{2}+\sigma_{2}(B)\left|\sum_{k=2}^{d+1} \alpha_{k}(\varepsilon)^{2}-1\right| . \tag{5.44}
\end{equation*}
$$

We can now combine (5.43) and (5.44), so to obtain

$$
\sum_{k \geq d+2}\left(\sigma_{k}(B)-\sigma_{2}(B)\right) \alpha_{k}(\varepsilon)^{2} \leq\left(c_{1}+c_{2}\right)(|d(\varepsilon)|+|Q(\varepsilon)|)+K \varepsilon^{2}
$$

Notice that

$$
1-\frac{\sigma_{2}(B)}{\sigma_{k}(B)}>0, \quad k \geq d+2
$$

since $\sigma_{2}(B)$ has multiplicty $d$ and this forms a non-decreasing sequence, then from the previous we can infer

$$
\sum_{k \geq d+2} \sigma_{k}(B) \alpha_{k}(\varepsilon)^{2} \leq C(|d(\varepsilon)|+|Q(\varepsilon)|)+K \varepsilon^{2}
$$

possibly with different constants $C$ and $K$, depending on the spectral gap $\sigma_{d+2}(B)-$ $\sigma_{2}(B)$, but not on $\varepsilon$. If we set

$$
\xi_{\varepsilon}=\sum_{k=2}^{d+1} \alpha_{k}(\varepsilon) \xi_{k}
$$

we have

$$
\left\|\varphi_{\varepsilon}-\xi_{\varepsilon}\right\|_{L^{2}(\partial B)}^{2} \leq \sigma_{d+2}(B)\left\|\nabla v_{\varepsilon}-\nabla \xi_{\varepsilon}\right\|_{L^{2}(B)}^{2}
$$

and

$$
\left\|\nabla \varphi_{\varepsilon}-\nabla \xi_{\varepsilon}\right\|_{L^{2}(B)}^{2}=\sum_{k \geq d+2} \sigma_{k}(B) \alpha_{k}(\varepsilon)^{2} \leq C(|d(\varepsilon)|+|Q(\varepsilon)|)+K \varepsilon^{2}
$$

which yields

$$
\left\|\varphi_{\varepsilon}-\xi_{\varepsilon}\right\|_{W^{1,2}(B)} \leq C \sqrt{|d(\varepsilon)|+|Q(\varepsilon)|}+K \varepsilon
$$

thanks to the fact that $u \mapsto\|u\|_{L^{2}(\partial B)}+\|\nabla u\|_{L^{2}(B)}$ is equivalent to the standard norm of $W^{1,2}(B)$. Finally, it is only left to observe that

$$
\left\|v_{\varepsilon}-\xi_{\varepsilon}\right\|_{W^{1,2}(B)} \leq\left\|\varphi_{\varepsilon}-\xi_{\varepsilon}\right\|_{W^{1,2}(B)}+\left\|v_{\varepsilon}-\varphi_{\varepsilon}\right\|_{W^{1,2}(B)}
$$

thus we have obtained (5.39).
We show how the previous Sobolev estimate (5.39) can be enhanced, replacing the $W^{1,2}(B)$ norm with the $C^{1}$ one.

Lemma 5.19. For every $\varepsilon \ll 1$, there exists an eigenfunction $\xi_{\varepsilon}$ relative to $\sigma_{2}(B)$ such that

$$
\begin{equation*}
\left\|v_{\varepsilon}-\xi_{\varepsilon}\right\|_{C^{1}(\bar{B})} \leq C_{7} \sqrt{|d(\varepsilon)|+|Q(\varepsilon)|}+C_{8} \varepsilon, \quad \text { for every } \varepsilon \ll 1 \tag{5.45}
\end{equation*}
$$

for some positive constants $C_{7}, C_{8}$ independent of $\varepsilon$.
Proof. First of all, let us write down the the Neumann boundary value problems solved by $v_{\varepsilon}$ and $\xi_{\varepsilon}$ : these are given respectively by
$\left\{\begin{array}{cccccc}\Delta v_{\varepsilon} & = & f_{\varepsilon}, & \text { in } \partial B \\ \left\langle\nabla v_{\varepsilon}, \nu\right\rangle & = & \sigma_{2}(B) g_{\varepsilon}, & \text { on } \partial B\end{array} \quad\right.$ and $\quad\left\{\begin{array}{clll}\Delta \xi_{\varepsilon} & = & 0, & \text { in } \partial B \\ \left\langle\nabla \xi_{\varepsilon}, \nu\right\rangle & = & \sigma_{2}(B) \xi_{\varepsilon}, & \text { on } \partial B\end{array}\right.$ where

$$
f_{\varepsilon}(x)=\Delta \widetilde{u}_{\varepsilon}(x), \quad x \in B
$$

and the boundary value $g_{\varepsilon}$ is given by (recall that $\nabla v_{\varepsilon}=\nabla \widetilde{u}_{\varepsilon}$ )

$$
\begin{aligned}
g_{\varepsilon}(x) & =v_{\varepsilon}(x)+\left[u_{\varepsilon}\left(\phi_{\varepsilon}(x)\right)-v_{\varepsilon}(x)\right] \\
& +\left(\frac{\sigma_{2}\left(E_{\varepsilon}\right)}{\sigma_{2}(B)}-1\right) u_{\varepsilon}\left(\phi_{\varepsilon}(x)\right) \\
& +\frac{1}{\sigma_{2}(B)}\left\langle\nabla \widetilde{u}_{\varepsilon}(x)-\nabla u_{\varepsilon}\left(\phi_{\varepsilon}(x)\right), \nu_{E_{\varepsilon}}\left(\phi_{\varepsilon}(x)\right)\right\rangle \\
& +\frac{1}{\sigma_{2}(B)}\left\langle\nabla \widetilde{u}_{\varepsilon}(x), \nu(x)-\nu_{E_{\varepsilon}}\left(\phi_{\varepsilon}(x)\right)\right\rangle=: v_{\varepsilon}(x)+\sum_{i=1}^{4} g_{\varepsilon, i}(x) \quad x \in \partial B .
\end{aligned}
$$

Thus in order to gain informations on the distance between $v_{\varepsilon}$ and $\xi_{\varepsilon}$, it suffices to estimate $f_{\varepsilon}$ and the boundary term $g_{\varepsilon}-\xi_{\varepsilon}$ : indeed, by standard Elliptic Regularity (see [109, Proposition 7.5]) and by the triangular inequality, for every $k \geq 1$ we have

$$
\begin{align*}
\left\|v_{\varepsilon}-\xi_{\varepsilon}\right\|_{W^{k, 2}(B)} & \leq C\left(\left\|v_{\varepsilon}-\xi_{\varepsilon}\right\|_{L^{2}(B)}+\left\|g_{\varepsilon}-\xi_{\varepsilon}\right\|_{W^{k-3 / 2,2}(\partial B)}+\left\|f_{\varepsilon}\right\|_{W^{k-2,2}(B)}\right) \\
& \leq C\left(\left\|v_{\varepsilon}-\xi_{\varepsilon}\right\|_{L^{2}(B)}+\left\|v_{\varepsilon}-\xi_{\varepsilon}\right\|_{W^{k-3 / 2,2}(\partial B)}\right.  \tag{5.46}\\
& \left.+\sum_{i=1}^{4}\left\|g_{\varepsilon, i}\right\|_{W^{k-3 / 2,2}(\partial B)}+\left\|f_{\varepsilon}\right\|_{W^{k-2,2}(B)}\right)
\end{align*}
$$

The first term on the right-hand side can be easily estimated as follows

$$
\left\|v_{\varepsilon}-\xi_{\varepsilon}\right\|_{L^{2}(B)} \leq\left\|v_{\varepsilon}-\xi_{\varepsilon}\right\|_{W^{1,2}(B)} \leq C \sqrt{|d(\varepsilon)|+|Q(\varepsilon)|}+K \varepsilon
$$

where we used (5.39) in the second inequality: then to obtain (5.45) it suffices to prove that

$$
\begin{gather*}
\left\|v_{\varepsilon}-\xi_{\varepsilon}\right\|_{W^{k-3 / 2,2}(\partial B)} \leq C \sqrt{|d(\varepsilon)|+|Q(\varepsilon)|}+K \varepsilon  \tag{5.47}\\
\sum_{i=1}^{4}\left\|g_{\varepsilon, i}\right\|_{W^{k-3 / 2,2}(\partial B)} \leq C \varepsilon  \tag{5.48}\\
\left\|f_{\varepsilon}\right\|_{W^{k-2,2}(B)} \leq C \varepsilon \tag{5.49}
\end{gather*}
$$

with $k=[d / 2]+2$. Indeed, using the Sobolev Imbedding Theorem, this would yield

$$
\left\|v_{\varepsilon}-\xi_{\varepsilon}\right\|_{C^{1}(\bar{B})} \leq C\left\|v_{\varepsilon}-\xi_{\varepsilon}\right\|_{W^{[d / 2]+2,2}(B)}
$$

and combining (5.46) and (5.47)-(5.49), we would conclude the proof.
We now begin to estimate the terms $g_{\varepsilon, i}$ : recalling that $u_{\varepsilon} \circ \phi_{\varepsilon}=\widetilde{u}_{\varepsilon} \circ \phi_{\varepsilon}$ on $\partial B$ and using (5.26) and the uniform estimates on $\widetilde{u}_{\varepsilon}$, we get that

$$
\left\|g_{\varepsilon, 1}\right\|_{W^{k-3 / 2,2}(\partial B)} \leq\left\|\widetilde{u}_{\varepsilon} \circ \phi_{\varepsilon}-\widetilde{u}_{\varepsilon}\right\|_{W^{k-3 / 2,2}(\partial B)}+\delta\left(\mathcal{H}^{d-1}(\partial B)\right)^{1 / 2}=O(\varepsilon)
$$

For the second, we use (5.22) and Lemma 5.15, to obtain

$$
\left\|g_{\varepsilon, 2}\right\|_{W^{k-3 / 2,2}(\partial B)} \leq K \frac{\left|\sigma_{2}\left(E_{\varepsilon}\right)-\sigma_{2}(B)\right|}{\sigma_{2}(B)} \leq K(|d(\varepsilon)|+|Q(\varepsilon)|)
$$

possibly with a different constant $K$, still not depending on $\varepsilon$. For the the third term, we just use a triangular inequality and the uniform estimates (5.22), (5.24)

$$
\begin{aligned}
\left\|g_{\varepsilon, 3}\right\|_{W^{k-3 / 2,2}(\partial B)} & \leq C\left\|\nabla \widetilde{u}_{\varepsilon}-\nabla u_{\varepsilon} \circ \phi_{\varepsilon}\right\|_{W^{k-3 / 2,2}(\partial B)} \\
& \leq C\left\|\nabla \widetilde{u}_{\varepsilon}-\nabla\left(u_{\varepsilon} \circ \phi_{\varepsilon}\right)\right\|_{W^{k-3 / 2,2}(\partial B)} \\
& +C\left\|\nabla\left(u_{\varepsilon} \circ \phi_{\varepsilon}\right)-\nabla u_{\varepsilon} \circ \phi_{\varepsilon}\right\|_{W^{k-3 / 2,2}(\partial B)} \leq C \varepsilon
\end{aligned}
$$

again thanks to the fact that $\widetilde{u}_{\varepsilon} \circ \phi_{\varepsilon}=u_{\varepsilon} \circ \phi_{\varepsilon}$ on $\partial B$. Finally, still using the uniform estimates (5.24) and (5.22), we have

$$
\left\|g_{\varepsilon, 4}\right\|_{W^{k-3 / 2,2}(\partial B)} \leq C\left\|\nu_{B}-\nu_{E_{\varepsilon}} \circ \phi_{\varepsilon}\right\|_{W^{k-3 / 2,2}(\partial B)}
$$

The term $\nu_{E_{\varepsilon}} \circ \phi_{\varepsilon}$ can be explicitly written as

$$
\nu_{E_{\varepsilon}}\left(\phi_{\varepsilon}(x)\right)=\frac{(1+\varepsilon \psi(x)) \nu_{B}(x)-\varepsilon \nabla_{\tau} \psi(x)}{\sqrt{(1+\varepsilon \psi(x))^{2}+\varepsilon^{2}\left|\nabla_{\tau} \psi(x)\right|^{2}}}, \quad x \in \partial B
$$

In this way

$$
\begin{aligned}
\nu_{B}(x)-\nu_{E_{\varepsilon}}\left(\phi_{\varepsilon}(x)\right) & =\nu_{B}(x)\left(1-\frac{1+\varepsilon \psi(x)}{\sqrt{(1+\varepsilon \psi(x))^{2}+\varepsilon^{2}\left|\nabla_{\tau} \psi(x)\right|^{2}}}\right) \\
& -\varepsilon \frac{\nabla_{\tau} \psi(x)}{\sqrt{(1+\varepsilon \psi(x))^{2}+\varepsilon^{2}\left|\nabla_{\tau} \psi(x)\right|^{2}}} .
\end{aligned}
$$

Then observe that

$$
\varphi_{1}(x)=1-\frac{1+\varepsilon \psi(x)}{\sqrt{(1+\varepsilon \psi(x))^{2}+\varepsilon^{2}\left|\nabla_{\tau} \psi(x)\right|^{2}}}, \quad x \in \partial B
$$

and

$$
\varphi_{2}(x)=\varepsilon \frac{\nabla_{\tau} \psi(x)(x)}{\sqrt{(1+\varepsilon \psi(x))^{2}+\varepsilon^{2}\left|\nabla_{\tau} \psi(x)\right|^{2}}}, \quad x \in \partial B
$$

are two $C^{\infty}$ applications on $\partial B$, such that for every $m \in \mathbb{N}$

$$
\left\|\varphi_{i}\right\|_{C^{m}(\partial B)} \leq C_{m} \varepsilon, \quad i=1,2,
$$

where $C_{m}$ is a constant depending on the $C^{m+1}(\partial B)$ norm of $\psi$, but not on $\varepsilon$. This permits to conclude the estimate on $g_{\varepsilon, 4}$ : we finally have

$$
\left\|g_{\varepsilon, 4}\right\|_{W^{k-3 / 2,2}(\partial B)} \leq C\left\|\nu_{B}-\nu_{E_{\varepsilon}} \circ \phi_{\varepsilon}\right\|_{W^{k-3 / 2,2}(\partial B)} \leq C \varepsilon,
$$

so collecting all these estimates we end up with (5.48), for any $k$.
Concerning the term $f_{\varepsilon}$, we have already seen that $\left\|f_{\varepsilon}\right\|_{L^{\infty}(B)} \leq C \varepsilon$ : repeating the argument (5.40) for every derivative and using that the $C^{[d / 2]+1}$ norm of $f_{\varepsilon}$ is uniformly bounded ${ }^{2}$, we obtain

$$
\left\|f_{\varepsilon}\right\|_{C^{k-2}(\bar{B})} \leq C \varepsilon
$$

for $k=[d / 2]+2$, so that the $W^{k-2,2}$ norm is estimated as follows

$$
\left\|f_{\varepsilon}\right\|_{W^{k-2,2}(B)} \leq C\left\|f_{\varepsilon}\right\|_{C^{k-2}(\bar{B})} \leq C \varepsilon .
$$

Finally, we aim to prove (5.47): by the trace inequality and (5.39) we have

$$
\left\|v_{\varepsilon}-\xi_{\varepsilon}\right\|_{W^{1 / 2,2}(\partial B)} \leq C\left\|v_{\varepsilon}-\xi_{\varepsilon}\right\|_{W^{1,2}(B)} \leq C \sqrt{|d(\varepsilon)|+|Q(\varepsilon)|}+K \varepsilon
$$

A first application of (5.46) with $k=2$, gives

$$
\left\|v_{\varepsilon}-\xi_{\varepsilon}\right\|_{W^{2,2}(E)} \leq C \sqrt{|d(\varepsilon)|+|Q(\varepsilon)|}+K \varepsilon
$$

and applying the trace inequality we obtain

$$
\left\|v_{\varepsilon}-\xi_{\varepsilon}\right\|_{W^{3 / 2,2}(\partial B)} \leq C\left\|v_{\varepsilon}-\xi_{\varepsilon}\right\|_{W^{2,2}(E)} \leq C \sqrt{|d(\varepsilon)|+|Q(\varepsilon)|}+K \varepsilon
$$

thus the validity of (5.47) with $k=3$. Finitely many repetitions of the previous argument give (5.47) with $k=[d / 2]+2$ and thus the proof is concluded.

[^5]
### 5.4.4 Step 4: conclusion

Thanks to Lemma 5.15, we know that

$$
\left|\sigma_{2}(B)-\sigma_{2}\left(E_{\varepsilon}\right)\right| \leq C_{1}(|d(\varepsilon)|+|Q(\varepsilon)|)+C_{2} \varepsilon^{2}
$$

First applying Lemma 5.19 and then Lemma 5.16 with $\omega(\varepsilon)=C_{7} \sqrt{|d(\varepsilon)|+|Q(\varepsilon)|}+$ $C_{8} \varepsilon$, we obtain

$$
\begin{equation*}
|d(\varepsilon)|+|Q(\varepsilon)| \leq \widetilde{C} \varepsilon \sqrt{|d(\varepsilon)|+|Q(\varepsilon)|}+\widetilde{C} \varepsilon^{2} . \tag{5.50}
\end{equation*}
$$

Let us set

$$
t(\varepsilon)=\frac{\varepsilon}{\sqrt{|d(\varepsilon)|+|Q(\varepsilon)|}},
$$

then from (5.50) we can infer

$$
\frac{1}{\widetilde{C}} \leq t(\varepsilon)+t(\varepsilon)^{2}
$$

which easily implies that $t(\varepsilon) \geq c$ for some costant $c>0$, i.e.

$$
\sqrt{|d(\varepsilon)|+|Q(\varepsilon)|} \leq \frac{\varepsilon}{c} .
$$

A further application of Lemma 5.15 finally shows that

$$
\left|\sigma_{2}(B)-\sigma_{2}\left(E_{\varepsilon}\right)\right| \leq C \varepsilon^{2},
$$

possibly with a different constant $C$, still independent of $\varepsilon$. Inserting this into (5.35), we can conclude the proof of Theorem 5.9.

## Chapter 6

## A reduction theorem for the stability of Gagliardo-Nirenberg-Sobolev inequalities

### 6.1 Introduction

In this chapter we address the problem of the stability of the Gagliardo-NirenbergSobolev inequalities (briefly: GNS). To lighten the notation, among the whole chapter we will drop the dependence on the set of integration when it is $\mathbb{R}^{d}$, and, analogously, we will always write $\|\cdot\|_{p}$ instead of $\|\cdot\|_{L^{p}\left(\mathbb{R}^{d}\right)}$ to indicate the $L^{p}$ norm of a function on $\mathbb{R}^{d}$.

The sharp Gagliardo-Nirenberg-Sobolev inequality in $\mathbb{R}^{d}$, with $d \geq 2$, takes the form, for a suitable $G=G(d, p, s, q)>0$,

$$
\begin{equation*}
G\|u\|_{q} \leq\|\nabla u\|_{p}^{\theta}\|u\|_{s}^{1-\theta} \tag{6.1}
\end{equation*}
$$

where the parameters $s, q, p$ satisfies

$$
\begin{align*}
& 1<p<d \\
& 1 \leq s<q<p^{\star}, \quad p^{\star}=\frac{d p}{d-p} \quad \text { and }  \tag{6.2}\\
& \frac{\theta}{p^{\star}}+\frac{1-\theta}{s}=\frac{1}{q}
\end{align*}
$$

and where $u$ is taken in $D^{p, s}\left(\mathbb{R}^{d}\right)$, that is the closure of $C_{c}^{\infty}\left(\mathbb{R}^{d}\right)$ under the norm $\|u\|_{D^{p, s}}=\|\nabla u\|_{p}+\|u\|_{s}$. Inequality (6.1) can be derived by combining the interpolation inequality between the $L^{s}, L^{q}$ and $L^{p^{\star}}$ norms on $\mathbb{R}^{d}$ with the Sobolev inequality:

$$
\begin{equation*}
S(n, p)\|u\|_{p^{\star}} \leq\|\nabla u\|_{p} \tag{6.3}
\end{equation*}
$$

## A reduction theorem for the stability of Gagliardo-Nirenberg-Sobolev

Here $S(n, p)$ is the optimal Sobolev constant, namely

$$
\begin{equation*}
S(n, p)=\inf \left\{\frac{\|\nabla u\|_{p}}{\|u\|_{p^{\star}}}: \quad u \in W^{1, p}\left(\mathbb{R}^{d}\right) \backslash\{0\}\right\} \tag{6.4}
\end{equation*}
$$

Explicit formulas for $S(n, p)$ and minimizers in (6.4) are known since the work of Aubin [6] and Talenti [109]. We stress that the same result is not available for the optimal constant and functions in (6.1) with the exception of the one-parameter family

$$
\begin{equation*}
p=2, \quad q=2 t, \quad s=t+1 \tag{6.5}
\end{equation*}
$$

see [48], [43]. What can be said in full generality, is that optimal functions in (6.1) exist and, according to the Pólya-Szegö inequality (see for instance [86, Chapter 3]), are non-negative, radially symmetric functions with decreasing profile (see [23], [106]). Furthermore, they are unique up to translations, rescaling and multiplication by (nonzero) constants. Now we want to address the quantitative stability of inequality (6.1). To state rigorously this problem, we introduce the (GNS) deficit $\delta(u)$ of a function $u \in D^{p, s}\left(\mathbb{R}^{d}\right)$ as

$$
\begin{equation*}
\delta(u)=\frac{\|\nabla u\|_{p}^{\theta}\|u\|_{s}^{1-\theta}}{G\|u\|_{q}}-1 \tag{6.6}
\end{equation*}
$$

and notice that inequality (6.1) reads, in terms of $\delta(u)$, as

$$
\begin{equation*}
\delta(u) \geq 0 \tag{6.7}
\end{equation*}
$$

Then, for a quantitative version of the GNS inequality, we mean an improvement of inequality (6.7) of the form

$$
\begin{equation*}
\delta(u) \geq \kappa_{0} \operatorname{dist}(\mathrm{u}, \mathrm{M})^{\alpha_{0}} \tag{6.8}
\end{equation*}
$$

$\kappa_{0}, \alpha_{0}$ are positive constants independent of $u$ and $\operatorname{dist}(\cdot, \mathrm{M})$ indicates an appropriate distance from $M$, the set of the optimizers for (6.1). The concept of distance we will adopt is the following (note the analogy with the Fraenkel asymmetry defined in Chapter 4 and Chapter 5):

$$
\begin{equation*}
\lambda(u)=\inf \left\{\frac{\|u-v\|_{q}^{q}}{\|u\|_{q}^{q}}: v \text { is optimal for }(6.1), \quad\|v\|_{q}=\|u\|_{q}\right\} \tag{6.9}
\end{equation*}
$$

Results in this direction have been recently obtained with some ad-hoc techniques valid for special classes of parameters among those in (6.2), see [34] and [35]. In particular, the parameters considered in [34] are those introduced in (6.5) (although the authors focus, because of their later applications, just on the particular case $p=2, q=6, s=4$ ), and the knowledge of minimizers is exploited in a crucial way. In [35] the authors address a class of parameters, $p=s=2, q>2$, for which the minimizers are not explicitly known and they follow a strategy developed by Bianchi and Egnell in [12], which heavily relies on the Hilbertian structure corresponding to $p=2$ and seems complicate to
generalize. Thus the above techniques seem adaptable to prove the stability of the GNS inequalities only for a particular class of parameters. In this chapter, thanks to a general symmetrization technique introduced by Cianchi, Fusco, Maggi and Pratelli in [39], we are able to prove a reduction principle which is valid for the whole class of parameters (6.2). Namely we reduce the problem to that of showing the stability just for radial symmetric functions, reducing the complexity of the task in its generality from a $d$-dimensional to a 1 -dimensional problem. Although this does not solve completely the issue, it offers a more simple way to attack it.

The main result we shall prove is the following
Theorem 6.1. Consider the functionals $\delta(\cdot)$ and $\lambda(\cdot)$ defined in (6.6) and (6.9) respectively. Suppose that there exist two positive constant $k_{0}$ and $\alpha_{0}$ such that the stability inequality

$$
\begin{equation*}
\delta(u) \geq \kappa_{0} \lambda(u)^{\alpha_{0}} \tag{6.10}
\end{equation*}
$$

holds for any radial non-increasing function $u \in D^{p, s}\left(\mathbb{R}^{d}\right)$. Then there are two positive constants $k_{1}$ and $\alpha_{1}$ such that the inequality

$$
\begin{equation*}
\delta(u) \geq \kappa_{1} \lambda(u)^{\alpha_{1}} \tag{6.11}
\end{equation*}
$$

holds true for any function in $D^{p, s}\left(\mathbb{R}^{d}\right)$.
The first step needed to prove Theorem 6.1 is a sort of continuity at 0 of the asymmetry $\lambda$ with respect to the deficit $\delta$. Namely we will prove, in Corollary 6.5 of Section 6.2, that given a sequence of functions $\left(u_{h}\right)_{h}$ such that $\delta\left(u_{h}\right)$ converges to 0 as $h$ goes to infinity, then also $\lambda\left(u_{h}\right)$ converges to 0 . This will be done by means of the compactness Theorem 6.2, where it is proved that a sequence of functions whose deficits are infinitesimal, up to be (suitably) rescaled and translated, is compact in $L^{q}\left(\mathbb{R}^{d}\right)$. Then we will pass to the proof of Theorem 6.1. Its proof is done in Section 6.3 and Section 6.4, each of them devoted to obtain a simplification of the class of the functions we deal with. In particular in Section 6.3 we prove a further reduction step, stating that if the stability inequality (6.10) holds for radial decreasing functions, it holds as well for $d$-symmetric functions, that is functions which are symmetric with respect to $d$ orthogonal hyperplanes. More precisely we prove that if there exist positive constants $\kappa_{0}$ and $\alpha_{0}$ such that for any radial decreasing function $u \in D^{p, s}\left(\mathbb{R}^{d}\right)$ inequality (6.10) holds true, then there exist positive constants $\widetilde{\kappa}_{0}$ and $\widetilde{\alpha}_{0}$ depending on $d, p, q$ and $s$ such that for any $d$-symmetric function $u \in D^{p, s}\left(\mathbb{R}^{d}\right)$ we have

$$
\begin{equation*}
\delta(u) \geq \widetilde{\kappa}_{0} \lambda(u)^{\widetilde{\alpha}_{0}} . \tag{6.12}
\end{equation*}
$$

Eventually, in Section 6.4, we prove that to get the stability of GNS inequality, it is not restrictive to consider only $d$-symmetric functions. Namely we prove the existence of two positive constants $\kappa_{2}$ and $\alpha_{2}$ such that for every function $u \in D^{p, s}\left(\mathbb{R}^{d}\right)$ there exists a $d$-symmetric function $\bar{u}$ such that the following reduction inequalities hold true

$$
\begin{equation*}
\lambda(u) \leq \kappa_{2} \lambda(\bar{u})^{\alpha_{2}}, \quad \delta(\bar{u}) \leq \kappa_{2} \delta(u)^{\alpha_{2}} . \tag{6.13}
\end{equation*}
$$

## A reduction theorem for the stability of Gagliardo-Nirenberg-Sobolev

It is then easy to see that combining (6.12) with the reduction inequalities (6.13) we prove the claim of Theorem 6.1.

### 6.2 Continuity of $\lambda$ with respect to $\delta$ via a compactness theorem

We devote this section to the proof of the following Theorem.
Theorem 6.2. Let $\left(u_{h}\right)_{h}$ be a sequence in $D^{p, s}\left(\mathbb{R}^{d}\right)$ such that $\delta\left(u_{h}\right)$ converges to 0 as $h \rightarrow \infty$. Then there exist $\left(\lambda_{h}\right)_{h} \subset(0,+\infty)$ and $\left(x_{h}\right)_{h} \subset \mathbb{R}^{d}$ such that the rescaled sequence

$$
\begin{equation*}
w_{h}(x)=\tau_{\lambda_{h}} u_{h}\left(x-x_{h}\right)=\lambda_{h}^{d / q} u_{h}\left(\lambda_{h}\left(x-x_{h}\right)\right) \tag{6.14}
\end{equation*}
$$

satisfies:
(i) $\delta\left(w_{h}\right)=\delta\left(u_{h}\right) ; \quad \lambda\left(w_{h}\right)=\lambda\left(u_{h}\right)$;
(ii) $\left\|w_{h}\right\|_{q}=\left\|u_{h}\right\|_{q} \quad \forall h \in N$;
(iii) there exist constants $C_{0}, C_{1}>0$, depending only on $d, p, q, s$, such that

$$
\frac{1}{C_{0}} \leq\left\|\nabla w_{h}\right\|_{p} \leq C_{0}, \quad \frac{1}{C_{1}} \leq\left\|w_{h}\right\|_{s} \leq C_{1}
$$

(iv) $w_{h} \rightarrow w$ strongly in $L^{q}\left(\mathbb{R}^{d}\right)$ as $h \rightarrow \infty$ with $w \in D^{p, s}\left(\mathbb{R}^{d}\right)$.

In this chapter the parameters $d, p, s, q$ and $\theta$ are always intended to satisfy conditions (6.11). We start our analysis defining the following functionals:

$$
\begin{equation*}
G(u)=\|\nabla u\|_{p}^{\theta}\|u\|_{s}^{1-\theta}, \quad F(u)=\int|\nabla u|^{p}+\int|u|^{s} \tag{6.15}
\end{equation*}
$$

defined for $u \in D^{p, s}\left(\mathbb{R}^{d}\right)$. Given $m>0$, we consider the infimum problems:

$$
\begin{equation*}
\psi(m)=\inf \left\{G(u):\|u\|_{q}^{q}=m\right\}, \quad \varphi(m)=\inf \left\{F(u):\|u\|_{q}^{q}=m\right\} \tag{6.16}
\end{equation*}
$$

Lemma 6.3. There exists $\eta_{0}=\eta_{0}(d, p, q, s)>0$ with the following property. For any $u \in D^{p, s}\left(\mathbb{R}^{d}\right)$ there exists $\lambda>0$ such that, if $\tau_{\lambda}(u)=\lambda^{d / q} u(\lambda x)$, then

$$
\begin{equation*}
F\left(\tau_{\lambda} u\right)=\eta_{0} G(u)^{k} \quad \text { where } \quad k=q \cdot \frac{n p+p s-n s}{n p+p q-n s}<q \tag{6.17}
\end{equation*}
$$

Proof. For $u \in D^{p, s}\left(\mathbb{R}^{d}\right)$ we have

$$
\begin{equation*}
\left\|\tau_{\lambda} u\right\|_{q}^{q}=\|u\|_{q}^{q}, \quad\left\|\tau_{\lambda} u\right\|_{s}^{s}=\lambda^{-d+d s / q}\|u\|_{s}^{s} \quad \text { and } \quad\left\|\nabla \tau_{\lambda} u\right\|_{p}^{p}=\lambda^{-d+p+\frac{d p}{q}}\|\nabla u\|_{p}^{p} \tag{6.18}
\end{equation*}
$$

hence

$$
F\left(\tau_{\lambda} u\right)=\lambda^{a} A+\lambda^{b} B=f(\lambda)
$$

where $A=\|\nabla u\|_{p}^{p}, B=\|u\|_{s}^{s}, a=-n+p+n p / q$ and $b=-n+n s / q$. The function $f$ attains its minimum at

$$
\lambda_{m}=\left(-\frac{b}{a}\right)^{\frac{1}{a-b}}\left(\frac{B}{A}\right)^{\frac{1}{a-b}}
$$

with the value

$$
f\left(\lambda_{m}\right)=\eta_{0}\left(A^{\theta / p}\right)^{q \nu}\left(B^{(1-\theta) / s}\right)^{q \nu}
$$

where $\eta_{0}=\eta_{0}(d, p, s, q)$ and $\nu=\frac{d p+p s-d s}{d p+p q-d s}$, that is, the claim of the lemma.
Lemma 6.4. There exists $\alpha=\alpha(d, p, q, s) \in(0,1)$ such that

$$
\varphi(m)=m^{\alpha} \varphi(1) \quad \forall m>0 .
$$

In particular $\varphi$ is strictly super-additive in $(0,1)$.
Proof. Let $u \in D^{p, s}\left(\mathbb{R}^{d}\right)$ be such that $\|u\|_{q}^{q}=m$. Let $v=u / m^{1 / q}$ and set $\tau_{\lambda} v(x)=$ $\lambda^{d / q} v(\lambda x)$. Setting $\lambda=m^{\frac{p-s}{d p+p q-d s}}$ we get, after some calculation analogous to those in Lemma 6.3,

$$
F\left(\tau_{\lambda} v\right)=m^{-\alpha} F(u)
$$

where $\alpha=\frac{d p+p s-d s}{d p+p q-d s}$. If we now consider a minimizing sequence $\left(u_{h}\right)_{h}$ for $\varphi(m)$ such that $\|u\|_{q}^{q}=m$ and as above $v_{h}=u_{h} / m^{1 / q}$, we obtain

$$
\varphi(m)=\lim _{h \rightarrow \infty} F\left(u_{h}\right)=m^{\alpha} \lim _{h \rightarrow \infty} F\left(\tau_{\lambda} v_{h}\right) \geq m^{\alpha} \varphi(1) .
$$

The opposite inequality can be proved with an analogous argument considering a sequence $\left(v_{h}\right)_{h}$ minimizing for $\varphi(1)$, and setting $u_{h}=m^{1 / q} v_{h}$.

We pass now to the proof of Theorem 6.2.
Proof of Theorem 6.2. Let $\varphi$ be the function defined in (6.16). We recall the Lions's Concentration-Compactness Theorem (see [89] or [107, Theorem 4.3 and Theorem 4.8]): given a non-negative sequence $\left(\rho_{h}\right)_{h}$ in $L^{1}\left(\mathbb{R}^{d}\right)$ with fixed $L^{1}$ norm, say 1 , there is a subsequence $\left(\rho_{h_{k}}\right)_{k}$ which satisfies one of the following properties:
(1) (concentration) there exists a sequence $\left(y_{k}\right)_{k} \in \mathbb{R}^{d}$ such that for every $\varepsilon>0$ there exists $R \in(0, \infty)$ such that $\int_{B_{R}\left(y_{k}\right)} \rho_{h_{k}} \geq 1-\varepsilon \quad$ for every $k \in \mathbb{N}$
(2) (vanishing) $\lim _{k \rightarrow \infty} \sup _{y \in \mathbb{R}^{d}} \int_{B_{R}(y)} \rho_{h_{k}}=0 \quad \forall \quad 0<R<\infty$
(3) (dichotomy) there exist $\alpha \in(0,1)$ and two sequences $R_{h} \rightarrow+\infty$ and $y_{h} \in \mathbb{R}^{d}$ such that $\int_{B_{R_{h_{k}}}\left(y_{h_{k}}\right)} \rho_{h_{k}} \rightarrow \alpha$ and $\int_{\mathbb{R}^{d} \backslash B_{2 R_{h_{k}}}\left(y_{h_{k}}\right)} \rho_{h_{k}} \rightarrow 1-\alpha$.

## A reduction theorem for the stability of Gagliardo-Nirenberg-Sobolev

Let $u_{h} \in D^{p, s}\left(\mathbb{R}^{d}\right)$ be such that $\delta\left(u_{h}\right) \rightarrow 0$. We can suppose that $\left\|u_{h}\right\|_{q}=1$. Consider $w_{h}\left(x-x_{h}\right)=\tau_{\lambda_{h}} u_{h}\left(x-x_{h}\right)$ where $\lambda_{h}>0$ is defined, for every $h$, as in Lemma 6.3. Thanks to formulas (6.18), each function $w_{h}$ satisfies statements (i) and (ii) of the theorem. Moreover, Lemma 6.3 provides us two positive constants $\eta_{0}$ and $k$ such that $F\left(w_{h}\right)=\eta_{0} G\left(u_{h}\right)^{k}$. Since $\delta\left(w_{h}\right)+1=G\left(w_{h}\right) / G$, and $\delta\left(w_{h}\right)$ tends to 0 , it follows that the sequence $F\left(w_{h}\right) \rightarrow G^{k}=\varphi(1)$ as $h$ tends to $\infty$, where $\varphi$ is defined in (6.16). In particular the sequence $\left(w_{h}\right)_{h}$ must satisfy statement (iii) of the theorem. In order to prove point (iv), we apply the Concentration-Compactness Theorem to the sequence $\left(\left|w_{h}\right|^{q}\right)_{h}$ aiming to exclude cases (2) and (3).

If the sequence vanishes, by Hölder inequality we would get vanishing also for the sequence $\left(\left|w_{h}\right|^{s}\right)_{h}$, since $s<q$. It is not difficult to see that these conditions, together to the equiboundedness of $\left(w_{h}\right)_{h}$ in $D^{p, s}\left(R^{n}\right)$, guarantees that $w_{h} \rightarrow 0$ strongly in $L^{q}$ as $h \rightarrow \infty$ (see for istance [90, Lemma $I .1]$ ). Since $\left\|w_{h}\right\|_{q}=\left\|u_{h}\right\|_{q}=1$, we would get a contradiction. So we can exclude case (2).

The dichotomy case is more complicated and requires a longer analysis. Suppose to have dichotomy for the sequence $\left(\left|w_{h}\right|^{q}\right)_{h}$. Then there exist $\alpha \in(0,1)$ and a sequence of positive numbers $R_{h} \rightarrow \infty$ as $h \rightarrow \infty$ such that

$$
\int_{B_{R_{h}}}\left|w_{h}\right|^{q} \rightarrow \alpha ; \quad \int_{B_{2 R_{h}}^{c}}\left|w_{h}\right|^{q} \rightarrow 1-\alpha ; \quad \int_{B_{2 R_{h}} \backslash B_{R_{h}}}\left|w_{h}\right|^{q} \rightarrow 0
$$

Let $f \in C_{c}^{1}(B(0,2) ;[0,1])$ such that $f=1$ on $B(0,1)$ and consider $f_{h}(x)=f\left(x / R_{h}\right) \in$ $C_{c}^{1}\left(B\left(0,2 R_{h}\right) ;[0,1]\right)$. Choose also $f$ such that $\left|\nabla f_{h}\right| \leq C / R_{h}$ for some $C>0$. Then we have

$$
\begin{aligned}
F\left(w_{h}\right)= & \int\left|\nabla w_{h}\right|^{p}+\int\left|w_{h}\right|^{s} \\
= & \int_{B_{R_{h}}}\left|\nabla\left(f_{h} w_{h}\right)\right|^{p}+\int_{B_{2 R_{h}}^{c}}\left|\nabla\left(\left(1-f_{h}\right) w_{h}\right)\right|^{p}+\int_{B_{2 R_{h}} \backslash B_{R_{h}}}\left|\nabla w_{h}\right|^{p} \\
& +\int_{B_{R_{h}}}\left|f_{h} w_{h}\right|^{s}+\int_{B_{2 R_{h}}^{c}}\left|\left(1-f_{h}\right) w_{h}\right|^{s}+\int_{B_{2 R_{h}} \backslash B_{R_{h}}}\left|w_{h}\right|^{s} \\
= & \int\left|\nabla\left(f_{h} w_{h}\right)\right|^{p}+\int\left|\nabla\left[\left(1-f_{h}\right) w_{h}\right]\right|^{p} \\
& +\int_{B_{2 R_{h}} \backslash B_{R_{h}}}\left[\left|\nabla w_{h}\right|^{p}-\left|\nabla\left(f_{h} w_{h}\right)\right|^{p}-\left|\nabla\left[\left(1-f_{h}\right) w_{h}\right]\right|^{p}\right] \\
& +\int\left[\left|f_{h} w_{h}\right|^{s}+\left|\left(1-f_{h}\right) w_{h}\right|^{s}\right] \\
& +\int_{B_{2 R_{h}} \backslash B_{R_{h}}}\left[\left|w_{h}\right|^{s}-f_{h}^{s}\left|w_{h}\right|^{s}-\left(1-f_{h}\right)^{s}\left|w_{h}\right|^{s}\right] .
\end{aligned}
$$

Since $f_{h}$ assume values only in $[0,1]$, the last integral is non-negative. Neglecting this
quantity and reordering the terms, we get

$$
\begin{align*}
F\left(w_{h}\right) & \geq F\left(f_{h} w_{h}\right)+F\left(\left(1-f_{h}\right) w_{h}\right)-\epsilon(h) \\
& \geq \varphi\left(\left\|f_{h} w_{h}\right\|_{q}^{q}\right)+\varphi\left(\left\|\left(1-f_{h}\right) w_{h}\right\|_{q}^{q}\right)-\epsilon(h) \tag{6.19}
\end{align*}
$$

where

$$
\epsilon(h)=\int_{B_{2 R_{h}} \backslash B_{R_{h}}}\left[\left|\nabla\left(f_{h} w_{h}\right)\right|^{p}+\left|\nabla\left[\left(1-f_{h}\right) w_{h}\right]\right|^{p}-\left|\nabla w_{h}\right|^{p}\right] .
$$

We claim that the error $\epsilon(h) \rightarrow 0$ is controlled from above by a quantity which converges to 0 as $h \rightarrow \infty$. Indeed

$$
\begin{align*}
\epsilon(h) & \left.=\left.\int_{B_{2 R_{h}} \backslash B_{R_{h}}}\left[\left|\nabla\left(f_{h} w_{h}\right)\right|^{p}+\mid \nabla\left[\left(1-f_{h}\right) w_{h}\right)\right]\right|^{p}-\left|\nabla w_{h}\right|^{p}\right] \\
& =\int_{B_{2 R_{h}} \backslash B_{R_{h}}}\left[\left|f_{h} \nabla w_{h}+w_{h} \nabla f_{h}\right|^{p}+\left|\left(1-f_{h}\right) \nabla w_{h}-w_{h} \nabla f_{h}\right|^{p}-\left|\nabla w_{h}\right|^{p}\right] \\
& \leq \int_{B_{2 R_{h}} \backslash B_{R_{h}}}\left[\left(\left|f_{h} \nabla w_{h}+w_{h} \nabla f_{h}\right|+\left|\left(1-f_{h}\right) \nabla w_{h}-w_{h} \nabla f_{h}\right|\right)^{p}-\left|\nabla w_{h}\right|^{p}\right]  \tag{6.20}\\
& \leq \int_{B_{2 R_{h}} \backslash B_{R_{h}}}\left[\left(\left|\nabla w_{h}\right|+2\left|w_{h}\right|\left|\nabla f_{h}\right|\right)^{p}-\left|\nabla w_{h}\right|^{p}\right] \\
& \leq \int_{B_{2 R_{h}} \backslash B_{R_{h}}}\left[\varepsilon C_{p}\left|\nabla w_{h}\right|^{p}+C_{\varepsilon}\left|\nabla f_{h}\right|^{p}\left|w_{h}\right|^{p}\right] \\
& \leq \varepsilon C_{p} \sup _{h \in \mathbb{N}} \int\left|\nabla w_{h}\right|^{p}+C_{\varepsilon} \int_{B_{2 R_{h}} \backslash B_{R_{h}}}\left|\nabla f_{h}\right|^{p}\left|w_{h}\right|^{p} .
\end{align*}
$$

where the first inequality is due to the super additivity of the map $t \mapsto t^{p}$ on $\mathbb{R}_{+}$, the second to the triangle inequality and the third one is the Young inequality of (suitable) parameters $\varepsilon>0$ and $C_{\varepsilon}$. We need to estimate the quantity

$$
g(h)=\int_{B_{2 R_{h}} \backslash B_{R_{h}}}\left|\nabla f_{h}\right|^{p}\left|w_{h}\right|^{p} .
$$

If $p>s$ then interpolating the $L^{p}$ norm of the $w_{h}$ 's between the $L^{s}$ norm and the $L^{p^{\star}}$ norm and recalling that $\left|\nabla f_{h}\right| \leq C / R_{h}$, we get

$$
\int_{B_{2 R_{h}} \backslash B_{R_{h}}}\left|\nabla f_{h}\right|^{p}\left|w_{h}\right|^{p} \leq \frac{C^{p}}{R_{h}^{p}}\left\|w_{h}\right\|_{p^{\star}}^{p \theta}\left\|w_{h}\right\|_{s}^{p(1-\theta)}
$$

Since we already know that $w_{h}$ satisfies the statement $(i i i)$ of the theorem, we get that $g(h) \rightarrow 0$ as $h \rightarrow \infty$. If $s \geq p$ we divide $g(h)$ into two terms:
$g(h)=\int_{\left(B_{2 R_{h}} \backslash B_{R_{h}}\right) \cap\left\{w_{h} \geq 1\right\}}\left|\nabla f_{h}\right|^{p}\left|w_{h}\right|^{p}+\int_{\left(B_{2 R_{h}} \backslash B_{R_{h}}\right) \cap\left\{w_{h}<1\right\}}\left|\nabla f_{h}\right|^{p}\left|w_{h}\right|^{p}=g_{1}(h)+g_{2}(h)$.

## A reduction theorem for the stability of Gagliardo-Nirenberg-Sobolev

Since $q>p$ we have

$$
g_{1}(h) \leq \frac{C^{p}}{R_{h}^{p}} \int_{B_{2 R_{h}} \backslash B_{R_{h}}}\left|w_{h}\right|^{q} \leq \frac{C^{p}}{R_{h}^{p}}
$$

and so $g_{1}(h) \rightarrow 0$ as $h \rightarrow \infty$. Moreover, by Hölder inequality of parameter $p^{\star} / p$, we get

$$
g_{2}(h) \leq \frac{C^{p}}{R_{h}^{p}}\left(\int_{\left(B_{2 R_{h}} \backslash B_{R_{h}}\right) \cap\left\{w_{h}<1\right\}}\left|w_{h}\right|^{p^{\star}}\right)^{p / p^{\star}}\left|\left(B_{2 R_{h}} \backslash B_{R_{h}}\right)\right|^{\frac{1}{\left(p / p^{\star}\right)^{\prime}}}
$$

where

$$
\left(p / p^{\star}\right)^{\prime}=\frac{p / p^{\star}}{-1+p / p^{\star}}=n / p
$$

Since $p^{\star}>q$, we obtain

$$
g_{2}(h) \leq\left(\omega_{d}\left(2^{d}-1\right)\right)^{p / d} C^{p}\left(\int_{\left(B_{2 R_{h}} \backslash B_{R_{h}}\right) \cap\left\{w_{h}<1\right\}}\left|w_{h}\right|^{q}\right)^{p / p^{\star}}
$$

where $\omega_{d}$ is the measure of the unit ball of $\mathbb{R}^{d}$. So also $g_{2}$ converges to 0 as $h \rightarrow \infty$. Thus, passing to the limit in (6.20), first in $h \rightarrow \infty$ and then in $\varepsilon \rightarrow 0$ we obtain that $\epsilon(h) \rightarrow 0$. Since $\left\|w_{h} f_{h}\right\|_{q}^{q}$ and $\left\|\left(1-f_{h}\right) w_{h}\right\|_{q}^{q}$ converges respectively to $\lambda$ and $1-\lambda$, we can conclude thanks to Lemma 6.4 that

$$
\varphi(1) \geq \varphi(\lambda)+\varphi(1-\lambda)>\varphi(1)
$$

obtaining a contradiction.
So we can exclude also the dichotomy phenomenon. Since $w_{h}$ is equibounded in $L^{q}\left(\mathbb{R}^{d}\right)$, we can consider its weak- $L^{q}$ limit $w$. This is also a strong limit in $L^{q}$. Indeed by concentration, up to translations and since $q>1$, we have

$$
1-\varepsilon \leq \lim _{h \rightarrow \infty} \int_{B_{R}}\left|w_{h}\right|^{q}=\int_{B_{R}}|w|^{q} \leq \int|w|^{q} \leq \underline{\lim }_{h \rightarrow \infty} \int\left|w_{h}\right|^{q}=1
$$

This conclude the proof of the theorem.
Corollary 6.5. Consider a sequence $\left(u_{h}\right)_{h} \subset D^{p, s}\left(\mathbb{R}^{d}\right)$ such that $\delta\left(u_{h}\right) \rightarrow \infty$ as $h \rightarrow 0$. Then also $\lambda\left(u_{h}\right) \rightarrow 0$ for $h \rightarrow \infty$.

Proof. As a consequence of Theorem 6.2, up to subsequence and to a rescaling of the form $\tau_{\lambda} u_{h}(x)=\lambda^{d / q} u(\lambda x)$, we can suppose that $u_{h} \rightarrow u$ strongly in $L^{q}\left(\mathbb{R}^{d}\right)$. Since the $\operatorname{map} u \mapsto \lambda(u)$ is strongly continuous in $L^{q}\left(\mathbb{R}^{d}\right), \lambda\left(u_{h}\right)$ converges to $\lambda(u)$. Furthermore, by the semicontinuity of the deficit (under $L^{q}$-convergence), $\delta(u)=0$. Hence $u$ is optimal for (6.1) and $\lambda(u)=0$.

### 6.3 Reduction to $d$-symmetric functions

In this section we will prove that under the hypothesis that inequality (6.10) holds true for radial decreasing functions, up to change the values of $\alpha_{0}$ and $k_{0}$, it holds also for $d$-symmetric functions (recall that a function is $k$-symmetric in $\mathbb{R}^{d}, d \geq k$, if it is symmetric with respect to $k$ mutually orthogonal hyperplanes). We begin with a brief overview of the strategy we shall adopt. Given a function $u \in D^{p, s}\left(\mathbb{R}^{d}\right)$, the natural radial symmetric function to look at is its spherical rearrangement $u^{\star}$ (see [86, Chapter 3] for its the definition and main properties). Suppose that inequality (6.11) holds true for radial symmetric decreasing functions (thus for $u^{\star}$ ). Then by the triangle inequality we get

$$
\lambda(u)^{1 / q} \leq\left\|u-u^{\star}\right\|_{q}+\lambda\left(u^{\star}\right)^{1 / q} \leq\left\|u-u^{\star}\right\|_{q}+\kappa_{0}^{1 / q} \delta\left(u^{\star}\right)^{\alpha_{0} / q} .
$$

We notice that by the Pólya-Szegö inequality we have that $\delta\left(u^{\star}\right) \leq \delta(u)$. But it is not clear if we can estimate the $L^{q}$ distance between $u$ and $u^{\star}$ in terms of $\delta(u)$. Indeed this turns out to be true only if a function is already d-symmetric. We shall prove, in Lemma 6.6, that $\delta(u)$ controls the Pólya-Szegö deficit, defined as

$$
\begin{equation*}
\delta_{P S}(u)=\frac{\|\nabla u\|_{p}-\left\|\nabla u^{\star}\right\|_{p}}{\left\|\nabla u^{\star}\right\|_{p}}, \tag{6.21}
\end{equation*}
$$

and then, in Lemma 6.7, we will obtain an estimate of the $L^{p^{\star}}$ distance between $u$ and $u^{\star}$ in terms of the $L^{p}$ distance between $|\nabla u|$ and $\left|\nabla u^{\star}\right|$.

Lemma 6.6. There exist two positive constants $\delta_{0}$ and $C_{0}$ such that for every $u \in$ $D^{p, s}\left(\mathbb{R}^{d}\right)$ such that $\|u\|_{q}=1$, with $\delta(u) \leq \delta_{0}$, up to the rescales (6.14), we have

$$
\begin{equation*}
\delta_{P S}(u) \leq C \delta(u)^{1 / \theta} \tag{6.22}
\end{equation*}
$$

where $\theta \in(0,1)$ is the parameter introduced in (6.2).
Proof. By Theorem 6.2, up to rescaling it we can suppose that $u$ satisfies properties (i) - (iii) in (6.14). If we choose $\delta(u) \leq 1 / G$, we obtain:

$$
\begin{equation*}
G \leq\left\|\nabla u^{\star}\right\|_{p}^{\theta}\left\|u^{\star}\right\|_{s}^{1-\theta} \leq\|\nabla u\|_{p}^{\theta}\|u\|_{s}^{1-\theta} \leq 1+G . \tag{6.23}
\end{equation*}
$$

Then,

$$
\begin{align*}
G \delta(u) & =\left(\|\nabla u\|_{p}^{\theta}-\left\|\nabla u^{\star}\right\|_{p}^{\theta}\right)\|u\|_{s}^{1-\theta}+G \delta\left(u^{\star}\right) \\
& \geq C_{1}^{\theta-1}\left(\|\nabla u\|_{p}^{\theta}-\left\|\nabla u^{\star}\right\|_{p}^{\theta}\right)+G \delta\left(u^{\star}\right) \tag{6.24}
\end{align*}
$$

where we used the fact that $\|u\|_{q}=1$ (statement (ii)) and $\|u\|_{s} \geq C_{1}^{-1}$ (statement (iii)). By (6.23) there exists a positive constant $c$ such that,

$$
\|\nabla u\|_{p}^{\theta}-\left\|\nabla u^{\star}\right\|_{p}^{\theta} \geq c\left(\|\nabla u\|_{p}-\left\|\nabla u^{\star}\right\|_{p}\right)^{\theta}
$$

Now the conclusion follows from (6.24) and definition (6.21), with $C=c^{1 / \theta}$.

## A reduction theorem for the stability of Gagliardo-Nirenberg-Sobolev

To obtain the desired estimate of the $L^{p}$ distance between $u$ and $u^{\star}$ we shall use the following result, whose proof can be found in [39, Theorem 3].

Lemma 6.7. Let $n \geq 2,1<p<n$ and $z=\max \{p, 2\}$. Then there exists a positive constant $C$ such that

$$
\begin{equation*}
\int\left|u-u^{\star}\right|^{p^{\star}} \leq C\left(\int|u|^{p^{\star}}\right)^{\frac{p}{n}}\left(\int\left|\nabla u^{\star}\right|^{p}\right)^{\frac{1}{z^{\prime}}}\left(\int|\nabla u|^{p}-\int\left|\nabla u^{\star}\right|^{p}\right)^{\frac{1}{z}} \tag{6.25}
\end{equation*}
$$

holds for every non-negative $u \in W^{1, p}\left(\mathbb{R}^{d}\right)$ which is symmetric with respect to the coordinate hyperplanes.

We are now able to proceed with the proof of (6.12) when $u$ is taken over the class of $d$-symmetric functions.

Theorem 6.8. Suppose that there exist positive constants $\kappa_{0}$ and $\alpha_{0}$ such that for any radial decreasing function $u \in D^{p, s}\left(\mathbb{R}^{d}\right)$ inequality (6.11) holds true. Then, there exist positive constants $\kappa_{1}$ and $\alpha_{1}$ depending on $d, p, q$ and $s$ such that for any $d-$ symmetric function $u \in D^{p, s}\left(\mathbb{R}^{d}\right)$ we have

$$
\begin{equation*}
\delta(u) \geq \widetilde{\kappa}_{0} \lambda(u)^{\widetilde{\alpha}_{0}} \tag{6.26}
\end{equation*}
$$

Proof. Since $\lambda(u) \leq 2^{q-1}$, if $\delta>0$ and $\delta(u) \geq \min \{\delta, 1 / G\}$, then we have $\lambda(u) \leq$ $\left(2^{q-1} / \delta\right) \delta(u)$ and so (6.12) holds true with $\kappa_{1}=2^{q-1} / \delta$ and $\alpha_{1}=1$. Hence we may assume that $\delta(u) \leq \min \{\delta, 1 / G\}$ for a suitably small $\delta$. Moreover, by Theorem 6.2 we can suppose $\|u\|_{q}=1$ and $\|u\|_{s} \in\left[1 / C_{0}, C_{0}\right]$ where $C_{0} \geq 1$ is a constant independent of $u$. Keeping in mind these remarks, we divide the proof into two steps:
Step 1: We first assume that $u \geq 0$. In this case we have, by an interpolation inequality, Lemma 6.7, and Theorem 6.2,

$$
\begin{align*}
\int\left|u-u^{\star}\right|^{q} & \leq\left(\int\left|u-u^{\star}\right|^{s}\right)^{(1-\theta) q / s}\left(\int\left|u-u^{\star}\right|^{p^{\star}}\right)^{\theta q / p^{\star}} \\
& \leq C^{\frac{\theta q}{p^{\star}}} 2^{\frac{(1-\theta) q(s-1)}{s}}\left(\int u^{s}\right)\left(\int|u|^{p^{\star}}\right)^{\frac{p \theta q}{n p^{\star}}}\left(\int\left|\nabla u^{\star}\right|^{p}\right)^{\frac{\theta q}{p^{\star} z^{\prime}}}\left(\int|\nabla u|^{p}-\int\left|\nabla u^{\star}\right|^{p}\right)^{\frac{\theta q}{p^{\star} z}} \\
& \leq C(d, p, q, s)\left(\int|\nabla u|^{p}-\int\left|\nabla u^{\star}\right|^{p}\right)^{\gamma} \tag{6.27}
\end{align*}
$$

where $C(d, p, q, s)$ and $\gamma$ are suitable positive constants depending on $d, p, s$ and $q$. Notice that we use the boundedness of the $L^{p}$ norm of $\nabla u$ and of the $L^{s}$ norm of $u$ granted by Theorem 6.2 (up to choose $\delta$ small enough). Moreover we exploited Lemma 6.7 and thus the assumption that $u \geq 0$. If we suppose $\delta(u) \leq \delta \leq G$, again by Theorem 6.2 we get

$$
G \leq C_{2}\left\|\nabla u^{\star}\right\|_{p}^{\theta} \leq C_{2}\|\nabla u\|_{p}^{\theta}=C_{3} G(\delta(u)+1) \leq C_{3} G(1+G)
$$

for suitable positive constants $C_{2}$ and $C_{3}$. Hence there exists $C_{4}>0$ such that

$$
\|\nabla u\|_{p}^{p}-\left\|\nabla u^{\star}\right\|_{p}^{p} \leq C_{4}\left(\|\nabla u\|_{p}-\left\|\nabla u^{\star}\right\|_{p}\right)
$$

By the triangle inequality, estimates (6.22) and (6.27), the Pólya-Szegö inequality and the assumption on radial functions that we have as hypotheses, we can find constants $C_{5}, C_{6}$ and $C_{7}$ independent of $u$ such that

$$
\begin{aligned}
\lambda(u)^{q} & \leq 2^{q-1}\left(\lambda\left(u^{\star}\right)+\left\|u-u^{\star}\right\|_{q}^{q}\right) \\
& \leq C_{5}\left[\delta(u)^{\alpha_{0}}+\left(\int|\nabla u|^{p}-\int\left|\nabla u^{\star}\right|^{p}\right)^{\gamma}\right] \\
& \left.\leq C_{6}\left[\delta(u)^{\alpha_{0}}+\delta(u)^{\gamma / \theta}\right]\right] \leq C_{7} \delta(u)^{\xi} .
\end{aligned}
$$

where $\xi=\min \left\{\alpha_{0}, \gamma / \theta\right\}$.
Step 2: We assume now that $u$ changes sign. In this case consider the positive and the negative part of $u: u^{+}=u 1_{\{u>0\}}$ and $u^{-}=-u 1_{\{u<0\}}$ (here $1_{A}$ denotes the characteristic function of the set $A$ ). By Lemma 6.3 we are provided a positive constant $\lambda=\lambda(u)$ such that

$$
\eta_{0} G(u)^{\kappa}=F\left(\tau_{\lambda} u\right) .
$$

Moreover we have that

$$
F\left(\tau_{\lambda} u^{ \pm}\right) \geq \inf _{\mu} F\left(\tau_{\mu} u^{ \pm}\right)=\eta_{0} G\left(u^{ \pm}\right)^{\kappa}
$$

where $F$ and $G$ are defined in (6.15). So we get

$$
\begin{aligned}
G^{\kappa}(\delta(u)+1)^{\kappa} & =\left(\|\nabla u\|_{p}^{\theta}\|u\|_{s}^{1-\theta}\right)^{\kappa}=\frac{1}{\eta_{0}}\left(\int\left|\nabla \tau_{\lambda} u\right|^{p}+\int\left|\tau_{\lambda} u\right|^{s}\right) \\
& =\frac{1}{\eta_{0}}\left(\int\left|\nabla \tau_{\lambda} u^{+}+\nabla \tau_{\lambda} u^{-}\right|^{p}+\left|\tau_{\lambda} u^{+}+\tau_{\lambda} u^{-}\right|^{s}\right) \\
& =\frac{1}{\eta_{0}}\left(\int\left|\nabla \tau_{\lambda} u^{+}\right|^{p}+\int\left|\nabla \tau_{\lambda} u^{-}\right|^{p}+\int\left|\tau_{\lambda} u^{+}\right|^{s}+\int\left|\tau_{\lambda} u^{-}\right|^{s}\right) \\
& \geq\left(\left\|\nabla u^{+}\right\|_{p}^{\theta}\left\|u^{+}\right\|_{s}^{1-\theta}\right)^{\kappa}+\left(\left\|\nabla u^{-}\right\|_{p}^{\theta}\left\|u^{-}\right\|_{s}^{1-\theta}\right)^{\kappa} \\
& \geq G^{\kappa}\left[\left\|u^{+}\right\|_{q}^{\kappa}+\left\|u^{-}\right\|_{q}^{\kappa}\right]=G^{\kappa}\left[\left\|u^{+}\right\|_{q}^{\kappa}+\left\|u^{-}\right\|_{q}^{\kappa}\right]
\end{aligned}
$$

The last equality is due to the fact that $\tau_{\lambda} u^{+}$and $\tau_{\lambda} u^{-}$have disjoint supports while in the last inequality we exploited the GNS inequality. Let us set $\int\left(u^{+}\right)^{q}=t$ and $\int\left(u^{-}\right)^{q}=1-t$. We can suppose $t \in(0,1)$, since $u$ changes sign. Then the previous formula takes the form

$$
f(t)=\left(t^{\kappa / q}+(1-t)^{\kappa / q}\right)^{1 / \kappa}-1 \leq \delta(u) .
$$

## A reduction theorem for the stability of Gagliardo-Nirenberg-Sobolev

The function $f$ symmetric in $[0,1]$, is null on 0 and 1 and since $\kappa<q$ (by Lemma 6.3) is concave. Thus there exists $a>0$ such that

$$
f(t) \geq \frac{1}{a} \min \{t, 1-t\}
$$

so that

$$
\begin{equation*}
\min \left\{\int\left(u^{+}\right)^{q}, \int\left(u^{-}\right)^{q}\right\} \leq a \delta(u) \tag{6.28}
\end{equation*}
$$

Suppose that the minimum in (6.28) is achieved by $\int\left(u^{-}\right)^{q}$ (being analogous the other case). Since $\delta(|u|)=\delta(u)$, we can conclude, thanks to the triangle inequality and to (6.28), that

$$
\lambda(u)^{1 / q} \leq \lambda(|u|)^{1 / q}+\left(\int\left|u-|u|^{q}\right)^{1 / q} \leq C_{8}\left(\delta(u)^{\xi / q}+\delta(u)^{1 / q}\right) \leq \widetilde{\kappa}_{0} \delta(u)^{\alpha / q}\right.
$$

where $\alpha=\min \{\xi, 1\}$ and $C_{8}$ a positive constant independent of $u$. The last inequality holds for $\delta(u)<1$. So (6.12) holds with $\widetilde{\alpha}_{0}=\alpha / q$. Eventually if $\delta(u) \geq 1$, then inequality (6.26) follows easily (possibly increasing the value of $\tilde{\kappa}_{0}$ ) since $\lambda(u)<2^{q-1}$.

### 6.4 Reduction inequalities

The goal of this section is to prove the reduction inequalities (6.13). Namely we will prove the following result.
Theorem 6.9. Assume the hypotheses of Theorem 6.8. Then there exist two positive constants $\kappa_{2}$ and $\alpha_{2}$ such that for every non-negative function $u \in D^{p, s}\left(\mathbb{R}^{d}\right)$ there exists a d-symmetric function $\bar{u}$ such that the following reduction inequalities hold true

$$
\begin{equation*}
\lambda(u) \leq \kappa_{2} \lambda(\bar{u})^{\alpha_{2}}, \quad \delta(\bar{u}) \leq \kappa_{2} \delta(u)^{\alpha_{2}} \tag{6.29}
\end{equation*}
$$

This will be done arguing similarly to the Sobolev case considered in [39], although some technical modifications shall be needed. We begin recalling that if $v$ is a optimal function for (6.1), then any other optimal function is of the form

$$
v_{a, b, x_{0}}(x)=\operatorname{av}\left(b\left(x-x_{0}\right)\right),
$$

where $a$ and $b$ are non-null constant and $x_{0} \in \mathbb{R}^{d}$. We define the relative asymmetry of a function on an affine subspace $S$ of $\mathbb{R}^{d}$ as

$$
\lambda(u \mid S)=\inf _{\left(a, b, x_{0}\right) \in \mathbb{R}^{2} \times S}\left\{\frac{\left\|u-v_{a, b, x_{0}}\right\|_{q}^{q}}{\|u\|_{q}^{q}}: \text { v optimal for }(6.1),\left\|v_{a, b, x_{0}}\right\|_{q}=\|u\|_{q}\right\}
$$

Next Lemma will show that the infimum in the definition of the relative asymmetry (and so of the asymmetry) is achieved.

Lemma 6.10. Let $u \in D^{p, s}\left(\mathbb{R}^{d}\right)$ and $S$ an affine space contained in $\mathbb{R}^{d}$. Then the infima in the definition of $\lambda(u)$ and $\lambda(u \mid S)$ are achieved.
Proof. Since the two cases are analogous, we show a proof just for the asymmetry. We can suppose without loss of generality that $\|u\|_{q}=1$. Let us start observing that $\lambda(u)<2$. Indeed, if $v$ is a competitor in the definition of $\lambda(u)$, then, up to translate the centre of symmetry of $v$, that $u$ and $v$ do not have disjointed supports. Then

$$
\begin{equation*}
\lambda(u) \leq \int|v-u|^{q}=\int_{\{u>v\}}(u-v)^{q}+\int_{\{v>u\}}(v-u)^{q}<\int u^{q}+\int v^{q}=2 \tag{6.30}
\end{equation*}
$$

Let now $v_{h}(x)=a_{h} v\left(b_{h}\left(x-x_{h}\right)\right)$ be a sequence of functions such that $\left\|v_{h}\right\|_{q}=1$ and $\left\|u-v_{h}\right\|_{q}^{q} \rightarrow \lambda(u)$ as $h \rightarrow \infty$. We want to show that, up to subsequences, $\left(a_{h}, b_{h}, x_{h}\right) \rightarrow$ $\left(a, b, x_{0}\right) \subset \mathbb{R}^{2} \times \mathbb{R}^{d}$ as $h \rightarrow \infty$. We have that

$$
\begin{equation*}
1=\int\left|v_{h}\right|^{q}=a_{h}^{q} \int\left|v\left(b_{h}\left(x-x_{h}\right)\right)\right|^{q}=\frac{a_{h}^{q}}{b_{h}^{n}} \int|v|^{q}=\frac{a_{h}^{q}}{b_{h}^{n}} \tag{6.31}
\end{equation*}
$$

so $a_{h}^{q}=b_{h}^{n}$. Since $v \in L^{q}\left(\mathbb{R}^{d}\right)$ there exists a function $\rho(\varepsilon)$ converging to 0 as $\varepsilon \rightarrow 0$ such that for each $z \in \mathbb{R}^{d}$ we have

$$
\begin{equation*}
\int_{B(z, \varepsilon)}|v|^{q} \leq \rho(\varepsilon), \quad \int_{B(0,1 / \varepsilon)}|v|^{q} \geq 1-\rho(\varepsilon) \tag{6.32}
\end{equation*}
$$

Set now $b_{-}=\underline{\lim }_{h \rightarrow \infty} b_{h}$ and $b_{+}=\overline{\lim }_{h \rightarrow \infty} b_{h}$. We claim that $b_{-}>0$ and that $b_{+}<\infty$. Suppose $b_{-}=0$; then, recalling that $v_{h}$ is a radial function, we have

$$
\begin{align*}
\int_{B(0,1 / \varepsilon)}\left|v_{h}\right|^{q} & \leq \int_{B\left(x_{h}, 1 / \varepsilon\right)}\left|v_{h}\right|^{q}=\frac{a_{h}^{q}}{b_{h}^{n}} \int_{B\left(0, b_{h} / \varepsilon\right)}|v(y)|^{q} d y  \tag{6.33}\\
& =\int_{B\left(0, b_{h} / \varepsilon\right)}|v(y)|^{q} d y
\end{align*}
$$

and the last quantity, up to pass to a subsequence, converge to 0 as $h \rightarrow \infty$. So we can suppose that for a fixed $\varepsilon$ with $h$ big enough,

$$
\begin{equation*}
\int_{B(0,1 / \varepsilon)}\left|v_{h}\right|^{q} \leq \varepsilon \tag{6.34}
\end{equation*}
$$

So, thanks to (6.34) and (6.32) we have

$$
\begin{aligned}
\left\|u-v_{h}\right\|_{q}^{q} & =\int_{B(0,1 / \varepsilon)}\left|u-v_{h}\right|^{q}+\int_{B(0,1 / \varepsilon)^{c}}\left|u-v_{h}\right|^{q} \\
& \geq\left|\left(\int_{B(0,1 / \varepsilon)}|u|^{q}\right)^{1 / q}-\left(\int_{B(0,1 / \varepsilon)}\left|v_{h}\right|^{q}\right)^{1 / q}\right|^{q} \\
& +\left|\left(\int_{B(0,1 / \varepsilon)^{c}}\left|v_{h}\right|^{q}\right)^{1 / q}-\left(\int_{B(0,1 / \varepsilon)^{c}}|u|^{q}\right)^{1 / q}\right|^{q} \\
& \geq\left[(1-\rho(\varepsilon))^{1 / q}-\varepsilon^{1 / q}\right]^{q}+\left[(1-\varepsilon)^{1 / q}-\rho(\varepsilon)^{1 / q}\right]^{q}
\end{aligned}
$$

## A reduction theorem for the stability of Gagliardo-Nirenberg-Sobolev

Passing to the limit in $h$ and then in $\varepsilon$ we obtain that $\lambda(u) \geq 2$, that is a contradiction. Suppose now that $b_{+}=\infty$. Then

$$
\begin{align*}
\int_{B\left(x_{h}, \varepsilon\right)^{c}}\left|v_{h}\right|^{q}= & \int_{B\left(x_{h}, \varepsilon\right)^{c}} a_{h}^{q}\left|v\left(b_{h}\left(x-x_{h}\right)\right)\right|^{q} d x=a_{h}^{q} \int_{B(0, \varepsilon)^{c}}\left|v\left(b_{h} x\right)\right|^{q} d x  \tag{6.35}\\
& =\frac{a_{h}^{q}}{b_{h}^{n}} \int_{B(0, \varepsilon)^{c}}|v(z)|^{q} d z=\int_{B\left(0, b_{h} \varepsilon\right)^{c}}|v|^{q}
\end{align*}
$$

and arguing as before we can suppose that $\int_{B\left(x_{h}, \varepsilon\right)^{c}}\left|v_{h}\right|^{q} \leq \varepsilon$ for $h$ big enough. By (6.32) we get

$$
\begin{aligned}
\left\|u-v_{h}\right\|_{q}^{q} & =\int_{B\left(x_{h}, \varepsilon\right)}\left|u-v_{h}\right|^{q}+\int_{B\left(x_{h}, \varepsilon\right)^{c}}\left|u-v_{h}\right|^{q} \\
& \geq\left|\left(\int_{B\left(x_{h}, \varepsilon\right)}\left|v_{h}\right|^{q}\right)^{1 / q}-\left(\int_{B\left(x_{h}, \varepsilon\right)}|u|^{q}\right)^{1 / q}\right|^{q} \\
& +\left|\left(\int_{B\left(x_{h}, \varepsilon\right)^{c}}|u|^{q}\right)^{1 / q}-\left(\int_{B\left(x_{h}, \varepsilon\right)^{c}}\left|v_{h}\right|^{q}\right)^{1 / q}\right|^{q} \\
& \geq(1-\varepsilon)-\rho(\varepsilon)+(1-\rho(\varepsilon))-\varepsilon
\end{aligned}
$$

so that we obtain again a contradiction. Suppose now that $\left(x_{h}\right)_{h}$ is not bounded and extract a subsequence (not relabelled) such that $\left|x_{h}\right| \rightarrow \infty$. Then given $N>0$, if $h$ is big enough we would get $\int_{B\left(x_{h}, N\right)}|u|^{q} \leq 1 / N$. If we choose $N$ such that $\int_{B\left(x_{h}, N\right)}\left|v_{h}\right| \geq \varepsilon$ for all $h \in \mathbb{N}$, we obtain

$$
\begin{aligned}
\left\|u-v_{h}\right\|_{q}^{q} & \geq \int_{B\left(x_{h}, N\right)}\left|u-v_{h}\right|^{q}+\int_{B\left(x_{h}, N\right)^{c}}\left|u-v_{h}\right|^{q} \\
& \geq\left|\left(\int_{B\left(x_{h}, N\right)}\left|v_{h}\right|^{q}\right)^{1 / q}-\left(\int_{B\left(x_{h}, N\right)}|u|^{q}\right)^{1 / q}\right|^{q} \\
& +\left|\left(\int_{B\left(x_{h}, N\right)^{c}}|u|^{q}\right)^{1 / q}-\left(\int_{B\left(x_{h}, N\right)}\left|v_{h}\right|^{q}\right)^{1 / q}\right|^{q} \\
& \geq\left[(1-\varepsilon)^{1 / q}-\frac{1}{N^{1 / q}}\right]^{q}+\left[\left(1-\frac{1}{N}\right)^{1 / q}-\varepsilon^{1 / q}\right]^{q} .
\end{aligned}
$$

and again we get a contradiction.

Clearly the asymmetry of a function controls its relative asymmetry. But if we consider a $d$-symmetric function is true also the opposite, as shown in next lemma.

Lemma 6.11. Let $u \in D^{p, s}\left(\mathbb{R}^{d}\right) k$-symmetric with respect to $k$ mutually orthogonal hyperplanes and let $S$ be the intersection of such hyperplanes. Then

$$
\lambda(u \mid S) \leq 3^{q} \lambda(u)
$$

Proof. Suppose as usual that $\|u\|_{q}=1$. Let $v_{a, b, x}$ a minimum for $\lambda(u)$. We consider now the orthogonal projection $x_{S}$ of $x$ on $S$ and $y$ the symmetric point of $x_{S}$ with respect to $S$. Notice that since $u$ is symmetric with respect to $S$, also $v_{a, b, y}$ is a minimum for $\lambda(u)$. Moreover, since the minima of the asymmetry are radial symmetric functions with decreasing profile we have

$$
\left\|v_{a, b, x}-v_{a, b, x_{S}}\right\|_{q} \leq\left\|v_{a, b, x}-v_{a, b, y}\right\|_{q}
$$

This observation and the triangle inequality imply

$$
\begin{aligned}
\lambda(u)^{1 / q} & \leq\left\|u-v_{a, b, x_{S}}\right\|_{q} \leq\left\|u-v_{a, b, x}\right\|_{q}+\left\|v_{a, b, x}-v_{a, b, x_{S}}\right\|_{q} \\
& =\lambda(u)^{1 / q}+\left\|v_{a, b, x}-v_{a, b, x_{S}}\right\|_{q} \leq \lambda(u)^{(1 / q)}+\left\|v_{a, b, x}-v_{a, b, y}\right\|_{q} \\
& \leq \lambda(u)^{1 / q}+\left\|v_{a, b, x}-u\right\|_{q}^{q}+\left\|u-v_{a, b, y}\right\|_{q}=3 \lambda(u)^{1 / q}
\end{aligned}
$$

and the conclusion follows.
Next result shows that the $d$-symmetry condition in the last Lemma can in some sense be relaxed:

Lemma 6.12. There exists a constant $C_{0}$ with the following property. Consider a function $u \in L^{q}\left(\mathbb{R}^{d}\right)$, $u \geq 0, H$ an hyperplane of $\mathbb{R}^{d}$ and $H^{+}$and $H^{-}$the half spaces in which $\mathbb{R}^{d}$ is divided by $H$. Suppose that

$$
\int_{H^{+}}|u|^{q}=\int_{H^{-}}|u|^{q}=\frac{1}{2} \int_{H}|u|^{q}
$$

then

$$
\begin{equation*}
\lambda(u \mid H) \leq C_{0} \lambda(u)^{1 / q} \tag{6.36}
\end{equation*}
$$

with a constant $C_{0}$ depending only on $q$ and $d$. Moreover, if $T_{H}$ denote the reflection with respect to $H$ of $\mathbb{R}^{d}$, it holds

$$
\begin{equation*}
\int\left|u \circ T_{H}-u\right|^{q} \leq C_{0}\|u\|_{q}^{q} \lambda(u)^{1 / q} \tag{6.37}
\end{equation*}
$$

Proof. Suppose without loss of generality that $\|u\|_{q}=1$ and let $v_{0}=v_{a, b, x_{0}}$ a minimum for $\lambda(u)$ centred at $x_{0}$. Suppose that $x_{0} \in H^{+}$, being the other case analogous, and let $\bar{x}$ the projection of $x_{0}$ on $H$ and $\bar{v}=v_{a, b, \bar{x}}$. Then

$$
\begin{equation*}
\lambda(u \mid H) \leq \int|u-\bar{u}|^{q} \leq 2^{q-1}\left(\lambda(u)+\int\left|v_{0}-\bar{v}\right|^{q}\right) \tag{6.38}
\end{equation*}
$$

## A reduction theorem for the stability of Gagliardo-Nirenberg-Sobolev

Consider the translated half spaces $K^{ \pm}=H^{ \pm}+\left(x_{0}-\bar{x}\right)$. Since $x_{0} \in H^{+}$it follows that $K^{+} \subseteq H^{+}$and $H^{-} \subseteq K^{-}$. We have that

$$
\frac{1}{2}=\int_{K^{ \pm}} v_{0}^{q}=\int_{H^{ \pm}} u^{q}=\int_{H^{ \pm}} \bar{v}^{q}
$$

and

$$
\int_{H^{-}}\left|v_{0}-\bar{v}\right|^{q}=\int_{K^{+}}\left|v_{0}-\bar{v}\right|^{q} \leq \int_{H^{+}}\left|v_{0}-\bar{v}\right|^{q}
$$

so that

$$
\begin{equation*}
\int\left|v_{0}-\bar{v}\right|^{q} \leq 2 \int_{H^{+}}\left|v_{0}-\bar{v}\right|^{q} \tag{6.39}
\end{equation*}
$$

Since $v_{0} \geq \bar{v}$ on $K^{+}$we get that $\left|v_{0}-\bar{v}\right|^{q} \leq v_{0}^{q}-\bar{v}^{q}$ on $K^{+}$. Then

$$
\begin{align*}
\int_{K^{+}}\left|v_{0}-\bar{v}\right|^{q} & \leq \int_{K^{+}} v_{0}^{q}-\int_{K^{+}} \bar{v}^{q}=\frac{1}{2}-\int_{H^{-}} v_{0}^{q}  \tag{6.40}\\
& \leq C\left(\|u\|_{L^{q}\left(H^{-}\right)}-\left\|v_{0}\right\|_{L^{q}\left(H^{-}\right)}\right) \leq C\left\|u-v_{0}\right\|_{q}=C \lambda(u)^{1 / q}
\end{align*}
$$

for a suitable positive constant $C$. Moreover

$$
\begin{aligned}
\int_{H^{+} \backslash K^{+}}\left|v_{0}-\bar{v}\right|^{q} & \leq 2^{q-1} \int_{H^{+} \backslash K^{+}}\left(v_{0}^{q}+\bar{v}^{q}\right)=2^{q} \int_{H^{+} \backslash K^{+}} v_{0}^{q} \\
& =2^{q}\left[\int_{H^{+}} v_{0}^{q}-\frac{1}{2}\right]=2^{q}\left[\int_{H^{+}} v_{0}^{q}-\int_{H^{+}} u^{q}\right]
\end{aligned}
$$

and reasoning as in (6.40) we obtain

$$
\begin{equation*}
\int_{H^{+} \backslash K^{+}}|v-\bar{v}|^{q} \leq C \lambda(u)^{1 / q} \tag{6.41}
\end{equation*}
$$

Inequality (6.36) is then a consequence of (6.38), (6.39), (6.40), and (6.41). We are left to show inequality (6.37). Let $\hat{u}$ be the minimal function for $\lambda(u \mid H)$. Then

$$
\begin{aligned}
\int_{H^{ \pm}}\left|u \circ T_{H}-u\right|^{q} & \leq 2^{q-1}\left(\int_{H^{ \pm}}\left|u \circ T_{H}-\hat{u}\right|^{q}+\int_{H^{ \pm}}|u-\hat{u}|^{q}\right) \\
& =2^{q-1} \int|u-\hat{u}|^{q}=2^{q-1} \lambda(u \mid H) \leq C_{0} \lambda(u)^{1 / q}
\end{aligned}
$$

Before passing to the proof of Theorem 6.9 we need another technical lemma which, roughly speaking, states that if two optimal functions for the GNS inequality are near in $L^{q}$ norm, then their $L^{q}$ distance on the whole $\mathbb{R}^{d}$ can be controlled just by that on a quarter of $\mathbb{R}^{d}$. Its proof is quite technical but it is essentially based on a Taylor expansion.

Lemma 6.13. Let $u$ be an optimal function for the GNS inequality of parameters $s, q, p$ centred in 0 with $\|u\|_{q}=1$, and set $u_{\alpha, z}(x)=\alpha^{d / q} u\left(\alpha(x-z)\right.$ ) (for simplicity, $u_{1, z}=u_{z}$ ). Consider two orthogonal half spaces $H$ and $K$ containing the origin in their boundaries. There exist two constants $k=k(d, s, q, p)>0$ and $\tilde{\rho} \ll 1$ such that if

$$
\int\left|u_{\lambda, x_{0}}-u_{\mu, y_{0}}\right|^{q} \leq \tilde{\rho}, \quad \int_{H \cap K}\left|u_{\mu, y_{0}}\right|^{q} \geq \frac{1}{8}
$$

then

$$
\begin{equation*}
\int_{H \cap K}\left|u_{\lambda, x_{0}}-u_{\mu, y_{0}}\right|^{q} \geq k \int\left|u_{\lambda, x_{0}}-u_{\mu, y_{0}}\right|^{q} \tag{6.42}
\end{equation*}
$$

Proof. Up to a rotation we can consider $H=\left\{e_{1} \geq 0\right\}$ and $K=\left\{e_{2} \geq 0\right\}$. Let us set $Q=H \cap K=\left\{e_{1} \geq 0, e_{2} \geq 0\right\}$. Consider the sets

$$
T_{k}=\{x: 1 / k<|\nabla u(x)|<M\} \quad M=\max _{\mathbb{R}^{d}}|\nabla u(x)|
$$

By the radial shape of $u$ we know that each $T_{k}$ is a radial set composed of a countable union of centred annuli, i.e. there exists a non-decreasing sequence of positive numbers $\left(r_{k, j}\right)_{j}$ such that

$$
T_{k}=\bigcup_{j \in \mathbb{N}}\left(B_{r_{k, j+1}} \backslash B_{r_{k, j}}\right)
$$

where $B_{r}$ is the ball centred at the origin of radius $r$. Let now $I_{k}=\left\{j \in \mathbb{N}: r_{k, j+1}-\right.$ $\left.r_{k, j} \geq 1 / k\right\}$ and set

$$
S_{k}=\bigcup_{j \in I_{k}}\left(B_{r_{k, j+1}} \backslash B_{r_{k, j}}\right)
$$

We consider, for $z \in \mathbb{R}^{d}$, the greater centred annulus contained in $S_{k} \cap\left(S_{k}+z\right)$ :

$$
\Sigma(k, z)=\left\{x \in S_{k} \cap\left(S_{k}+z\right): \partial B_{|x|} \subseteq S_{k} \cap\left(S_{k}+z\right)\right\}
$$

Notice, that $\Sigma\left(k, z_{1}\right)=\Sigma\left(k, z_{2}\right)$ whenever $\left|z_{1}\right|=\left|z_{2}\right|$ so we may define $\Sigma(k, r)=\Sigma(k, z)$ if $|z|=r$. Clearly $\bigcup_{k \in \mathbb{N}} S_{k}=\mathbb{R}^{d}$. Moreover, since $|\nabla u|$ is continuous, we have that

$$
\begin{equation*}
\bigcup_{|z|>0} \Sigma(k, z)=S_{k} \tag{6.43}
\end{equation*}
$$

Indeed, if $x \in S_{k}$, there exists $r=r(x)>0$ such that $|\nabla u|(y)>1 / k$ for every $y \in B_{r(x)}(x)$, that is $\overline{B_{r(x)}(x)} \subseteq S_{k}$. Thus it is sufficient to choose $|z|<\operatorname{dist}\left(x, \partial S_{k}\right)$ and we get that $\partial B_{|x|} \subseteq S_{k} \cap\left(S_{k}+z\right)$, i.e $x \in \Sigma(k, z)$. With these definitions in mind, we pass to prove inequality (6.42). Up to a change of variables we can consider just the case $\lambda=1+l>1, \mu=1$. For any Borel set $A$ we have:

## A reduction theorem for the stability of Gagliardo-Nirenberg-Sobolev

$$
\begin{align*}
& \int_{A}\left|\lambda^{d / q} u\left(\lambda\left(x-x_{0}\right)\right)-u\left(x-y_{0}\right)\right|^{q}=\int_{A+y_{0}}\left|(1+l)^{d / q} u\left((1+l)\left(x+y_{0}-x_{0}\right)\right)-u(x)\right|^{q} \\
& =\int_{A+y_{0}}\left|(1+l)^{d / q}\left[u(x)+\left\langle\nabla u(x), l x+(1+l)\left(y_{0}-x_{0}\right)\right\rangle\right]-u(x)+R(x)\right|^{q} \\
& =\int_{A+y_{0}} \left\lvert\,\left\langle\nabla u(x), l x+(1+l)\left(y_{0}-x_{0}\right)\right\rangle+\frac{n}{q} l\left\langle\nabla u(x), l x+(1+l)\left(y_{0}-x_{0}\right)\right\rangle\right. \\
& +\frac{n}{q} l u(x)+\left.R(x)\right|^{q} \\
& =\int_{A+y_{0}}|\nabla u(x)|\left\langle\frac{x}{|x|}, y_{0}-x_{0}\right\rangle+l\left[|x||\nabla u(x)|+\frac{n}{q} u(x)\right]+\left.R(x)\right|^{q} \\
& :=\int_{A+y_{0}}\left|E_{x_{0}, y_{0}, l}(x)+R(x)\right|^{q} \tag{6.44}
\end{align*}
$$

where the last inequality is due to the radial symmetry of $\nabla u$ and the error term $R(x)$ is given by

$$
R(x)=(1+l)^{d / q}\left[\frac{\left|\nabla^{2} u(x)\right|}{|x|^{2}}(x \otimes x)\left(l x+(1+l)\left(y_{0}-x_{0}\right)\right)^{2}\right]=O\left(l^{2}\right)+O\left(\left|l x+(1+l)\left(y_{0}-x_{0}\right)\right|^{2}\right)
$$

We notice that there exists $\tilde{\rho}_{1}$ such that for every $x \in S_{k}$

$$
\begin{equation*}
|R(x)| \leq \frac{1}{2}\left|E_{x_{0}, y_{0}, l}(x)\right| \tag{6.45}
\end{equation*}
$$

if $\left|x_{0}\right|+\left|y_{0}\right|+l \leq \tilde{\rho}_{1}$, since $R$ contains any terms of $E_{x_{0}, y_{0}, l}$ with an higher power (so it is of higher order). We aim now to find $\tilde{\rho}$ such that, for $\left|x_{0}\right|+\left|y_{0}\right|+l \leq \tilde{\rho}$, the following chain of inequalities hold true:
$\int\left|E_{x_{0}, y_{0}, l}+R\right|^{q} \leq c_{1} \int_{\Sigma_{k, y_{0}}}\left|E_{x_{0}, y_{0}, l}+R\right|^{q} \leq c_{2} \int_{\Sigma_{k, y_{0} \cap Q}}\left|E_{x_{0}, y_{0}, l}+R\right|^{q} \leq c_{2} \int_{Q}\left|E_{x_{0}, y_{0}, l}+R\right|^{q}$
for suitable constants $c_{1}, c_{2}$ and $k$. This would immediately imply inequality (6.42). We start remarking that

$$
\lim _{k \rightarrow \infty} \int_{S_{k}}\left|E_{x_{0}, y_{0}, l}+R\right|^{q}=\int\left|E_{x_{0}, y_{0}, l}+R\right|^{q}
$$

so there exists $\bar{k}$ such that for $k \geq \bar{k}$ we have

$$
\int_{S_{k}}\left|E_{x_{0}, y_{0}, l}+R\right|^{q} \geq \frac{1}{2} \int\left|E_{x_{0}, y_{0}, l}+R\right|^{q}
$$

Moreover, in view of (6.43) we get,

$$
\int_{\Sigma\left(k, y_{0}\right)}\left|E_{x_{0}, y_{0}, l}+R\right|^{q} \geq \frac{1}{2} \int_{S_{k}}\left|E_{x_{0}, y_{0}, l}+R\right|^{q}
$$

So the first inequality in (6.46) holds with $c_{1}=1 / 4$. Since the last inequality in (6.46) is trivial, we are left to prove the central one. Consider now, by contradiction, a sequence $\left(x_{h}, y_{h}, l_{h}\right) \rightarrow 0 \in \mathbb{R}^{d} \times \mathbb{R}^{d} \times \mathbb{R}$ as $h \rightarrow \infty$ such that (6.46) does not hold. Integrating on the whole strip we get, for $h$ big enough, the following estimate:

$$
\begin{align*}
& \int_{\Sigma\left(k, y_{h}\right)}\left|E_{x_{h}, y_{h}, l_{h}}+R\right|^{q} \leq(3 / 2)^{q} \int_{\Sigma\left(k, y_{h}\right)}\left|E_{x_{h}, y_{h}, l_{h}}\right|^{q} \\
& \leq(3 / 2)^{q} \int_{\Sigma\left(k, y_{h}\right)}\left|M\left\langle\frac{x}{|x|}, y_{h}-x_{h}\right\rangle+l(|x||\nabla u(x)|+(d / q) u(x))\right|^{q} \\
& \leq(3 / 2)^{q} \int_{\Sigma\left(k, y_{h}\right)}|M| y_{h}-x_{h}\left|+l_{h}(|x||\nabla u(x)|+(d / q) u(x))\right|^{q}  \tag{6.47}\\
& \leq(3 / 2)^{q} \int_{\Sigma\left(k, y_{h}\right)}|M| v_{h}\left|+c_{0}\right|^{q}
\end{align*}
$$

where we set $v_{h}=y_{h}-x_{h}$ and

$$
c_{0}=\varlimsup_{\lim _{h \rightarrow \infty}} \max _{\Sigma\left(k, y_{h}\right)}|x||\nabla u(x)|+\frac{n}{q} u(x)>0 .
$$

So there are constant $k_{0}, k_{1}$ such that

$$
\begin{equation*}
\int_{\Sigma\left(k, y_{h}\right)}\left|E_{x_{h}, y_{h}, l_{h}}+R\right|^{q} \leq k_{0} \int_{\Sigma\left(k, y_{h}\right)}\left|k_{1} l+\left|v_{h}\right|^{q}\right. \tag{6.48}
\end{equation*}
$$

By the other hand we get

$$
\begin{aligned}
& \left.\int_{\Sigma\left(k, y_{h}\right) \cap Q}\left|E_{x_{h}, y_{h}, l_{h}}+R\right|^{q} \geq \frac{1}{2^{q}} \int_{\Sigma\left(k, y_{h}\right) \cap Q} \right\rvert\, E_{x_{h}, y_{h},\left.l_{h}\right|^{q}} \\
& =\frac{1}{2^{q}} \int_{\Sigma\left(k, y_{h}\right)}| | \nabla u(x)\left|\left\langle\frac{x}{|x|}, \frac{v_{h}}{\left|v_{h}\right|}\right\rangle\right| v_{h}|+l(|x||\nabla u(x)|+(d / q) u(x))|^{q} .
\end{aligned}
$$

We consider now three possible situation: $\left|v_{h}\right| \ll l_{h}, l_{h} \ll\left|v_{h}\right|$ or $l_{h} \simeq\left|v_{h}\right|$. In the first case we have, thanks to (6.48), that

$$
\int_{\Sigma\left(k, y_{h}\right)}\left|E_{x_{h}, y_{h}, l_{h}}+R\right|^{q} \leq k_{1} \int_{\Sigma\left(k, y_{h}\right)} l^{q}
$$

for a suitable $k_{1}$. Moreover it is easy to find positive constants $k_{2}$ and $k_{3}$ independent of $l_{h}$ and $v_{h}$ such that

$$
\int_{\Sigma\left(k, y_{h}\right) \cap Q}\left|E_{x_{h}, y_{h}, l_{h}}+R\right|^{q} \geq k_{2} \int_{\Sigma\left(k, y_{h}\right) \cap Q} l^{q} \geq k_{3} \int_{\Sigma\left(k, y_{h}\right)} l^{q} .
$$

So in this case (6.46) holds with $c_{2}=\frac{k_{2}}{k_{3}} c_{1}$ (or, in other terms, we get a contradiction). The second case, $l_{h} \ll\left|v_{h}\right|$ can be treated with the same argument, with the only

## A reduction theorem for the stability of Gagliardo-Nirenberg-Sobolev

observation that, for the estimate from above, we must restrict furtherly the set of integration to the set $\left.U=\left\{x \in Q:|\langle x /| x|, v_{0} /\left|v_{0}\right|\right\rangle \mid \geq 1 / 10\right\}$. We are left to study the case where $\left|v_{h}\right| \simeq l_{h}$. If $\underline{\lim }_{h}\left|v_{h}\right| / l_{h} \geq \widetilde{c}>0$, we have that

$$
\int_{\Sigma\left(k, y_{h}\right)}\left|E_{x_{h}, y_{h}, l_{h}}+R\right|^{q}=l_{h}^{q} \int_{\Sigma\left(k, y_{h}\right)}\left|c_{0} \frac{l_{h}}{\left|v_{h}\right|}+M\right|^{q} .
$$

Let us define $\left.V=\left\{x \in Q:|\langle x /| x|, v_{h} /\left|v_{h}\right|\right\rangle \mid \leq \alpha\right\}$ where $\alpha$ is a constant (depending on $\widetilde{c}$ ) that will be chosen later. We have:

$$
\begin{aligned}
& \int_{\Sigma\left(k, y_{h}\right) \cap Q}\left|E_{x_{h}, y_{h}, l_{h}}+R\right|^{q} \geq \int_{\Sigma\left(k, y_{h}\right) \cap V}\left|E_{x_{h}, y_{h}, l_{h}}+R\right|^{q} \\
& \geq \kappa_{0} l^{q} \int_{\Sigma\left(k, y_{h}\right) \cap V}\left|c_{1} \frac{l}{\left|v_{h}\right|}-M \alpha\right|^{q} .
\end{aligned}
$$

Choosing $\alpha$ small enough, since $l_{h} /\left|v_{h}\right| \gg 0$, we can find constant independent of $l_{h}$ and $v_{h}$ such that

$$
\int_{\Sigma\left(k, y_{h}\right) \cap Q}\left|E_{x_{h}, y_{h}, l_{h}}+R\right|^{q} \geq \kappa_{1} l^{q} \int_{\Sigma\left(k, y_{h}\right) \cap V}\left|c_{0} \frac{l_{h}}{\left|v_{h}\right|}+M\right|^{q} \geq \kappa_{2} l^{q} \int_{\Sigma\left(k, y_{h}\right)}\left|c_{0} \frac{l_{h}}{\left|v_{h}\right|}+M\right|^{q},
$$

so again (6.46) holds with $c_{2}=k_{2} / c_{1}$ (thus again a contradiction). Eventually, if $\underline{\lim }_{h} l_{h} /\left|v_{h}\right| \geq \widetilde{c}>0$, we repeat a similar argument integrating, in the estimate from above, over $\left.V_{2}=\left\{x \in Q:|\langle x /| x|, v_{h} /\left|v_{h}\right|\right\rangle \mid \geq 1 / 10\right\}$.

We pass now to prove Theorem 6.9. For the sake of clearness we divide its proof into two parts. We first prove a proposition which provides us a method to transform a generic function in $D^{p, s}\left(\mathbb{R}^{d}\right)$ in an $(d-1)$-symmetric function which satisfies the reduction inequalities (6.29). Then we will see how to obtain the last symmetry.

Proposition 6.14. There exists a positive constant $C$ such that for every function $u \in D^{p, s}\left(\mathbb{R}^{d}\right)$ there is an $(d-1)$-symmetric function $\tilde{u}$ such that

$$
\begin{equation*}
\lambda(u) \leq C \lambda(\tilde{u}), \quad \delta(\tilde{u}) \leq 2^{d-1} \delta(u) \tag{6.49}
\end{equation*}
$$

Proof. As usual, by homogeneity of the deficit and the asymmetry, we can consider $\|u\|_{q}=1$. Moreover we can suppose that $\delta(u)<\bar{\delta}$ for an arbitrary small $\bar{\delta}$. Indeed, if $\delta(u) \geq \bar{\delta}$, let $v$ be a radial (and so $d$-symmetric!) function such that $0<\delta(v)<2^{d-1} \bar{\delta}$. Then

$$
\lambda(u) \leq 2^{q}=\frac{2^{q}}{\lambda(v)} \lambda(v) \leq \bar{C} \lambda(v), \quad \delta(v) \leq 2^{d-1} \bar{\delta} \leq 2^{d-1} \delta(u)
$$

Consider, for $k=1, \ldots d$, the $d$ hyperplanes orthogonal to the coordinate axis such that

$$
\int_{H_{k}^{+}} u^{q}=\int_{H_{k}^{-}} u^{q}=\frac{1}{2}
$$

where $H_{k}^{ \pm}$are the two half spaces in which $\mathbb{R}^{d}$ is divided by $H_{k}$. Denoting $T_{k}$ the reflection with respect to $H_{k}$, we define

$$
\begin{align*}
& u_{k}^{+}(x)=\left\{\begin{array}{ccc}
u(x) & \text { if } & x \in H_{k}^{+} \\
u\left(T_{k}(x)\right) & \text { if } & x \in H_{k}^{-}
\end{array}\right. \\
& u_{k}^{-}(x)=\left\{\begin{array}{cll}
u(x) & \text { if } & x \in H_{k}^{-} \\
u\left(T_{k}(x)\right) & \text { if } & x \in H_{k}^{+}
\end{array}\right. \tag{6.50}
\end{align*}
$$

By construction $u_{k}^{ \pm}$are symmetric with respect to $H_{k}$. We observe now that

$$
\int u^{s}=\frac{\int\left(u_{k}^{+}\right)^{s}+\int\left(u_{k}^{-}\right)^{s}}{2}, \quad \int|\nabla u|^{p}=\frac{\int\left|\nabla u_{k}^{+}\right|^{p}+\int\left|\nabla u_{k}^{-}\right|^{p}}{2}
$$

and since $t \rightarrow t^{1 / p}$ and $t \rightarrow t^{1 / s}$ are concave functions, we have that

$$
\begin{equation*}
\|u\|_{s} \geq \frac{\left\|u_{k}^{+}\right\|_{s}+\left\|u_{k}^{-}\right\|_{s}}{2}, \quad\|\nabla u\|_{p} \geq \frac{\left\|\nabla u_{k}^{+}\right\|_{p}+\left\|\nabla u_{k}^{-}\right\|_{p}}{2} \tag{6.51}
\end{equation*}
$$

By the definition of $\delta(u)$ and since $(x, y) \mapsto x^{\theta} y^{1-\theta}$ is concave on $\mathbb{R}_{+}^{2}$ and strictly increasing in $x$ and $y$, we get that

$$
\begin{aligned}
G \delta(u) & \geq\left(\frac{\left\|\nabla u_{k}^{+}\right\|_{p}+\left\|\nabla u_{k}^{-}\right\|_{p}}{2}\right)^{\theta}\left(\frac{\left\|u_{k}^{+}\right\|_{s}+\left\|u_{k}^{-}\right\|_{s}}{2}\right)^{1-\theta}-G \\
& \geq \frac{G}{2} \delta\left(u_{k}^{+}\right)+\frac{G}{2} \delta\left(u_{k}^{-}\right)
\end{aligned}
$$

and so

$$
\delta(u) \geq \frac{\delta\left(u_{k}^{+}\right)+\delta\left(u_{k}^{-}\right)}{2}
$$

In particular for every $k=1, \ldots, d$

$$
\max \left\{\delta\left(u_{k}^{+}\right), \delta\left(u_{k}^{-}\right)\right\} \leq 2 \delta(u)
$$

Let $v_{k}^{+}$and $v_{k}^{-}$be the functions which minimize $\lambda\left(u_{k}^{ \pm} \mid H_{k}\right)$. Then, by triangle inequality and Lemma 6.11 we have

$$
\begin{aligned}
\lambda(u) & \leq \int\left|u-v_{k}^{+}\right|^{q}=\int_{H_{k}^{+}}\left|u_{k}^{+}-v_{k}^{+}\right|^{q}+\int_{H_{k}^{-}}\left|u_{k}^{-}-v_{k}^{+}\right|^{q} \\
& \leq 2^{q-1}\left(\frac{\lambda\left(u_{k}^{+} \mid H_{k}\right)+\lambda\left(u_{k}^{-} \mid H_{k}\right)}{2}+\int_{H_{k}^{-}}\left|v_{k}^{+}-v_{k}^{-}\right|^{q}\right) \\
& \leq 2^{q-2} 3^{q}\left(\lambda\left(u_{k}^{+}\right)+\lambda\left(u_{k}^{-}\right)+\int_{H_{k}^{-}}\left|v_{k}^{+}-v_{k}^{-}\right|^{q}\right)
\end{aligned}
$$

## A reduction theorem for the stability of Gagliardo-Nirenberg-Sobolev

We claim now that for $\delta(u)$ is small enough, chosen any couple of indexes $1 \leq i<j \leq d$, for $k_{1}=i$ or $k_{1}=j$ the following inequality holds:

$$
\begin{equation*}
\int_{H_{k_{1}}^{-}}\left|v_{k_{1}}^{+}-v_{k_{1}}^{-}\right|^{q} \leq C\left(\int_{H_{k_{1}}^{+}}\left|u_{k_{1}}^{+}-v_{k_{1}}^{+}\right|^{q}+\int_{H_{k_{1}}^{-}}\left|u_{k_{1}}^{-}-v_{k_{1}}^{-}\right|^{q}\right) \tag{6.52}
\end{equation*}
$$

Let us show how this concludes: by (6.51) and (6.52) we would have

$$
\lambda(u) \leq C \max \left\{\lambda\left(u_{k_{1}}^{+}\right), \lambda\left(u_{k_{1}}^{-}\right)\right\}, \quad \max \left\{\delta\left(u_{k_{1}}^{+}\right), \delta\left(u_{k_{1}}^{-}\right)\right\} \leq 2 \delta(u)
$$

So we would obtain, among $u_{k_{1}}^{ \pm}$, a 1 -symmetric function with $L^{q}$ norm equal to 1 , say $u_{k_{1}}^{+}$. We can now iterate this procedure exploiting two hyperplanes between the $(d-1)$ we did not use yet. We would then obtain a 2 -symmetric function which would satisfy the reductions inequalities (6.49) (with respect to $u_{k_{1}}^{+}$). We can iterate such construction until we have just one hyperplane left. But then we would have an $(d-1)$-symmetric function which satisfies inequalities (6.49). Thus we are only left to prove (6.52). We divide such proof into two further steps:

Step 1 There exists a positive constant $C_{0}$ such that for all $1 \leq i<j \leq d, \sigma, \tau \in\{+,-\}$ it holds

$$
\begin{equation*}
\int\left|v_{i}^{\sigma}-v_{j}^{\tau}\right|^{q} \leq C_{0} \int_{H_{i}^{\sigma} \cap H_{j}^{\tau}}\left|v_{i}^{\sigma}-v_{j}^{\tau}\right|^{q} \tag{6.53}
\end{equation*}
$$

Step 2 (6.53) implies (6.52).
To prove (6.53) (thus Step 1) we notice that, thanks to Lemma 6.13, it is verified if there are two positive constants $\rho$ and $C_{1}$ such that:
(i) $\int\left(v_{i}^{\sigma}\right)^{q}=\int\left(v_{j}^{\tau}\right)^{q}=1$;
(ii) $H_{i}^{\sigma}$ ed $H_{j}^{\tau}$ are two orthogonal half spaces which contains on their boundary the centre of symmetry of $v_{i}^{\sigma}$ e $v_{j}^{\tau}$;
(iii) $\int\left|v_{i}^{\sigma}-v_{j}^{\tau}\right|^{q} \leq \rho$.

So we are left to check (iii). We have

$$
\begin{equation*}
\left\|v_{i}^{\sigma}-v_{j}^{\tau}\right\|_{q} \leq\left\|v_{i}^{\sigma}-u_{i}^{\sigma}\right\|_{q}+\left\|u_{i}^{\sigma}-u\right\|_{q}+\left\|u-u_{j}^{\tau}\right\|_{q}+\left\|u_{j}^{\tau}-v_{j}^{\tau}\right\|_{q} \tag{6.54}
\end{equation*}
$$

Thanks to Lemma 6.12

$$
\begin{equation*}
\int\left|u_{i}^{\sigma}-u\right|^{q}=\frac{1}{2} \int\left|u \circ T_{i}-u\right|^{q} \leq C \lambda(u)^{1 / q} \tag{6.55}
\end{equation*}
$$

Moreover

$$
\begin{equation*}
\int\left|v_{i}^{\sigma}-u_{i}^{\sigma}\right|^{q} \leq 2^{q-1}\left(\lambda\left(u \mid H_{i}\right)+\left\|u-u_{i}^{\sigma}\right\|_{p}^{p}\right) \leq C \lambda(u)^{1 / q} \tag{6.56}
\end{equation*}
$$

Clearly the same estimate holds for the two addends not yet considered in (6.54). Putting together (6.55) and (6.56) we obtain (iii) and this conclude the proof of Step 1. Let us prove now Step 2. Suppose to fix the ideas that $i=1$ and $j=2$. For $k=1,2$ we set

$$
h_{k}=v_{k}^{+} 1_{H_{k}^{+}}+v_{k}^{-} 1_{H_{k}^{-}} .
$$

Thanks to (6.53),

$$
\int\left|h_{1}-h_{2}\right|^{q} \geq \int_{H_{1}^{+} \cap H_{2}^{+}}\left|h_{1}-h_{2}\right|^{q}=\int_{H_{1}^{+} \cap H_{2}^{+}}\left|v_{1}^{+}-v_{2}^{+}\right|^{q} \geq \frac{1}{C} \int\left|v_{1}^{+}-v_{2}^{+}\right|^{q}
$$

With an similar argument, using $H_{1}^{-} \cap H_{2}^{+}$instead of $H_{1}^{+} \cap H_{2}^{+}$we get

$$
\int\left|h_{1}-h_{2}\right|^{q} \geq \frac{1}{C} \int\left|v_{1}^{-}-v_{2}^{+}\right|^{q} .
$$

Hence

$$
\begin{equation*}
\int\left|v_{1}^{+}-v_{1}^{-}\right|^{q} \leq 2^{q} C \int\left|h_{1}-h_{2}\right|^{q} . \tag{6.57}
\end{equation*}
$$

Similarly we can see that

$$
\begin{equation*}
\int\left|v_{2}^{+}-v_{2}^{-}\right|^{q} \leq 2^{q} C \int\left|h_{1}-h_{2}\right|^{q} \tag{6.58}
\end{equation*}
$$

Furthermore we have

$$
\begin{align*}
\int\left|h_{1}-h_{2}\right|^{q} & \leq 2^{q-1}\left(\int\left|h_{1}-u\right|^{q}+\int\left|h_{2}-u\right|^{q}\right) \\
& =2^{q-1}\left(\int_{H_{1}^{+}}\left|v_{1}^{+}-u_{1}^{+}\right|^{q}+\int_{H_{1}^{-}}\left|v_{1}^{-}-u_{1}^{-}\right|^{q}\right. \\
& \left.+\int_{H_{2}^{+}}\left|v_{2}^{+}-u_{2}^{+}\right|^{q}+\int_{H_{2}^{-}}\left|v_{2}^{-}-u_{2}^{-}\right|^{q}\right) \\
& \leq 2^{q} \max \left\{\int_{H_{1}^{+}}\left|v_{1}^{+}-u_{1}^{+}\right|^{q}+\int_{H_{1}^{-}}\left|v_{1}^{-}-u_{1}^{-}\right|^{q},\right.  \tag{6.59}\\
& \left.\int_{H_{2}^{+}}\left|v_{2}^{+}-u_{2}^{+}\right|^{q}+\int_{H_{2}^{-}}\left|v_{2}^{-}-u_{2}^{-}\right|^{q}\right\} \\
& =2^{q}\left(\int_{H_{k}^{+}}\left|v_{k}^{+}-u_{k}^{+}\right|^{q}+\int_{H_{k}^{-}}\left|v_{k}^{-}-u_{k}^{-}\right| q^{q}\right) .
\end{align*}
$$

putting together (6.57), (6.58) and (6.59) we obtain the claim of Step 2.

## A reduction theorem for the stability of Gagliardo-Nirenberg-Sobolev

We stress here that we cannot symmetrize directly our function once again. A formal argument which shows a problem that may occur is the following: consider a function $v$ such that $\delta(v)=\lambda(v)=0$ and construct a new function $u$ as follows:

$$
u(x)=v(x) 1_{\left\{x_{1} \geq 0\right\}}(x)+2^{d / q} v(2 x) \mathbf{1}_{\left\{x_{1}<0\right\}}(x)
$$

Such a function is $(d-1)$-symmetric with respect to $H_{k}=\left\{x_{k}=0\right\}$ for $k \neq 1$. if we try to symmetrize such function with respect to $H_{1}$ we would obtain $u_{+}(x)=$ $2^{d / q}(v(2 x))$ and $u_{-}(x)=v(x)$. Clearly none of those functions satisfy the first inequality in (6.29). However we are going to see that a more refined symmetrization can bypass this problem.

Proof of Theorem 6.9. Assume as usual that $\|u\|_{q}=1$. We can assume, thanks to Proposition 6.14, that $u$ is an $(d-1)$-symmetric function and that $\delta(u)<\bar{\delta}$ with $\bar{\delta}$ arbitrarily small. Up to a rotation and a translation, we can consider $u$ to be symmetric with respect to the coordinate axes $\left\{x_{k}=0\right\}$ for $k=2, \ldots, d$ and such that

$$
\int_{\left\{x_{k}>0\right\}} u^{q}=\frac{1}{2}=\int_{\left\{x_{k}<0\right\}} u^{q}
$$

Let $u^{ \pm}$be the two symmetrizations of $u$ with respect to the hyperplane $\left\{x_{1}=0\right\}$, constructed as in Proposition 6.14. We have that $\max \left\{\delta\left(u^{+}\right), \delta\left(u^{-}\right)\right\} \leq 2 \delta(u)$. So, if $\min \left\{\lambda\left(u^{+}\right), \lambda\left(u^{-}\right)\right\} \geq C_{0} \lambda(u)$, we would be done. Thus we suppose that

$$
\begin{equation*}
\max \left\{\lambda\left(u^{+}\right), \lambda\left(u^{-}\right)\right\}<\varepsilon \lambda(u) \tag{6.60}
\end{equation*}
$$

for some constant $\varepsilon$ to be chosen. Consider $Q=\left\{\left|x_{1}\right| \leq x_{2}\right\}, Q^{+}=Q \cap\left\{x_{1}>0\right\}$ and $Q^{-}=Q \cap\left\{x_{1}<0\right\}$ and define a function $\hat{u}$ as follows:

$$
\hat{u}(x)= \begin{cases}u(x) & \text { if } \quad x \in Q \\ u\left(R_{1} x\right) & \text { if } \quad x \in R_{1}(Q) \\ \hat{u}\left(R_{2} x\right) & \text { if } \quad x \in R_{2}\left(Q \cup R_{1}(Q)\right)\end{cases}
$$

where $R_{1}$ and $R_{2}$ are the reflection with respect to $\left\{x_{1}=x_{2}\right\}$ and $\left\{x_{1}=-x_{2}\right\}$ respectively. The function $\hat{u}$ satisfies all the symmetries of $u$ with the exception of the one related to the hyperplane $\left\{x_{2}=0\right\}$, but by construction it is symmetric also with respect to $\left\{x_{1}= \pm x_{2}\right\}$. So it is $d$-symmetric. It remains to show that $\hat{u}$ satisfies the reduction inequalities (6.29). Let us start with the one relative to the asymmetry. To this aim we will denote $\hat{v}, v^{+}$and $v^{-}$as the functions who achieve the minima of $\lambda(\hat{u} \mid\{0\}), \lambda\left(u^{+} \mid\{0\}\right)$ and $\lambda\left(u^{-} \mid\{0\}\right)$ respectively. Since $\int \hat{u}^{q} \leq 4$, we get that

$$
\begin{aligned}
3^{q} \lambda(\hat{u}) & \geq \lambda(\hat{u} \mid\{0\})=\frac{\int|\hat{u}-\hat{v}|^{q}}{\int \hat{u}^{q}}=\frac{4}{\int \hat{u}^{q}} \int_{Q}|u-\hat{v}|^{q} \\
& \geq \int_{Q^{+}}\left|u^{+}-\hat{v}\right|^{q}+\int_{Q^{-}}\left|u^{-}-\hat{v}\right|^{q} \\
& =\int_{Q^{+}}\left|u^{+}-\hat{v}\right|^{q}+\int_{Q^{+}}\left|u^{-}-\hat{v}\right|^{q} \geq \frac{1}{2^{q-1}} \int_{Q^{+}}\left|u^{+}-u^{-}\right|^{q}
\end{aligned}
$$

The first inequality in (6.29) is then true if we can estimate $\lambda(u)$ in terms of $\int_{Q^{+}} \mid u^{+}{ }^{-}$ $\left.u^{-}\right|^{q}$. To this aim we observe that

$$
\begin{align*}
\left\|u^{+}-u^{-}\right\|_{L^{q}\left(Q^{+}\right)} & =\frac{1}{2}\left\|u^{+}-u^{-}\right\|_{L^{q}(Q)} \\
& =\frac{1}{2}\left\|\left(v^{+}-v^{-}\right)-\left(v^{+}-u^{+}\right)-\left(u^{-}-v^{-}\right)\right\|_{L^{q}(Q)}  \tag{6.61}\\
& \geq \frac{1}{2}\left[\left\|v^{+}-v^{-}\right\|_{L^{q}(Q)}-\left\|u^{+}-v^{+}\right\|_{L^{q}(Q)}-\left\|u^{-}-v^{-}\right\|_{L^{q}(Q)}\right] .
\end{align*}
$$

Moreover

$$
\begin{equation*}
\int_{Q}\left|u^{ \pm}-v^{ \pm}\right|^{q} \leq \int\left|u^{ \pm}-v^{ \pm}\right|^{q}=\lambda\left(u^{ \pm} \mid\{0\}\right) \leq 3^{q} \lambda\left(u^{ \pm}\right) \leq \varepsilon 3^{q} \lambda(u) \tag{6.62}
\end{equation*}
$$

where we used Lemma 6.11 and Lemma 6.12 and the fact that $u^{ \pm}$are $d$-symmetric functions. Thanks to (6.61) and (6.62) we obtain

$$
\begin{equation*}
\left\|u^{+}-u^{-}\right\|_{L^{q}\left(Q^{+}\right)} \geq \frac{1}{2}\left[\left\|v^{+}-v^{-}\right\|_{L^{q}(Q)}-2\left(\frac{3^{q}}{C(\varepsilon)} \lambda(u)\right)^{1 / q}\right] \tag{6.63}
\end{equation*}
$$

where $C(\varepsilon)$ is a suitable positive constant. Furthermore, always thanks to (6.62) we have

$$
\begin{aligned}
\lambda(u) & \leq \int\left|u-v^{+}\right|^{q}=\frac{1}{2} \int\left|v^{+}-u^{+}\right|^{q}+\frac{1}{2} \int\left|v^{+}-u^{-}\right|^{q} \\
& \leq \varepsilon \frac{3^{q}}{2} \lambda(u)+2^{q-2}\left(\varepsilon 3^{q} \lambda(u)+\int\left|v^{+}-v^{-}\right|^{q}\right)
\end{aligned}
$$

and so

$$
\begin{equation*}
\int\left|v^{+}-v^{-}\right|^{q} \geq \frac{1}{2^{q-2}}\left[\lambda(u)-\varepsilon \frac{3^{q}}{2} \lambda(u)-\varepsilon 3^{q} \lambda(u)\right] \geq \frac{1}{2^{q}} \lambda(u) \tag{6.64}
\end{equation*}
$$

where last inequality is justified choosing $\varepsilon$ small enough. Summarying we obtained

$$
\left\|u^{+}-u^{-}\right\|_{L^{q}\left(Q^{+}\right)} \geq C_{0} \lambda(u)^{1 / q}
$$

where $C_{0}$ is a suitable positive constant independent of $u$. Let us consider now the second inequality in (6.29), that is the one concerning the deficit. We stress here that we cannot say that $\|\hat{u}\|_{q}=1$. We have

$$
\left|\|\hat{u}\|_{L^{q}\left(Q^{+}\right)}-(1 / 8)^{1 / q}\right|=\left|\left\|u^{+}\right\|_{L^{q}\left(Q^{+}\right)}-\left\|v^{+}\right\|_{L^{q}\left(Q^{+}\right)}\right| \leq\left\|u^{+}-v^{+}\right\|_{L^{q}\left(Q^{+}\right)}
$$

and, since $\left|s^{q}-t^{q}\right| \leq C_{1}|s-t|$ for a suitable $C_{1}$ if $s$ and $t$ are in $[0,1]$,

$$
\begin{align*}
\left|\int_{Q^{+}} \hat{u}^{q}-\frac{1}{8}\right| & \leq C_{2}\left\|u^{+}-v^{+}\right\|_{q}=C_{2} \lambda\left(u^{+} \mid\{0\}\right)^{1 / q}  \tag{6.65}\\
& \leq 3^{1 / q} C_{3} \lambda\left(u^{+}\right)^{1 / q} \leq 3^{1 / q} C_{4} \delta\left(u^{+}\right)^{\alpha / q} \leq 2 C_{4} \delta(u)^{\alpha / q}
\end{align*}
$$

## A reduction theorem for the stability of Gagliardo-Nirenberg-Sobolev

where we used Lemma 6.11, the fact that $u$ is $d$-symmetric and that $\delta\left(u^{+}\right) \leq 2 \delta(u)$ and where $C_{2}, C_{3}$ and $C_{4}$ are positive constants. An analogous estimate holds on $Q^{-}$, on $U^{+}=\left\{x_{1}>0, x_{2}>0\right\} \backslash Q$ and on $U^{-}=\left\{x_{1}<0, x_{2}>0\right\} \backslash Q$. Then, by triangle inequality we obtain that

$$
\begin{equation*}
\left|\int \hat{u}^{q}-1\right| \leq C \delta(u)^{\alpha / q} \tag{6.66}
\end{equation*}
$$

for a suitable $\alpha>0$. Let us recall the definition of the functionals $F$ and $G$ given in (6.15):

$$
\begin{equation*}
G(u)=\|\nabla u\|_{p}^{\theta}\|u\|_{s}^{1-\theta}, \quad F(u)=\int|\nabla u|^{p}+\int|u|^{s} . \tag{6.67}
\end{equation*}
$$

By Lemma 6.3 we know that there exist two positive constants $\kappa$ and $\eta_{0}$ such that for every $u$ there exists $\lambda>0$ such that $F\left(\tau_{\lambda} u\right)=\eta_{0} G(u)^{\kappa}$ where $\tau_{\lambda} u(x)=\lambda^{d / q} u(\lambda x)$. Furthermore such $\lambda$ minimize the function $\mu \mapsto F\left(\tau_{\mu} u\right)$ in $\mathbb{R}^{+}$. So we get

$$
\begin{equation*}
G^{\kappa}\|\hat{u}\|_{q}^{\kappa}(1+\delta(\hat{u}))^{\kappa}=G(\hat{u})^{\kappa} \leq\left(\frac{1}{\eta_{0}} \int\left|\nabla \tau_{\lambda} \hat{u}\right|^{p}+\frac{1}{\eta_{0}} \int\left|\tau_{\lambda} \hat{u}\right|^{s}\right) \tag{6.68}
\end{equation*}
$$

for all $\lambda>0$. Then

$$
\begin{align*}
& \frac{1}{\eta_{0}}\left(\int\left|\nabla \tau_{\lambda} \hat{u}\right|^{p}+\int\left|\tau_{\lambda} \hat{u}\right|^{s}\right)=\frac{4}{\eta_{0}}\left(\int_{Q^{+} \cup Q^{-}}\left|\nabla \tau_{\lambda} \hat{u}\right|^{p}+\int_{Q^{+} \cup Q^{-}}\left|\tau_{\lambda} \hat{u}\right|^{s}\right)  \tag{6.69}\\
& =\frac{4}{\eta_{0}}\left(\int_{\left\{x_{2}>0\right\}}\left(\left|\nabla \tau_{\lambda} u\right|^{p}+\left|\tau_{\lambda} u\right|^{s}\right)-\int_{U^{+} \cup U^{-}}\left(\left|\nabla \tau_{\lambda} u\right|^{p}+\left|\tau_{\lambda} u\right|^{s}\right)\right) .
\end{align*}
$$

Choosing $\lambda>0$ such that $F\left(\tau_{\lambda} u\right)=\eta_{0} G(u)^{\kappa}$, since $u$ is symmetric with respect to $\left\{x_{2}=0\right\}$, we get

$$
\begin{align*}
\int_{\left\{x_{2}>0\right\}}\left(\left|\nabla \tau_{\lambda} u\right|^{p}+\left|\tau_{\lambda} u\right|^{s}\right) & =\frac{1}{2} F\left(\tau_{\lambda} u\right) \\
& =\frac{\eta_{0}}{2} G(u)^{\kappa}=\frac{\eta_{0} G^{\kappa}}{2}(\delta(u)+1)^{\kappa}  \tag{6.70}\\
& \leq \frac{\eta_{0} G^{\kappa}}{2}\left(1+C_{5} \delta(u)\right),
\end{align*}
$$

where the last passage is true for $\delta(u)$ small enough. Let us consider now the function

$$
v(x)= \begin{cases}\tau_{\lambda} u(x) & x \in U^{+} \\ \tau_{\lambda} u\left(R_{1} x\right) & x \in R_{1} U^{+} \\ v\left(S_{1} x\right) & x \in\left\{x_{1}<0, x_{2}>0\right\} \\ v\left(S_{2} x\right) & x \in\left\{x_{2}<0\right\}\end{cases}
$$

where $S_{i}$ is the symmetrization with respect to $\left\{x_{i}=0\right\}$. We have

$$
\begin{align*}
\int_{U^{+}}\left(\left|\nabla \tau_{\lambda}(u)\right|^{p}+\left|\tau_{\lambda}(u)\right|^{s}\right) & =\frac{1}{8}\left(\int\left(|\nabla v|^{p}+|v|^{s}\right)\right) \geq \frac{1}{8} \min _{\mu \in \mathbb{R}^{+}} F\left(\tau_{\mu} v\right) \\
& =\frac{\eta_{0}}{8} G(v)^{\kappa} \geq \frac{\eta_{0} G^{\kappa}}{8}\left(\int v^{q}\right)^{\kappa / q}=\frac{\eta_{0} G^{\kappa}}{8} \cdot 8^{\kappa / q}\left(\int_{U^{+}} u^{q}\right)^{\kappa / q} \\
& \geq \eta_{0} G^{\kappa} 8^{-1+\kappa / q}\left(\frac{1}{8}-C_{6} \delta(u)^{\alpha}\right)^{\kappa / q} \geq \eta_{0} G^{\kappa}\left(\frac{1}{8}-C_{7} \delta(u)^{\alpha}\right) . \tag{6.71}
\end{align*}
$$

with suitable constants $C_{6}$ and $C_{7}$ and for $\delta(u)$ small enough. Notice that in the last passage we used (6.65). An analogous estimate can be obtained on $U^{-}$. By (6.68), (6.69), (6.70) and (6.71) we get

$$
\begin{align*}
G^{\kappa}\|\hat{u}\|_{q}^{\kappa}(1+\delta(\hat{u}))^{\kappa} & \leq \frac{4}{\eta_{0}}\left[\frac{\eta_{0} G^{\kappa}}{2}\left(1+C_{5} \delta(u)\right)-2 \eta_{0} G^{\kappa}\left(\frac{1}{8}-C_{7} \delta(u)\right)^{\beta}\right]  \tag{6.72}\\
& =G^{\kappa}\left[1+C_{8} \delta(u)^{\alpha}\right]
\end{align*}
$$

for a suitable $C_{8}>0$ and where $\alpha$ is the minumum between 1 and $\beta$. By (6.72) and thanks to (6.66) we conclude that

$$
\delta(\hat{u}) \leq \frac{1+C_{9} \delta(u)^{\alpha}}{1-C_{9} \delta(u)^{\alpha}}-1 \leq C \delta(u)^{\alpha}
$$

that is true again for $\delta(u)$ small enough. We conclude setting $\alpha_{2}=\alpha$ and $\kappa_{2}=C$.

## Chapter 7

## Estimate of the dimension of the singular set of the MS functional: a short proof

### 7.1 Introduction

In this chapter we report a short note written in collaboration with Camillo De Lellis and Matteo Focardi. In this we observe that a classical result proved in [3] about the Hausdorff dimension, $\operatorname{dim}_{\mathcal{H}}$, of the singular set of a minimizer of the MumfordShah functional, can be easily derived from the fact that the blow-up of such a set is a Caccioppoli partition (see Section 7.2), as recently proved in [50]. This work is unrelated to the rest of the Thesis and we enclose it without properly introducing all the technical preliminaries exploited, which are, anyway, mentioned and recalled whenever needed.

The (localized) Mumford-Shah energy on a bounded open subset $\Omega \subset \mathbb{R}^{d}$ is given by

$$
\begin{equation*}
\operatorname{MS}(v, A)=\int_{A}|\nabla v|^{2} d x+\mathcal{H}^{d-1}\left(S_{v} \cap A\right), \quad \text { for } v \in S B V(\Omega) \text { and } A \subseteq \Omega \text { open } \tag{7.1}
\end{equation*}
$$

In the previous definition, the space $S B V(\Omega)$ (Special functions of Bounded Variation) is that of the classical $B V$ functions (functions of Bounded Variation) which Cantor part vanishes, while the set $S_{v}$ denotes the jump set of the function $v \in S B V(\Omega)$, that is the set where the derivative of $v$ is not absolutely continuous with respect to the Lebesgue measure. A comprehensive account on the subject is [2, Chapter 4]. In what follows if $A=\Omega$ we shall drop the dependence on the set of integration.

We recall the following result due to L. Ambrosio, N. Fusco and J. E. Hutchinson [3, Theorem 5.6].

## Estimate of the dimension of the singular set of the MS functional: a short 152 <br> proof

Theorem 7.1. Let $u$ be a local minimizer of the Mumford-Shah energy, i.e. any function $u \in S B V(\Omega)$ with $\operatorname{MS}(u)<\infty$ and such that

$$
\operatorname{MS}(u) \leq \operatorname{MS}(w) \quad \text { whenever }\{w \neq u\} \subset \subset \Omega
$$

Let $\Sigma_{u} \subseteq \overline{S_{u}}$ be the set of points out of which $\overline{S_{u}}$ is locally regular, and let

$$
\Sigma_{u}^{\prime}:=\left\{x \in \Sigma_{u}: \lim _{\rho \rightarrow 0} \rho^{1-d} \int_{B_{\rho}(x)}|\nabla u|^{2}=0\right\}
$$

Then, $\operatorname{dim}_{\mathcal{H}} \Sigma_{u}^{\prime} \leq d-2$.
The main interest in establishing such an estimate on the set $\Sigma_{u}^{\prime}$, the so-called subset of triple-junctions, is related to the understanding of the Mumford-Shah conjecture (see [2, Chapter 6] for a related discussion, see also [50, Section 7]).

Indeed, Theorem 7.1, together with the higher integrability property of the approximate gradients enjoyed by minimizers as established in 2 -dimensions by [50] and more recently in any dimension by [52], imply straightforwardly an analogous estimate on the full singular set $\Sigma_{u}$. More precisely, in view of [50, Theorem 1] and [52, Theorem 1.1] any local minimizer $u$ of the Mumford-Shah energy is such that $|\nabla u| \in L_{\text {loc }}^{p}(\Omega)$ for some $p>2$, therefore [3, Corollary 5.7] yields that

$$
\operatorname{dim}_{\mathcal{H}} \Sigma_{u} \leq \max \{d-2, d-p / 2\}
$$

where $\operatorname{dim}_{\mathcal{H}}(E)$ is the Hausdorff dimension of the set $E$. A characterization (of a suitable version) of the Mumford-Shah conjecture in $2-$ dimensions in terms of a refined higher integrability property of the gradient in the finer scale of weak Lebesgue spaces has been recently established in [50, Proposition 5].

Our proof of Theorem 7.1 rests on a compactness result proved by C. De Lellis and M. Focardi (see [50, Theorem 13]) showing that the blow-up limits of the jump set $S_{u}$ in points in the regime of small gradients, i.e. in points of $\Sigma_{u}^{\prime}$, are minimal Caccioppoli partitions. The original approach in [3], instead, relies on the notion of Almgren's area mimizing sets, for which an involved analysis of the composition of $S B V$ functions with Lipschitz deformations (not necessarily one-to-one) and a revision of the regularity theory for those sets are needed (cp. with [3, Sections 2, 3 and 4]).

Given [50, Theorem 13], the regularity theory of minimal Caccioppoli partitions developed in [94], [87] and [88], and standard arguments in geometric measure theory yield the conclusion, thus bypassing the above mentioned technical complications.

We describe briefly the plan of this chapter: in Section 7.2 we introduce the necessary definitions and recall some well-known facts about Caccioppoli partitions. In Section 7.3 we prove our main result, Theorem 7.1.

### 7.2 Caccioppoli partitions

In what follows $\Omega \subset \mathbb{R}^{d}$ will denote a bounded open set.
Definition 7.2. A Caccioppoli partition of $\Omega$ is a countable partition $\mathscr{E}=\left\{E_{i}\right\}_{i=1}^{\infty}$ of $\Omega$ in sets of (positive Lebesgue measure and) finite perimeter with $\sum_{i=1}^{\infty} P\left(E_{i} ; \Omega\right)<\infty$.

For each Caccioppoli partition $\mathscr{E}$ we define its set of interfaces as

$$
J_{\mathscr{E}}:=\bigcup_{i \in} \partial^{*} E_{i} .
$$

Here $\partial^{*} E$ denotes the essential boundary of the set $E$ :

$$
\partial^{*} E=\left\{x \in \mathbb{R}^{d}: \lim _{\rho \rightarrow 0^{+}} \frac{D \mathbf{1}_{E}\left(B_{\rho}(x)\right)}{\left|D \mathbf{1}_{E}\right|\left(B_{\rho}(x)\right)}:=\nu_{E}(x) \text { exists and satisfies }\left|\nu_{E}(x)\right|=1\right\},
$$

where $D \mathbf{1}_{E}$ is the distributional derivative of the characteristic function of the set $E$ while $\left|D \mathbf{1}_{E}\right|$ is its total variation. Also about the main properties of $\partial^{*} E$ we refer to [2] and [92]. The partition $\mathscr{E}$ is said to be minimal if

$$
\mathcal{H}^{d-1}\left(J_{\mathscr{E}}\right) \leq \mathcal{H}^{d-1}\left(J_{\mathscr{F}}\right)
$$

for all Caccioppoli partitions $\mathscr{F}$ for which there exists an open subset $\Omega^{\prime} \subset \subset \Omega$ with

$$
\sum_{i=1}^{\infty}\left|\left(F_{i} \triangle E_{i}\right) \cap\left(\Omega \backslash \Omega^{\prime}\right)\right|=0
$$

Definition 7.3. Given a Caccioppoli partition $\mathscr{E}$ we define its singular set $\Sigma_{\mathscr{E}}$ as the set of points for which the approximate tangent plane to $J_{\mathscr{E}}$ does not exist. That is, the complementary of $\Sigma_{\mathscr{E}}$ is the set of points $x$ such that $(\mathscr{E}-x) / \rho$ has a weak limit point as $\rho \rightarrow 0^{+}$, as a Radon measure (called approximate tangent space). See [2, Section 2.11].

A characterization of the singular set $\Sigma_{\mathscr{E}}$ for minimal Caccioppoli partitions in the spirit of $\varepsilon$-regularity results is provided in the ensuing statement (cp. with [88, Corollary 4.2.4] and [92, Theorem III.6.5] ).
Theorem 7.4. Let $\Omega$ be an open set and $\mathscr{E}=\left\{E_{i}\right\}_{i \in}$ a minimal Caccioppoli partition of $\Omega$.

Then, there exists a dimensional constant $\varepsilon=\varepsilon(d)>0$ such that

$$
\begin{equation*}
\Sigma_{\mathscr{E}}=\left\{x \in \Omega \cap \overline{J_{\mathscr{E}}}: \inf _{B_{\rho}(x) \subset \subset \Omega} e(x, \rho) \geq \varepsilon\right\}, \tag{7.2}
\end{equation*}
$$

where e $(x, \rho)$ denotes the spherical excess of $\mathscr{E}$ at the point $x \in J_{\mathscr{E}}$ at the scale $\rho>0$, that is

$$
e(x, \rho):=\min _{\nu \in \mathbb{S}^{d-1}} \frac{1}{\rho^{d-1}} \int_{B_{\rho}(x) \cap J_{\mathscr{E}}} \frac{\left|\nu_{\mathscr{E}}(y)-\nu\right|^{2}}{2} d \mathcal{H}^{d-1}(y) .
$$

## Estimate of the dimension of the singular set of the MS functional: a short 154 <br> proof

We recall next a result that is probably well-known in literature.
Theorem 7.5. Let $\mathscr{E}$ be a minimal Caccioppoli partition in $\Omega$, then $\operatorname{dim}_{\mathcal{H}} \Sigma_{\mathscr{E}} \leq d-2$. If, in addition, $d=2$, then $\Sigma_{\mathscr{E}}$ is locally finite.
Proof. We apply the abstract version of Federer's reduction argument in [105, Theorem A.4] with the set of functions

$$
\mathcal{F}=\left\{\mathbf{1}_{J_{\mathscr{E}}}: \mathscr{E} \text { is a minimal Caccioppoli partition }\right\}
$$

endowed with the convergence

$$
\mathbf{1}_{J_{\mathscr{E}_{h}}} \rightarrow \mathbf{1}_{J_{\mathscr{E}}} \Longleftrightarrow \lim _{h \rightarrow \infty} \int_{J_{\mathcal{E}_{h}}} g d \mathcal{H}^{d-1}=\int_{J_{\mathscr{E}}} g d \mathcal{H}^{d-1}, \quad \text { for all } g \in C_{c}^{1}(\Omega) .
$$

and singularity map $\operatorname{sing}\left(\mathbf{1}_{\mathscr{E}}\right)=\Sigma_{\mathscr{E}}$.
It is easy to see that condition $A .1$ (closure under scaling) and $A .3(2)$ hold true. Moreover, the blow-ups of a minimal Caccioppoli partition converge to a minimizing cone (see [87, Theorem 3.5], or [88, Theorem 4.4.5 (a)]), so that A. 2 holds as well. About $A .3(1)$, we notice that the singular set of an hyperplane is empty. Eventually, if a sequence $\left(\mathbf{1}_{J_{\mathscr{E}_{h}}}\right)_{h \in \mathbb{N}} \subseteq \mathcal{F}$ converges to $\mathbf{1}_{J_{\mathscr{E}}}$ and $\left(x_{h}\right)_{h \in \mathbb{N}}$ converges to $x$, with $x_{h} \in \Sigma_{\mathscr{C}_{h}}$ for all $h$, then by the continuity of the excess and the characterization in (7.2), $x \in \Sigma_{\mathscr{E}}$, so that condition $A .3(3)$ is satisfied as well.

To conclude, we recall that $\left[105\right.$, Theorem A.4] itself ensures that the set $\Sigma_{\mathscr{E}}$ is locally finite being in this setting $\operatorname{dim}_{\mathcal{H}} \Sigma_{\mathscr{E}}=0$.

### 7.3 Proof of the main result

We are now ready to prove the main result of the chapter following the approach exploited in [3, Theorem 5.6]. To this aim we recall that in [2, Theorems 8.1-8.3] it is characterized alternatively the singular set $\Sigma_{u}$ as follows

$$
\begin{equation*}
\Sigma_{u}=\left\{x \in \overline{S_{u}}: \varliminf_{\rho \rightarrow 0}(\mathscr{D}(x, \rho)+\mathscr{A}(x, \rho)) \geq \varepsilon_{0}\right\}, \tag{7.3}
\end{equation*}
$$

where $\varepsilon_{0}$ is a dimensional constant, and the scaled Dirichlet Energy and the scaled mean-flatness are respectively defined as

$$
\mathscr{D}(x, \rho):=\rho^{1-d} \int_{B_{\rho}(x)}|\nabla u|^{2} d y, \quad \mathscr{A}(x, \rho):=\rho^{-1-d} \min _{T \in \Pi} \int_{S_{u} \cap B_{\rho}(x)} \operatorname{dist}^{2}(y, T) d \mathcal{H}^{d-1}(y),
$$

with $\Pi$ the set of all affine $(d-1)$-hyperplanes in $\mathbb{R}^{d}$.
Proof of Theorem 7.1. We argue by contradiction: suppose that there exists $s>d-2$ such that $\mathcal{H}^{s}\left(\Sigma_{u}^{\prime}\right)>0$. From this, we infer that $\mathcal{H}_{\infty}^{s}\left(\Sigma_{u}^{\prime}\right)>0$, (here $\mathcal{H}_{\infty}^{s}$ is the preHausdorff measure, see [2, Definition 2.46]) and moreover that for $\mathcal{H}^{s}$-a.e. $x \in \Sigma_{u}^{\prime}$ it holds

$$
\begin{equation*}
\overline{\lim }_{\rho \rightarrow 0^{+}} \frac{\mathcal{H}_{\infty}^{s}\left(\Sigma_{u}^{\prime} \cap B_{\rho}(x)\right)}{\rho^{s}} \geq \frac{\omega_{s}}{2^{s}} \tag{7.4}
\end{equation*}
$$

where $\omega_{s}$ is the $s$-dimensional Hausdorff measure of the unitary ball (see for instance [2, Theorem 2.56 and formula (2.43)] or [92, Lemma III.8.15]). Without loss of generality, suppose that (7.4) holds at $x=0$, and consider a sequence $\rho_{h} \rightarrow 0$ for which

$$
\begin{equation*}
\mathcal{H}_{\infty}^{s}\left(\Sigma_{u}^{\prime} \cap B_{\rho_{h}}\right) \geq \frac{\omega_{s}}{2^{s+1}} \rho_{h}^{s} \quad \text { for all } h \in \mathbb{N} . \tag{7.5}
\end{equation*}
$$

[50, Theorem 13] provides a subsequence, not relabeled for convenience, and a minimal Caccioppoli partition $\mathscr{E}$ such that

$$
\begin{equation*}
\mathcal{H}^{d-1}\left\llcorner\rho _ { h } ^ { - 1 } S _ { u } \stackrel { * } { \rightharpoonup } \mathcal { H } ^ { d - 1 } \left\llcorner J_{\mathscr{E}}, \text { and } \quad \rho_{h}^{-1} \overline{S_{u}} \rightarrow \overline{J_{\mathscr{E}}}\right.\right. \text { locally Hausdorff. } \tag{7.6}
\end{equation*}
$$

In turn, from the latter we claim that if $\mathcal{F}$ is any open cover of $\Sigma_{\mathscr{E}} \cap \bar{B}_{1}$, then for some $h_{0} \in \mathbb{N}$

$$
\begin{equation*}
\rho_{h}^{-1} \Sigma_{u}^{\prime} \cap \bar{B}_{1} \subseteq \bigcup_{F \in \mathcal{F}} F \quad \text { for all } h \geq h_{0} \tag{7.7}
\end{equation*}
$$

Indeed, if this is not the case we can find a sequence $x_{h_{j}} \in \rho_{h_{j}}^{-1} \Sigma_{u}^{\prime} \cap \bar{B}_{1}$ converging to some point $x_{0} \notin \Sigma_{\mathscr{E}}$. If $\pi_{x_{0}}^{\mathscr{E}}$ is the approximate tangent plane to $J_{\mathscr{E}}$ at $x_{0}$, that exists by the very definition of $\Sigma_{\mathscr{E}}$, then for some $\rho_{0}$ we have

$$
\rho^{1-d} \int_{B_{\rho}\left(x_{0}\right) \cap J_{\mathscr{E}}} \operatorname{dist}^{2}\left(y, \pi_{x_{0}}^{\mathscr{E}}\right) d \mathcal{H}^{d-1}<\varepsilon_{0}, \quad \text { for all } \rho \in\left(0, \rho_{0}\right)
$$

In turn, from the latter inequality it follows for $\rho \in\left(0, \rho_{0} \wedge 1\right)$

$$
\varlimsup_{j \rightarrow \infty} \int_{B_{\rho}\left(x_{h_{j}}\right) \cap \rho_{h_{j}}^{-1} S_{u}} \operatorname{dist}^{2}\left(y, \pi_{x_{0}}^{\mathscr{E}}\right) d \mathcal{H}^{d-1}<\varepsilon_{0} .
$$

Therefore, as $x_{h_{j}} \in \rho_{h_{j}}^{-1} \Sigma_{u}^{\prime}$, we get for $j$ large enough

$$
\overline{\lim }_{\rho \rightarrow 0}\left(\mathscr{D}\left(x_{h_{j}}, \rho\right)+\mathscr{A}\left(x_{h_{j}}, \rho\right)\right)<\varepsilon_{0}
$$

a contradiction in view of the characterization of the singular set in (7.3).
To conclude, we note that by (7.7) we get

$$
\mathcal{H}_{\infty}^{s}\left(\Sigma_{\mathscr{E}} \cap \overline{B_{1}}\right) \geq \overline{\lim }_{h \rightarrow \infty} \mathcal{H}_{\infty}^{s}\left(\rho_{h}^{-1} \Sigma_{u}^{\prime} \cap \overline{B_{1}}\right) ;
$$

given this, (7.5) and (7.6) yield that

$$
\mathcal{H}^{s}\left(\Sigma_{\mathscr{E}} \cap \overline{B_{1}}\right) \geq \mathcal{H}_{\infty}^{s}\left(\Sigma_{\mathscr{E}} \cap \overline{B_{1}}\right) \geq \overline{\lim }_{h \rightarrow \infty} \mathcal{H}_{\infty}^{s}\left(\rho_{h}^{-1} \Sigma_{u}^{\prime} \cap \overline{B_{1}}\right) \geq \frac{\omega_{s}}{2^{s+1}}
$$

thus contradicting Theorem 7.5.

## Bibliography

[1] E. Acerbi, N. Fusco, M. Morini: Minimality via second variation for a nonlocal isoperimetric problem, Comm. Math. Phys., 322 (2013), 515-557.
[2] L. Ambrosio, N. Fusco, D. Pallara: Functions of bounded variation and free discontinuity problems, Oxford Mathematical Monographs. The Clarendon Press, Oxford University Press, New York, 2000.
[3] L. Ambrosio, N. Fusco, J.E. Hutchinson: Higher integrability of the gradient and dimension of the singular set for minimisers of the Mumford-Shah functionals, Calc. Var. Partial Differential Equation, 16 (2003), 187-215.
[4] L.Ambrosio, P.Tilli: Topics on Analysis in Metric Spaces, Oxford Lecture Series in Mathematics and its Applications, Oxford University Press, Oxford (2004).
[5] M.S. Ashbaugh, E.M. Harrell: Maximal and minimal eigenvalues and their associated nonlinear equations, J. Math. Phys., 28 (1987), 1770-1786.
[6] T. Aubin: Problèmes isopérimétriques et espaces de Sobolev, J. Differential Geometry, 11 (1976), 573-598.
[7] A. I. Ávila: Stability results for the first eigenvalue of the Laplacian on domains in space forms, J. Math. Anal. Appl., 267 (2002), 760-774.
[8] M. Barchiesi, F. Cagnetti, N. Fusco: Stability of the Steiner symmetrization of convex sets, J. Eur. Math. Soc., 15 (2013), 1245-1278.
[9] M. F. Betta, F. Brock, A. Mercaldo, M. R. Posteraro: A weighted isoperimetric inequality and applications to symmetrization, J. of Inequal. \& Appl., 4 (1999), 215-240.
[10] T. Bhattacharya: Some observations on the first eigenvalue of the p-Laplacian and its connections with asymmetry, Electron. J. Differential Equations, 35 (2001), 1-15.
[11] T. Bhattacharya, A. Weitsman: Estimates for Green's function in terms of asymmetry, Contemp. Math. Series, 221 (1999), 31-58.
[12] G. Bianchi, H. Egnell: A note on the Sobolev inequality, J. Funct. Anal., 100 (1991), 18-24.
[13] V. Bögelein, F. Duzaar, N. Fusco: A quantitative isoperimetric inequality on the sphere, preprint at http://cvgmt.sns.it/paper/2146/ (2013).
[14] V. Bögelein, F. Duzaar, N. Fusco: A sharp quantitative isoperimetric inequality in higher codimension, preprint at http://cvgmt.sns.it/paper/1865/ (2012).
[15] L. Brasco, G. De Philippis, B. Velichkov: Faber-Krahn inequalities in sharp quantitative form, preprint at http://cvgmt.sns.it/paper/2161/ (2013).
[16] L. Brasco, G. De Philippis, B. Ruffini: Spectral optimization for the Stekloff-Laplacian: the stability issue, J. Funct. Anal., 262 (2012), 4675-4710.
[17] L. Brasco, G. Franzina: On the Hong-Krahn-Szego inequality for the pLaplace operator, Manuscripta Math., 141 (2013), 537-557.
[18] L. Brasco, A. Pratelli: Sharp stability of some spectral inequalities, Geom. Func. An., 22 (2012), 107-135.
[19] H. Brezis: Functional analysis, Sobolev spaces and partial differential equations, Universitext, Springer, New York, 2011.
[20] F. Brock: An isoperimetric inequality for eigenvalues of the Stekloff problem, ZAMM Z. Angew. Math. Mech., 81 (2001), 69-71.
[21] F. Brock, A. Chiacchio, M. R. Mercaldo: A class of degenerate elliptic equations and a Didos problem with respect to a measure, J. Math. Anal. Appl., 348 (2008), 356-365.
[22] Brock, Chiacchio, Mercaldo: Weighted isoperimetric inequalities in cones and applications, Nonlinear Analysis: Theory, Methods and Applications, 75 (2012), 5737-5755.
[23] J. E. Brothers, W. P. Ziemer: Minimal rearrangements of Sobolev functions, J. Reine Angew. Math., 384 (1988), 153-179.
[24] D. Bucur: Minimization of the $k$-th eigenvalue of the Dirichlet Laplacian, Arch. Ration. Mech. Anal., 206 (2012), 1073-1083.
[25] D. Bucur, G. Buttazzo: Variational Methods in Shape Optimization Problems, Progress in nonlinear Differential Equations Vol. 65, Birkhuser Boston, Inc., Boston, MA, 2005.
[26] D. Bucur, G. Buttazzo: On the characterization of the compact embedding of Sobolev spaces, Calc. Var., 44 (2012), 455-475.
[27] G. Buttazzo: Semicontinuity, relaxation and integral representation in the calculus of variations, Pitman Research Notes in Mathematics Series, 207. Longman Scientific \& Technical, Harlow, copublished in the United States with John Wiley \& Sons, Inc., New York, 1989.
[28] G. Buttazzo: Spectral optimization problems, Rev. Mat. Complut., 24 (2011), 277-322.
[29] G. Buttazzo, G. Dal Maso: Shape optimization for Dirichlet problems: relaxed formulation and optimality conditions, Appl. Math. Optim., 23 (1991), 17-49.
[30] G. Buttazzo, G. Dal Maso: An existence result for a class of shape optimization problems, Arch. Rational Mech. Anal., 122 (1993), 183-195.
[31] G. Buttazzo, A. Gerolin, B. Ruffini, B. Velichkov: Optimal potential for Schrödinger operators, preprint at http://cvgmt.sns.it/paper/2140/ (2013).
[32] G. Buttazzo, B. Ruffini, B.Velichkov: Shape Optimization Problems for Metric Graphs, to appear in ESAIM: Control, Optimisation and Calculus of Variations, (2013).
[33] G. Buttazzo, H. Soubairi, N. Varchon: Optimal measures for elliptic problems, Ann. Mat. Pura Appl., 185 (2006), 207-221.
[34] E. Carlen, A. Figalli: Stability for a GNS inequality and the log-HLS inequality, with application to the critical mass Keller-Segel equation, Duke Math. J., 162 (2013), 579-625.
[35] E. Carlen, R. Frank, E. Lieb: Stability estimates for the lowest eigenvalue of a Schrdinger operator, preprint at http://arxiv.org/abs/1301.5032 (2012).
[36] E. A. Carlen, M. Loss: Sharp Constant in Nash's Inequality, Sc. Math. Geor. Inst. Tech., 7 (1993), 213-215.
[37] J. Cheeger: Differentiability of Lipschitz functions on metric measure spaces Geom. Funct. Anal., 9 (1999), 428-517.
[38] A. Cianchi, N. Fusco: Functions of Bounded Variation and Rearrangements, Arch. Rational Mech. Anal., 165 (2002) 1-40.
[39] A. Cianchi, N. Fusco, F. Maggi, A. Pratelli: The sharp Sobolev inequality in quantitative form. J. Eur. Math. Soc., 11 (2009), 1105-1139.
[40] M. Cicalese, G. Leonardi A selection principle for the sharp quantitative isoperimetric inequality, Arch. Ration. Mech. Anal., 206 (2012), 617-643.
[41] M. Cicalese E.N. Spadaro: Droplet Minimizers of an Isoperimetric Problem with long-range interactions, to appear in Commun. Pure Appl. Math., (2013).
[42] R. Courant, D. Hilbert: Methods of Mathematical Physics, Wiley, vol. 1 and 2, New York, 1953 and 1962.
[43] D. Cordero-Erausquin, B. Nazaret, C. Villani: A Mass-Transportation Approach to Sobolev and Nirenberg Inequalities, Adv. Math., 182 (2004), 307332.
[44] G. Dal Maso: An Introduction to $\Gamma$-convergence, Progress in Nonlinear Differential Equations and their Applications, Vol. 8, Birkhäuser Boston, Inc., Boston MA, 1993.
[45] G. Dal Maso, U. Mosco: Wiener's criterion and $\Gamma$-convergenceAppl. Math. Optim., 15 (1987), 15-63.
[46] G. David: Singular Sets of Minimizers for the Mumford-Shah Functional, Progress in Mathematics, vol. 233, Birkhäuser Verlag, Basel, (2005).
[47] E. De Giorgi: Sulla proprietà isoperimetrica dellipersfera, nella classe degli insiemi aventi frontiera orientata di misura finita, Atti Accad. Naz. Lincei. Mem. Cl. Sci. Fis. Mat. Nat. Sez. I, 5 (1958), 33-44.
[48] M. Del Pino, J. Dolbeault: Best constants for Gagliardo-Nirenberg inequalities and applications to nonlinear diffusions, J. Math. Pures Appl., 81 (2002) 847-875.
[49] M.C. Delfour and J.P. Zolésio: Shapes and geometries. Metrics, analysis, differential calculus, and optimization, Second edition, Advances in Design and Control, 22. Society for Industrial and Applied Mathematics (SIAM), Philadelphia, PA, 2011.
[50] C. De Lellis, M. Focardi: Higher integrability of the gradient for minimizers of the $2 d$ Mumford-Shah energy, J. Math. Pures Appl., 100 (2013), 391-409.
[51] C. De Lellis, M. Focardi, B. Ruffini: A note on the Hausdorff dimension of the singular set for minimizers of the Mumford-Shah energy, preprint at http://cvgmt.sns.it/paper/2178/ (2013).
[52] G. De Philippis A. Figalli: Higher integrability for minimizers of the Mumford-Shah functional, preprint at http://arxiv.org/abs/1303.1196 (2013).
[53] G. De Philippis, F. Maggi, Sharp stability inequalities for the Plateau problem, preprint at http://cvgmt.sns.it/paper/1715/ (2013).
[54] J. Dolbeault G. Toscani: Improved Sobolevs inequalities, relative entropy and fast diffusion equations, preprint at http://arxiv.org/abs/1110.5175, (2011)
[55] H. Egnell: Extremal properties of the first eigenvalue of a class of elliptic eigenvalue problems, Ann. Sc. Norm. Super. Pisa Cl. Sci., 14 (1987), 1-48.
[56] L. Evans: Partial Differential Equations, American Mathematical Society, Providence, RI 1998.
[57] L. Evans, R. Gariepy: Measure theory and fine properties of functions, Studies in Advanced Mathematics, CRC Press, Boca Raton, FL, 1992.
[58] A. Figalli, F. Maggi, A. Pratelli: A mass transportation approach to quantitative isoperimetric inequalities, Invent. Math., 182 (2010), 167-211.
[59] L. Friedlander: Extremal properties of eigenvalues for a metric graph, Ann. Inst. Fourier, 55 (2005), 199-211.
[60] B. Fuglede: Stability in the isoperimetric problem for convex or nearly spherical domains in Rn. Trans. Amer. Math. Soc., 314 (1989), 619-638.
[61] N. Fusco, V. Julin: A strong form of the quantitative isoperimetric inequality, preprint at http://cvgmt.sns.it/paper/1692/ (2011).
[62] N. Fusco, F. Maggi, A. Pratelli: The sharp quantitative isoperimetric inequality. Ann. of Math., 168 (2008), 941-980.
[63] N. Fusco, F. Maggi, A. Pratelli: The sharp quantitative Sobolev inequality for functions of bounded variation, J. Funct. Anal., 244 (2007), 315-341.
[64] N. Fusco, F. Maggi, A. Pratelli: Stability estimates for certain FaberKrahn, Isocapacitary and Cheeger inequalities, Ann. Sc. Norm. Super. Pisa Cl. Sci., 8 (2009), 51-71.
[65] N. Fusco, F. Maggi, A. Pratelli: On the isoperimetric problem with respect to a mixed Euclidean-Gaussian density, J. Funct. Anal., 260 (2011), 3678-3717.
[66] M. Fontelos, A. Friedman: Symmetry-Breaking Bifurcations of Charged Drops, Arch. Rational Mech. Anal., 172 (2004), 267-294.
[67] D. Gilbarg, N. Trudinger: Elliptic partial differential equations of second order, Classics in Mathematics, Springer-Verlag, Berlin, 2001.
[68] E. Giusti: Minimal surfaces and functions of bounded variation, Monographs in Mathematics, 80, Birkhäuser Verlag, Basel, 1984.
[69] M. Goldman, M. Novaga: Volume-constrained minimizers for the prescribed curvature problem in periodic media, Calc. Var. Partial Differential Equations, 44 (2012), 297-318.
[70] M. Goldman, M. Novaga, B. Ruffini: A non-local model for charged droplets, in preparation (2013).
[71] W. Hansen, N. Nadirashvili Isoperimetric inequalities in potential theory, Potential Anal., 3 (1994), 1-14.
[72] R. R. Hall: A quantitative isoperimetric inequality in $n$-dimensional, J. Reine Angew. Math., 428 (1992), 161-176.
[73] R. R. Hall, W. K. Hayman, A. W. Weitsman: On asymmetry and capacity, J. d'Analyse Math., 56 (1991), 87-123.
[74] E.M. Harrell: Hamiltonian operators with maximal eigenvalues, J. Math. Phys., 25 (1984), 48-51.
[75] A. Henrot: Extremum problems for eigenvalues of elliptic operators, Frontiers in Mathematics, Birkhäuser Verlag, Basel, 2006.
[76] A. Henrot, G. A. Philippin, A. Safoui: Some isoperimetric inequalities with application to the Stekloff problem, J. Convex Anal., 15 (2008), 581-592.
[77] A. Henrot, M. Pierre: Variation et optimisation de formes. Une analyse géométrique. Mathématiques \& Applications, 48, Springer, Berlin, 2005.
[78] J. Hersch, L. E. Payne, M. M. Schiffer: Some inequalities for Stekloff eigenvalues, Arch. Rat. Mech. Anal., 57 (1954), 99-114.
[79] G. N. Hile, Z. Xu: Inequalities for sums of reciprocals of eigenvalues, J. Math. Anal. Appl., 180 (1993), 412-430.
[80] H. Knüpfer, C. B. Muratov: On an isoperimetric problem with a non-local term. I. The planar case, Comm. Pure Appl. Math., 66 (2013), 1129-1162.
[81] H. Knüpfer, C. B. Muratov: On an isoperimetric problem with a non-local term. II. The general case, preprint at http://arxiv.org/abs/1206.7078 (2012).
[82] A. V. Kolesnikov, R. I. Zhdanov: On isoperimetric sets of radially symmetric measures, Contemp. Math., 545 (2011), 123-154.
[83] P. Kuchment: Quantum graphs: an introduction and a brief survey AMS Proc. Symp. Pure. Math. 77, (2008), 291-312.
[84] N. S. Landkof: Foundations of Modern Potential Theory, Die Grundlehren der mathematischen Wissenschaften, 180, Springer-Verlag, New York-Heidelberg, 1972.
[85] A. Lemenant: Regularity of the singular set for Mumford-Shah minimizers in $\mathbb{R}^{3}$ near a minimal cone, Ann. Sc. Norm. Super. Pisa Cl. Sci., 10 (2011), 561-609.
[86] H. Lieb, M. Loss: Analysis, Graduate Studies in Mathematics, 14 American Mathematical Society, Providence, RI 2000.
[87] G.P. Leonardi: Blow-Up of Oriented Boundaries, Rend. Sem. Mat. Univ. Padova, 103 (2000), 211-232.
[88] G.P. Leonardi: Optimal Subdivisions of $d$-dimensional Domains, PhD thesis, Università di Trento, 1998.
[89] P.L. Lions: The concentration-compactness principle in the Calculus of Variation. The Locally compact case, part 1, Annales de l'I. H. P., section C, tome 1, 4 (1984), 109-145.
[90] P.L. Lions: The concentration-compactness principle in the Calculus of Variation. The Locally compact case, part 2, Annales de l'I. H. P., section C, tome 1, 14 (1984), 223-283.
[91] P.L. Lions: The Concentration-Compactness Principle in the Calculus of Variations. The limit case, part 1, Revista Matemática iberoamericana Vol. 1 (1985), 145-201.
[92] F. Maggi: Sets of Finite Perimeter and geometric variational problems. An introduction to geometric measure theory, Cambridge Studies in Advanced Mathematics 135, Cambridge University Press, Cambridge 2012.
[93] F. Maggi: Some methods for studying stability in isoperimetric type problems, Bull. Amer. Math. Society, 45 (2008), 367-408.
[94] U. Massari, I. Tamanini: Regularity properties of optimal segmentations, J. Reine Angew. Math., 420 (1991), 61-84.
[95] P. Mattila: Geometry of sets and measures in Euclidean spaces. Fractals and rectifiability, Cambridge Studies in Advanced Mathematics, 44, Cambridge University Press, Cambridge, 1995.
[96] D. Mazzoleni, A. Pratelli: Existence of minimizers for spectral problems, J. Math. Pures Appl., 100 (2013), 433-453.
[97] A. Melas: The stability of some eigenvalue estimates, J. Differential Geom., $\mathbf{3 6}$ (1992), 19-33.
[98] R.L. Moore: Concerning triods in the plane and the junction points of plane continua, Proceedings of the National Academy of Sciences of the United States of America, 14 (1928), 85-88.
[99] C. MüLler: Analysis of spherical symmetries in Euclidean spaces, Applied Mathematical Sciences, 129, Springer-Verlag, New York, 1998.
[100] N. Nadirashvili: Conformal maps and isoperimetric inequalities for eigenvalues of the Neumann problem,Proceedings of the Ashkelon Workshop on Complex Function Theory, 11 (1996), 197-201, Israel Math. Conf. Proc., Bar-Ilan Univ., Ramat Gan, 1997.
[101] C. Pommerenke: Boundary behaviour of conformal maps. Grundlehren der Mathematischen Wissenschaften, 299, Springer-Verlag, Berlin 1992.
[102] T. Povel: Confinement of Brownian motion among Poissonian obstacles in $\mathbb{R}^{d}$, $d \geq 3$, Probab. Theory Relat. Fields, 114 (1999), 177-205.
[103] E.B. Saff, V. Totik: Logarithmic potentials with external fields, Grundlehren der Mathematischen Wissenschaften, 316, Springer-Verlag, Berlin, 1997.
[104] J. Serrin, M. Tang: Uniqueness of ground states for quasilinear elliptic equations, Indiana Univ. Math. J., 49 (2000), 897-923.
[105] L. Simon: Lectures on geometric measure theory, Proceedings of the Centre for Mathematical Analysis, Australian National University, 3. Australian National University, Centre for Mathematical Analysis, Canberra, (1983).
[106] J. Serrin, M. Tang: Uniqueness of ground states for quasilinear elliptic equations, Indiana Univ. Math. J., 49 (2000), 897-923.
[107] M. Struwe: Variational Methods. Applications to nonlinear partial differential equations and Hamiltonian systems, fourth edition, Ergebnisse der Mathematik und ihrer Grenzgebiete, A Series of Modern Surveys in Mathematics, 34, SpringerVerlag, Berlin, 2008.
[108] A. S. Sznitman: Fluctuations of principal eigenvalues and random scales: Comm. Math. Phys., 189 (1997), 337-363.
[109] G. Talenti: Best constants in Sobolev inequality, Ann. Mat Pura Appl., 110 (1976), 353-372.
[110] C. Villani: Optimal transport. Old and new, Grundlehren der Mathematischen Wissenschaften, 338, Springer-Verlag, Berlin, 2009.
[111] R. Weinstock: Inequalities for a classical eigenvalue problem, J. Rational Mech. Anal., 3 (1954), 745-753.
[112] Y. Xu: The first nonzero eigenvalue of Neumann problem on Riemannian manifolds, J. Geom. Anal., 5 (1995), 151-165.
[113] G.S. Young: A generalization of Moore's theorem on simple triods, Bull. Amer. Math. Soc., 50 (1944), 714.
[114] B. White: Stratification of minimal surfaces, mean curvature flows, and harmonic maps, J. Reine Angew. Math., 488 (1997), 1-35.


[^0]:    ${ }^{1}$ Here as well as in the rest of the thesis, lim indicates the limit inferior

[^1]:    ${ }^{2}$ A proof of this fact can be easily performed by means of the theory about Functions of Bounded Variation, see for instance [2]. Anyway it is worth having a look to the original De Giorgi proof in [47].

[^2]:    ${ }^{3}$ Actually Weinstock prescribes the perimeter of the set, implying so, in view of the isoperimetric inequality, a stronger version for the Brock-Weinstock inequality in dimension 2.

[^3]:    ${ }^{1}$ Here we follow the classic notation about Riesz potentials. Although they share such a notation, the concept of potential discussed in this chapter is not related to the Schrödinger potentials of Chapter 1.

[^4]:    ${ }^{1}$ The choice of $k$ will be clear in the proof of Lemma 5.19.

[^5]:    ${ }^{2}$ This is the reason why we choose $\widetilde{u}_{\varepsilon}$ to be a $C^{k+1}$ extension of $u_{\varepsilon}$ with $k=[d / 2]+2$.

